

# Air Shower Simulations

---

Johannes Knapp,  
Physics & Astronomy  
U of Leeds, UK

3<sup>rd</sup> School on  
Cosmic Rays and Astrophysics  
Arequipa, Peru  
2008

- Part 1: Astroparticle Physics, Air Showers and Simulations
- Part 2, 3: Hadronic & Nuclear Models, CORSIKA, Performance and Limitations
- Part 4: Selected Aspects of EAS Simulations, Simulation Techniques

# Air Shower Simulations

---

Johannes Knapp,  
Physics & Astronomy  
U of Leeds, UK

3<sup>rd</sup> School on  
Cosmic Rays and Astrophysics  
Arequipa, Peru  
2008

- Part 1: Astroparticle Physics, Air Showers and Simulations
- Part 2, 3: Hadronic & Nuclear Models, CORSIKA, Performance and Limitations
- Part 4: Selected Aspects of EAS Simulations, Simulation Techniques

# The Atmosphere

All other components  
(e.g.  $\text{CO}_2$ ,  $\text{H}_2\text{O}$ , ...) are  $\ll 1\%$   
and can be neglected.

composition: 78.0%  $\text{N}_2$   
21.5%  $\text{O}_2$   
0.5% Ar

barometric formula:  $\rho(h) = \rho(0) \exp(-h/h_0)$

$\rho(0) = 0.00123 \text{ g/cm}^3$  at sea level  
 $h_0 \sim 7 \text{ km}$

Total (vertical) thickness  $T$  of atmosphere:  $\sim 1000 \text{ g/cm}^2$

Pressure at sea level: 1 bar    1 mbar = 1 hPa

$$P \sim T \cdot g = 1000 \text{ g/cm}^2 \cdot 9.81 \text{ m/s}^2 \sim 10^5 \text{ N/m}^2 = 10^5 \text{ Pa}$$

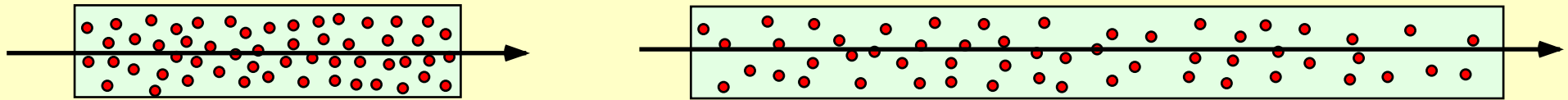
Chacaltaya (5200 m a.s.l.):

$$r = 0.59 r_0$$

$$T = 538 \text{ g/cm}^2$$

(temperature complicates things....  
"adiabatic atmosphere")

The natural scale for the shower development is  
the atmospheric thickness  $t$  (in  $\text{g}/\text{cm}^2$ ), not the height (in  $\text{cm}$ ).



(the number of scattering centers matters, not the distance travelled.)

Characteristic lengths:  $x_0$  radiation length,  $\lambda_0$  interaction length

Atmospheric Thickness (=mass overburden)

$$\tau(h) = \int_h^{\infty} \rho(h) dh = \rho(0) h_0 \exp(-h/h_0)$$

Shower development:  $N(t) \sim c \cdot t^a \cdot \exp(-bt)$

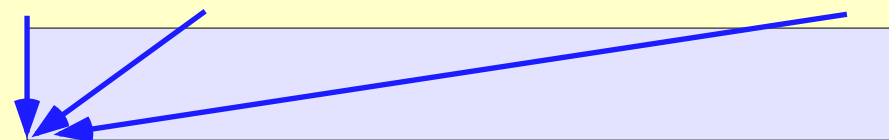
$T$  increases with zenith angle  $\theta$

$$T(\theta) \sim T(0) / \cos \theta \quad \text{for } \theta < 70^\circ$$

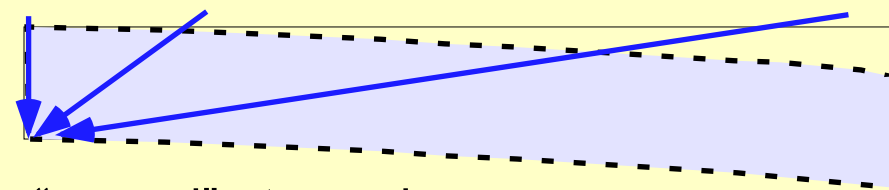
$$T(\theta) \sim \infty \quad \text{for } \theta = 90^\circ$$

$$T(\theta) < T(0) / \cos \theta \quad \text{for } \theta < 70^\circ$$

$$T(\theta) \sim 36 T(0) \quad \text{for } \theta = 90^\circ$$



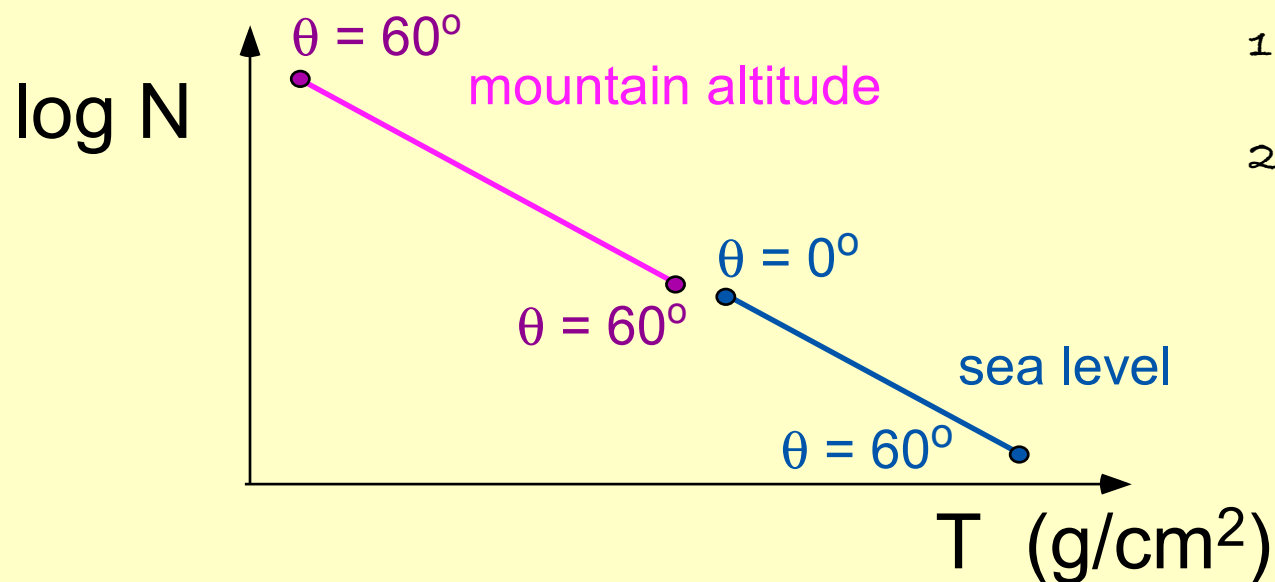
“flat” atmosphere



“curved” atmosphere

A vertical  $10^{20}$  eV shower has its maximum near sea level ( $T_{\max} \sim 1000 \text{ g/cm}^2$ ).

A  $10^{20}$  eV shower at  $>70^\circ$  ( $T > 3000 \text{ g/cm}^2$ ) is almost completely absorbed in the atmosphere (just muons survive).



1. Compare experiments at different altitudes.
2. Study the longitudinal shower development by varying  $q$  and look for signatures with the same frequency (constant intensity method)

# The Atmosphere in CORSIKA

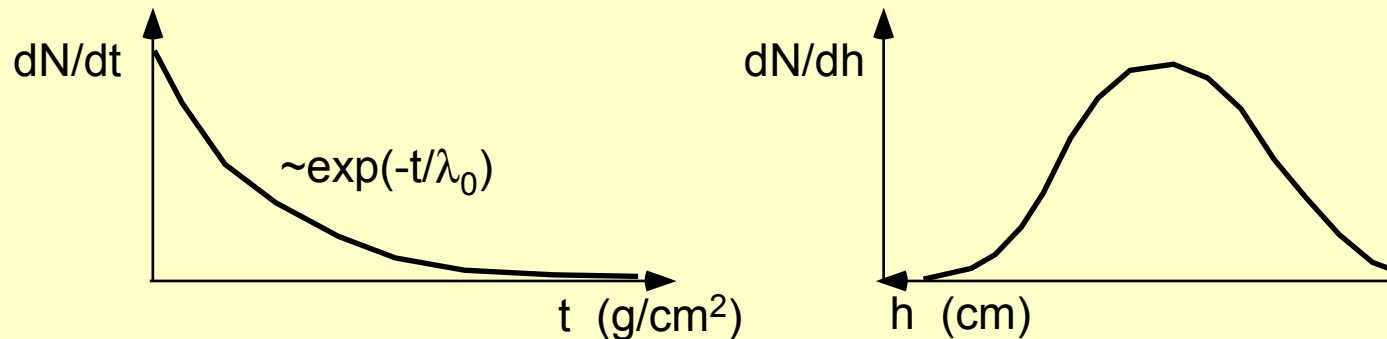
US Standard Atmosphere (parameterized by J. Linsley in 5 layers)

> 100 km	linear		
40 ... 100 km	} exponential density variation	78.0%	N <sub>2</sub>
10 ... 40 km		21.5%	O <sub>2</sub>
4 ... 10 km		0.5%	Ar
0 ... 4 km			

$\rho(h)$   
 $T(h)$   
 $h(T)$ 
} implemented as functions

$\rho(0) = 0.00123 \text{ g/cm}^3$   
 $T(0) = 1036.1 \text{ g/cm}^2$   
 $h(T=0) = 112.8 \text{ km}$

Position of next interaction follows exponential in  $t$   
 but a small  $dt$  may correspond to huge  $dh$  since  $T \sim \exp(h/h_0)$



# Decay or Interaction ?

Position of **next interaction** of a particle depends on the matter traversed (i.e.  $T$ ).

$$dN/dx \sim \exp(-x/\lambda_0)$$

Position of **decay** of a particle depends on time past (i.e.  $t$ ) or distance travelled ( $s$ ).

$$dN/dt \sim \exp(-t/T) \quad \text{and since } t = s/c = x/\rho c \quad dN/dx \sim \exp(-x/\rho c T)$$

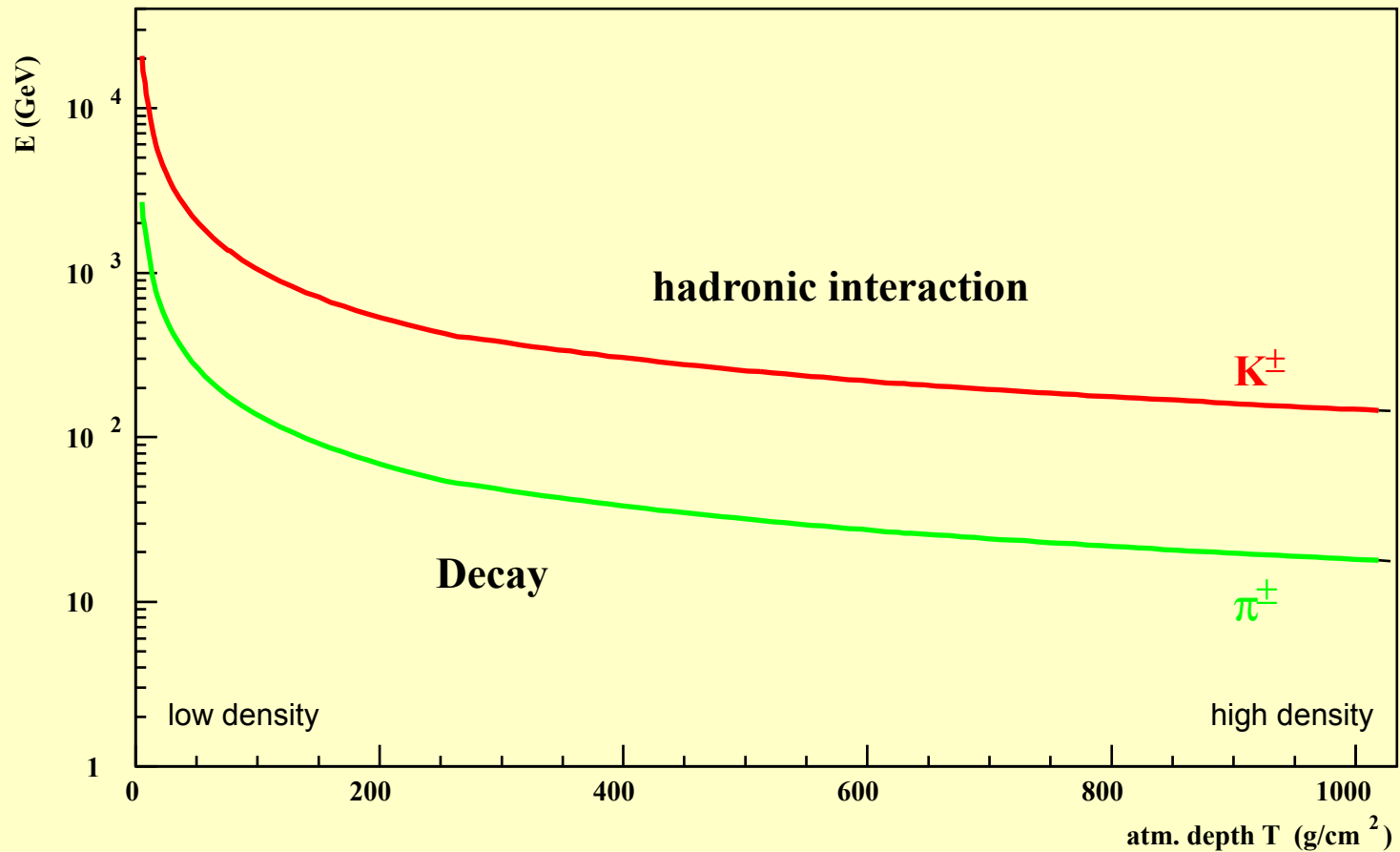
What if interaction and decay compete, e.g. for  $\pi^\pm$  ?

in MC very simple: Calculate independently a decay point and an interaction point.  
Select the mechanism that happens first.

- 1.) draw a random path  $x_i$  from the distribution  $\exp(-x/\lambda_0)$  to get interaction point
- 2.) draw a random time  $t_d$  from the distribution  $\exp(-t/T)$  to get decay point
- 3.) convert time  $t_d$  into a distance  $x_d = t_d c \rho$ .
- 4.) If  $x_i < x_d$ , then an interaction happens, otherwise a decay.

This is actually how nature does it.

# Decay or Interaction ?



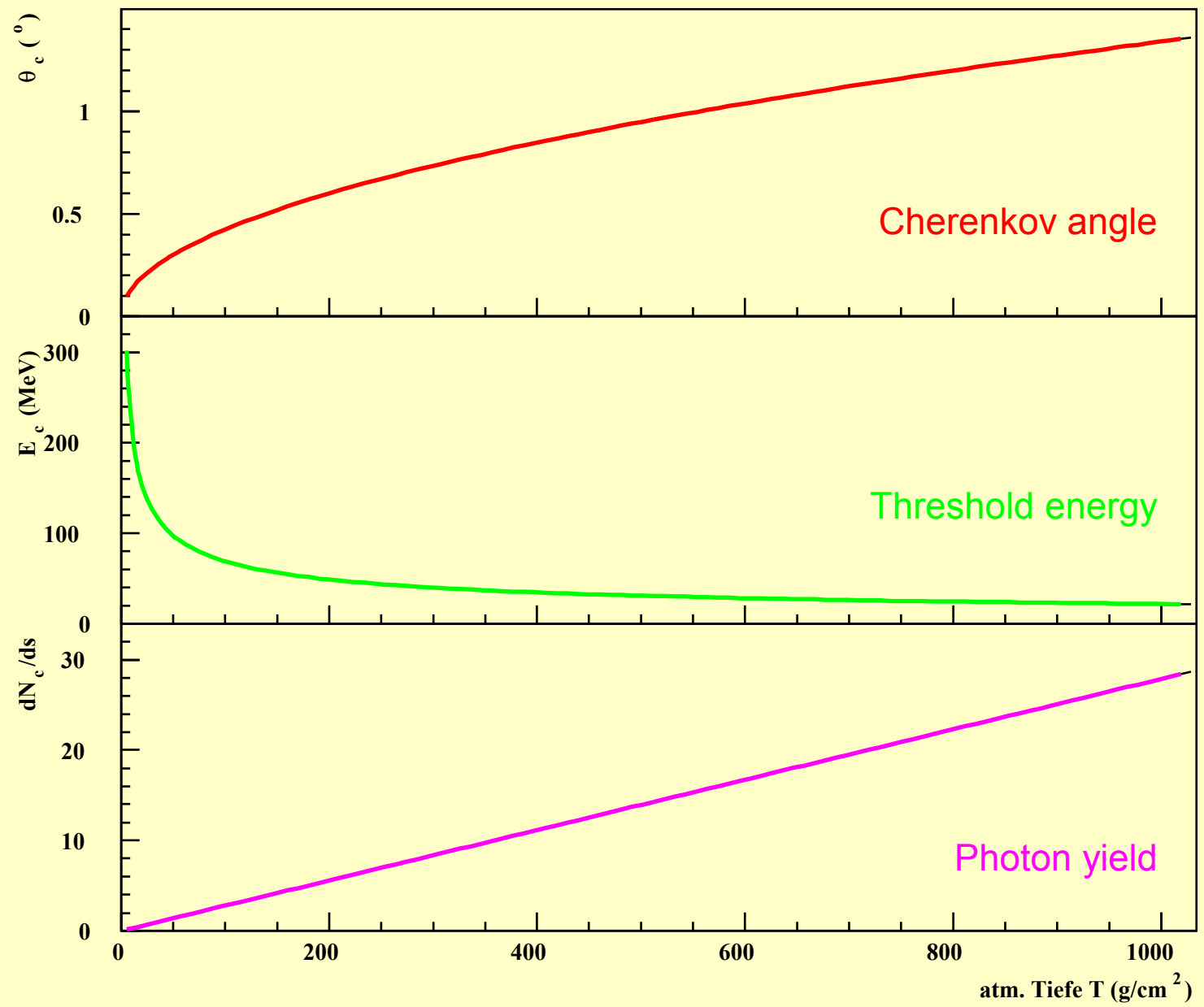
decay:  $P_d(t) \sim \exp(t/\tau)$        $\tau = \text{lifetime}$

interaction:  $P_i(x) \sim \exp(x/\lambda)$        $\lambda = \text{interaction length}$

analytic:  $P_{\text{tot}}^{-1} = P_d^{-1} + P_i^{-1}$       difficult, since exponentially varying atmospheric density



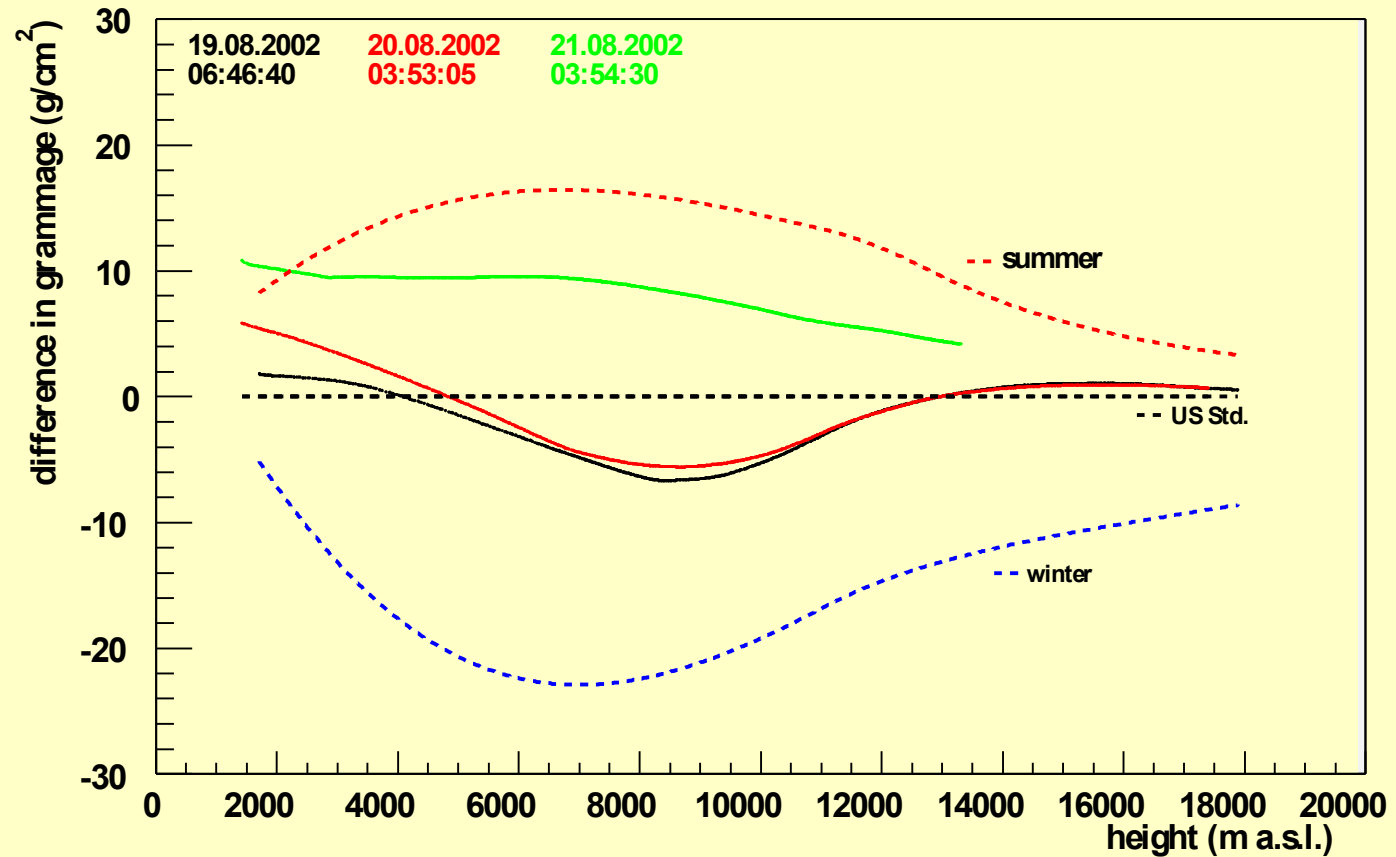
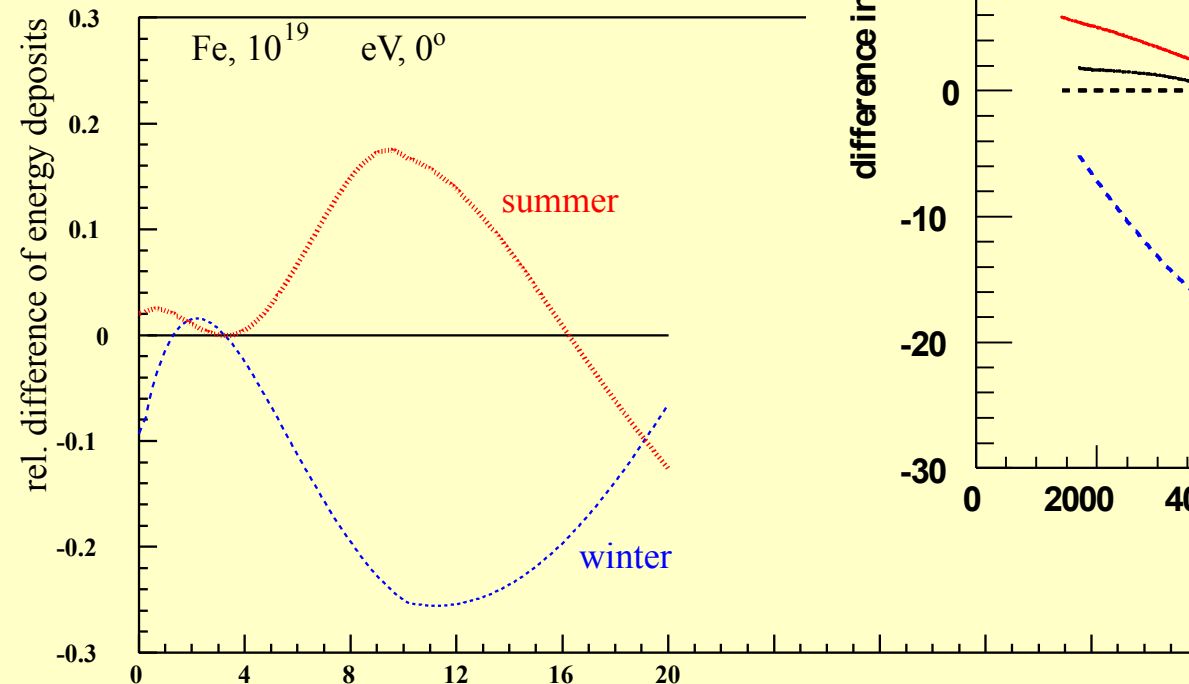
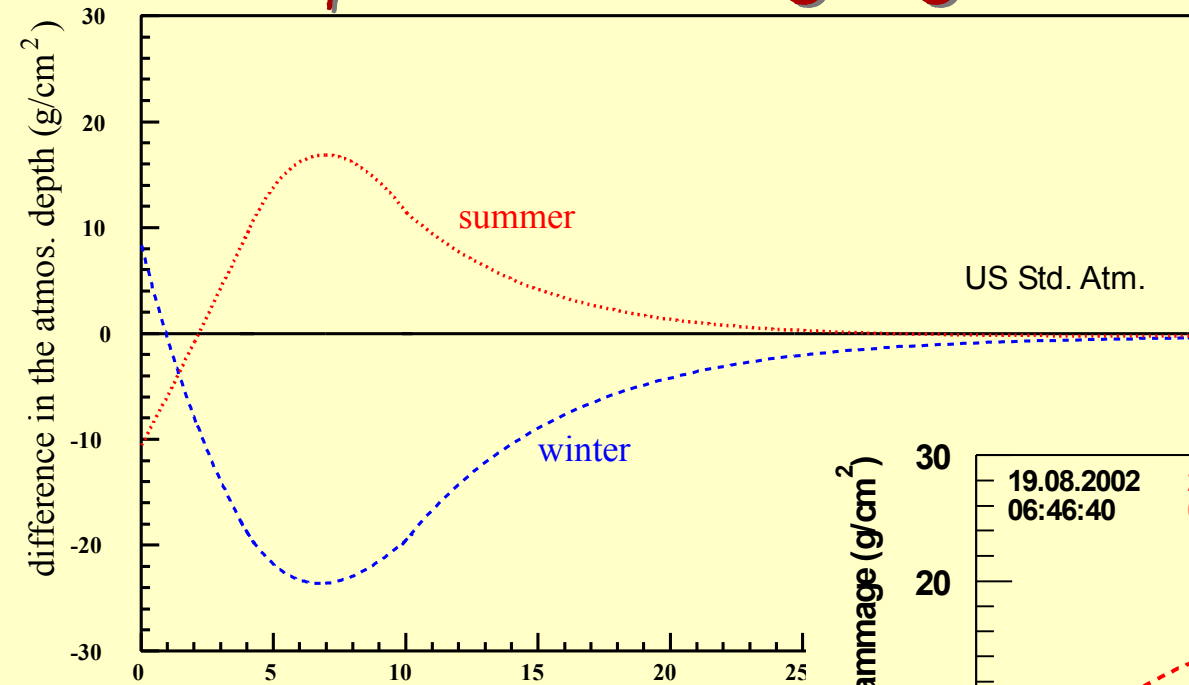
# Cherenkov Light Production in Air



needs to be folded  
with the  
particle distribution

# Atmosphere is changing all the time ....

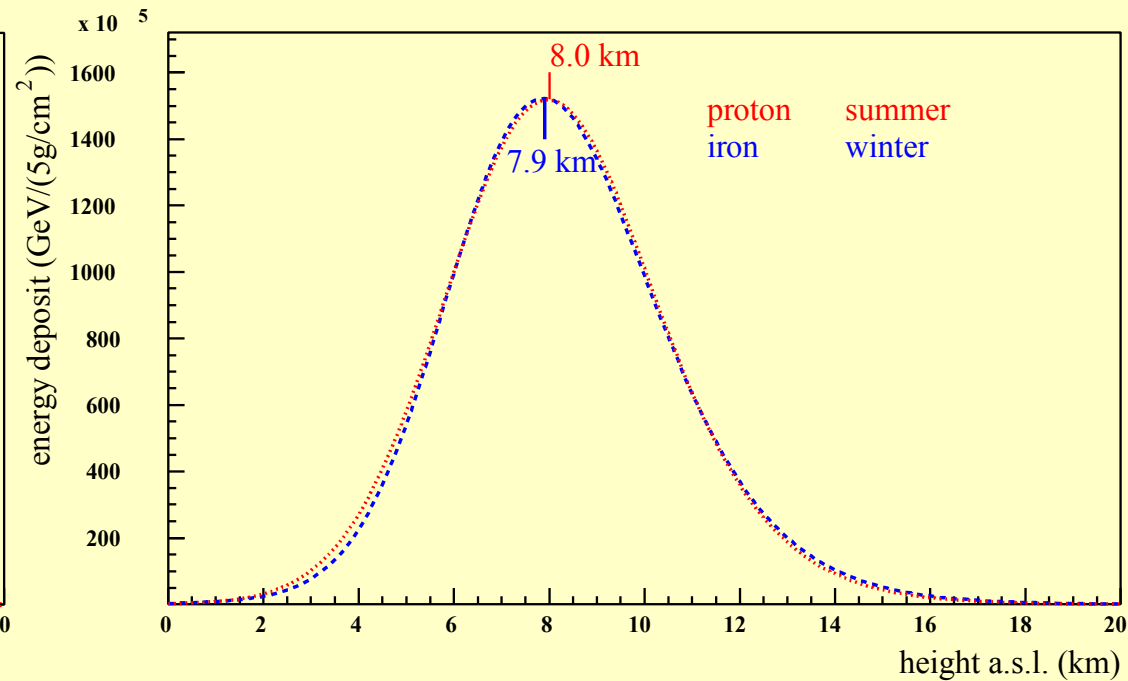
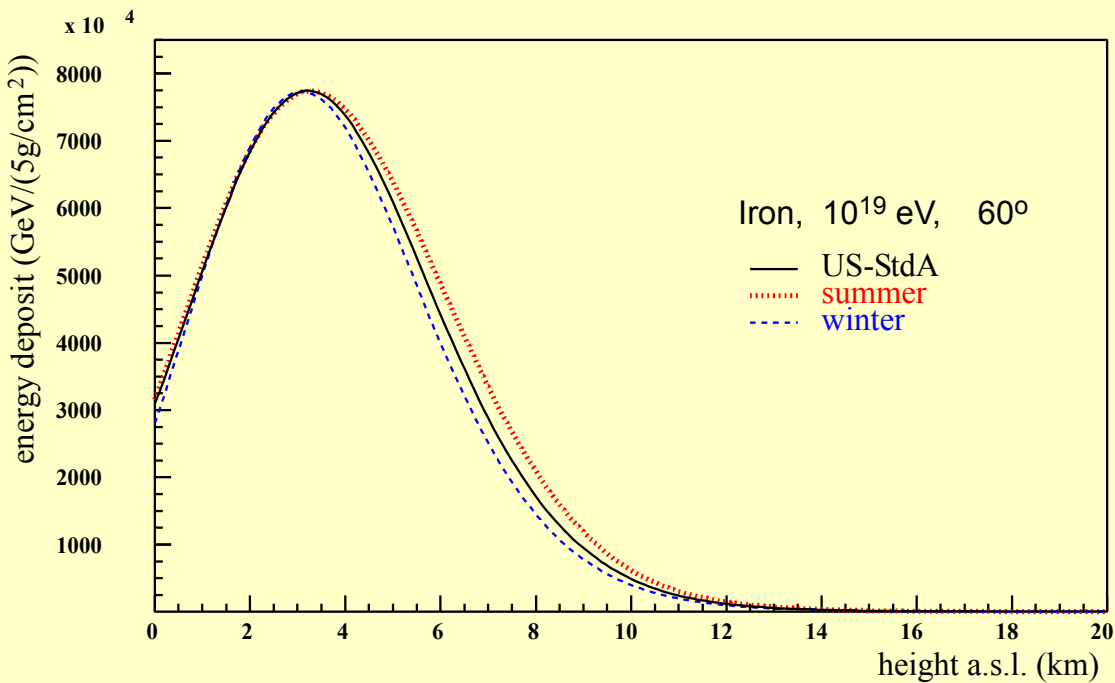
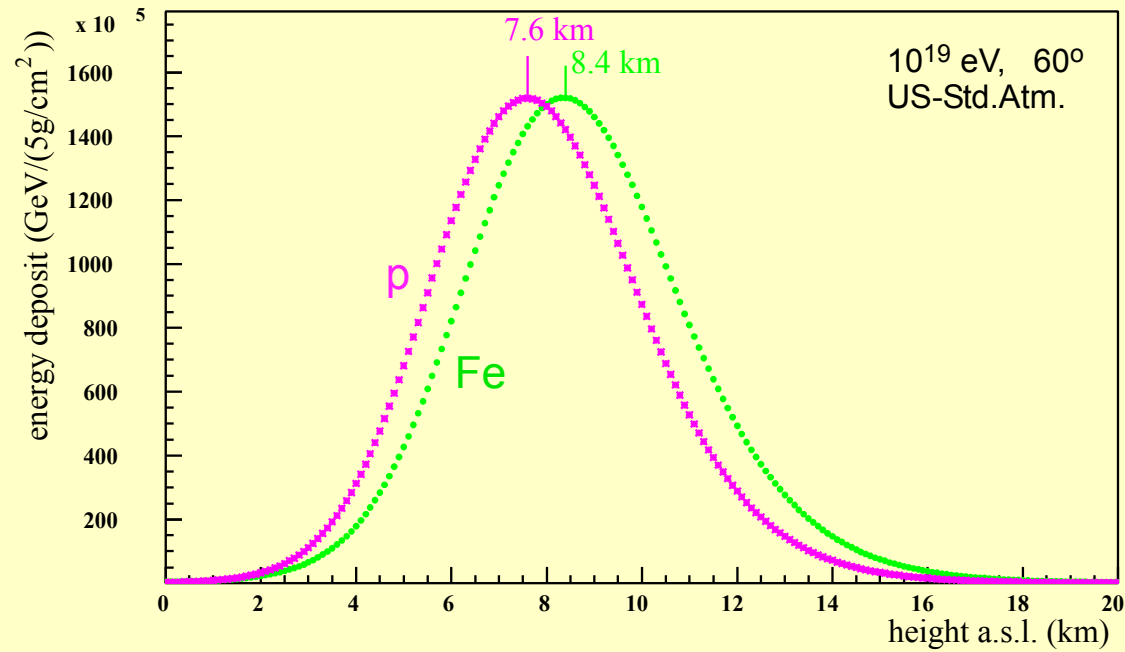
Total thickness of absorber varies as well as the layering.  
 variation of atmospheric conditions:  
 (to be measured by balloon sondes at the site,  
 e.g. after each  $> 10^{20}$  ev event?)



height a.s.l. (km)

B. Keilhauer

# $X_{max}$ and the atmosphere



# Simulation Speed-up

Computing time  $\sim 1 \text{ h} \times E/10^{15} \text{ eV}$

Disc space  $\sim 300 \text{ MB} \times E/10^{15} \text{ eV}$  per shower.

At  $10^{20} \text{ eV}$ : more than  $10^{11}$  secondaries !!!

excessive resources needed per shower:  $\sim 10^5 \text{ h} = 11 \text{ years}$   $\sim 30 \text{ Tera Bytes}$

No way (nor need?) to follow them all: statistical sampling "Thinning"

Standard method:

(by A.M. Hillas)

define the thinning threshold  $E_{\text{th}} = \epsilon_{\text{th}} E_0$

$E > E_{\text{th}}$  follow each particle

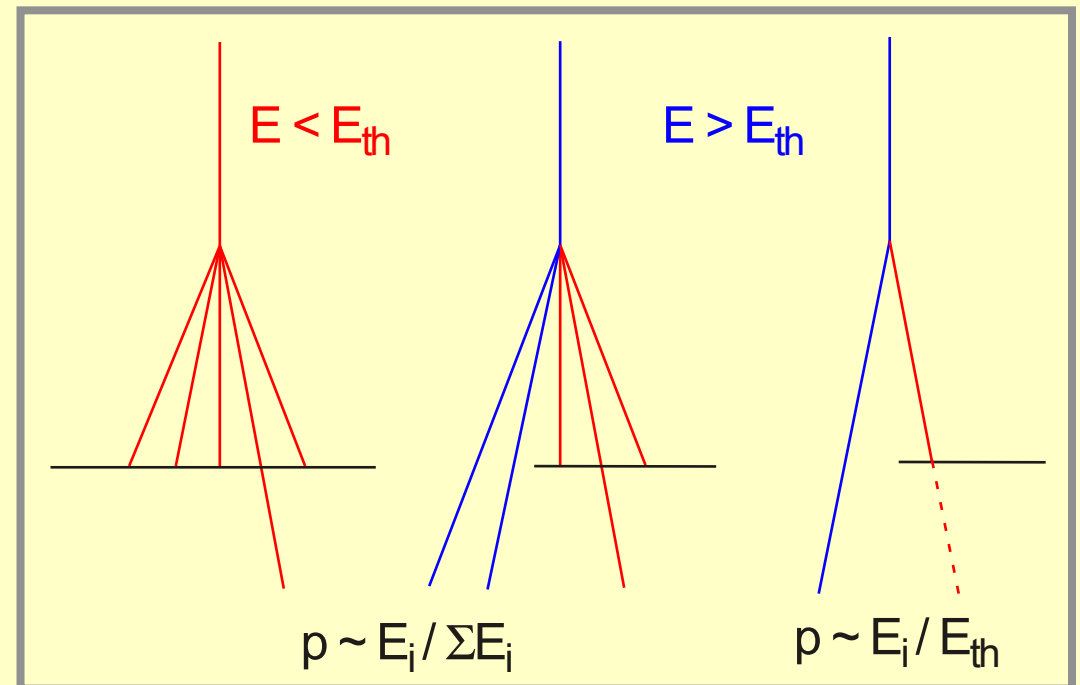
$E < E_{\text{th}}$  follow only one (or few),

but give it a weight to account for  
the discarded particles:  $w_i' \sim w_i 1/p$

+ energy is conserved

+ mean  $N_{e,\gamma,\mu}$  are preserved

- but: fluctuations are enlarged artificially



Added weight increases output per particle (8 words instead of 7),

but computing time and number of particles to be written out are greatly reduced.

# Which is the right thinning level?

longitudinal development  $N(t)$ :

very many particles in the shower core, i.e.

low thinning level is sufficient ( $E_{th} \sim 10^{-4} E_0$ )

The smaller  $E_{th}$ ,

the better the shower is modelled

§ the larger the computing time.

particles far from shower core, e.g.  $S(r)$  in Auger:

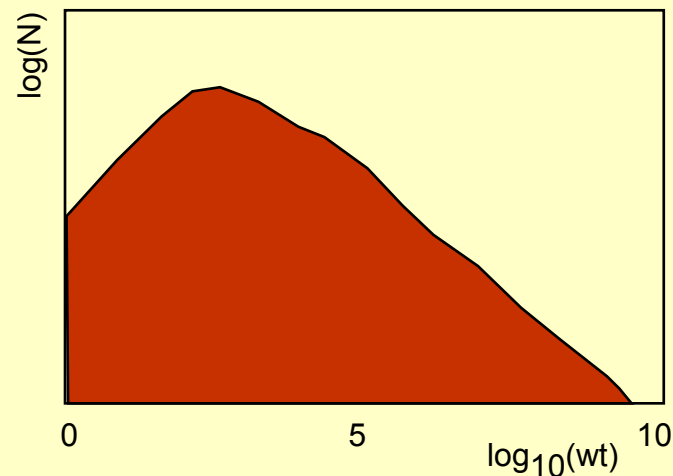
small particle density requires good thinning ( $E_{th} < 10^{-7} E_0$ )

Artificial fluctuations due to thinning should be smaller than the intrinsic shower fluctuations  
(to be checked for the variable in question...)

Computing time and disk space are reduced, but do still grow proportional to  $E_0$ .

Particle weights can go up to  $w_{max} = E_{th} / E_{min} \sim 10^{20} \times 10^{-6} / 10^5 = 10^9$

↑ start    ↑ end of thinning ( $E_{min} =$  low energy cut-off)



High weights are problematic:

Is there a way to avoid them?

# Weight Limitation and Optimum Thinning

Avoid too high weights by setting a maximum allowed weight  $w_{max}$  (e.g.  $10^5$ ).  
If weight gets larger, follow all particles again.

Of course, weight limitation increases run time again.

Which is the best setting of  $\epsilon_{th}$  and  $w_{max}$ ?

i.e. which minimizes the statistical error for a given run time?

Optimum thinning for a given  $\epsilon_{th}$

$$w_{max} = E_0 \text{ (in GeV)} \cdot \epsilon_{th}$$

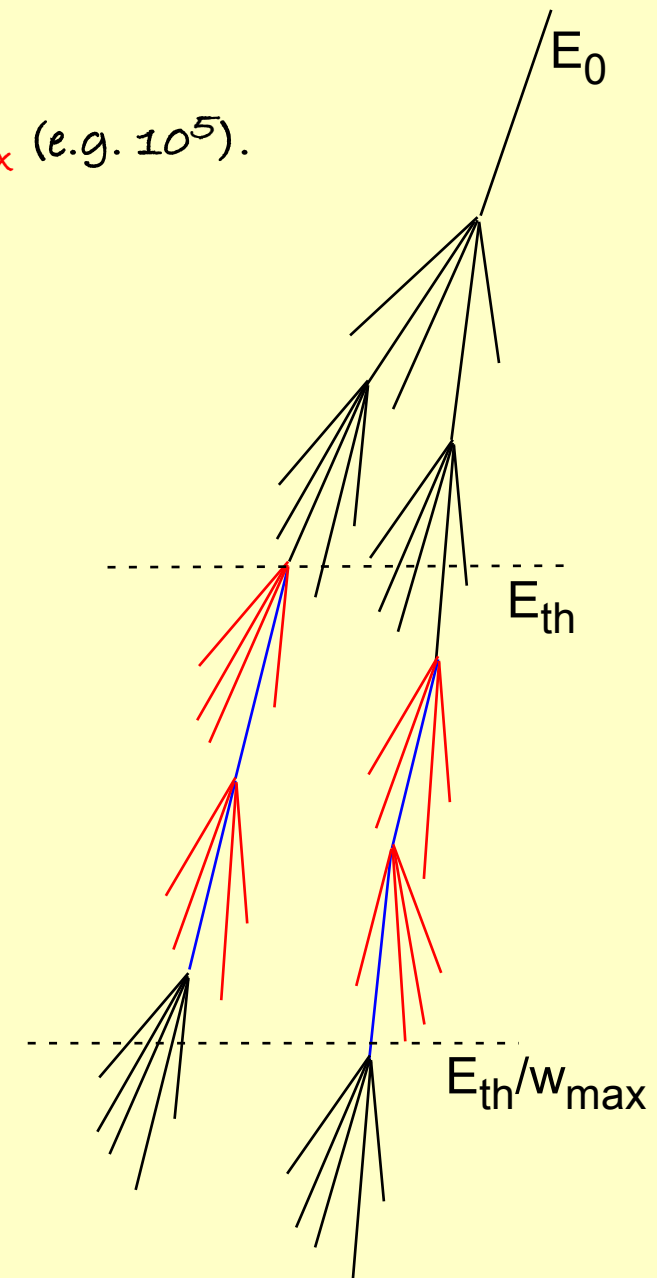
$$10^{18} \text{ eV} \quad \epsilon_{th} = 10^{-6} \quad w_{max} = 10^3$$

$$10^{19} \text{ eV} \quad \epsilon_{th} = 10^{-6} \quad w_{max} = 10^4$$

$$10^{20} \text{ eV} \quad \epsilon_{th} = 10^{-6} \quad w_{max} = 10^5$$

The run time is **only dependent on  $\epsilon_{th}$** , no longer on energy.

$10^{-5}$  optimum thinning is about as good as  $10^{-7}$  thinning without weight limitation.

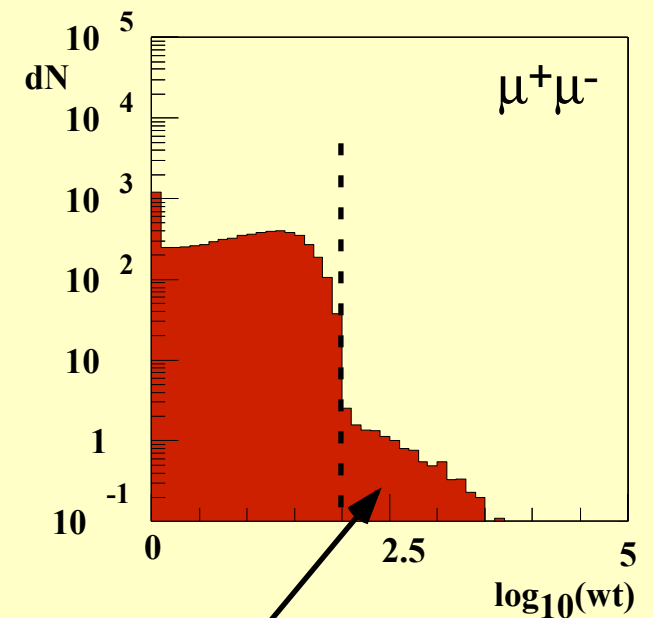
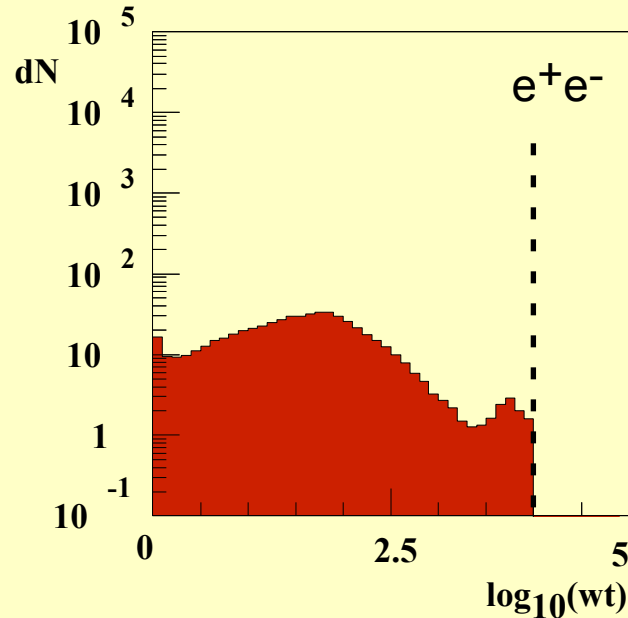
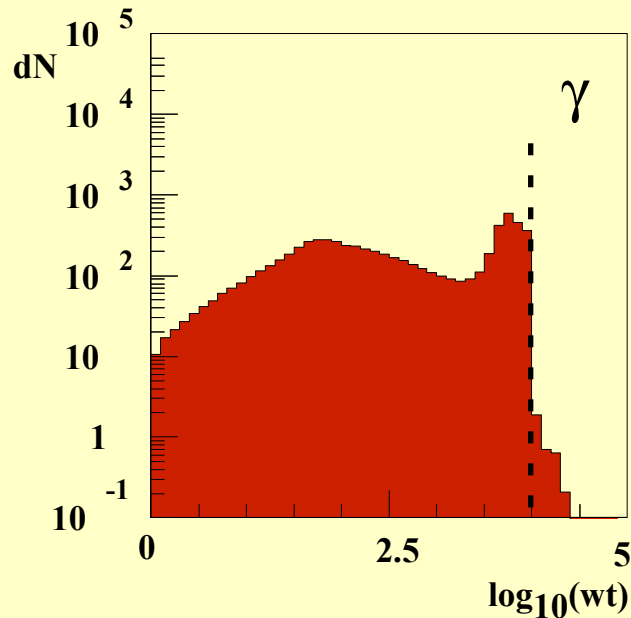


Since  $g$  and  $e^+e^-$  are much more abundant ( $\times 100$ ) than  $\mu^+\mu^-$  or hadrons we set different  $\mathcal{E}_{th}$  and  $w_{max}$  for  $\gamma$  &  $e^+e^-$  than for  $\mu^+\mu^-$  or hadrons.

e.g.: proton  $10^{19}$  eV  $37^\circ$

$$\gamma, e^+e^-: \quad \mathcal{E}_{th} = 10^{-6} \quad w_{max} = 10^4$$

$$\mu^+\mu^-, h: \quad \mathcal{E}_{th} = 10^{-8} \quad w_{max} = 10^2$$



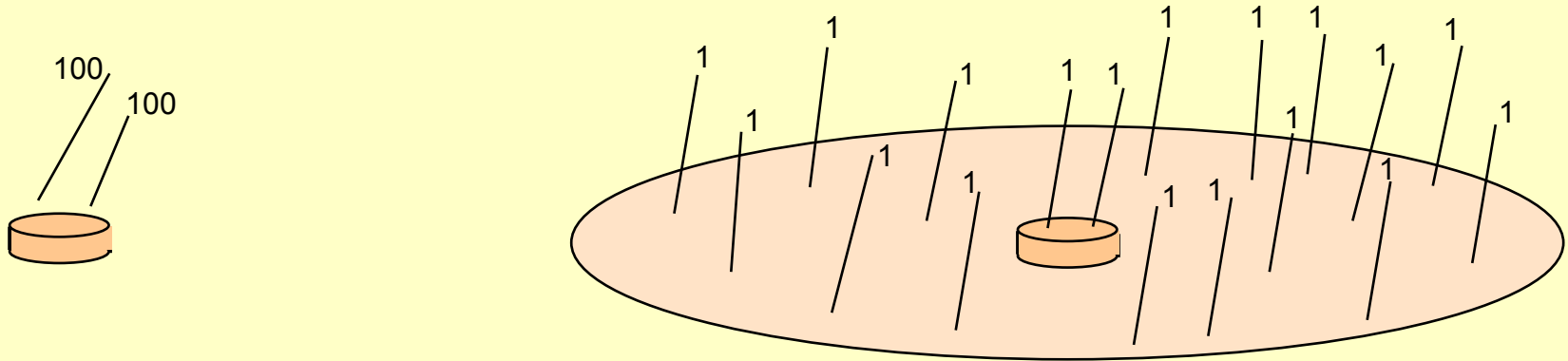
1 particles with weight  $10^5$  stands for  
 10 particles with weight  $10^4$  or  
 $10^5$  particles with weight 1

muons produced  
 in photo production  
 from photons with weight  $> 10^2$

## How to treat a particle with a high weight once it hits a detector ???

A particle with weight stands for many others with different masses, energies, angles, impact points.

Some sort of un-thinning is needed (e.g. method used in Auger)



Instead of  $N$  particles with high weights  $wt$  on a small area  $A$ , use  $N \cdot wt$  particles, each with weight 1, from a larger area  $A \cdot wt$ .

Problem: For  $wt > 10^5$  the area becomes so large ( $\sim$  distance between detectors) that particle densities, energies, ... will change from one side to the other.

There is no ideal method to get rid of the weights again.  
Some of the information has been lost in the thinning process.



## Other speed-up tricks?

All tricks come at a cost !

Speed-up mechanisms cause biases and distortions.

Careful checks are needed that they don't invalidate the results.

Example: early KASCADE simulations:

Low-energy electromagnetic subshowers created high in the atmosphere were discarded (x3 speed-up, with bias of less than 3%)

When higher energies were simulated,

this speed-up caused a large bias which led to wrong results.

Massive simulations had to be repeated after the problem was found.

## Use parallel processing?

In principle possible: do first interaction & distribute secondaries on different processors

But we usually need more than one shower to evaluate fluctuations.

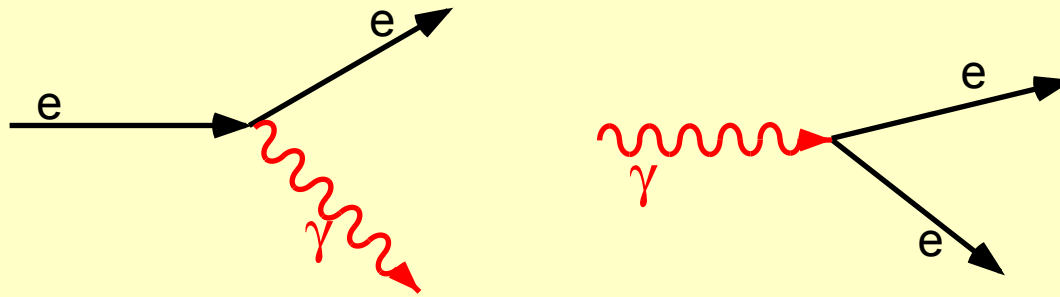
Best way to parallelize: run on each processor one shower  
to avoid any inter-processor communication.

e.g. processor farm in Lyon (~1000 processors + lots of storage space)

70,000 showers for Auger, ~150 processor years, 20 TB

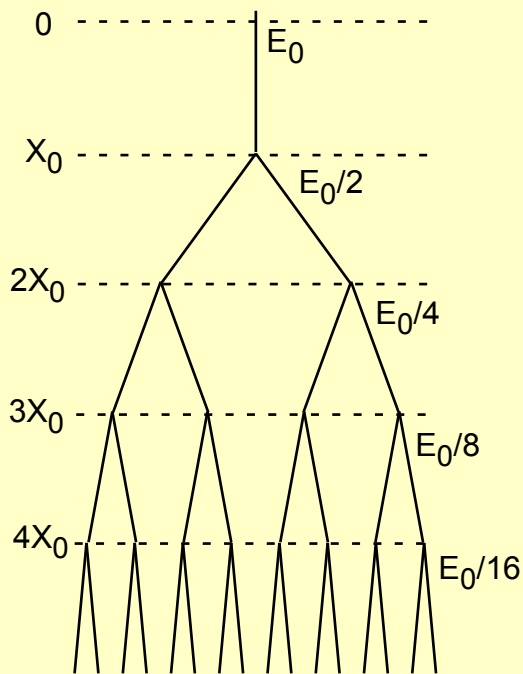
# Electromagnetic Showers: from Toy Model to EGS4

basic reactions:    photons:    pair production  
                           electrons:    bremsstrahlung



Both reactions have the same scale length ( $X_0$ ) and have two outgoing particles per incoming particle.

**Toy Model** (one-dimensional, very simplified, yet qualitatively correct):



particle multiplication ( $\times 2$ ) in each step ( $X_0$ ) until  $E < E_{crit}$ ,  
 then particle losses due to ionisation dominant.

$$t = k X_0, \quad k = 1, 2, \dots$$

$$N = 2^k \quad E = E_0 / N$$

$$k_{max}: \quad E_0 / 2^{k_{max}} = E_{crit} \quad k_{max} = \ln(E_0 / E_{crit}) / \ln(2)$$

grows only logarithmically with  $E_0$

$$t_{max} = k_{max} \cdot X_0$$

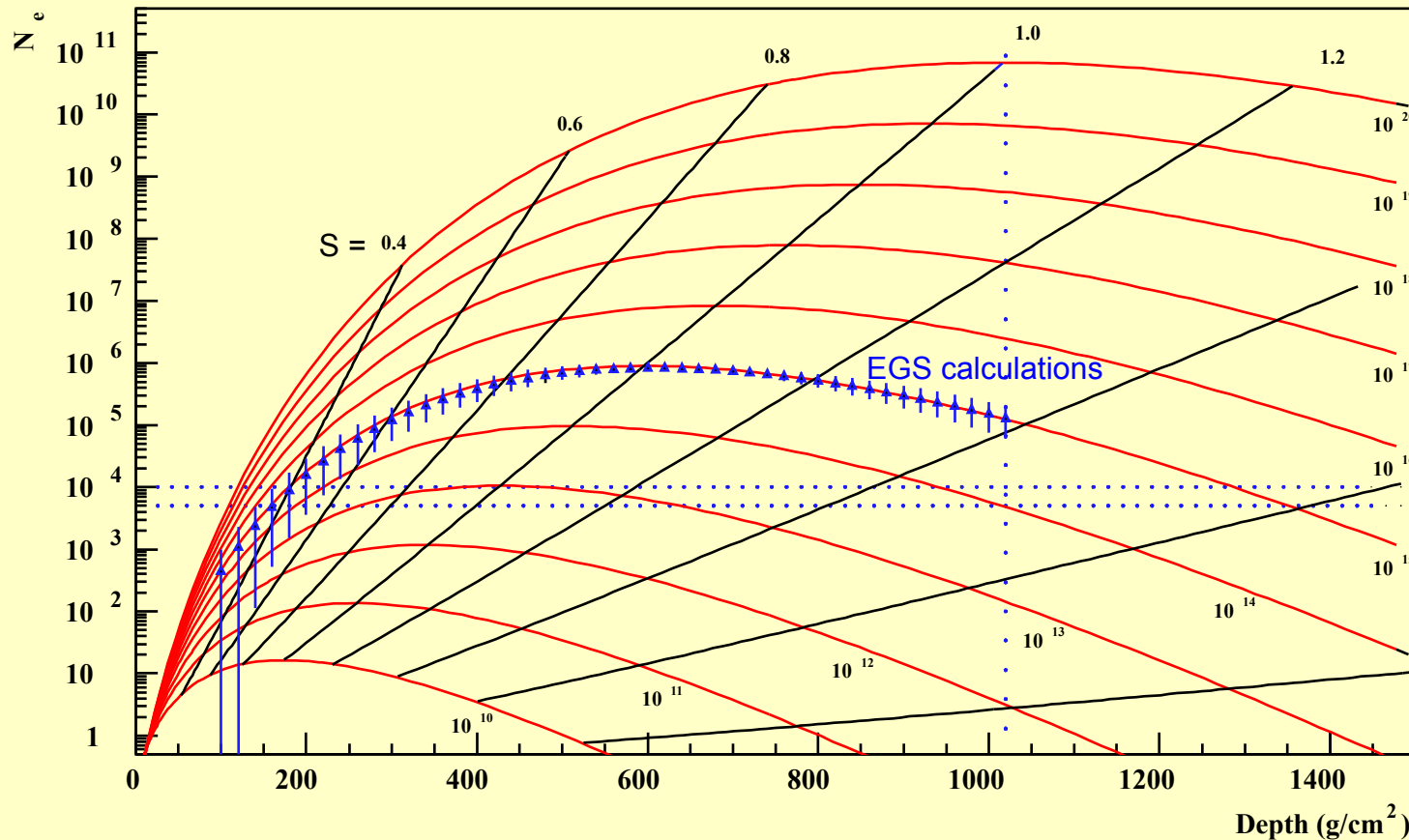
$$N_{max} = E_0 / E_{crit}$$

Measure  $t_{max}$  or  $N_{max}$   
 and estimate  $E_0$ .

# Nishimura Kamata Greisen (NKG): Longitudinal Shower Development

analytic description of purely electromagnetic showers:

$$N_e = \frac{0.31 \exp(t(1-1.5 \ln s))}{\sqrt{\ln(E_0/E_{\text{crit}})}} \quad s = \frac{3t}{t + 2 \ln(E_0/E_{\text{crit}})}$$



$N_e$ :  
number of electrons  
down to energy 0 ?  
(unphysical)

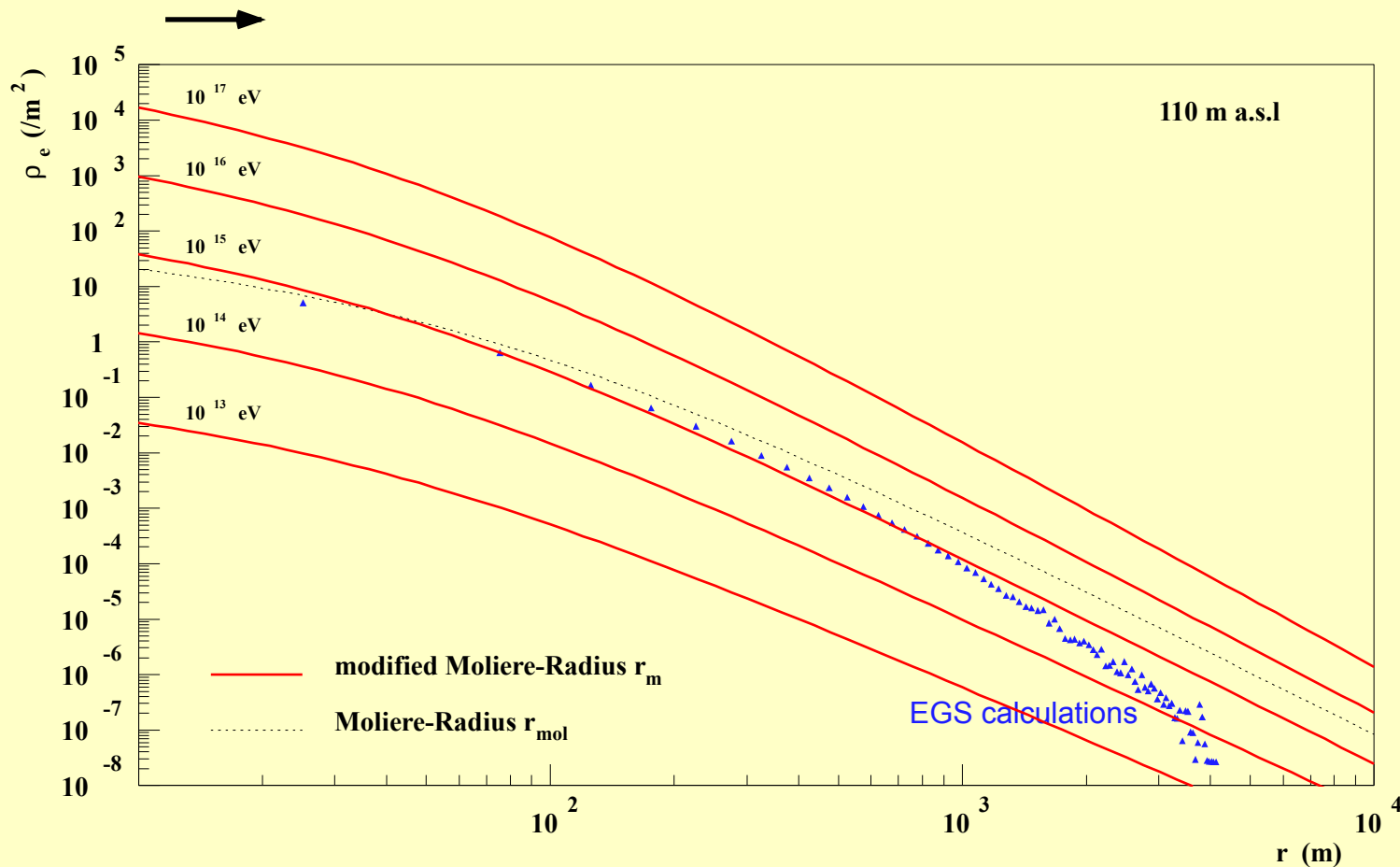
Just average,  
no fluctuations.

# Nishimura Kamata Greisen: lateral shower development

purely electromagnetic showers:

$$\rho_e = \frac{Ne}{2\pi r_m^2} \frac{\Gamma(4.5-s)}{\Gamma(s) \Gamma(4.5-2s)} \left(\frac{r}{r_m}\right)^{s-2} \left(1 + \frac{r}{r_m}\right)^{s-4.5}$$

$$r_m = (0.78 - 0.21 s) r_{mol}$$



$$r_{mol} = X_0 E_s / E_{crit}$$

$$\sim 9.6 \text{ g/cm}^2$$

$$\sim 78 \text{ m at sea level}$$

$$E_s = m_e c^2 (4p/a)^{1/2} \sim 21 \text{ MeV}$$

A cylinder around the shower axis with radius  $r_{mol}$  contains 90% of the shower energy.

NKG formalism allows a **fast**, semi-analytical simulation of electromagnetic sub-showers.

# Limitations of NKG:

$N_e$ : number of electrons down to energy 0 is unphysical.

When realistic detector thresholds are used up to 30% difference.

Also purely electromagnetic showers contain some muonic / hadronic component.

$\gamma \longrightarrow p$        $\gamma \longrightarrow \mu^+ \mu^-$  (suppressed, but nevertheless there)

no fluctuations

no info on gammas

no info on particle times and directions

# Full 4-dim simulation with EGS

Electron Gamma Shower Code      Nelson et al. ~1970

ALL processes of electrons and gammas are included.

bremsstrahlung, ionisation, d-electrons, Bhabha & Moeller scattering,  
multiple scattering, annihilation, ...

$e^+e^-$  pair production, Compton effect, photo effect, Rayleigh scattering, ...

based on QED calculations and is very well checked and verified.

extended by LPM effect ( $> \text{TeV}$  in dense materials;  $> 10^{18} \text{ eV}$  in atmosphere)

EGS gives precise predictions of all sorts of electromagnetic interactions in materials.

negative: 40 x slower than NKG version,  
but the quality of the results is worth the while.

# Random Numbers

... are of central importance in Monte Carlo methods.

Use random physical processes to create real random numbers?  
(e.g. radio active decay, electronic noise, ...)

In principle yes, but ....

MC programs must be reproducible (e.g. for bug fixing).

**Pseudo-Random** numbers:

they are produced by a predictable algorithm  
but behave in all respects like random numbers.

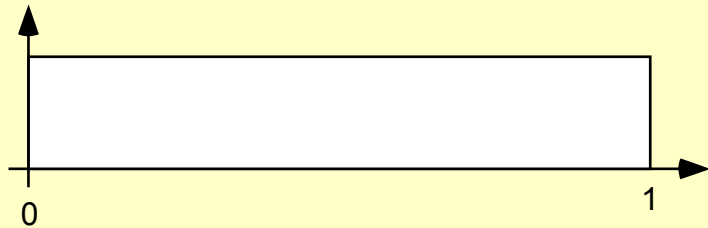
(all digits, all combinations of digits appear with equal probability,  
there are no correlations within the sequence)

Computers are deterministic machines,

i.e. computer generated random number sequences are not really random  
and computer sequences have a finite sequence length (period)

It is an art to produce good random numbers !

# Uniform Random Number Generator



uniform probability in range 0, ... 1.

Simple example: linear congruent generator

$$l_j = a (l_{j-1} + c) \bmod m$$

$$u_j = l_j / m$$

3 parameters:  $a, c, m$

$l_0$ : seed

$l_1, l_2, \dots$  are integers between 0 and  $m-1$

$u_j$  is a real number between 0 and 1

The maximum period is  $m$ , but real period depends on  $a$  and  $c$ . typical:  $10^5 \dots 10^{11}$   
not enough for serious applications.

Random numbers from this generator are not uncorrelated:

$k$ -tuples of random numbers lie in  $k$ -dim space on  $(k-1)$ -dim hyper planes

Less significant bits are usually less random



better: Fibonacci generator

$$u_n = (u_{n-24} + u_{n-55}) \bmod 1 \quad n > 55 \quad \text{period : } > 2^{55} \sim 10^{18}$$

to be initialised with  $u_1 \dots u_{55}$  still not sufficient.

modern generators combine simpler methods:

- 1) Combine two random numbers (from different generators) with "+", "-", or "exclusive OR".
- 2) A sequence of random numbers from generator 1 is stored in a memory. A random number of generator 2 is used as address of the next random number in the memory.

used in CORSIKA: RANMAR (CERNLIB)

32-bit floating point numbers between 0 and 1.

900.000.000 different sequences of  $\sim 2^{144} = 10^{43}$  period length

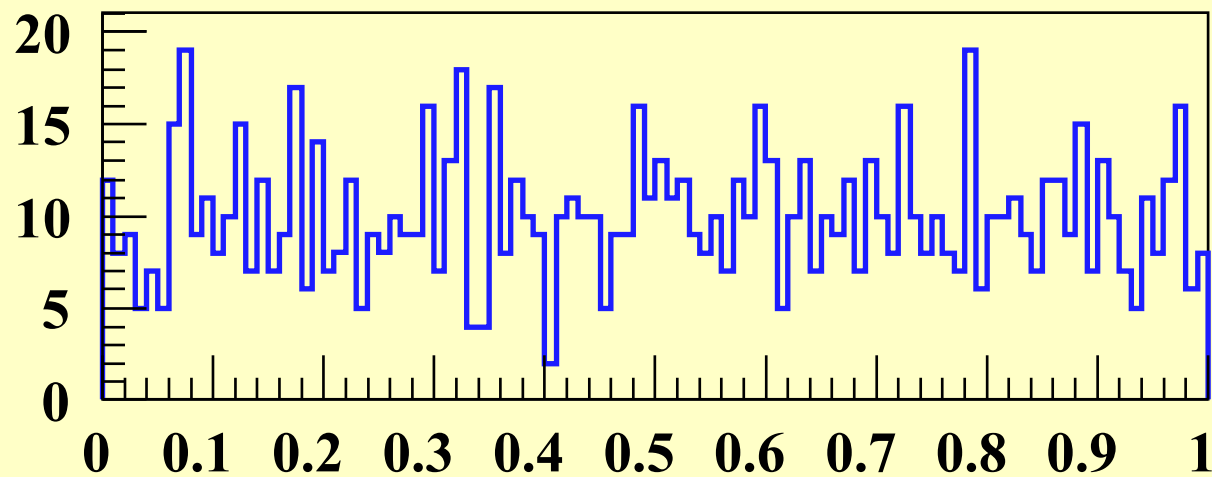
There are even better ones:

... but the better the random number generator, the slower it is.

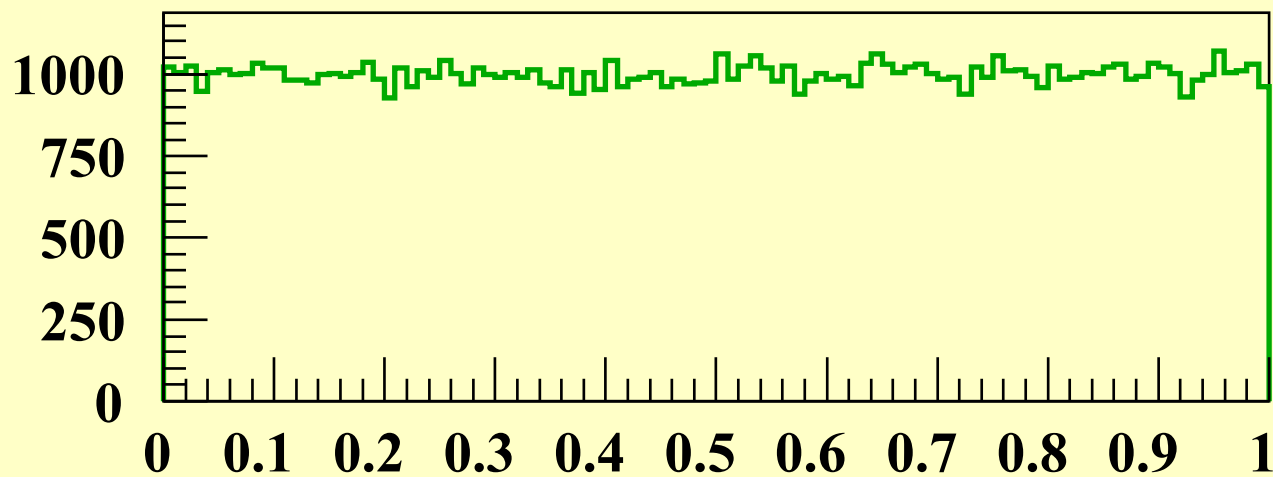
About 30% of the computing time of CORSIKA goes into calculation of random numbers !

uniform

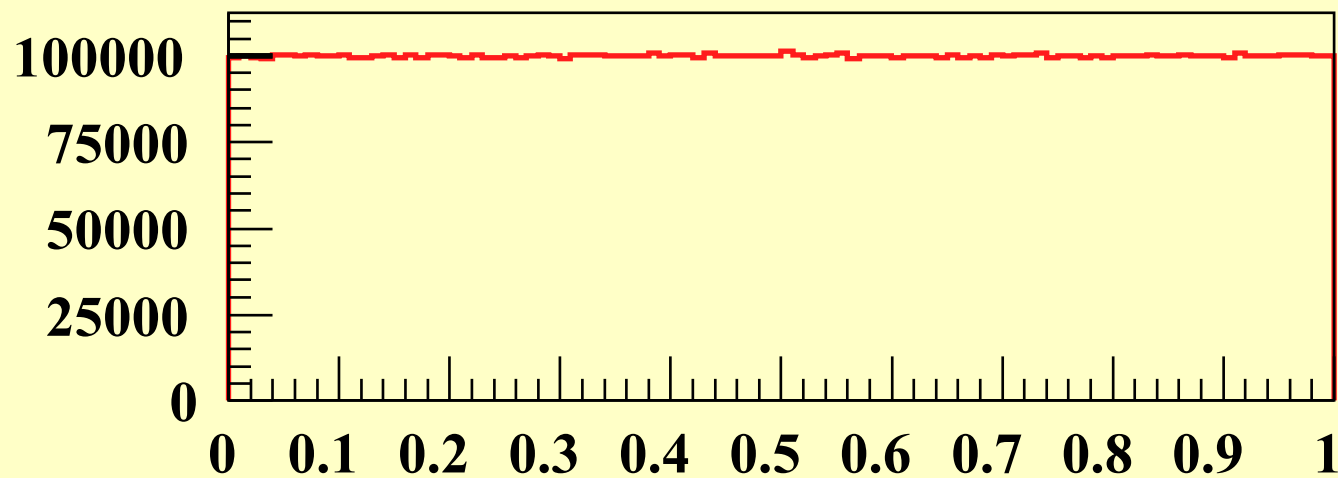
$10^3$  events

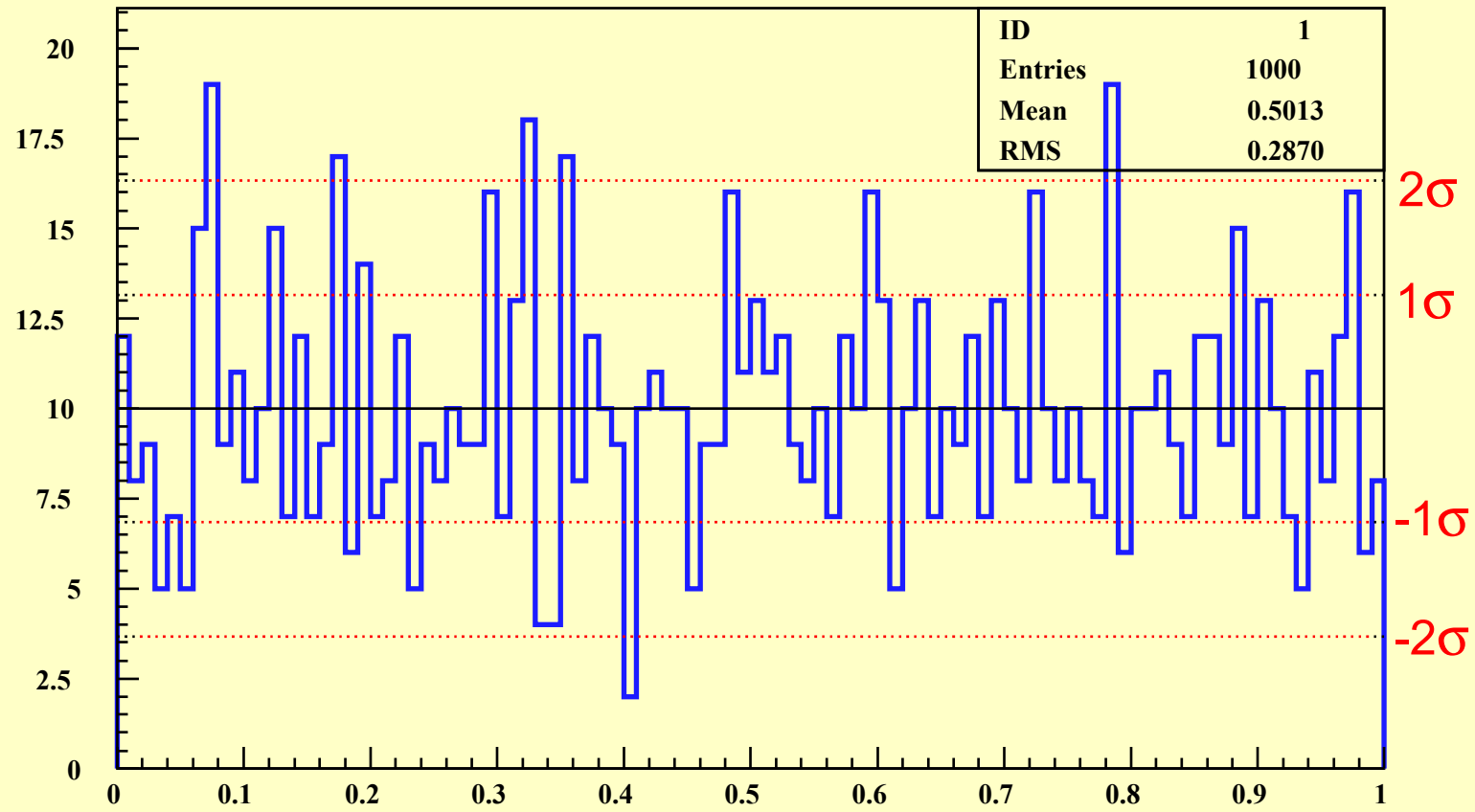


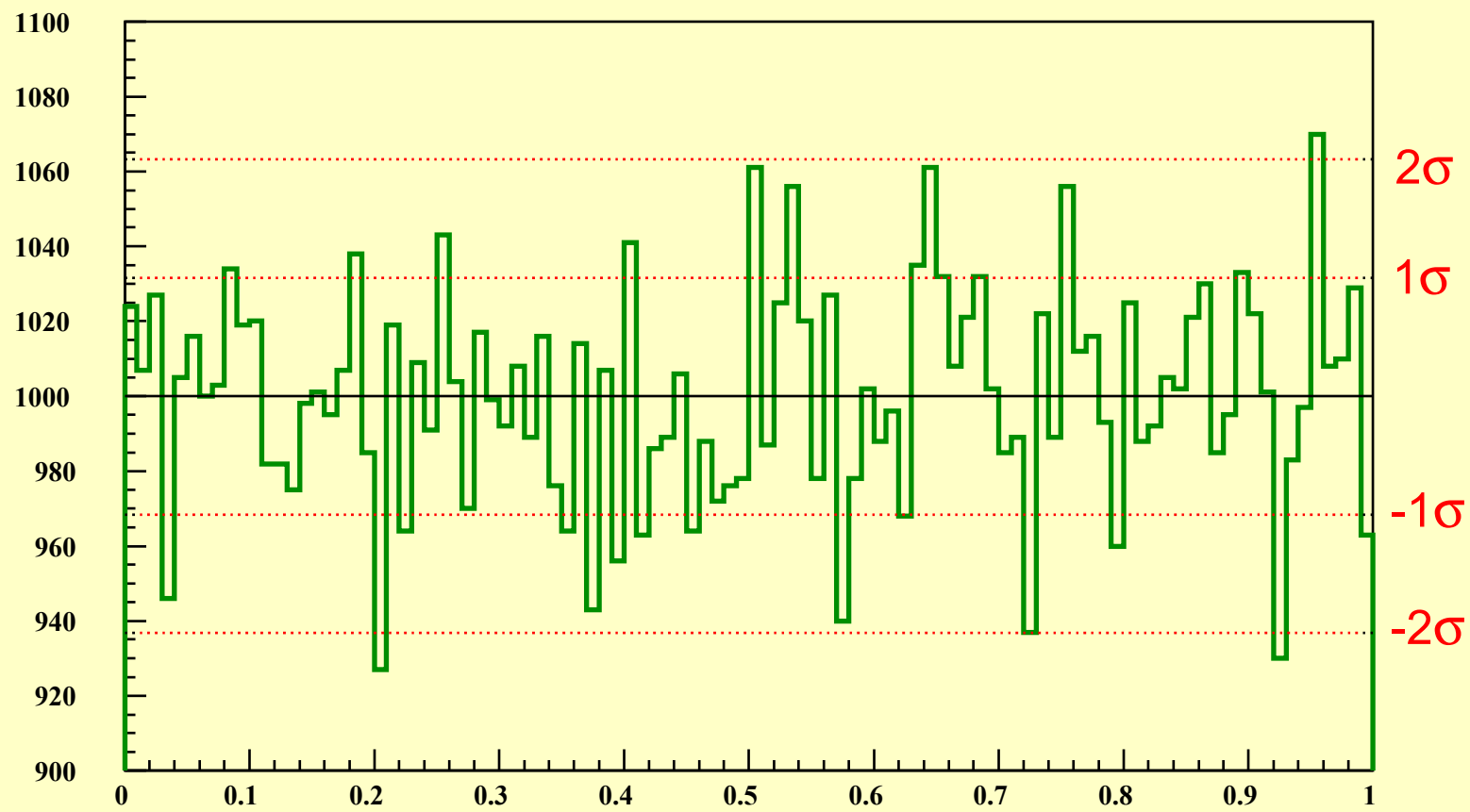
$10^5$  events



$10^7$  events





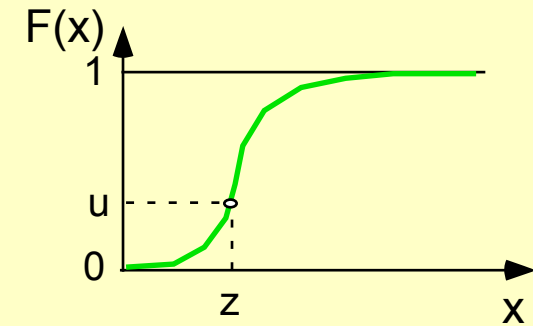
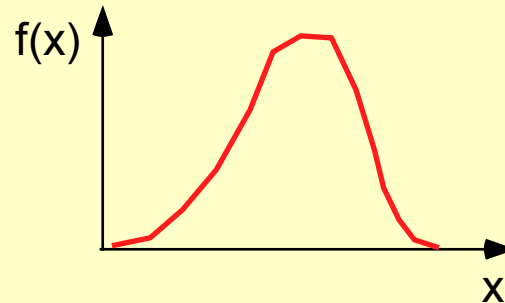


# Generators for Arbitrary Distributions

create random numbers according to the probability distribution  $f(x)$

Method 1: inverse integral method

$$F(x) = \int_{-\infty}^x f(x) dx$$



by construction  $F(x)$  ranges between 0 and 1, if  $f(x)$  is a probability density.

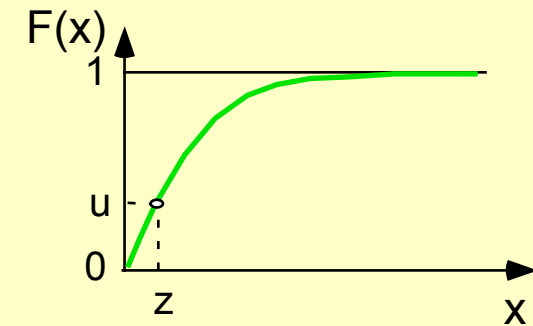
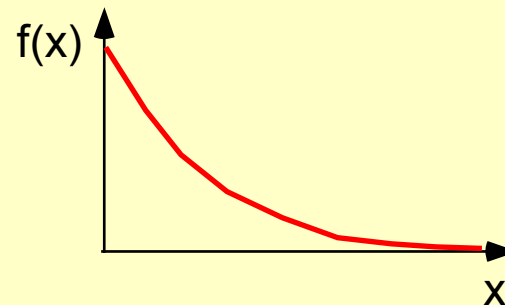
if we draw  $u = F(x)$  at random from a uniform distribution and find  $z = F^{-1}(u)$  then  $z$  is distributed like  $f(x)$ .

e.g. **exponential distribution**:  $f(x) = \lambda \exp(-\lambda x)$

$$u = F(x) = 1 - \exp(-\lambda x)$$

$$z = -\ln(1-u) / \lambda \quad \text{or}$$

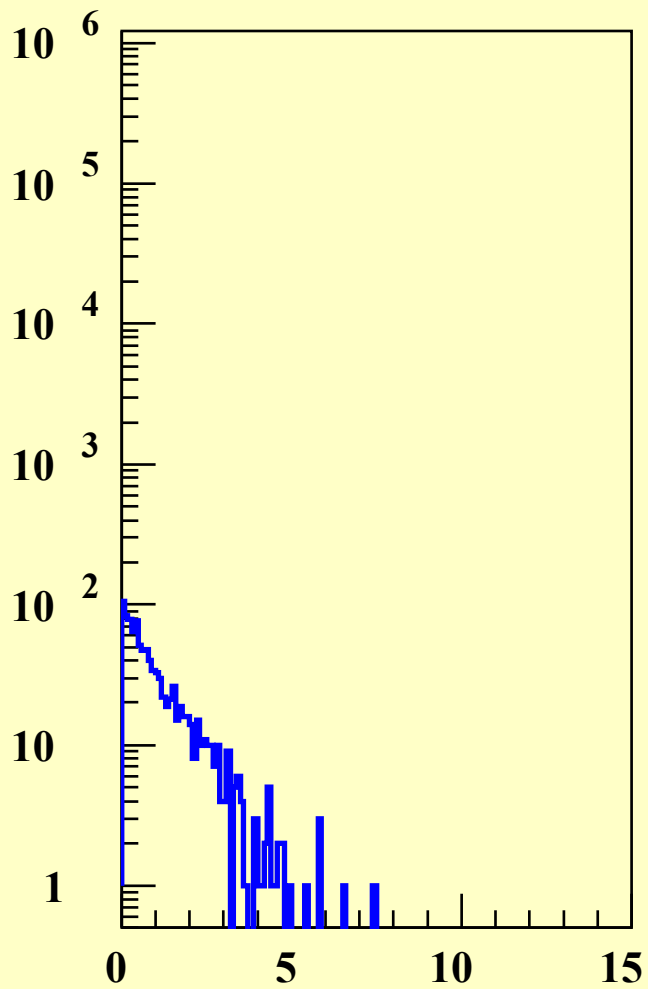
$$z = -\ln(u) / \lambda$$



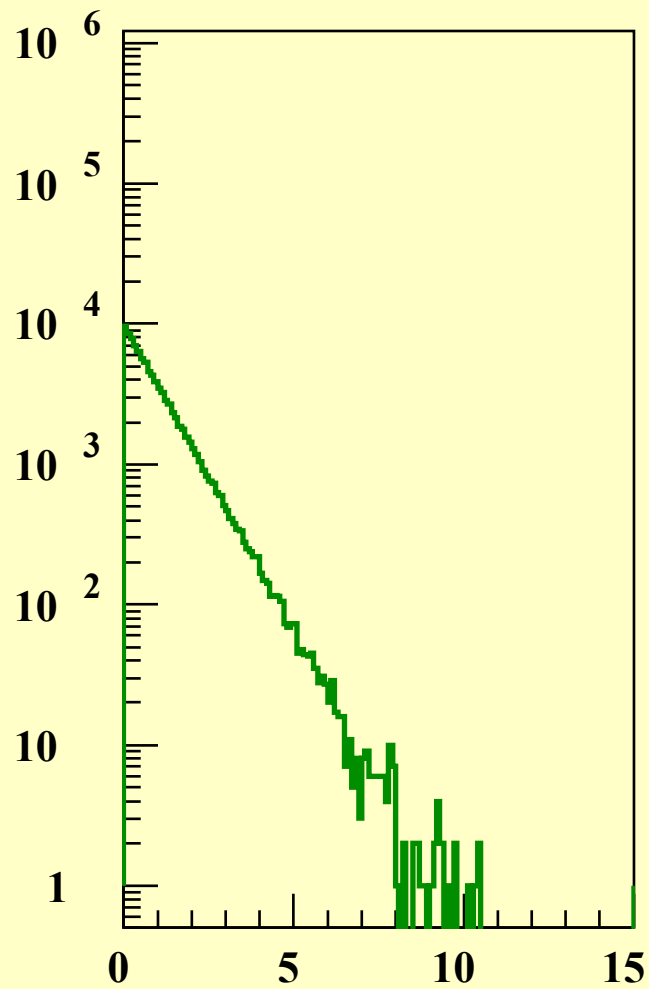
Only one call to uniform generator gives  $u$   
then  $z$  is exponentially distributed.

# Exponential

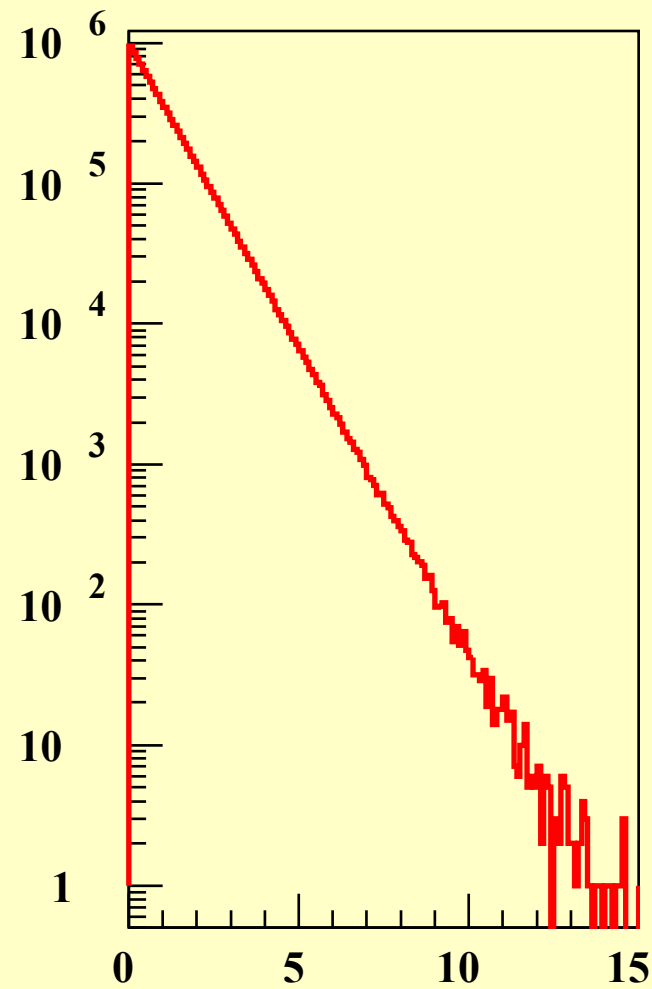
$10^3$  Events



$10^5$  Events



$10^7$  Events



# What if we cannot compute $F^{-1}(x)$ ?

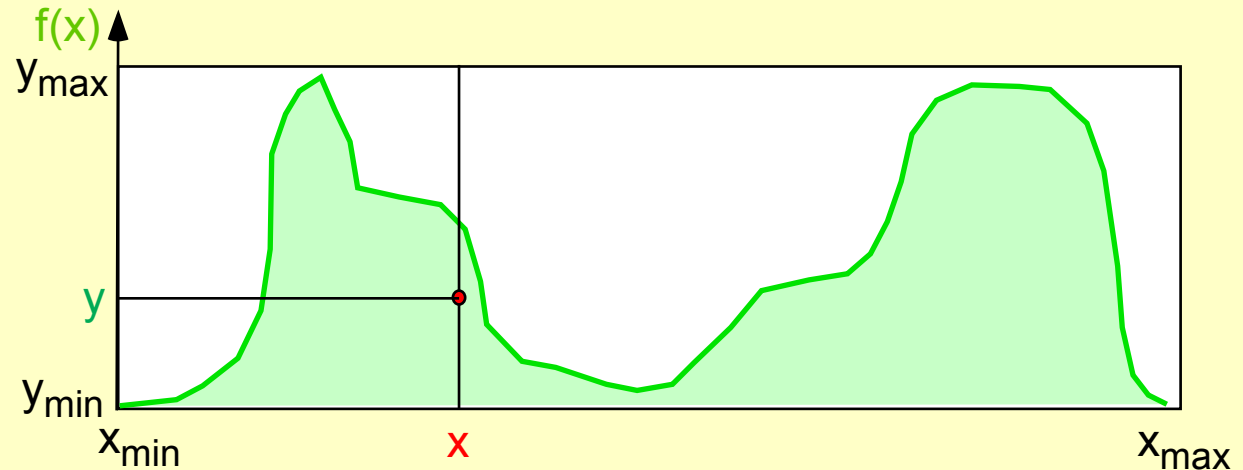
## Method 2: brute-force method

draw  $x$  at random from a uniform distribution between  $(x_{\min}, x_{\max})$

draw  $y$  at random from a uniform distribution between  $(y_{\min}, y_{\max})$

if  $y < f(x)$  accept  $x$  as random number

if  $y > f(x)$  reject  $x$  and start over again



$x$  will be distributed like  $f(x)$ , but at least 2 calls for uniformly distributed random numbers are necessary.

The rejection makes this method even more inefficient.

The larger the white area the more rejections will happen.

# Gaussian Distribution

$$f(x) = \frac{1}{\sqrt{2\pi} \sigma} \exp\left(-\frac{(x-\mu)^2}{2\sigma^2}\right)$$

Gaussian with mean  $\mu$  and standard deviation  $\sigma$

... is very important: if many independent little variations contribute to one variable then it is following a Gaussian distribution

i.e.  $z = \sum_{j=1,12} u_j - 6$  will produce a distribution of  $z$  that is about Gaussian, but needs 12 calls for uniform random numbers.

A better way:

1) create 2 uniformly distributed random numbers  $u_1$  and  $u_2$

2)  $v_1 = 2u_1 - 1$  uniformly distributed in  $(-1,1)$

$$v_2 = 2u_2 - 1$$

3)  $r^2 = v_1^2 + v_2^2$

if  $r > 1$  then goto 1)

if  $r < 1$  continue

$$4) z_{1/2} = v_{1/2} \sqrt{\frac{-2 \ln r^2}{r^2}}$$

are two independent and Gaussian distributed random numbers with mean 0 and standard deviation 1

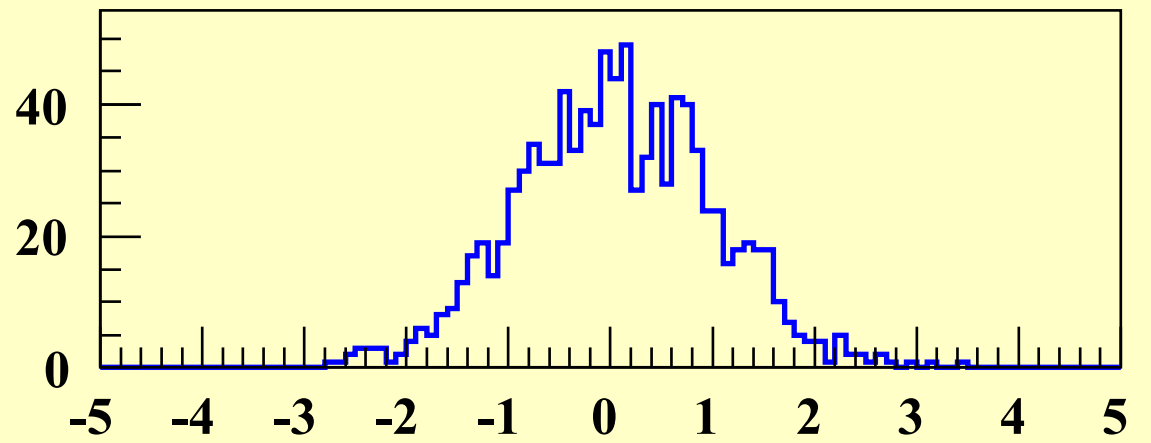
To get a Gaussian with mean  $m$  and standard deviation  $s$

$$\text{compute } Z_{1/2} = z_{1/2} \cdot s + m$$

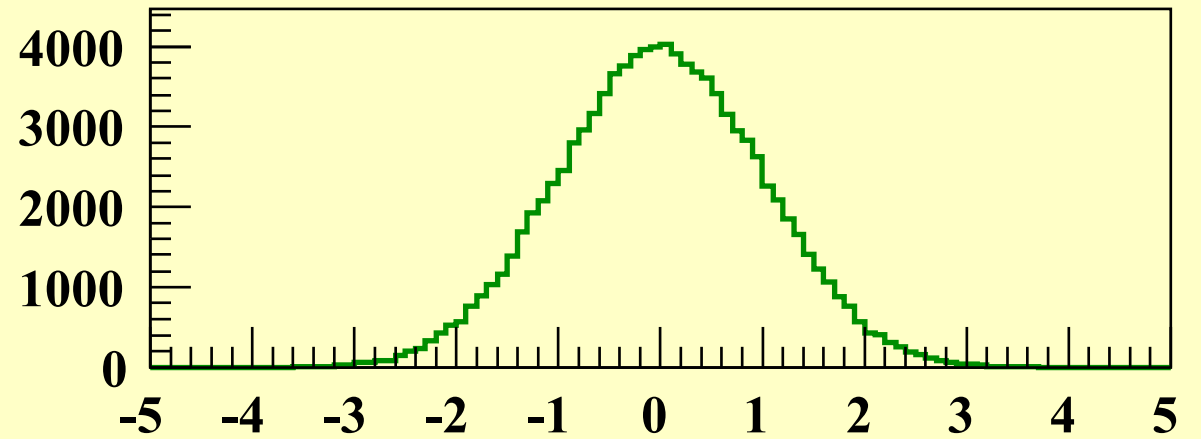


# Gaussian

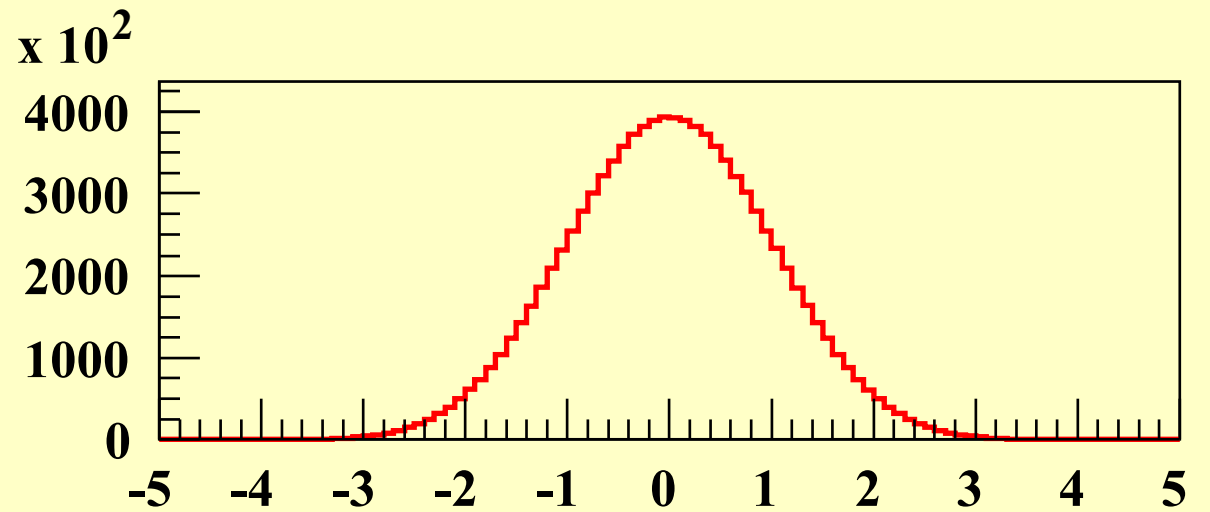
$10^3$  events



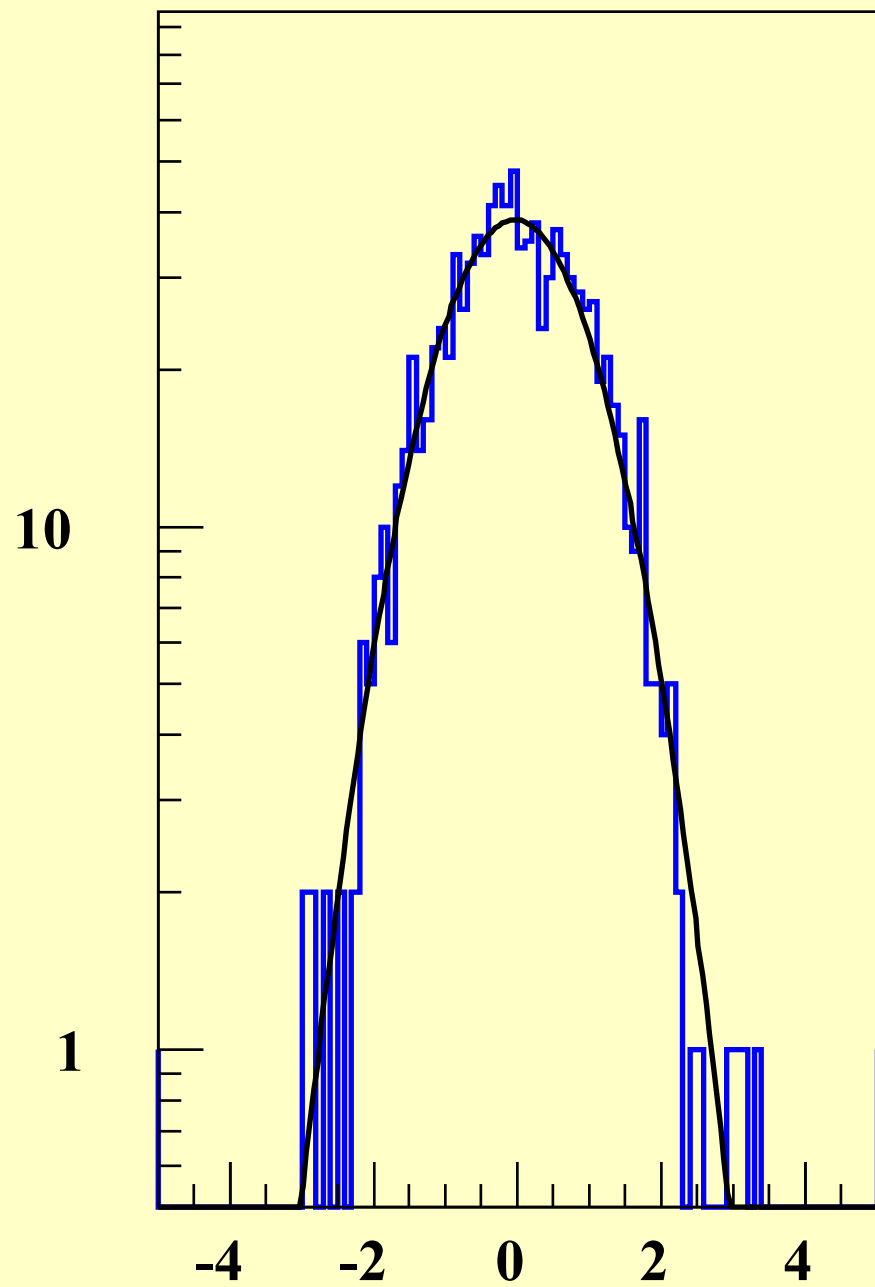
$10^5$  events



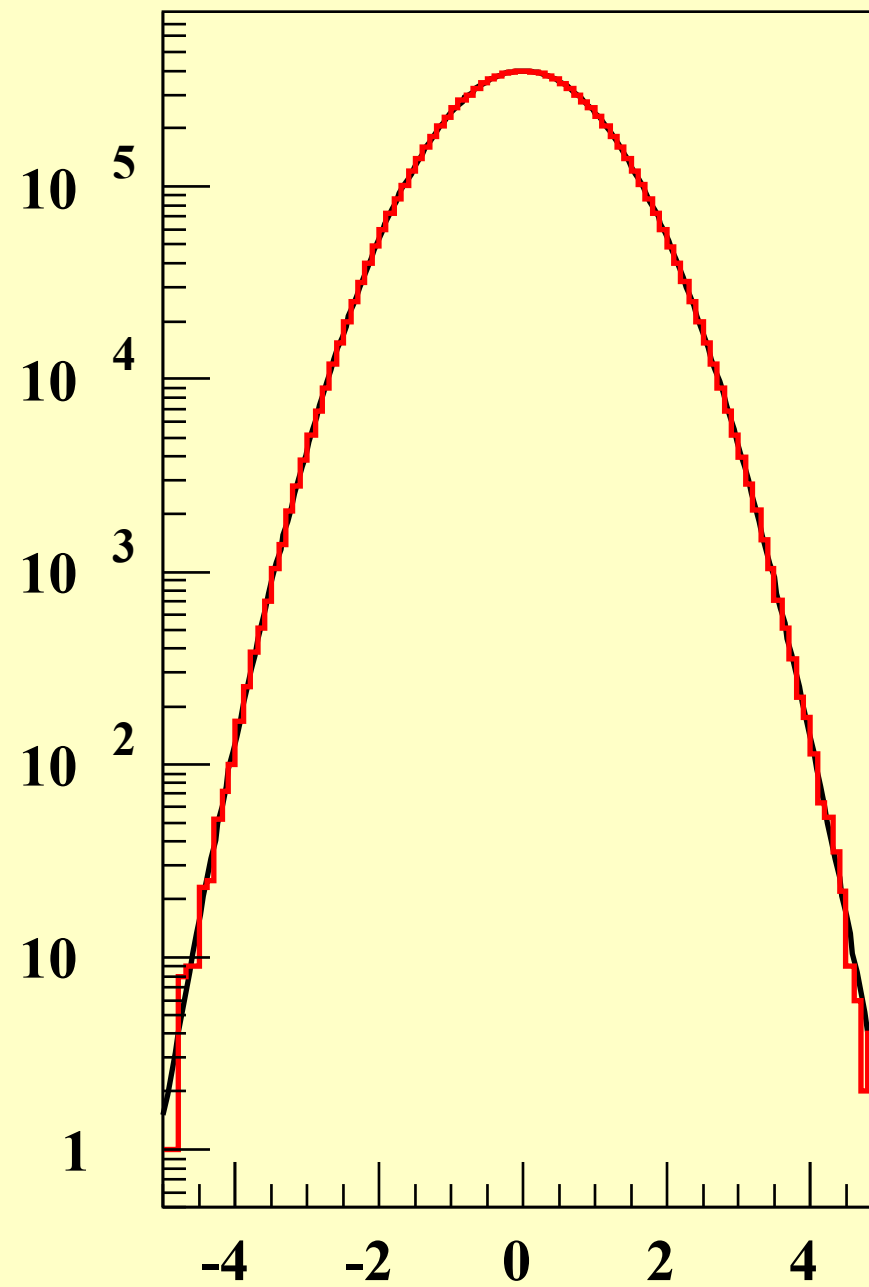
$10^7$  events



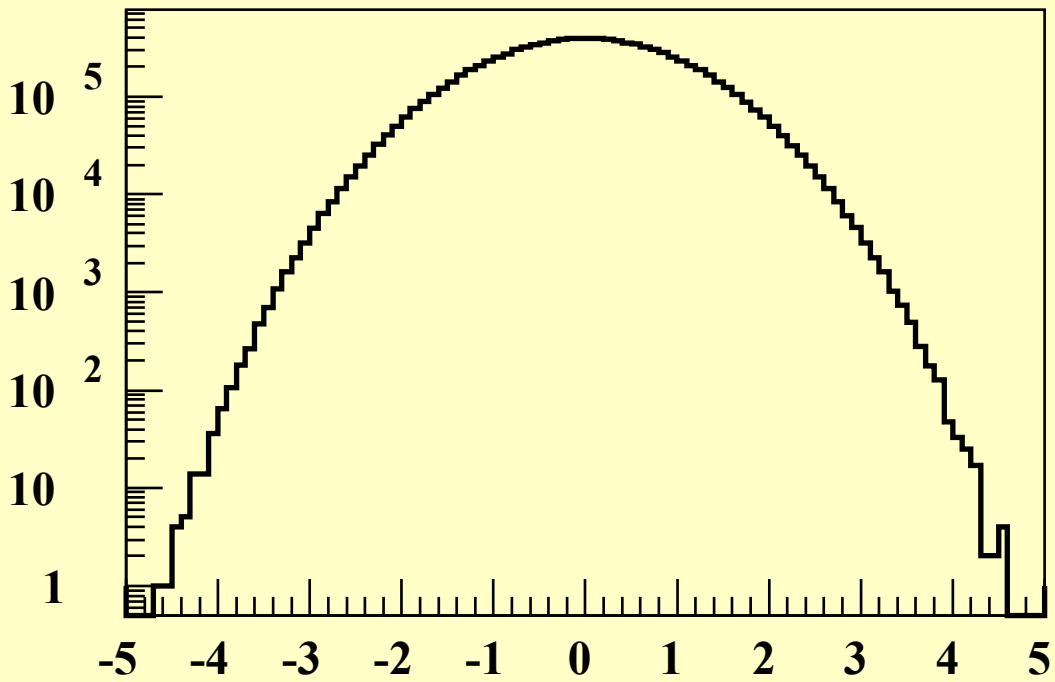
$10^3$  Events



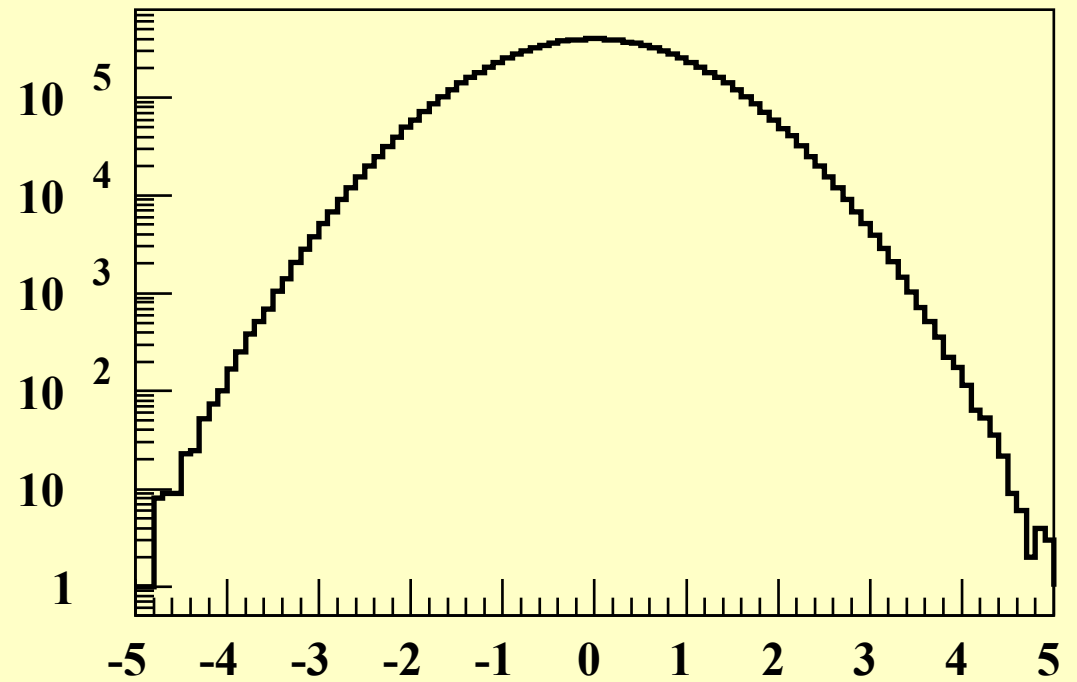
$10^7$  Events



*Simple Gaussian*



*Advanced Gaussian*

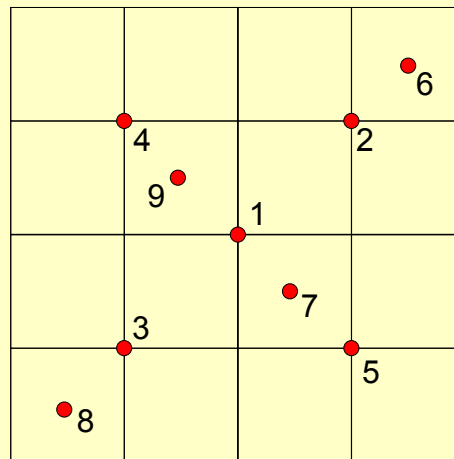


# Quasi-Random Number Generators

Sometimes one wants to map out a value range faster and more uniform than with pseudo-random numbers.

Quasi-Random numbers are **constructed** to fill a given space as uniformly as possible. (i.e. to avoid clusters)

e.g. Sobol Generator:

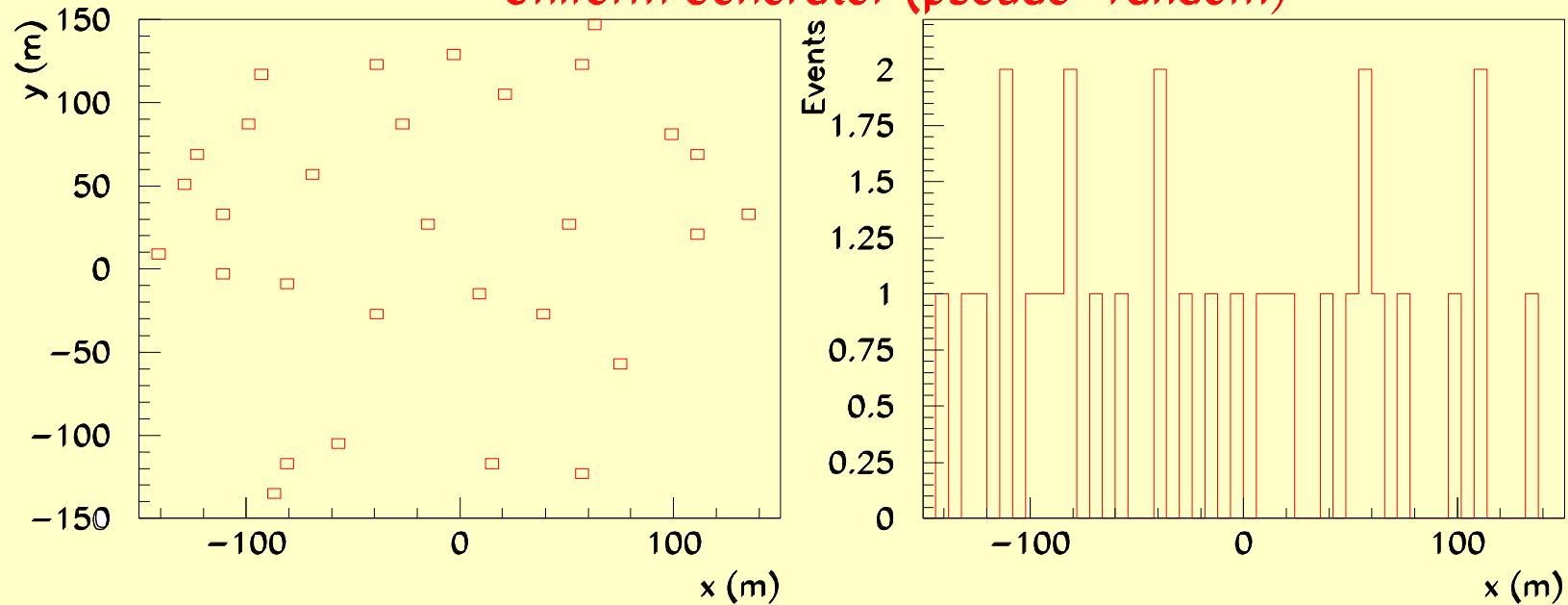


Exhibits reduced fluctuations as compared to pseudo-random generators.

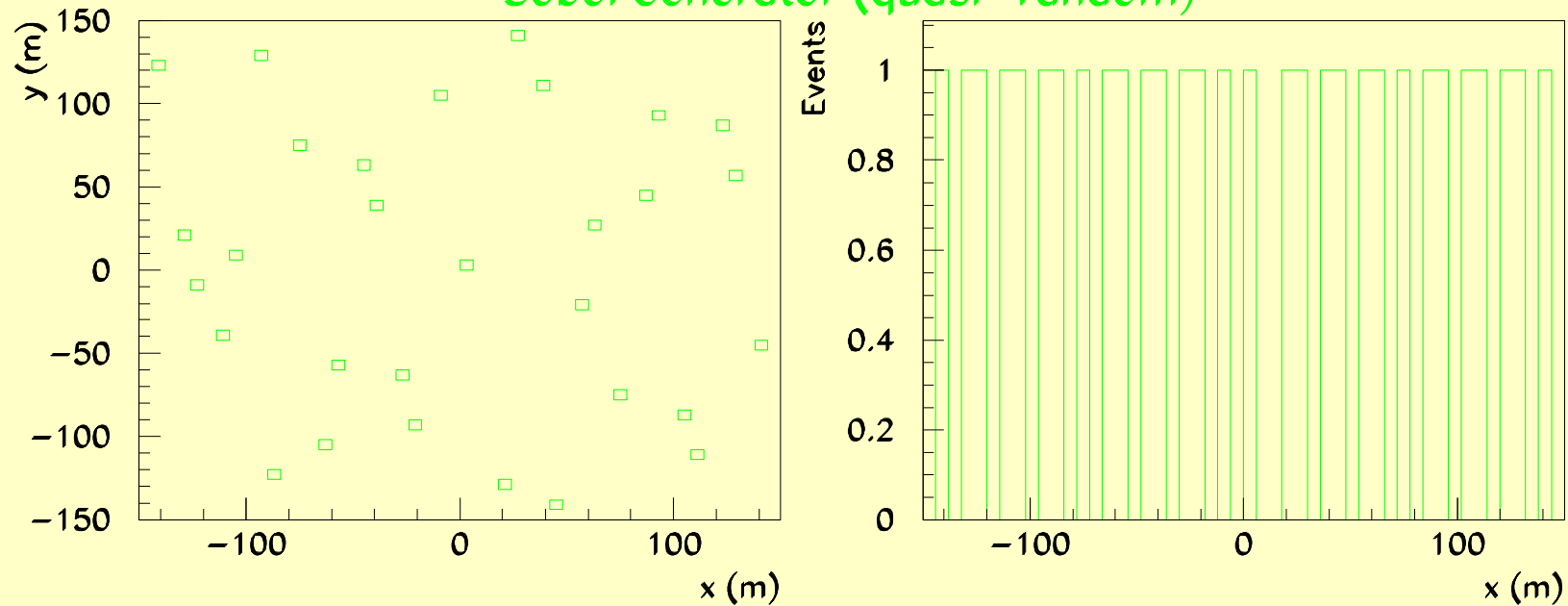
Therefore usually not used often in Monte Carlo.

# Sobol Quasi-Random Generator vs Pseudo-Random Generator: 30 evts

## Uniform Generator (pseudo-random)

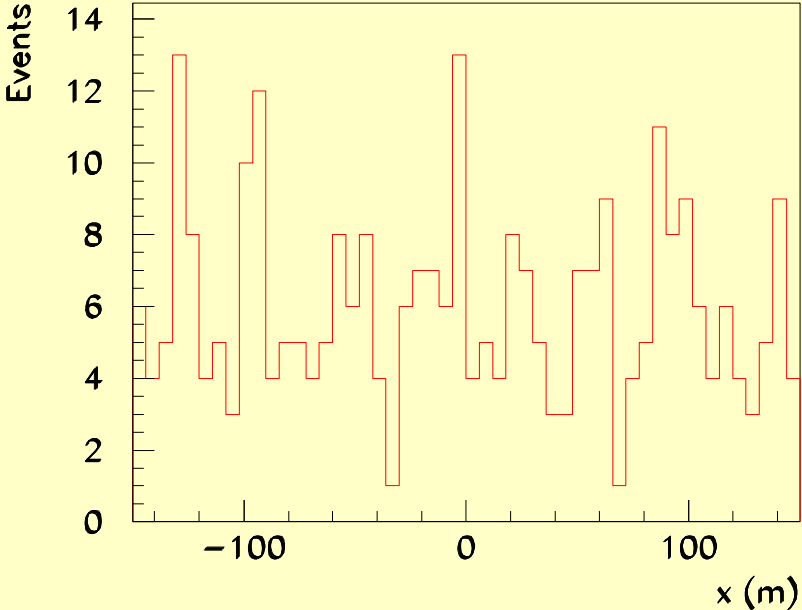
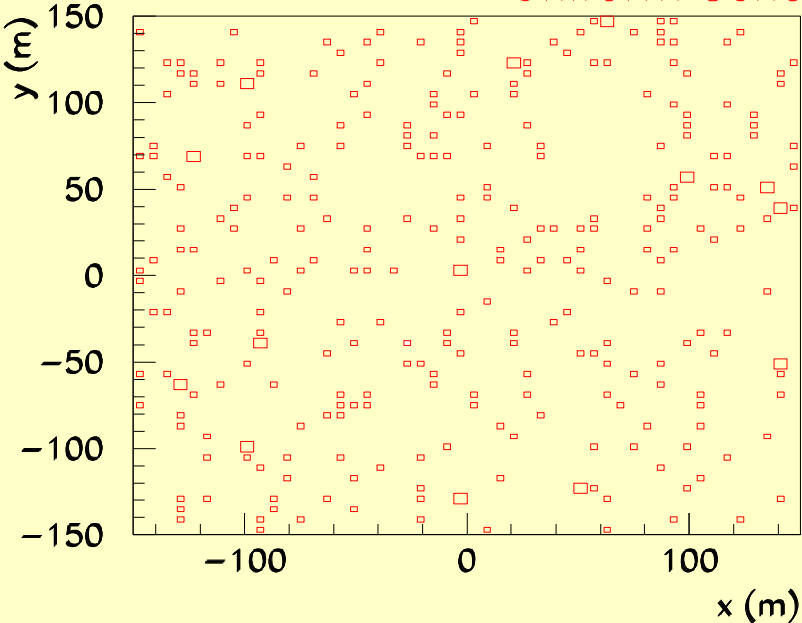


## Sobol Generator (quasi-random)

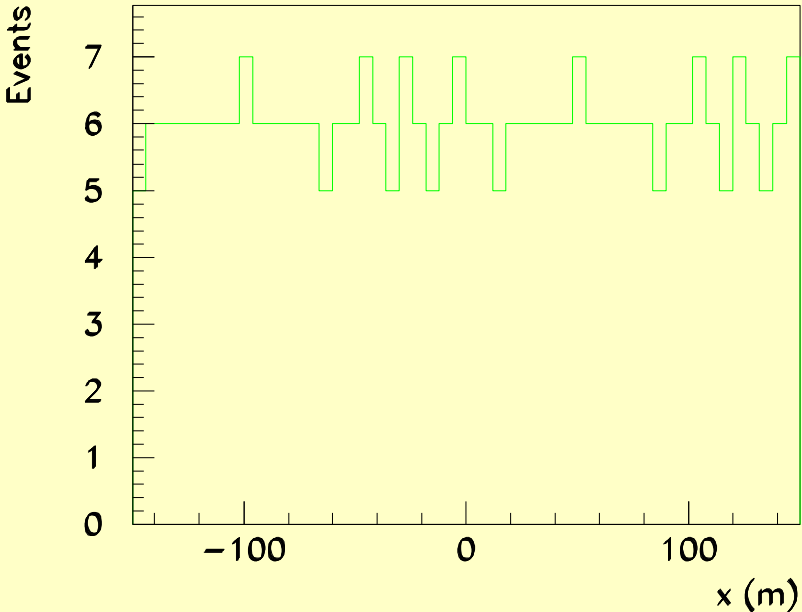
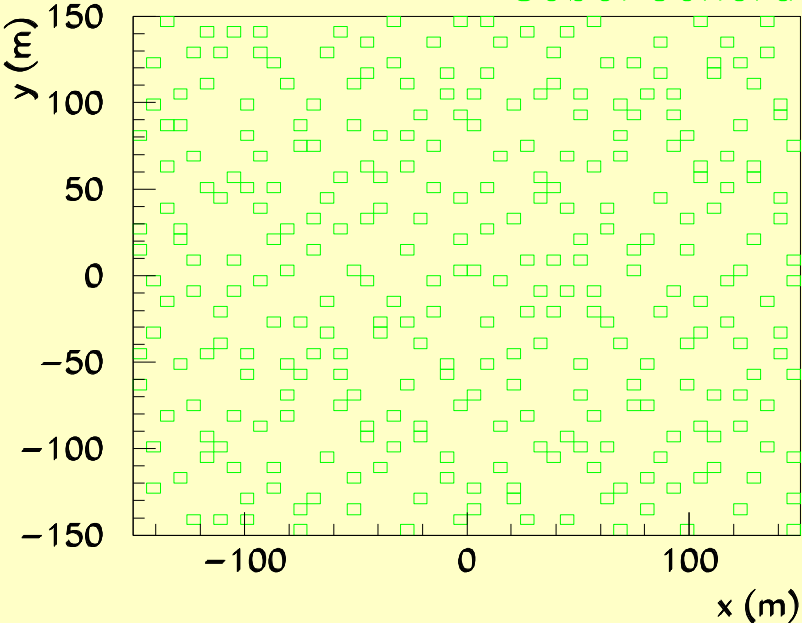


# Sobol Quasi-Random Generator vs Pseudo-Random Generator: 300 evts

Uniform Generator (pseudo-random)

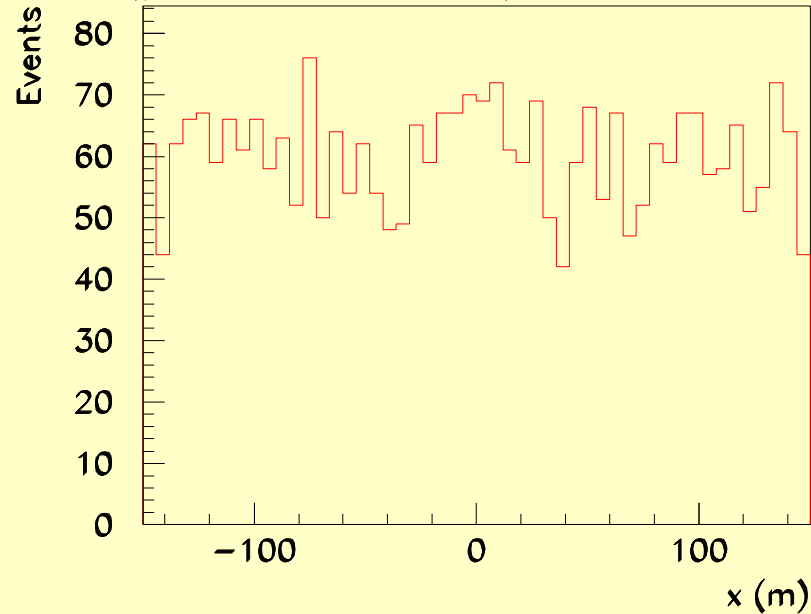
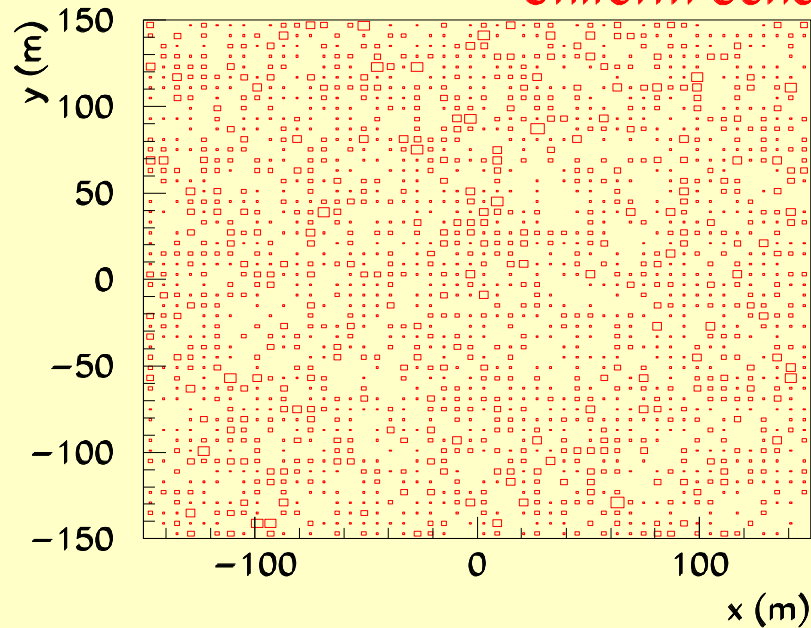


Sobol Generator (quasi-random)

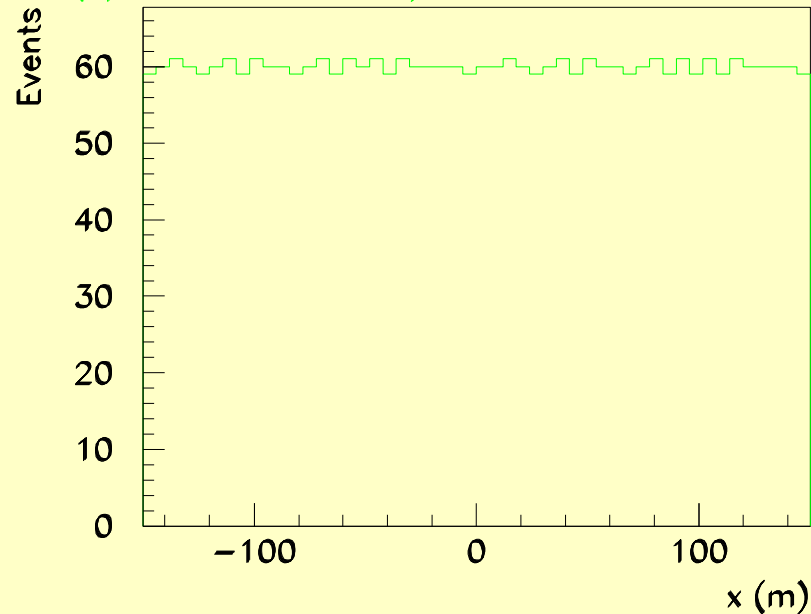
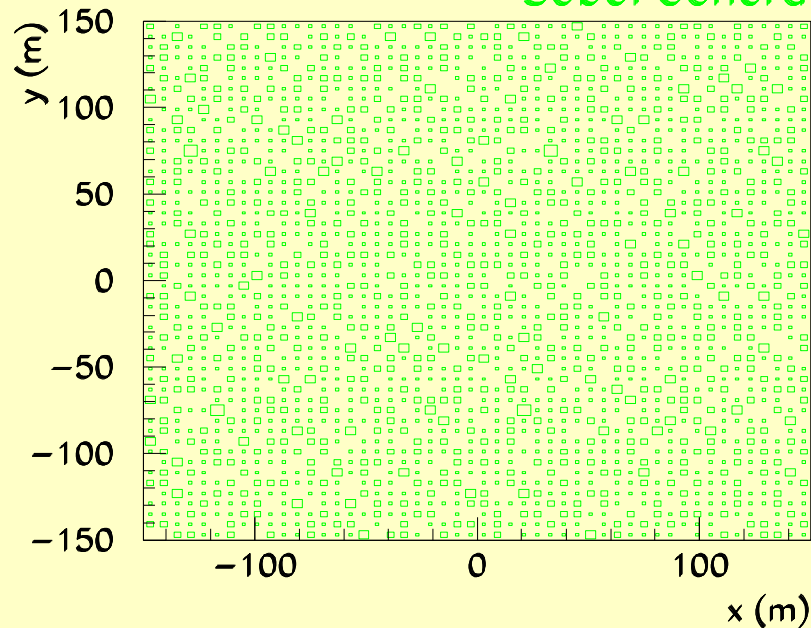


# Sobol Quasi-Random Generator vs Pseudo-Random Generator: 3000 evts

## Uniform Generator (pseudo-random)



## Sobol Generator (quasi-random)



# Monte Carlo Pitfalls

beware of bad random number generators

not random, too short sequence, correlations

rounding errors

e.g. emission angles at high energies

steeply falling distributions



## Summary:

Monte Carlo Simulations provide a powerful tool for many parts of science and engineering, especially when statistical processes are involved.

MC Methods are prone to subtle errors, due to random numbers and finite numerical precision. Therefore, they need careful testing.

Nevertheless, MCs are **undispensible** for Air Shower analysis and experiment design.