

Detectors of Cosmic Rays Gamma-Rays and Neutrinos

(I)

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Third School on Cosmic Rays and Astrophysics
August 25 to September 5, 2008
Arequipa - Perú

(I)

- Introduction
- Something about historical detectors
- Space detectors

Our aims:

PARTICLE IDENTIFICATION

ENERGY MEASUREMENTS

ARRIVAL DIRECTIONS

TIME VARIATIONS

Wide range of energies:

- Dark matter 10^3 eV
- Solar/SN neutrinos 10^6 eV
- Gamma-rays $10^3 - 10^{12}$ eV ... 10^{21} eV
- Charged CR: 10^6 eV – 10^{21} eV
- C. neutrinos ... 10^{21} eV

Wide quality of detectors

- $-dE/dx$
- Secondary/cascade measurements
- Magnetic deflection
- Cherenkov
- Acoustic
- Radio

Wide range of dimensions

- dark matter / space detectors ($< \text{m}$)

LE neutrinos (10 – 100 m)

- “sampling arrays”

Acceptance area $\gg$ sensitive area

(Auger: 3000 km^2 vs 16000 m^2 ... factor 200000!)

Shower detectors

- Extensive Air Showers (HE CRs, Gamma rays)
- Underwater/ice neutrino (km)

Wide range of locations

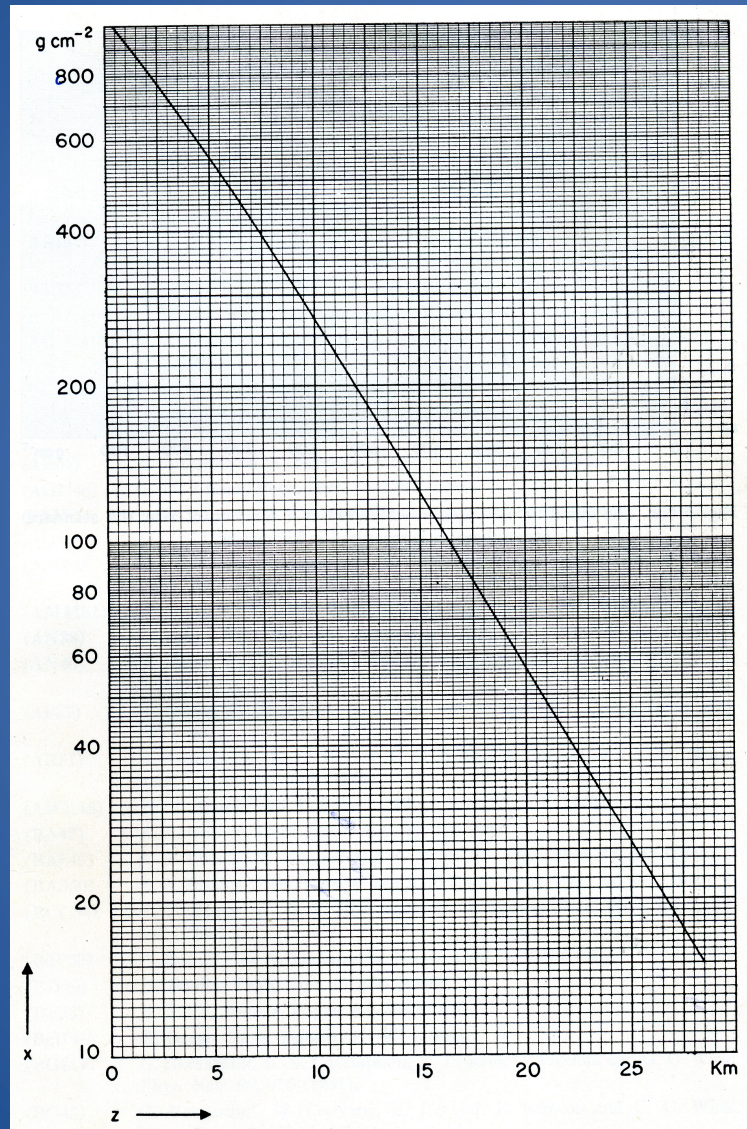
- Space
 - Ground based
 - ... sea level... mountain... “deserts”
 - Underground
 - Underwater
 - Under-ice / on-ice
- + long duration, stable observations
- + sensitivity to bursts

DIRECT/NON-DIRECT observations

- DIRECT : the particle under study interacts in your detector: you can measure the particle itself or its products (as for photons)
- NON-DIRECT: the secondaries produced by the particle are detected: typical of ground based charged particle experiments

THE EARTH ATMOSPHERE

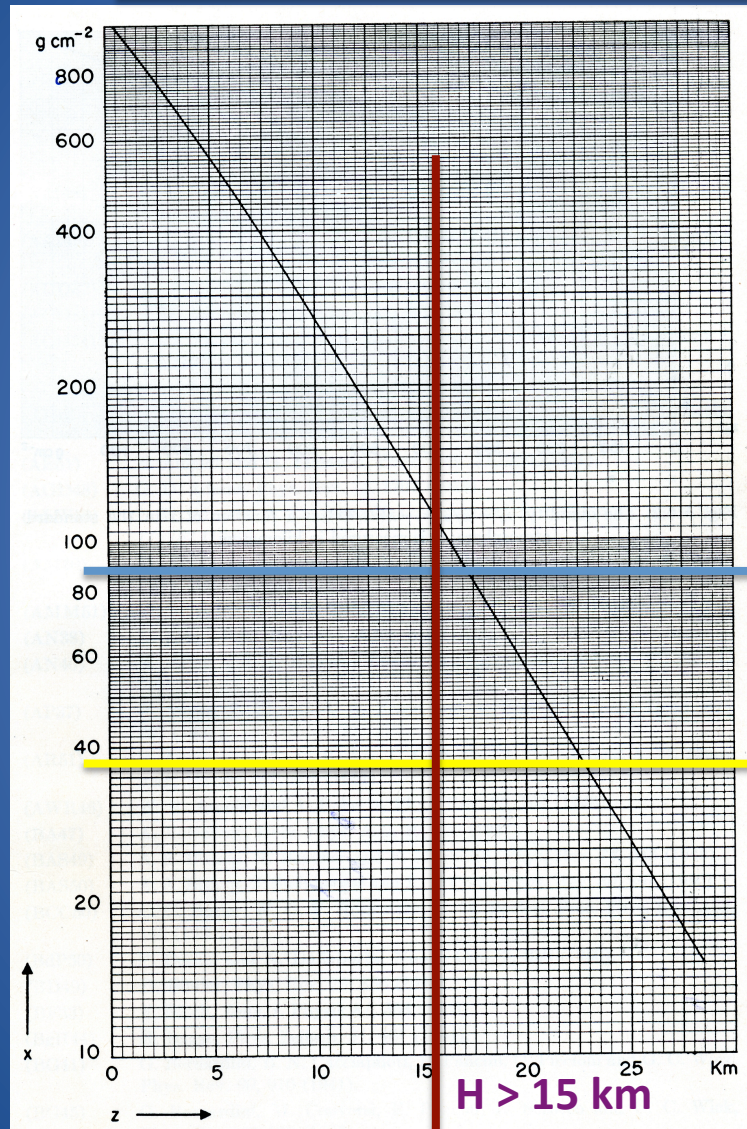
ATMOSPHERIC
PRESSURE



ALTITUDE A.S.L.

THE EARTH ATMOSPHERE

ATMOSPHERIC
PRESSURE



$\times 10^{10}$ TeV neutrinos m.f.p.

Proton m.f.p. (hadronic int.)

Photon m.f.p. (pair production)

ALTITUDE A.S.L.

A few words on history...

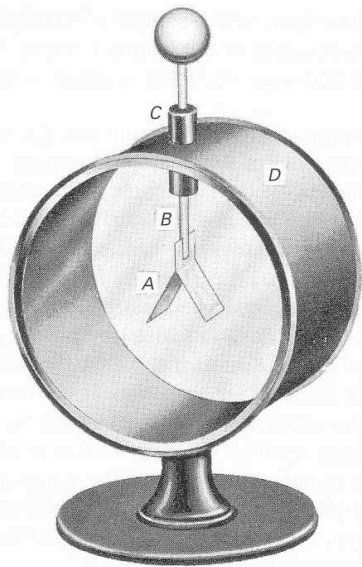
- Detectors and methods

Why...

- to start with the basis of our detectors
- and also to stress the “globality” of CR detectors and remind that these methods are ready to be applied on different scales

THE DISCOVERY OF COSMIC RAYS

The electroscope



COLLECTIVE DETECTOR

When electroscope is charged, the leaves (A) are pushed apart

The ionization of the gas inside discharge the electroscope and the leaves move towards each other. The rate at which the leaves came together measured the amount of ionization.

Spontaneous discharge of the electroscope !

1901 Wilson: the discharge is identical on the ground and in a tunnel

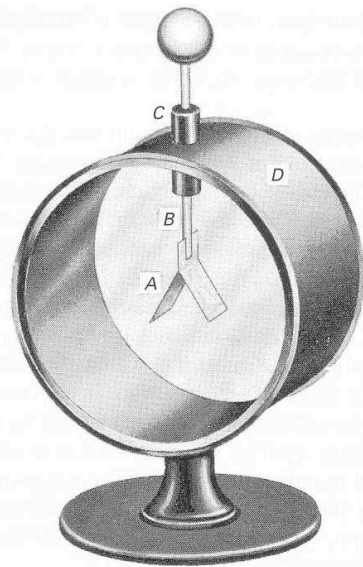
X-rays and radioactive substances discharged the electroscopes

Rutherford: this is due to the natural radioactivity

1912 V. Hess: COSMIC RAYS

THE DISCOVERY OF COSMIC RAYS

The electroscope

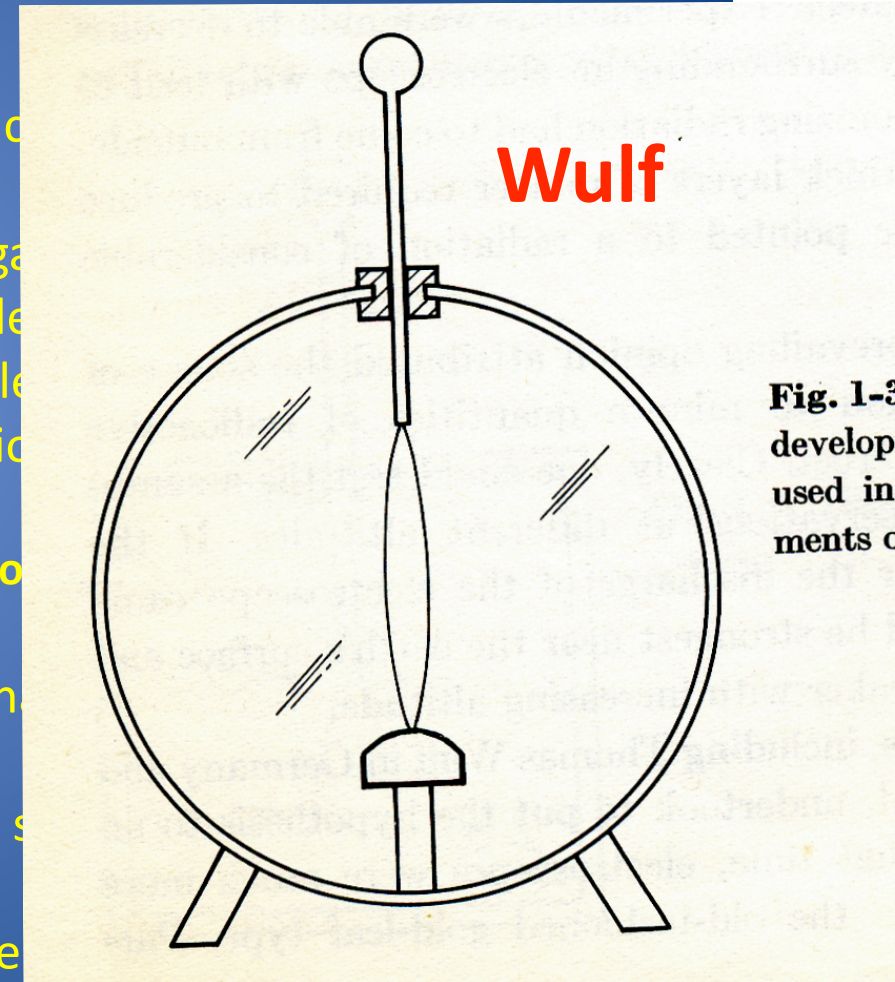


COLLECTIVE DETECTOR

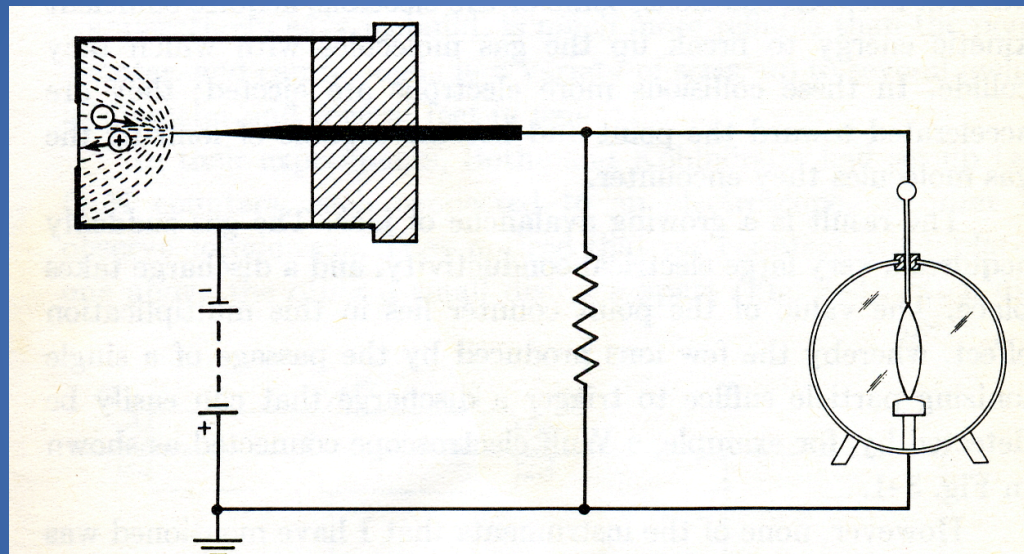
When electroscope is charged, the leaves are pushed apart.
The ionization of the gas inside the electroscope and the leakage of charge from the top plate.
The rate at which the leaves collapse is a measure of the amount of ionization.

Spontaneous discharge of electroscope

1901 Wilson: the discharge of an electroscope in a dark room and in a tunnel.
X-rays and radioactive substances cause the discharge of electrosopes.
Rutherford: this is due to alpha particles.
1912 V. Hess: COSMIC RAYS



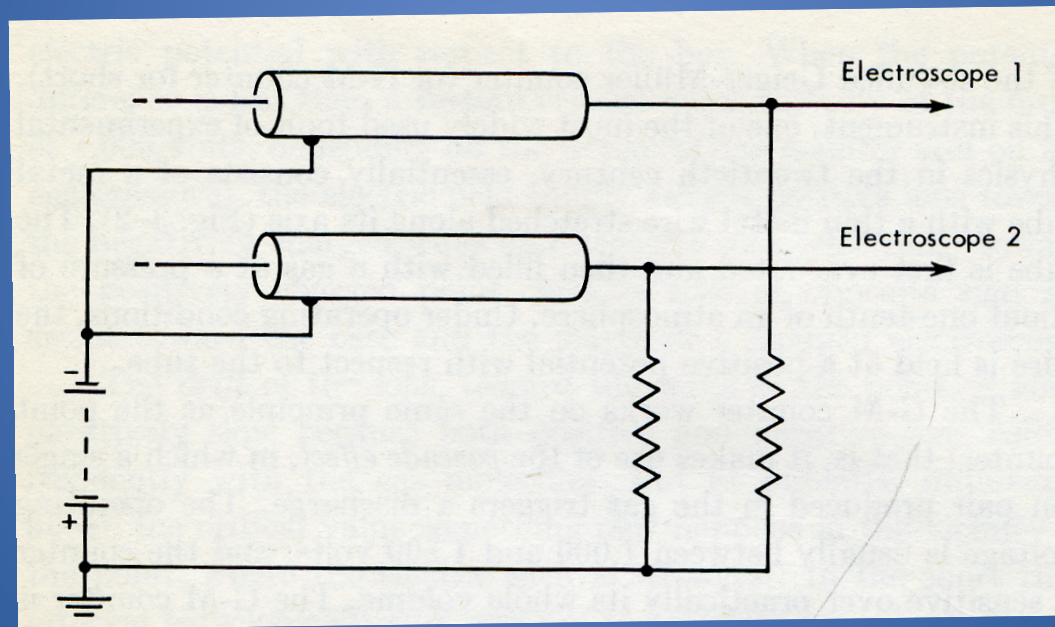
Geiger: the *point counter*



1929 ==> Geiger-Muller counter, avalanche

SINGLE PARTICLE DETECTOR

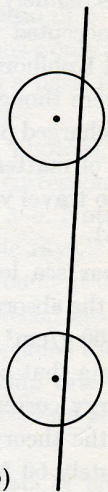
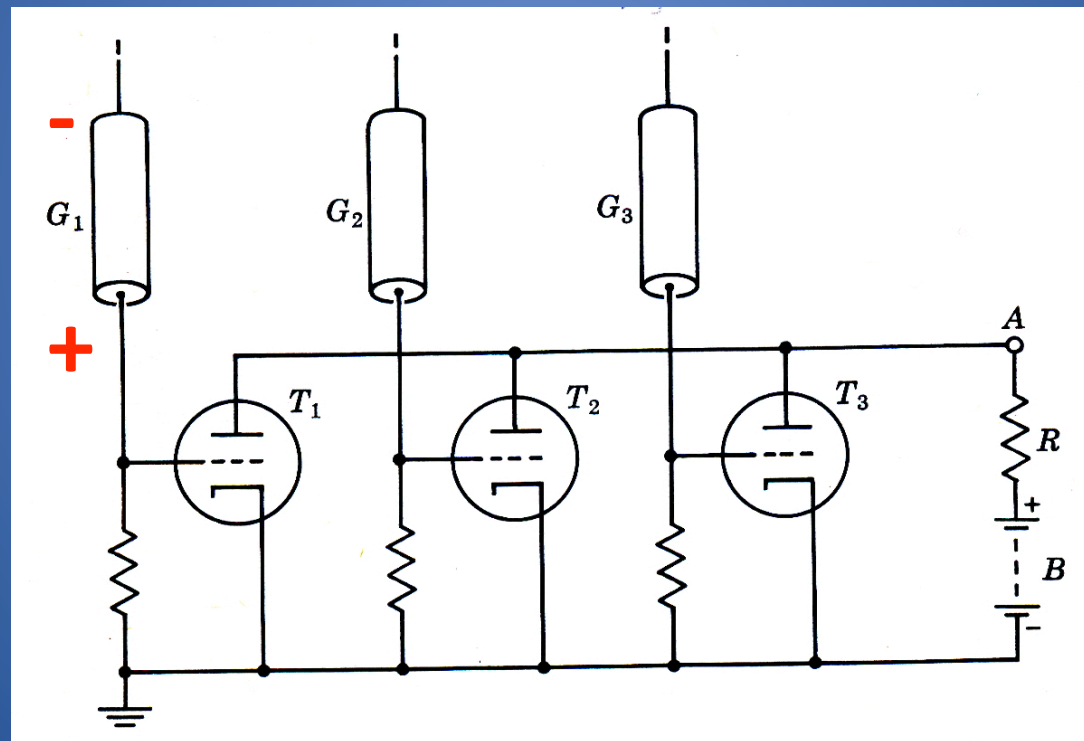
Coincidence: Bothe and Kohlhoerster



(Field of views, directions, particle penetration, trigger)

The coincidence circuit

B. Rossi



(b)

d-particle coincidences. (a) A photon produces a

Earth magnetic field

$$E \text{ [eV]} = 300 B R_c \text{ [gauss cm]}$$

B geomagnetic field

R_c radius of Earth

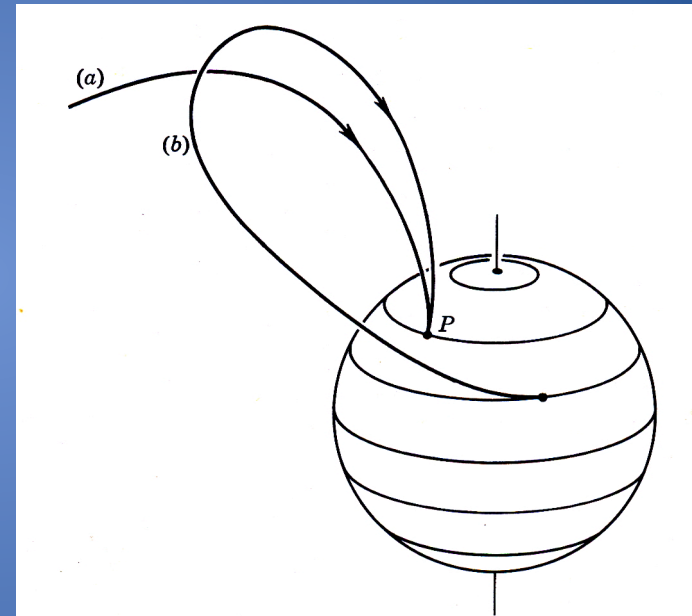
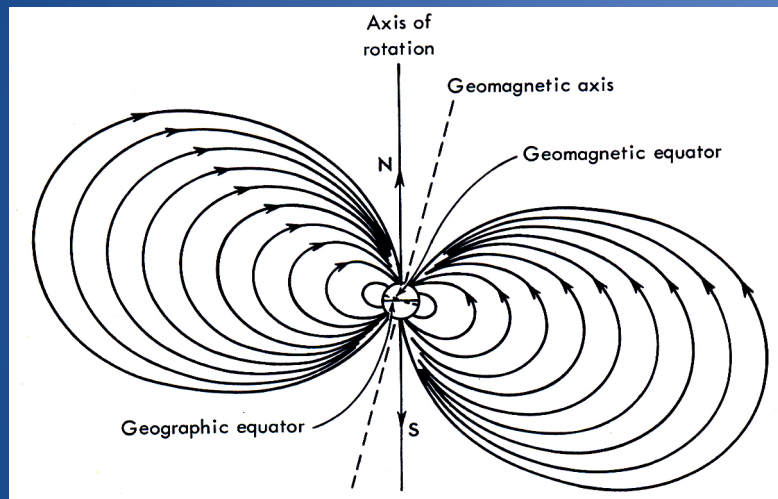
$$(B R_c)_E = 0.32 \times 6.4 \times 10^8 = 2 \times 10^8 \text{ gauss cm}$$

$$E_E = 300 \times 2 \times 10^8 \text{ [eV]} = 60 \text{ GeV}$$

Particles with $E < E_E$ can be affected by Earth

$$\gamma_E = 300 (Z_e/M_{eV}) B R_c \text{ [gauss cm]}$$

Earth environment as a magnetic analyzer



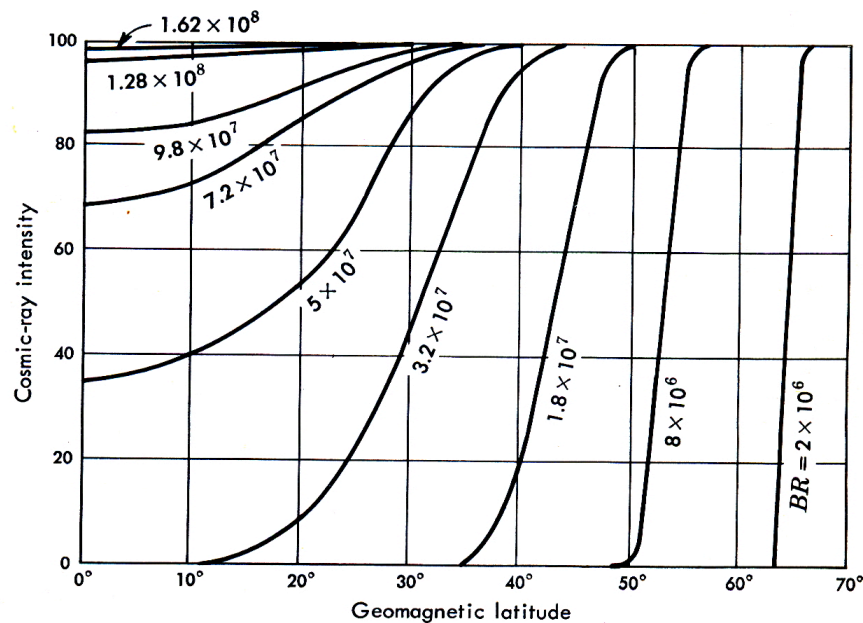
Allowed and forbidden directions

Atmosphere as a magnetic analyzer

latitude effect

CR charged particles

Allowed and forbidden directions



Clay 1927 Leiden-Java Compton

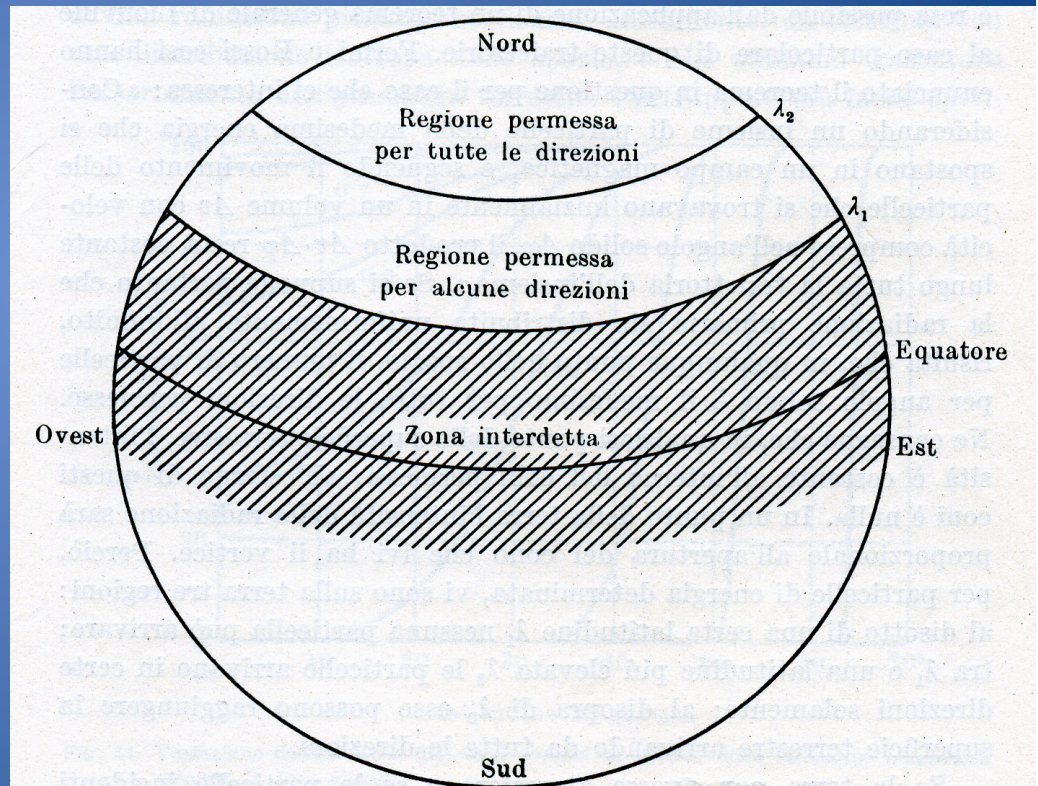


Fig. 50. Date delle particelle di una certa energia, esse non potranno raggiungere la Terra che alle latitudini maggiori di λ_1 . Al di sopra di λ_2 tutte potranno raggiungere la Terra, qualunque sia la loro direzione. Tra λ_1 e λ_2 solamente quelle la cui direzione è contenuta in un certo cono (aperto verso l'ovest per le positive, e verso l'est per le negative) potranno raggiungere la Terra.

Atmosphere as a magnetic analyzer

CR dominated by positive charges

E-W effect 1933

Johnson+Alvarez

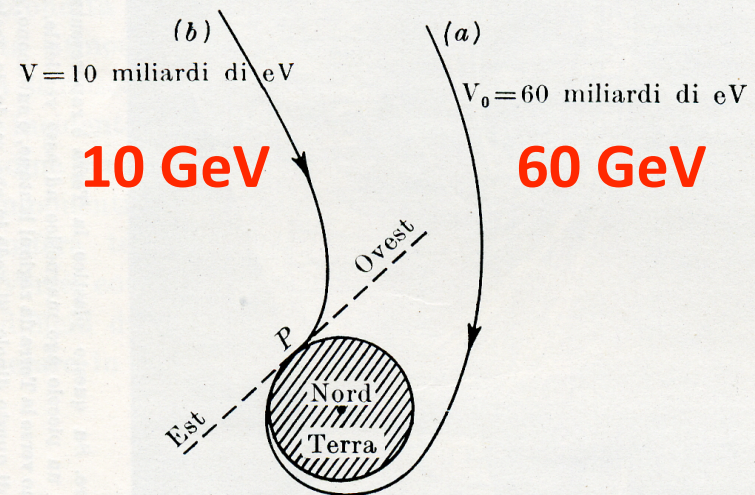
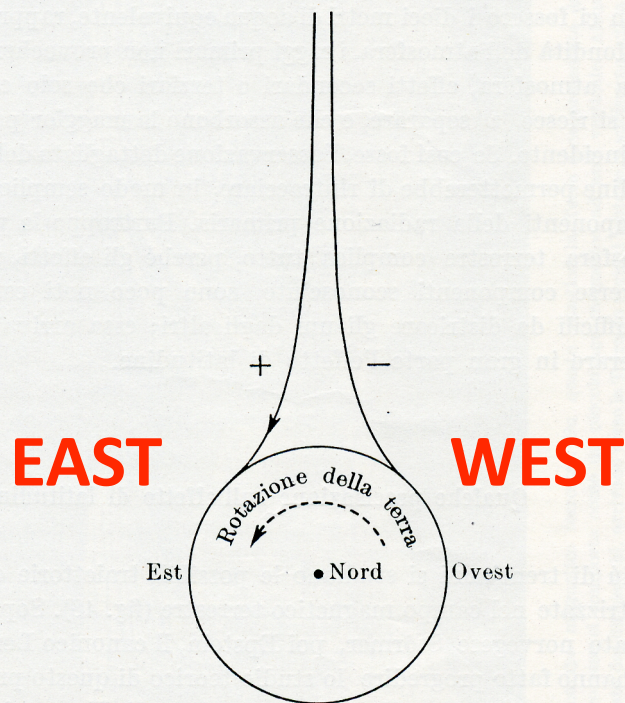
Mexico City 29 deg North

10% @ 45 deg

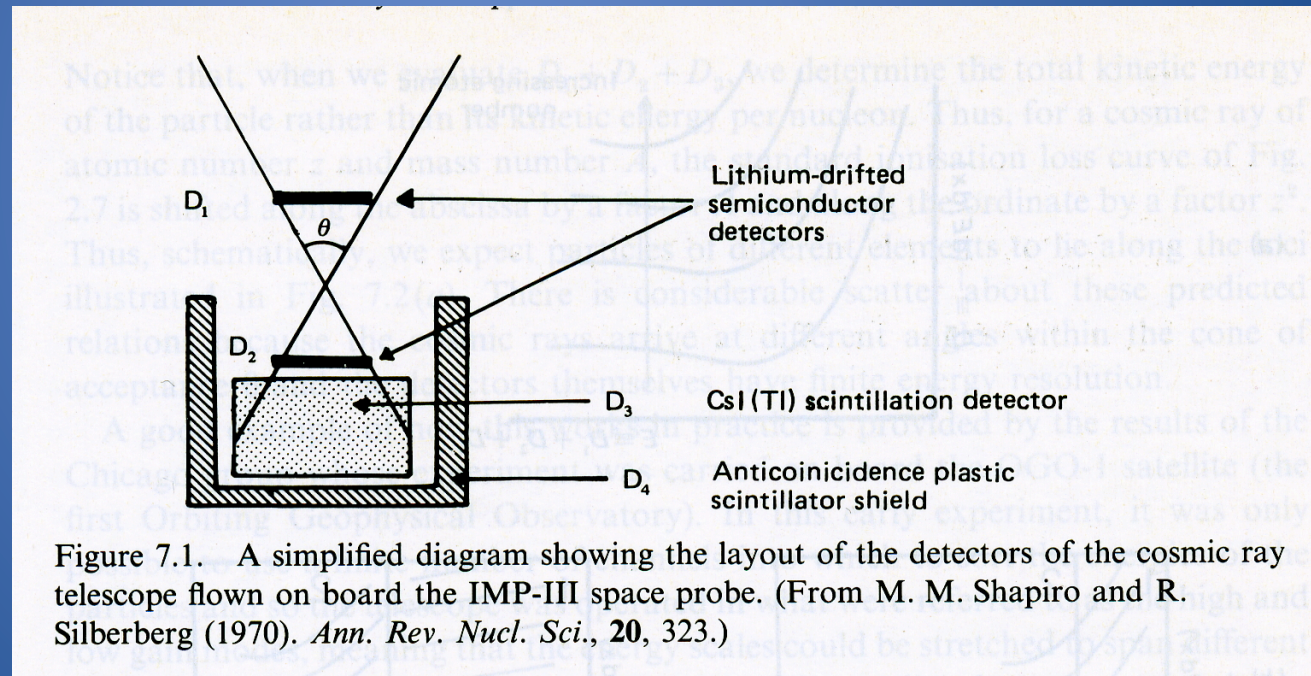
Rossi + De Benedetti

Asmara 11 dg North 26%

Excess from West ==> positive



DIRECT Space “classical” energy/particle identification Experiment (E vs $-dE/dx$)



Space “classical” energy/particle identification Experiment (E vs $-dE/dx$)

Geometry

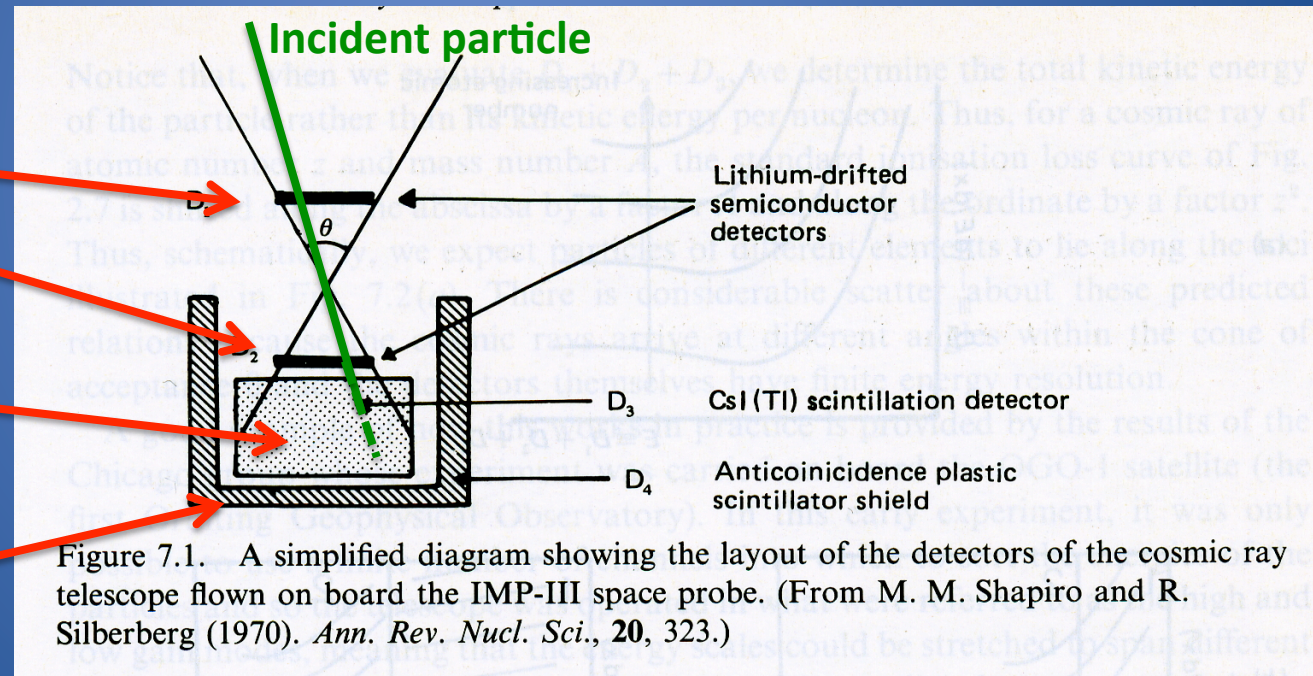
$-dE/dx$

Energy

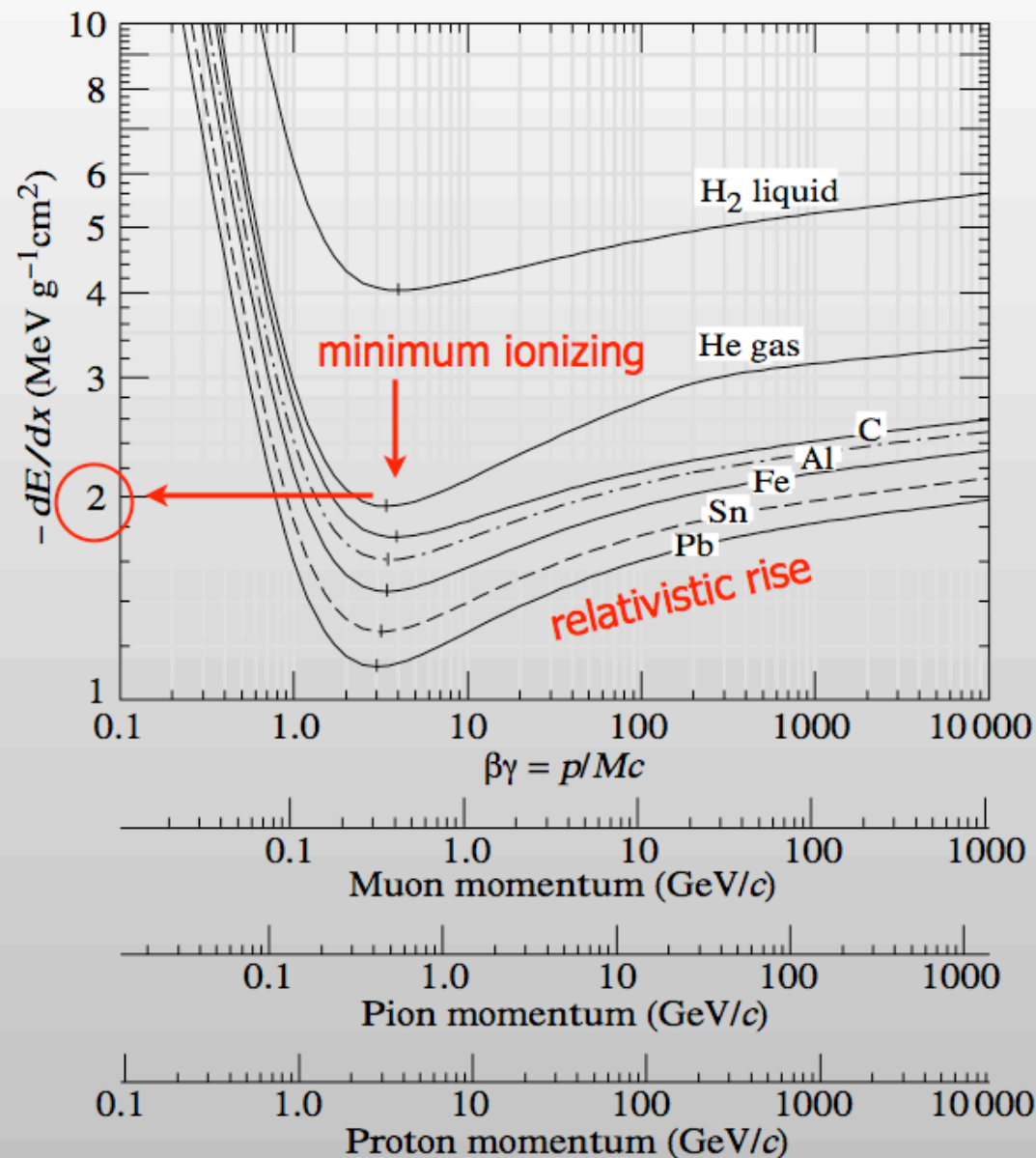
Calorimetry

AC for

Containment



Passage of ch. particles through matter (I)

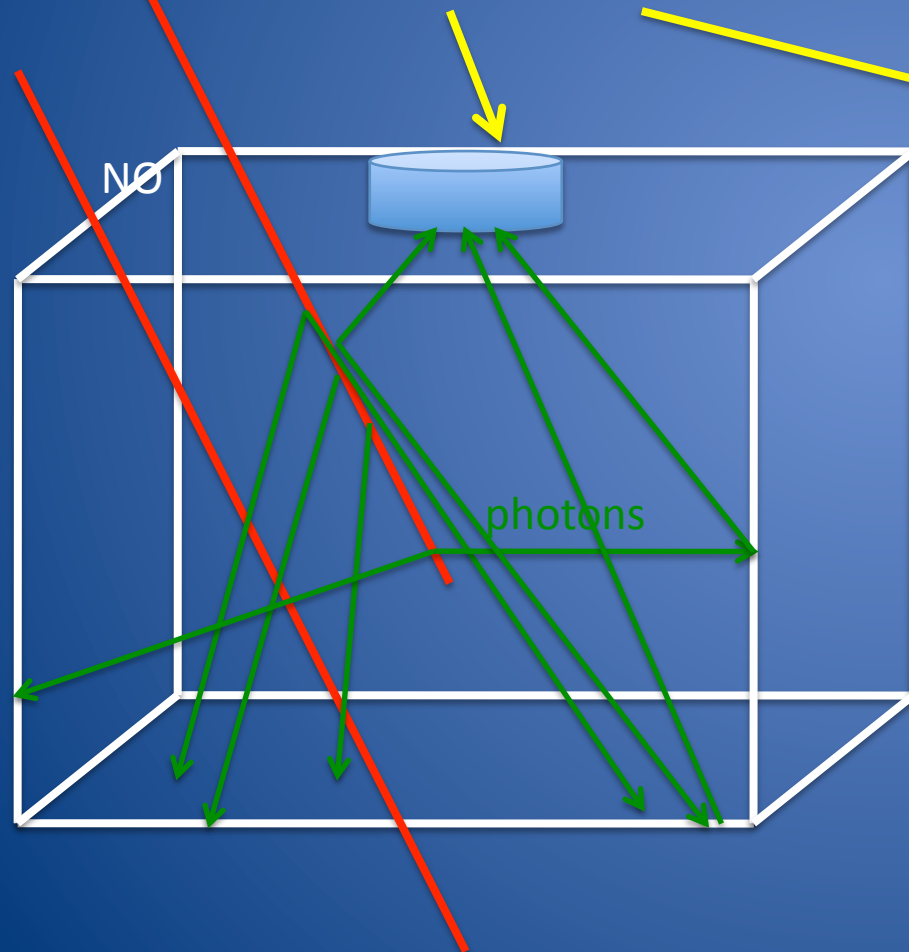


$$(dE/dx)_{\min} \approx 2 \text{ MeV}/(\text{g cm}^2)$$

A “calorimeter” light collection

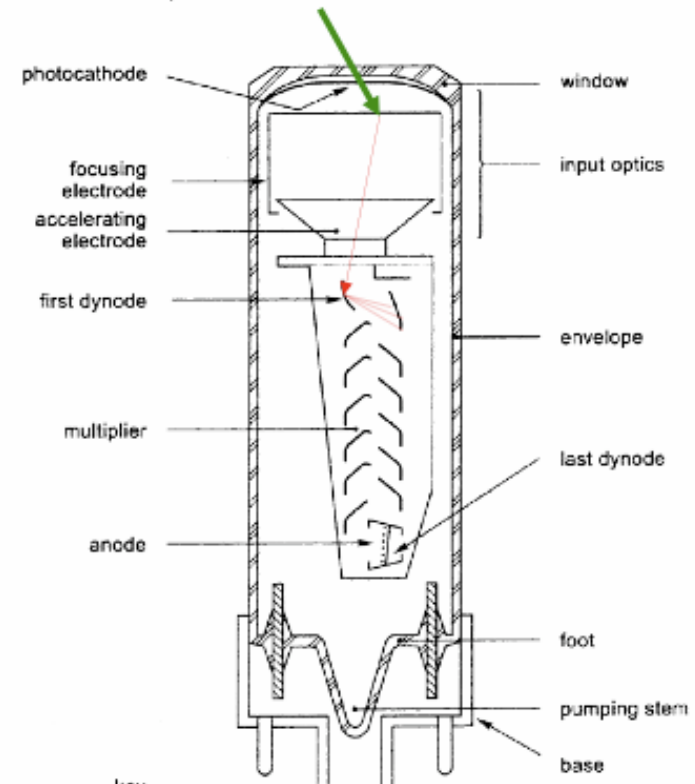
Incident particle

YES PHOTOMULTIPLIER



$$I \sim n_{ph} \sim E_{loss} \sim E_o$$

Most Common Device:
Photomultiplier Tube



Calorimetry

Energy is converted into “quantum” (photons, pairs...
scintillator, semiconductors...)

n of quantum detected = $n_0 \cdot \varepsilon$

n_0 = quantum produced = E_0 / E_q (E_q = energy necessary to produce a
quantum: 100 eV, 1 eV...)

ε = quantum recording efficiency

Resolution dominated by statistical fluctuations:

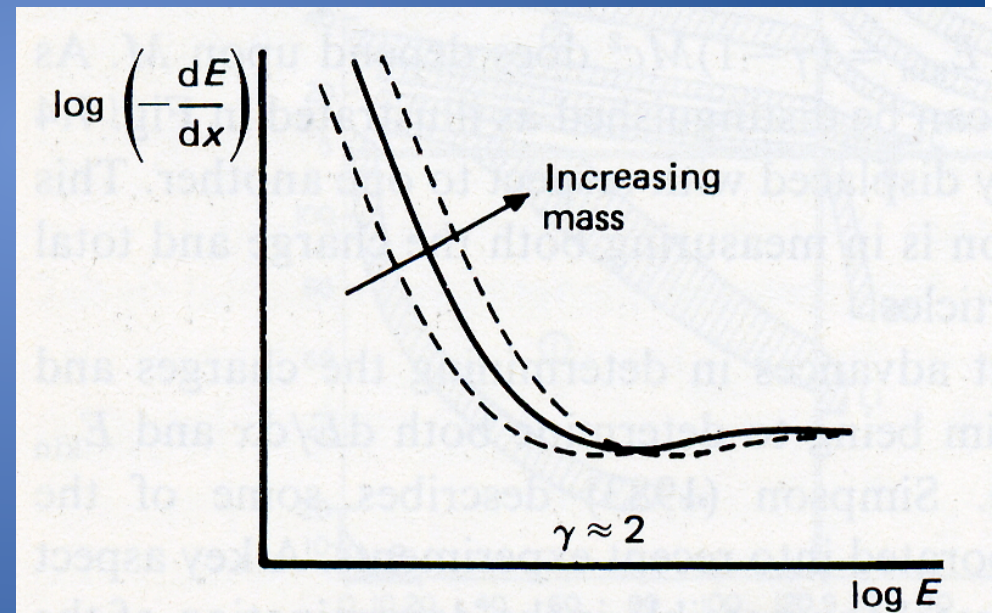
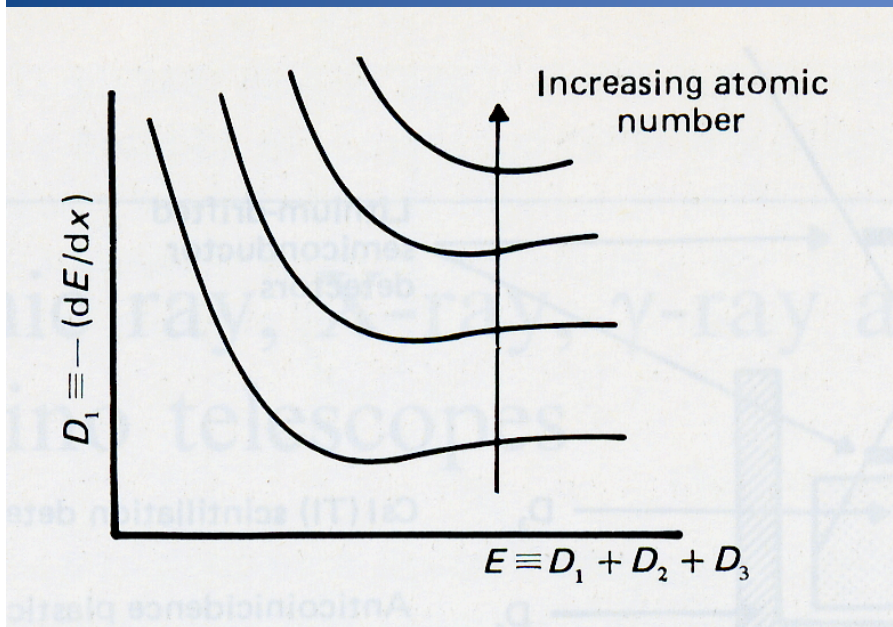
$$\sigma(E_0) / E_0 = \sqrt{n} / n = 1 / \sqrt{n} = \sqrt{E_q / \varepsilon E_0} = c / \sqrt{E_0}$$

$$\sigma(E_0) / E_0 \text{ about } " 10 \% / \sqrt{E_0 \text{ (MeV)}} "$$

With scintillator and ε about 1 % (ε assumed including e.g. q.e.)

Energy losses of charged particles & space detectors

$$-\frac{dE}{dx} = KZ^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 T_{\max}}{I^2} - \beta^2 - \frac{\delta}{2} \right]$$



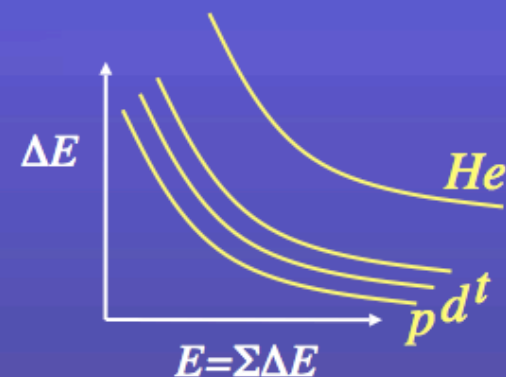
$$1/v^2 \approx M/E$$

CRIS: Cosmic Ray Isotope Spectrometer

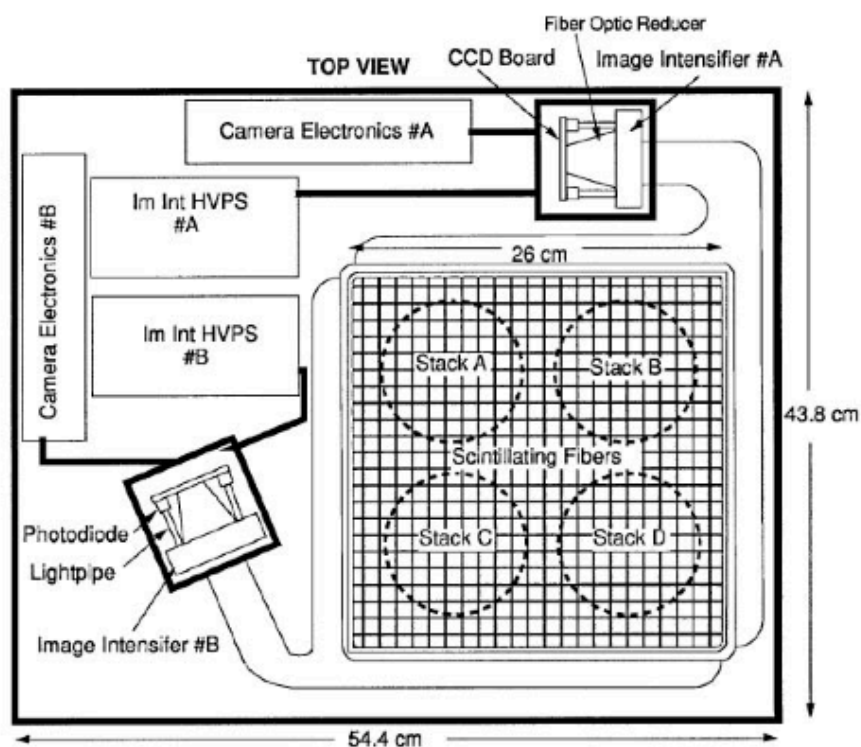
launched August '97

Bethe Bloch:

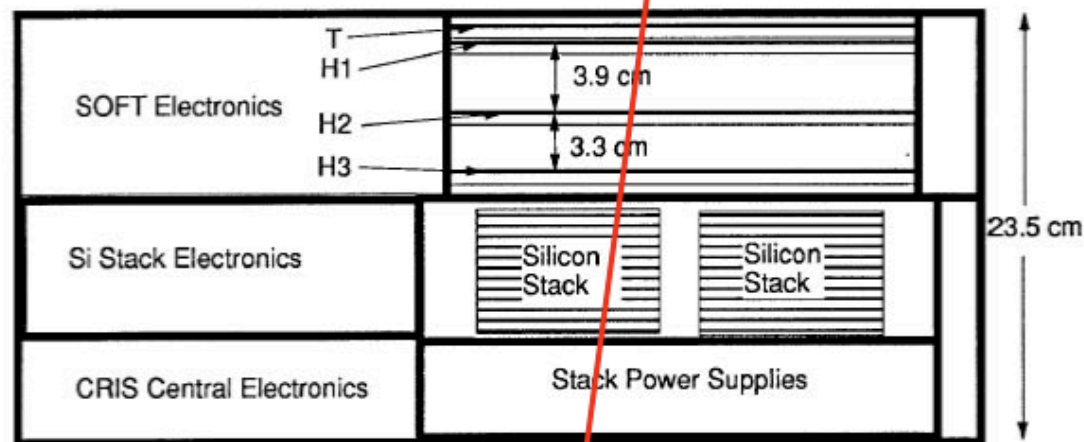
$$\Delta E \propto \frac{z^2 \cdot A}{E}$$



Multiple ΔE measurements
in each particle track

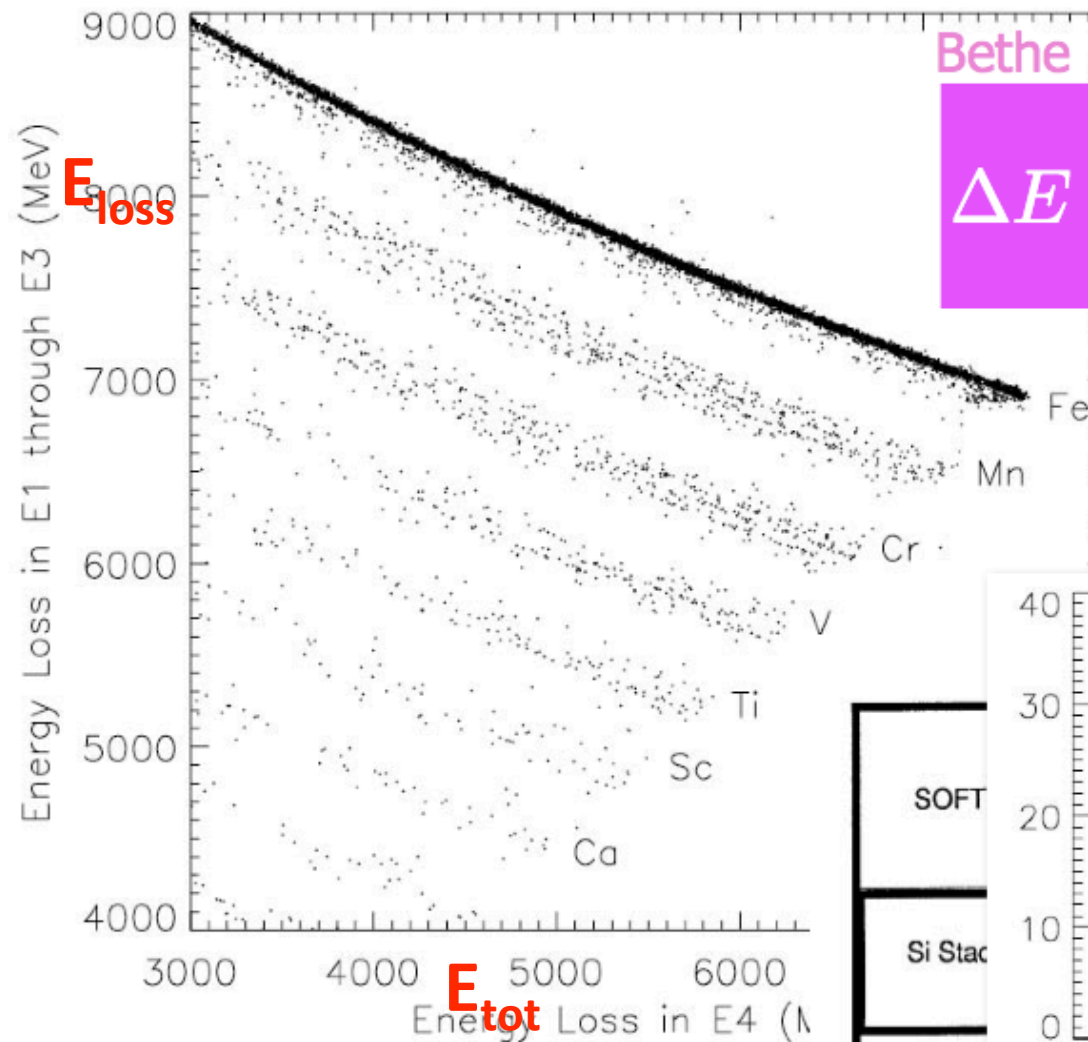


SIDE VIEW



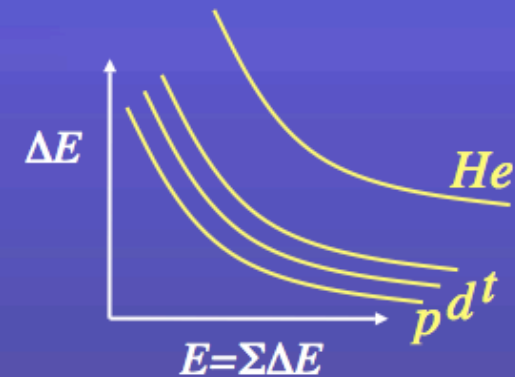
CRIS: Cosmic Ray Isotope Spectrometer

launched August '97

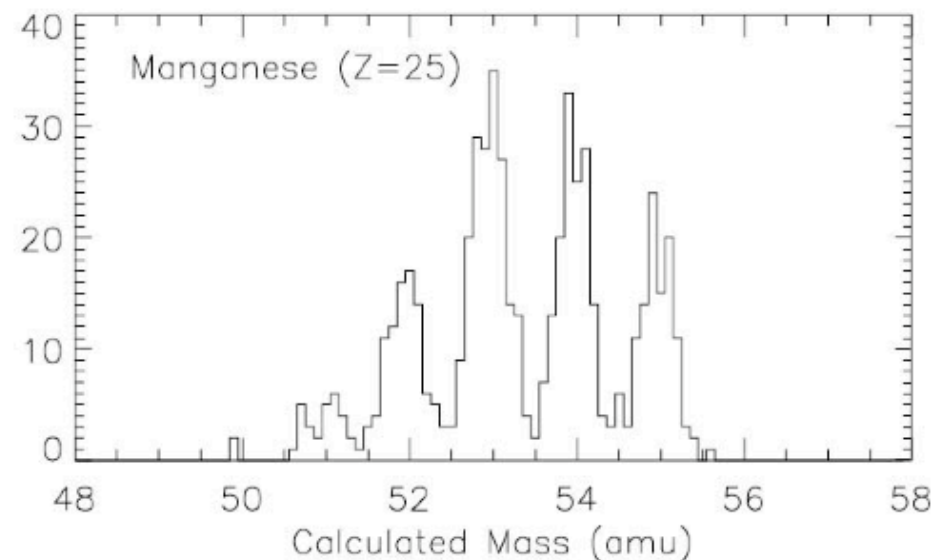


Bethe Bloch:

$$\Delta E \propto \frac{z^2 \cdot A}{E}$$

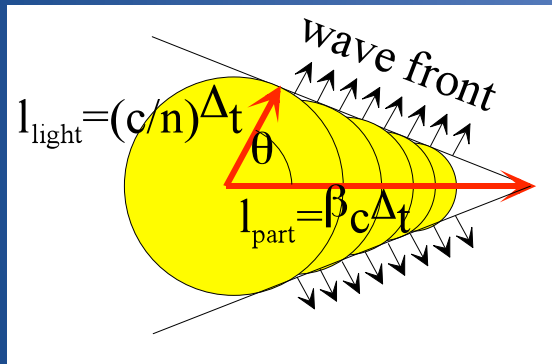


Multiple ΔE measurements
in each particle track

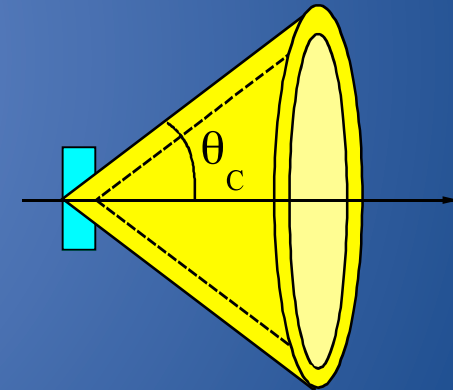


Cherenkov light

Qualitatively: polarization of the molecules of the medium due to the electric field of the incident particle. The depolarization composes destructively in all directions except, for ultra-relativistic particles, in direction θ_c (as for a shock wave from a supersonic aircraft)



$$\cos \theta_C = \frac{1}{n\beta} \quad \text{with } \underline{n = n(\lambda) \geq 1}$$



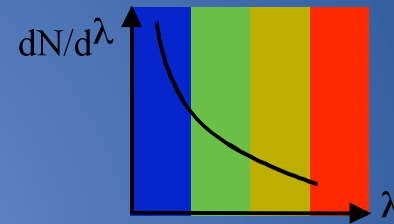
Threshold velocity: $\beta_s = 1/n \rightarrow \theta_c \sim 0$

Maximum angle: $\theta_{\max} = \arccos(1/n)$

Cherenkov light

Number of photons emitted
per path-length and
wavelength unit.

It decreases with increasing λ .



$$\frac{d^2 N}{dx d\lambda} = \frac{2\pi z^2 \alpha}{\lambda^2} \left(1 - \frac{1}{\beta^2 n^2} \right) = \frac{2\pi z^2 \alpha}{\lambda^2} \sin^2 \theta_C$$

$$\frac{d^2 N}{dx d\lambda} \propto \frac{1}{\lambda^2} \quad \text{with} \quad \lambda = \frac{c}{\nu} = \frac{hc}{E} \quad \frac{d^2 N}{dx dE} = \text{const.}$$

Cherenkov light

$$\frac{d^2 N}{dx d\lambda} = \frac{2\pi z^2 \alpha}{\lambda^2} \left(1 - \frac{1}{\beta^2 n^2} \right) = \frac{2\pi z^2 \alpha}{\lambda^2} \sin^2 \theta_C$$

$$\frac{d^2 N}{dx d\lambda} \propto \frac{1}{\lambda^2} \quad \text{with} \quad \lambda = \frac{c}{\nu} = \frac{hc}{E} \quad \frac{d^2 N}{dx dE} = \text{const.}$$

medium	n	$\theta_{\max} (\beta=1)$	$N_{\text{ph}} (\text{eV}^{-1} \text{ cm}^{-1})$
air	1.000283	1.36	0.208
isobutane	1.00127	2.89	0.941
water	1.33	41.2	160.8
quartz	1.46	46.7	196.4

deg

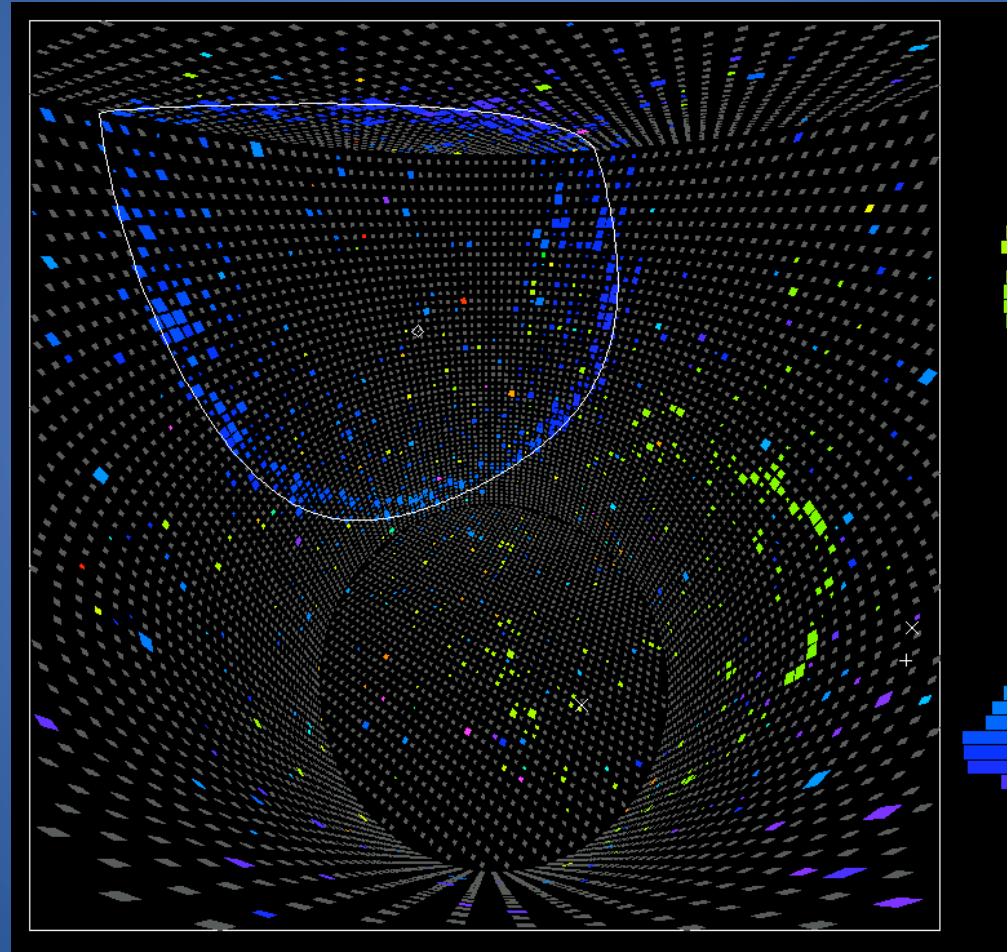
$$\Delta E = 1 \text{ eV} \sim \Delta \lambda = 300 - 600 \text{ nm}$$

Cherenkov light

- Energy measurement and particle identification
- Threshold
- Angle
- Particle identification:
 β + momentum (e.g. through a magnetic field)
=> mass

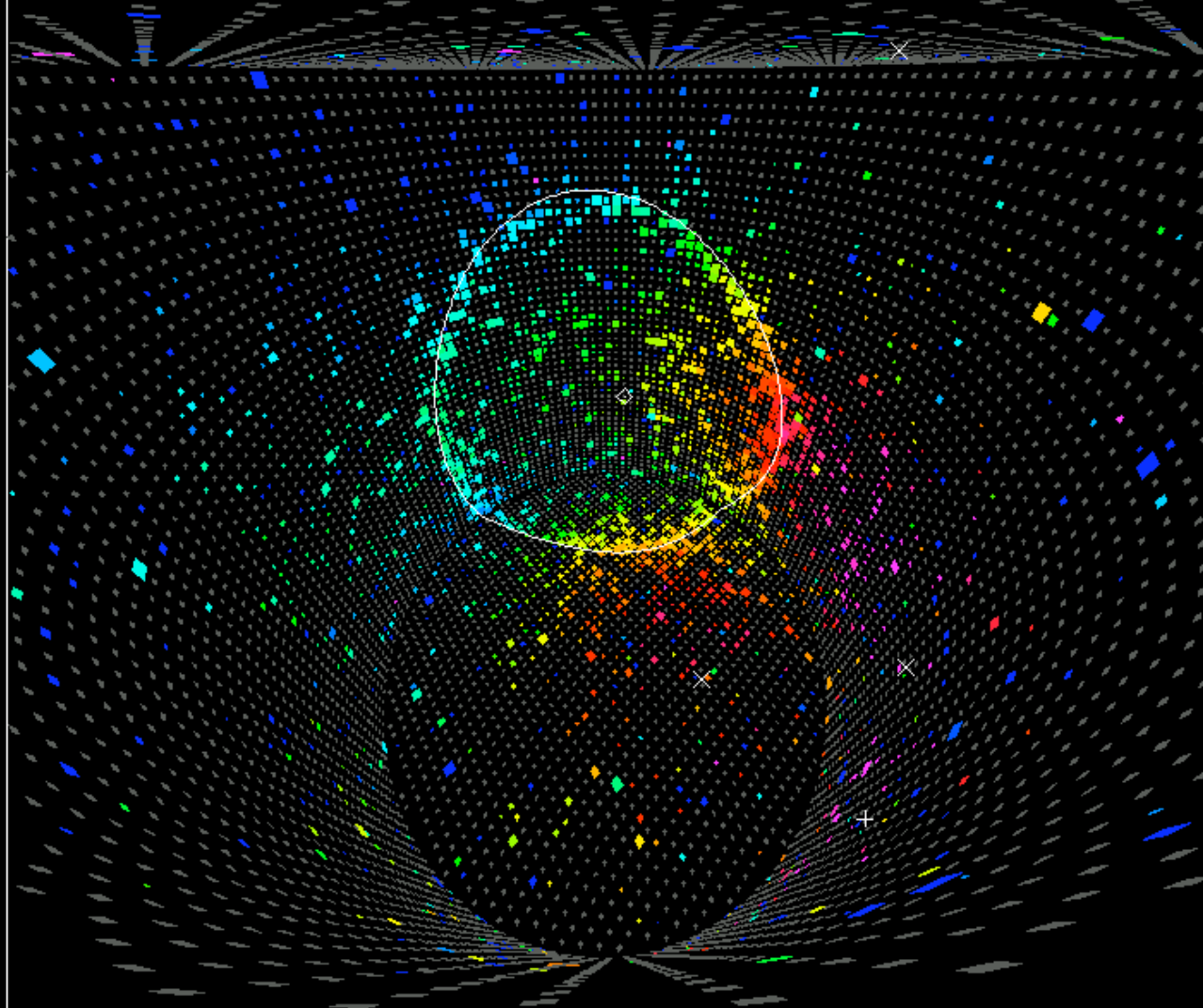
SUPER-KAMIOKANDE

CHERENKOV LIGHT MUON DECAY



SUPER-KAMIOKANDE: huge underground water Cherenkov light detector

SUPER-KAMIOKANDE CHERENKOV LIGHT E.M. SHOWER



Cherenkov detectors HEAO-C3

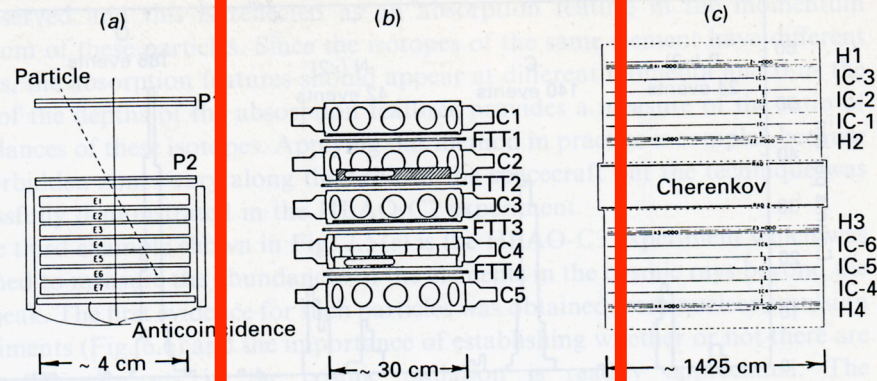


Figure 7.5. Examples of three important cosmic ray satellite telescopes. (From J. A. Simpson (1983). *Ann. Rev. Nucl. Part. Sci.*, **33**, 351.)
 (a) A schematic diagram of an ISEE class telescope for the measurement of cosmic ray isotopes. (W. E. Althouse, A. C. Cummings, T. L. Garrard, R. A. Meawaldt, E. C. Stone and R. E. Vogt (1978). *IEEE Trans. Geosci. Electron.*, **GE15**, 204.)
 (b) The HEAO-C2 experiment of the Danish-French collaboration. (Copenhagen-Saclay Collaboration (1981). *Adv. Space Res.*, **1**, 173.)
 (c) The HEAO-C3 ultraheavy nuclei telescope for the study of cosmic rays with $Z \geq 30$. (W. R. Binns, M. H. Israel, J. Klarmann, W. R. Scarlett, E. C. Stone and C. J. Waddington (1981). *Nucl. Instrum. Methods*, **185**, 415.)

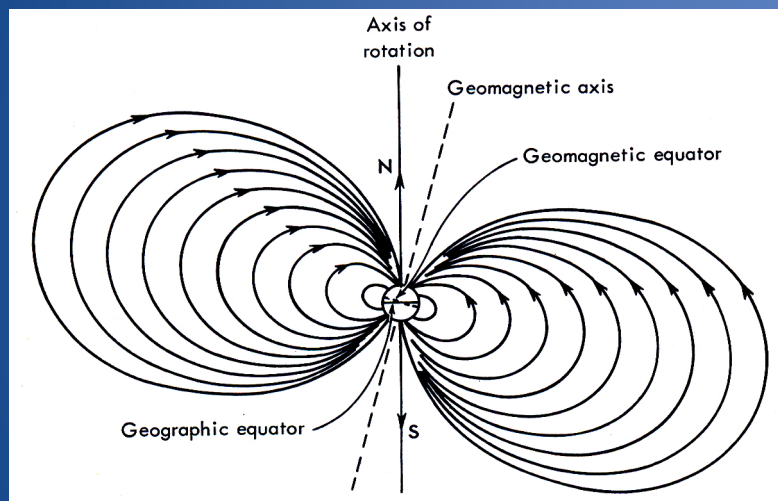
$$n = 1.015 \quad E_{th} = 4.5 \text{ GeV/nucleon}$$

1.33

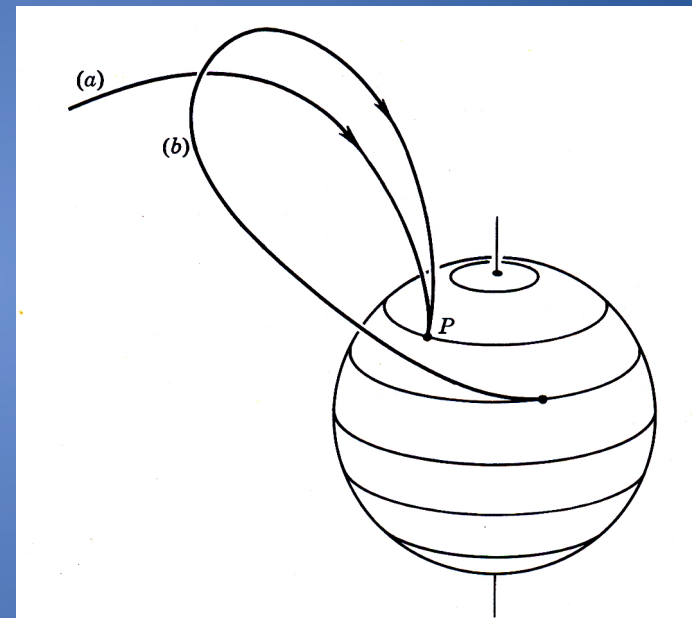
0.5

isotopes identification
 through arrival directions
 and selection of forbidden
 zones in the Earth magnetic
 field

Atmosphere as a magnetic analyzer

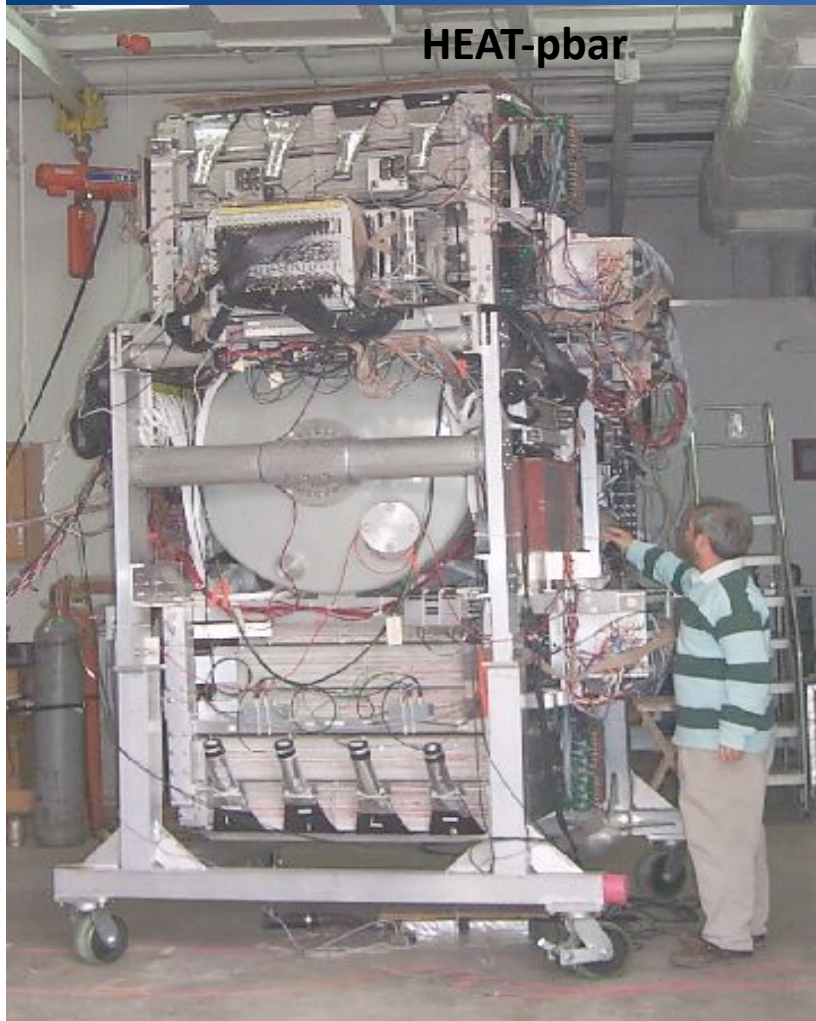


$$\gamma_E = 300 (Z_e/M_{eV}) B R_c \text{ [gauss cm]}$$

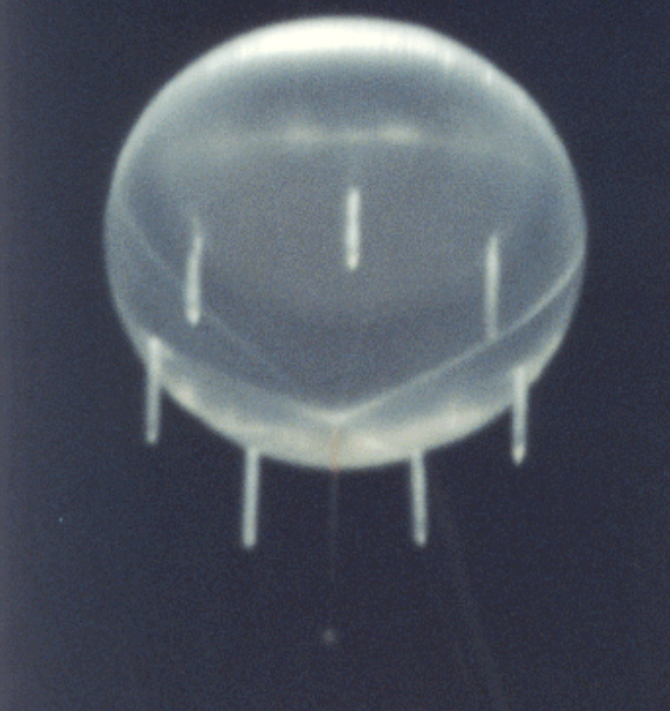


Allowed and forbidden directions
given γ (measured by CL)
depend on M

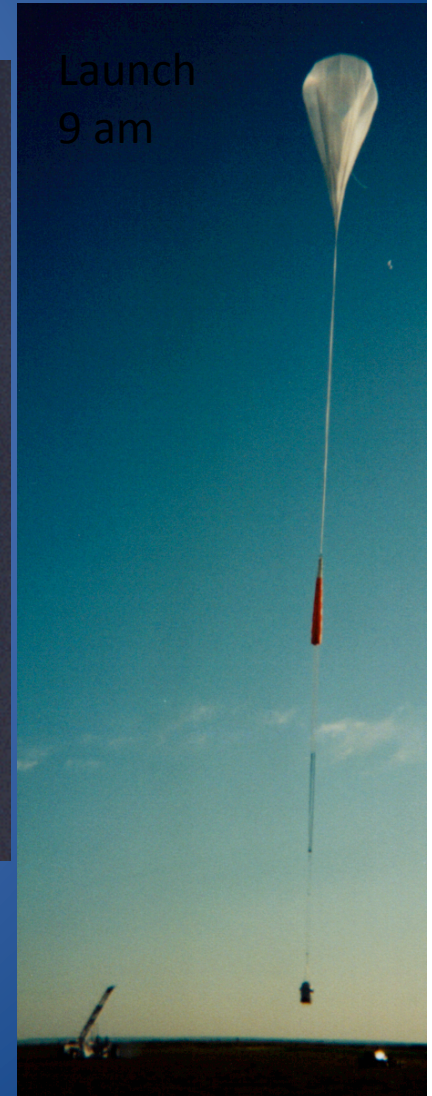
HEAT-pbar (High Energy Antimatter Telescope)



Balloon at float, 12 pm



Launch
9 am



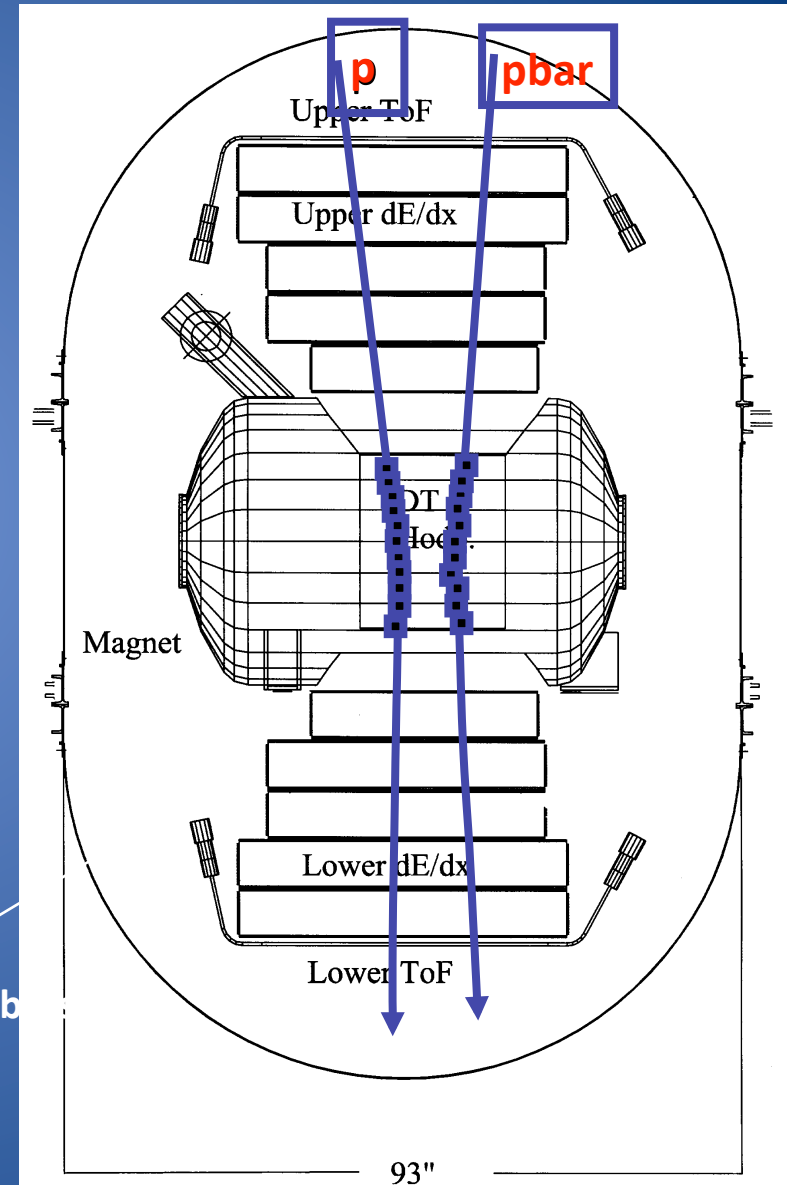
Superconducting Magnet Spectrometer with Drift Tube Hodoscope (DTH), Multiple Ionization (dE/dx) Detector and Time-of-Flight (TOF) system.

Identifying Antiprotons with HEAT-pbar

- DTH:
 - p from amount of bending in B=1T
 - Sign of Z from direction

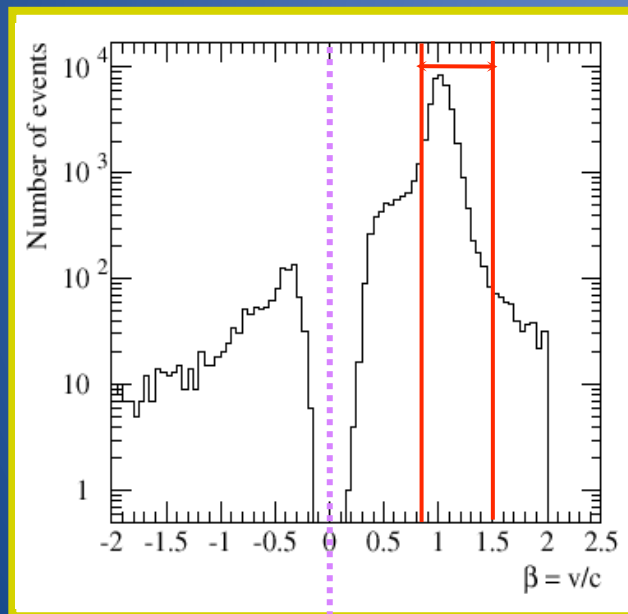
$$R = pc/Z_e,$$
$$R_{\text{max}} = 170 \text{ GV}$$

Xe Gaz chamb



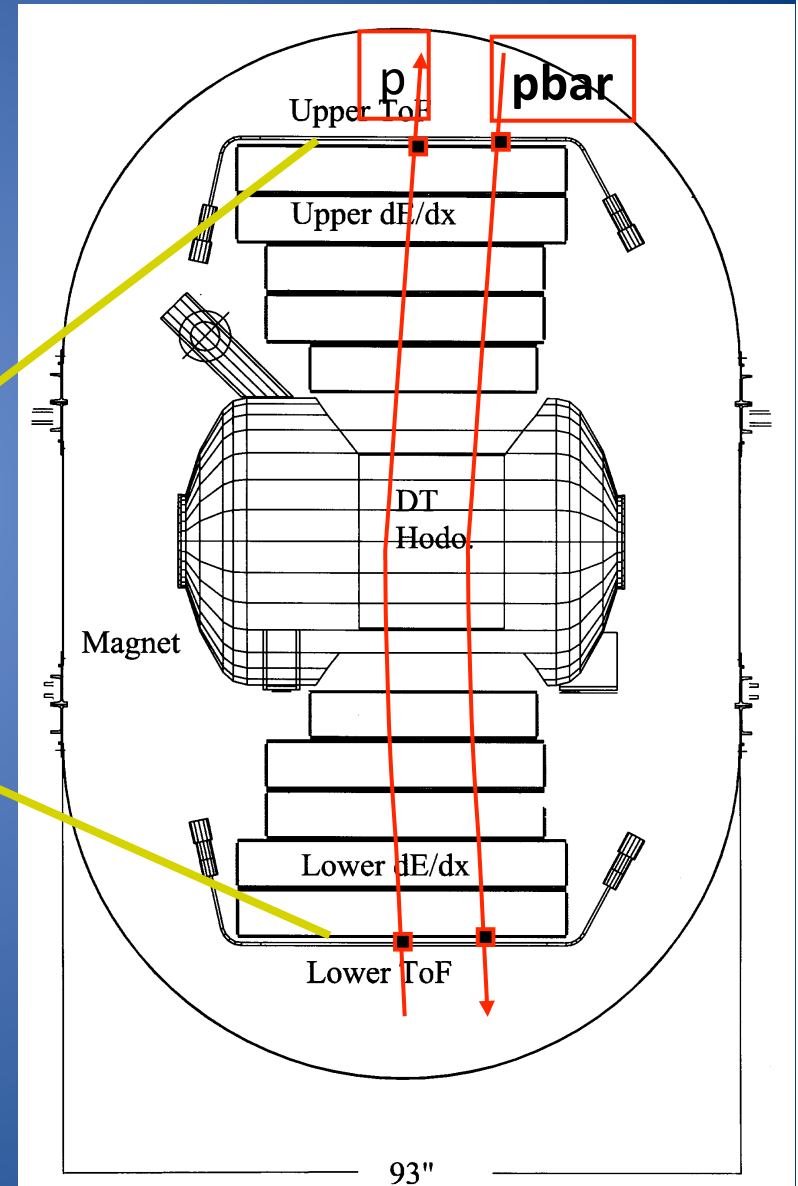
Identifying Antiprotons with HEAT-pbar

Up going proton (albedo particles) looks like down going antiproton -> Need to know start and stop in the time-of-flight



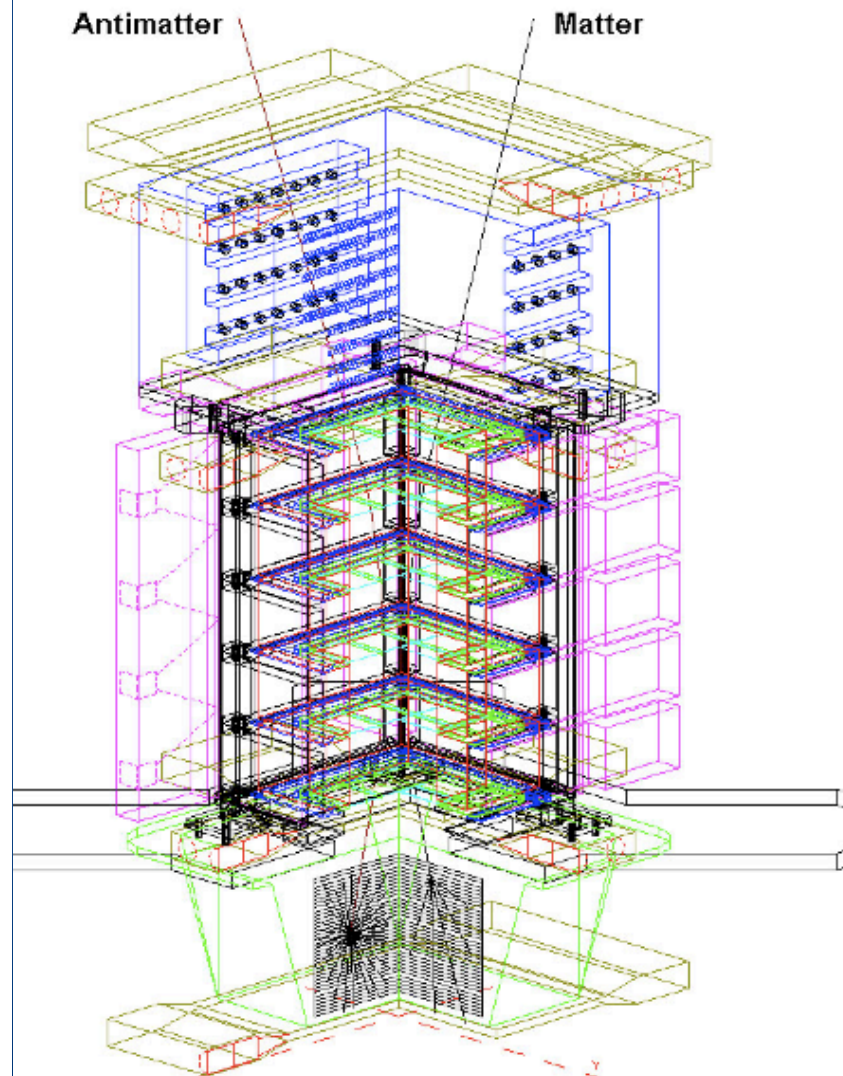
Upgoing

Downgoing

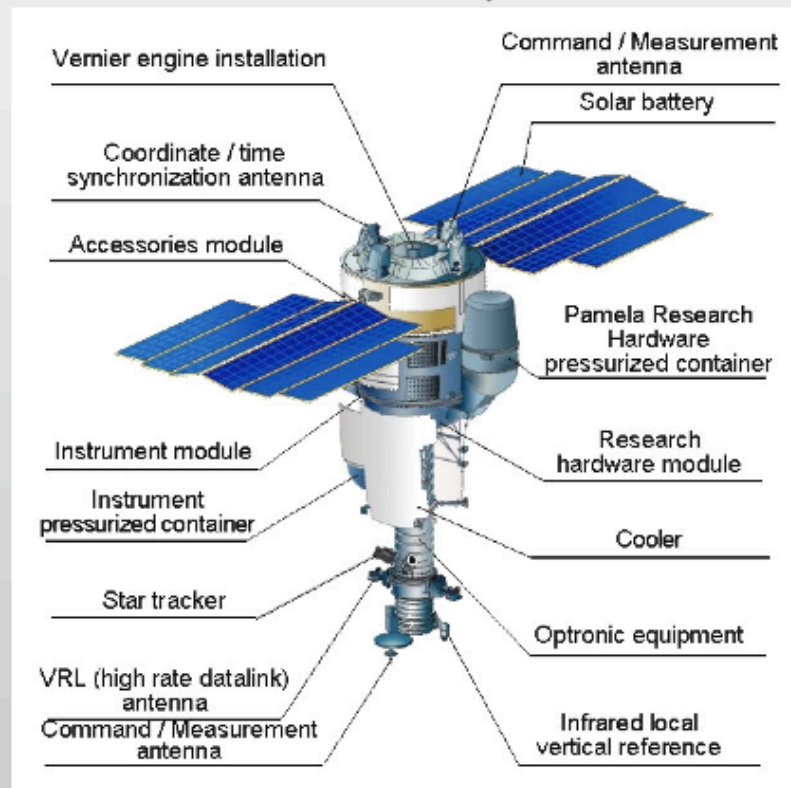


New major project: PAMELA

Pamela (launched 15.5.06)



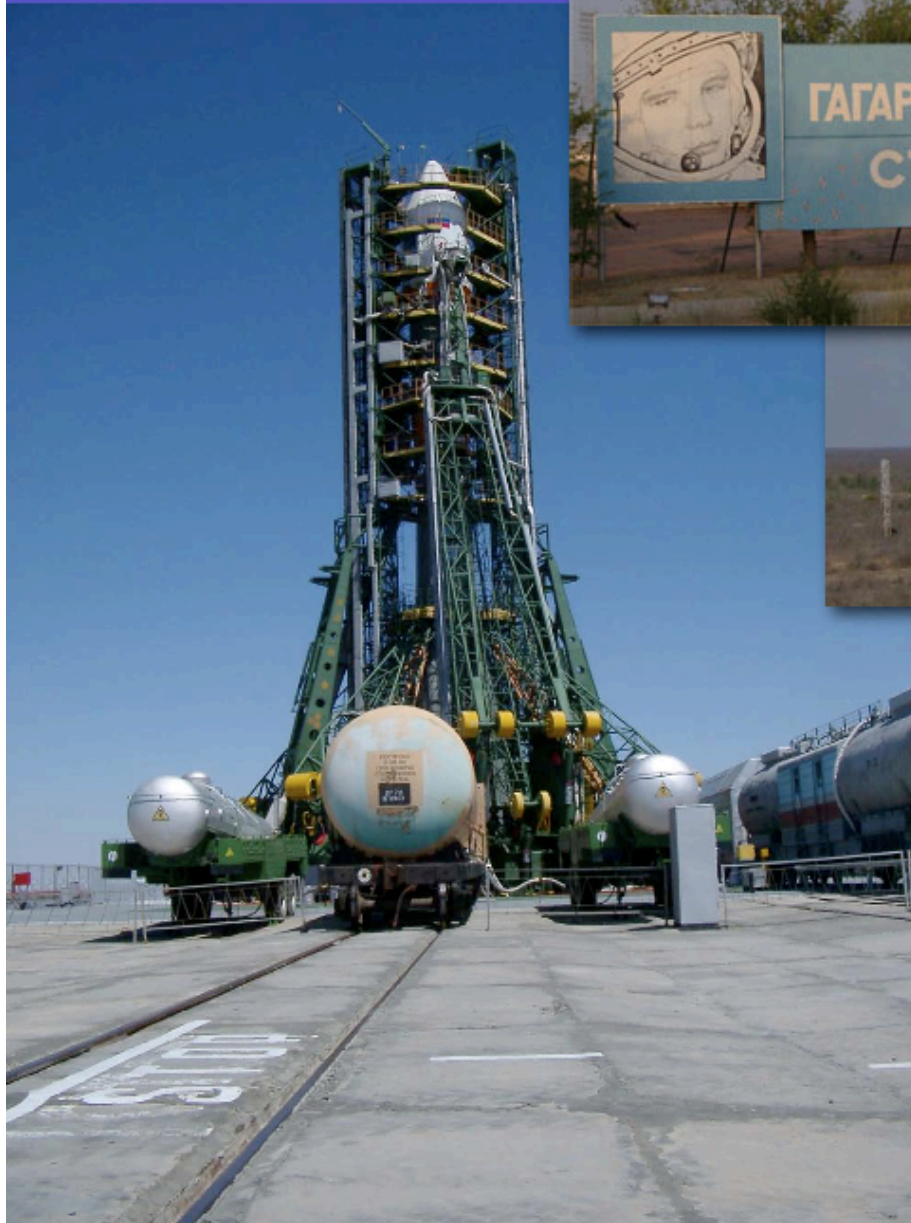
on-board of the Russian
Resurs-DK1 satellite by Sojuz rocket



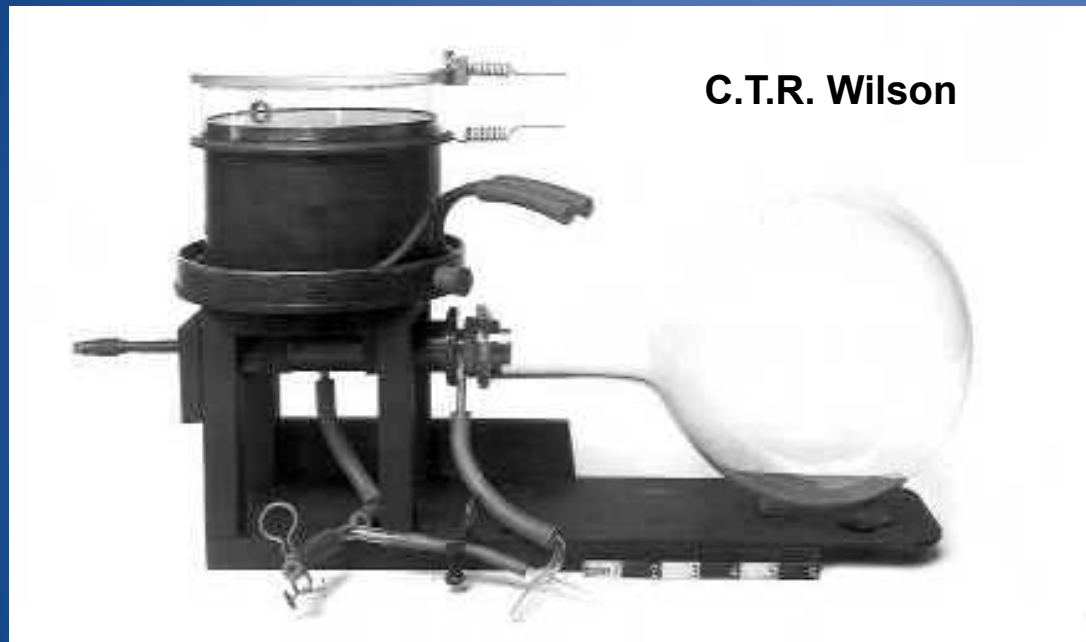
magnetic spectrometer with permanent
0.4 T magnet, calorimeter, ToF, TRD, ...
size: 120 x 40 x 45 cm³; weight: 380 kg

PAMELA launch

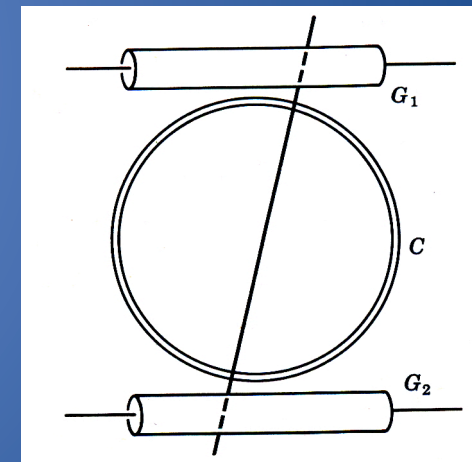
15.5.2006, Baikonur Cosmodrome
Kazakhstan



The cloud chamber (1929)



Controlled by external trigger



Fast expansion of gas in order
to decrease temperature and
over saturate the gas
Condensation of drops on the
ions produced by particles

Blacket Occhialini (1931 - 1934)

E.M. SHOWERS

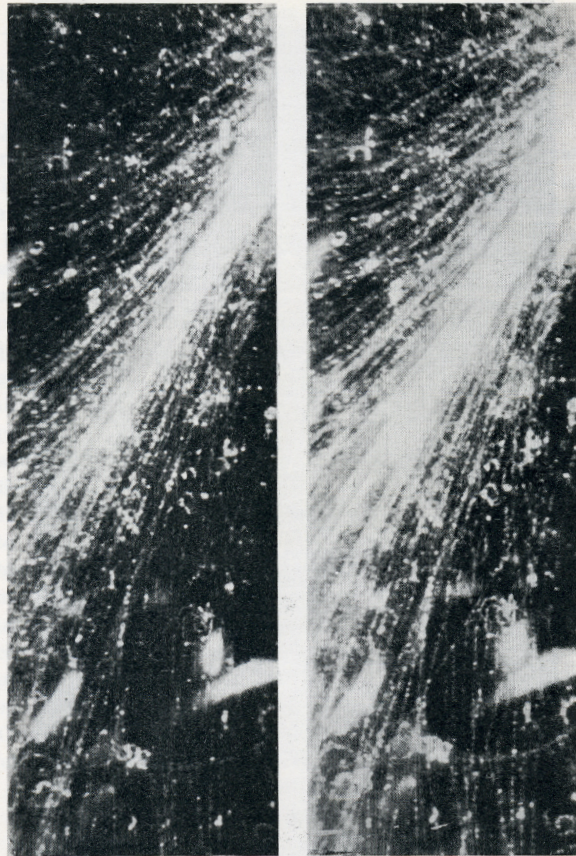


Fig. 45. Grande sciame di elettroni fotografato al laboratorio di Largentière con una camera di 80 cm di altezza e un campo magnetico di 3500 gauss. Vi si vedono un gran numero di elettroni, tanto positivi che negativi, di grande energia: nella parte inferiore della foto si vedono traiettorie circolari di elettroni, secondari di fotoni, prodotti anch'essi nello sciame.

LHERITIER, PEYROU, LAGARRIGUE, laboratorio di Largentière La Bessée (Delfinato).

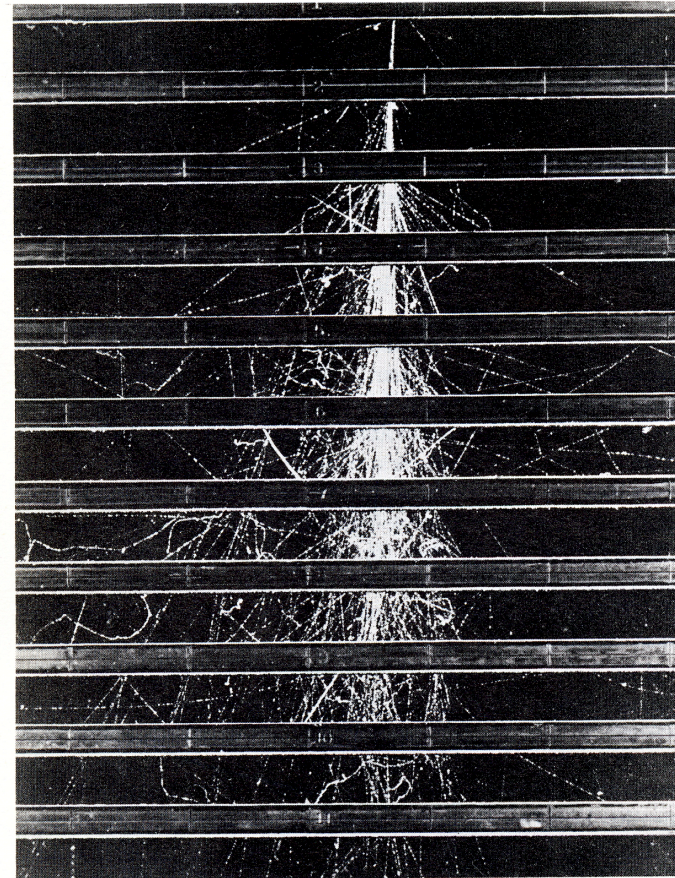


Fig. 7-5 A shower developing through a number of brass plates 1.25 cm thick placed across a cloud chamber. The shower was initiated in the top plate by an incident high-energy electron or photon. The photograph was taken by the MIT cosmic-ray group.

CALORIMETRY AT HIGH ENERGIES: E.M. CASCADES

$$N(t) = 2^t$$

$$\underline{E}(t) = E_0 / N(t) = E_0 / 2^t$$

up to $\underline{E}(t) > E_{\text{crit}}$

for $\underline{E}(t) < E_{\text{crit}}$

→ $-(dE/dx)_{\text{coll}}$ → absorption

$\underline{E}(t) = E_{\text{crit}}$ → $N(t)$ MAX

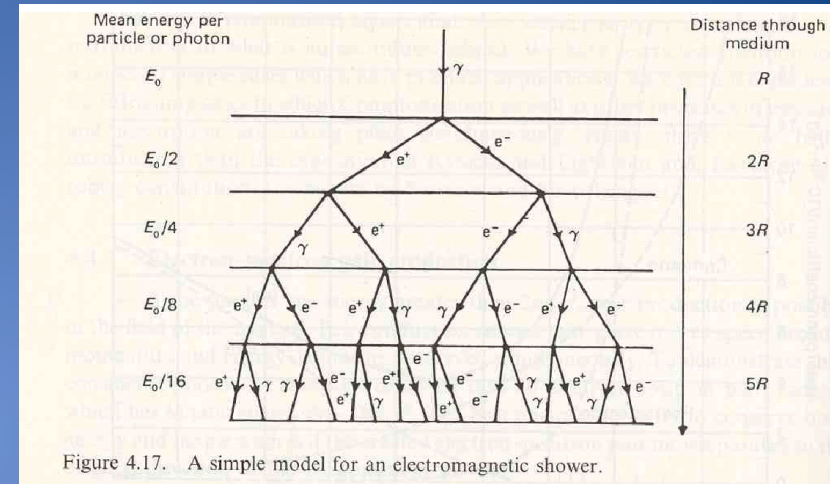
$$E_{\text{crit}} = E_0 / 2^{t_{\text{max}}} \quad 2^{t_{\text{max}}} = E_0 / E_{\text{crit}}$$

$$t_{\text{max}} = \ln(E_0 / E_{\text{crit}}) / \ln 2$$

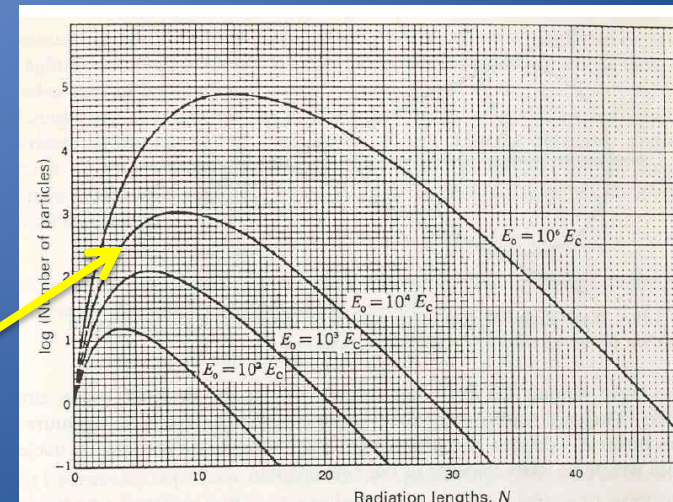
$$N(t) \text{ MAX} = N(t_{\text{max}}) = E_0 / E_{\text{crit}}$$



π^+, π^- decay, interact



1 c.u. in carbon ~ 15 cm



MEASUREMENT OF INTEGRAL BELOW THE RELATIVE CURVE

WHEN ENERGY CANNOT BE CONTAINED: SAMPLING CALORIMETERS

$$\Delta E \sim X_{\text{sens}}/X_{\text{tot}}$$

$$\sim 10 / 500 \sim 2 \%$$

sens:

$$1 \text{ l r} \sim 15 \text{ cm}$$

$$\rho \sim 1 \text{ gr cm}^{-3}$$

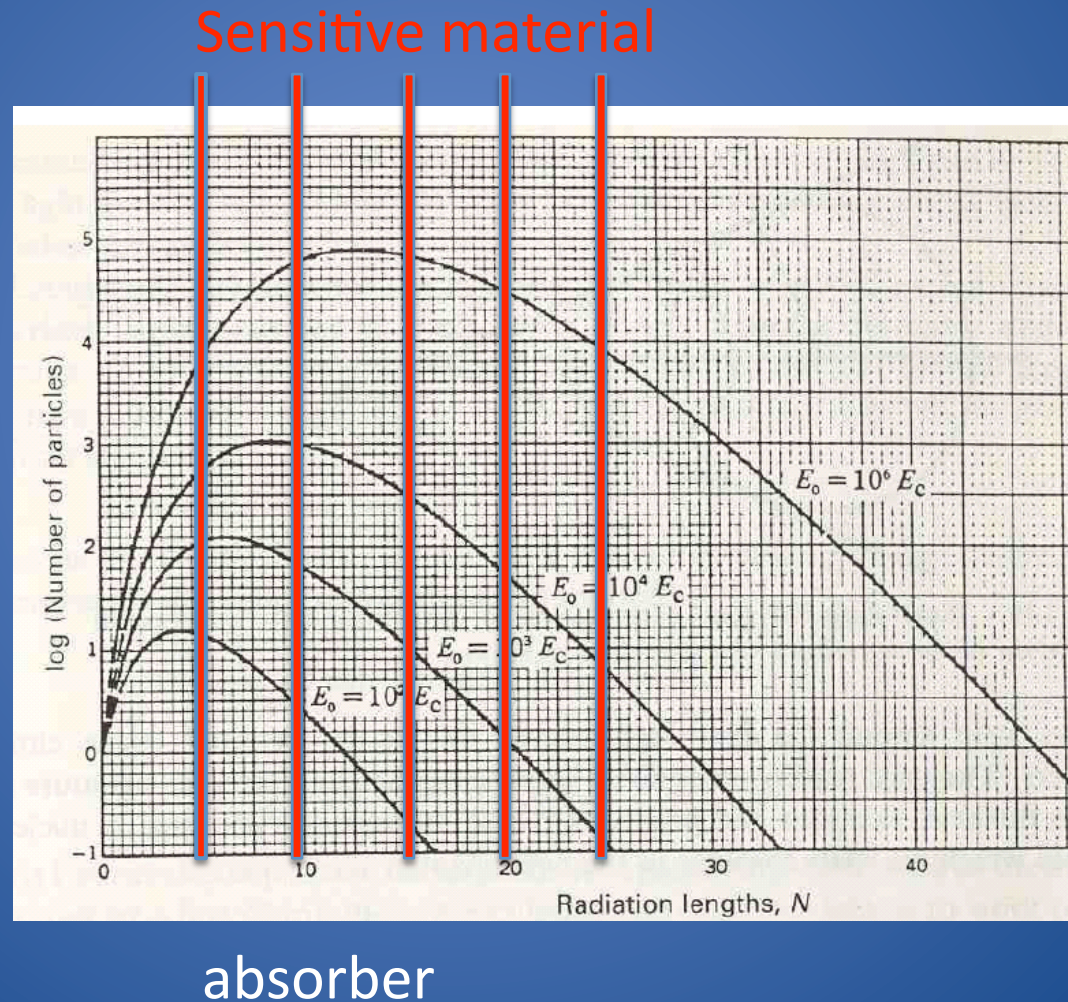
$$E_c \sim 80 \text{ MeV}$$

abs:

$$1 \text{ l r} \sim 1.8 \text{ cm}$$

$$\rho \sim 7.8 \text{ gr cm}^{-3}$$

$$E_c \sim 20.7 \text{ MeV}$$



Sampling Calorimeters

Grigorov PROTON satellites

Fe 5 cm
Scint. 1.5 cm

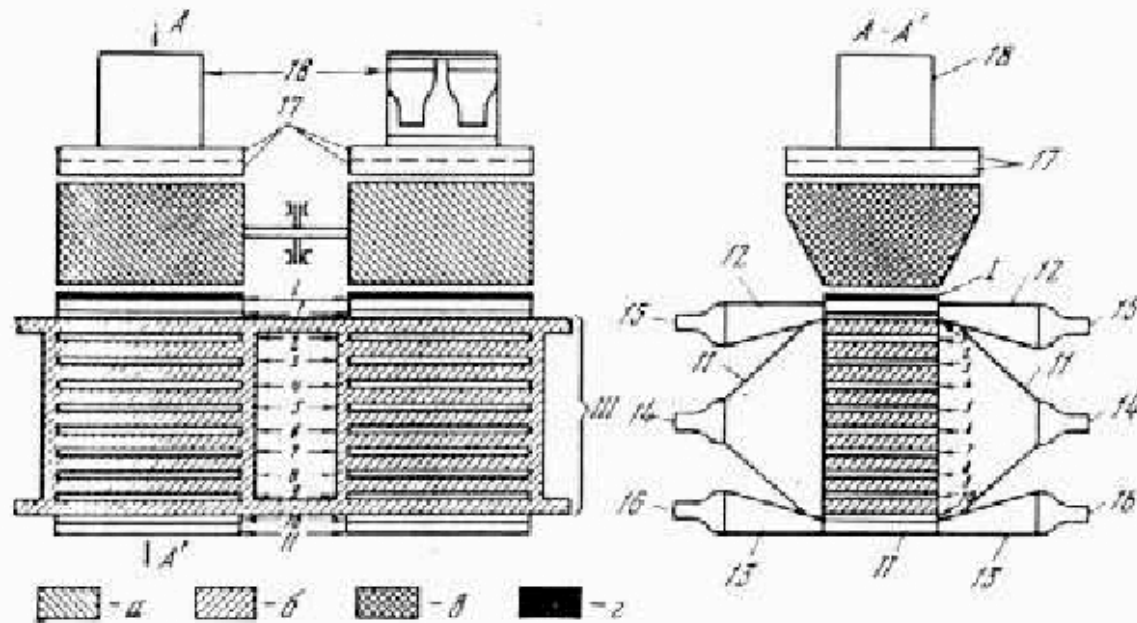


Рис. 9.1 Схематическое изображение прибора СЗЗ-14 I - детектор взаимодействия; II - нижний сцинтилляционный счетчик; III - ионизационный калориметр; 1-10 - сцинтилляционные детекторы энергии; 11 - диффузор детектора энергии; 12 - диффузор детектора взаимодействия; 14-16 - фотоумножители; 17 - детектор заряда (сдвоенный пропорциональный счетчик); 18 - детектор направления; а - полиэтилен; б - железо; в - углерод; г - свинец.

PROTON satellite calorimeter

Telescope:

$$A = 1369 \text{ cm}^2$$

Cerenkov counters (plexiglas):

Direction + charge

Calorimeter depth (Fe) 386 g/cm^2

$$X \sim 2.7 \lambda_{\text{Fe}}$$

50% energy released

$$E_o = 3.9 \cdot E_r \quad (E_{\text{th}} = 29 \text{ GeV})$$

$$\Delta E/E \sim 20 \% \text{ (my estimate)}$$

Fraction of energy released inside the calorimeter

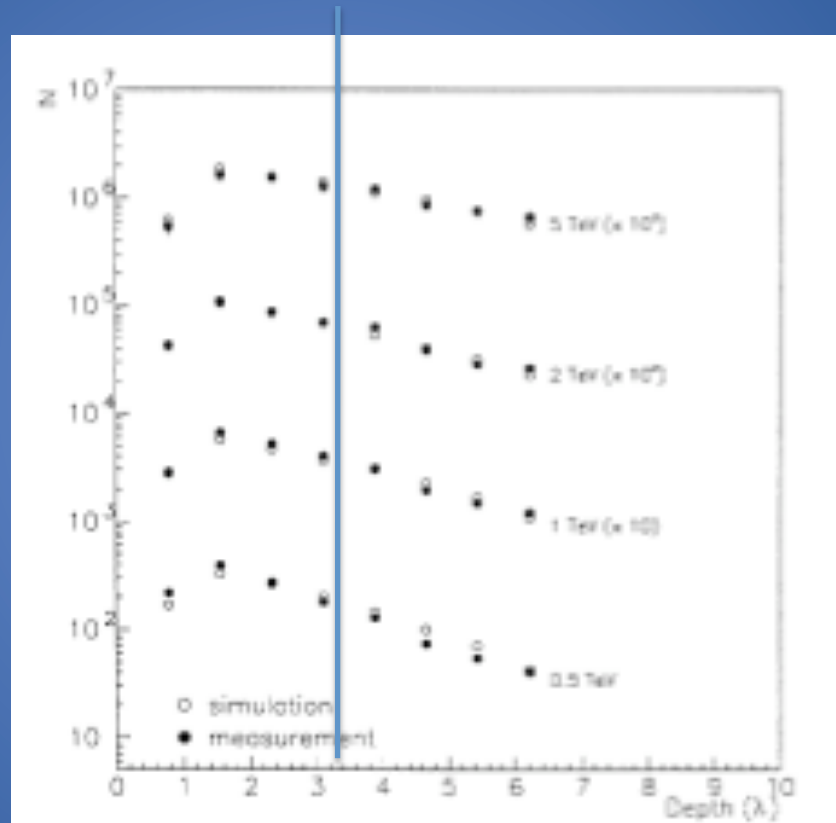


Fig. 11. Mean longitudinal transition curves for hadrons. Full circles: data; empty circles: simulation.

Hadron sampling calorimeters: resolutions

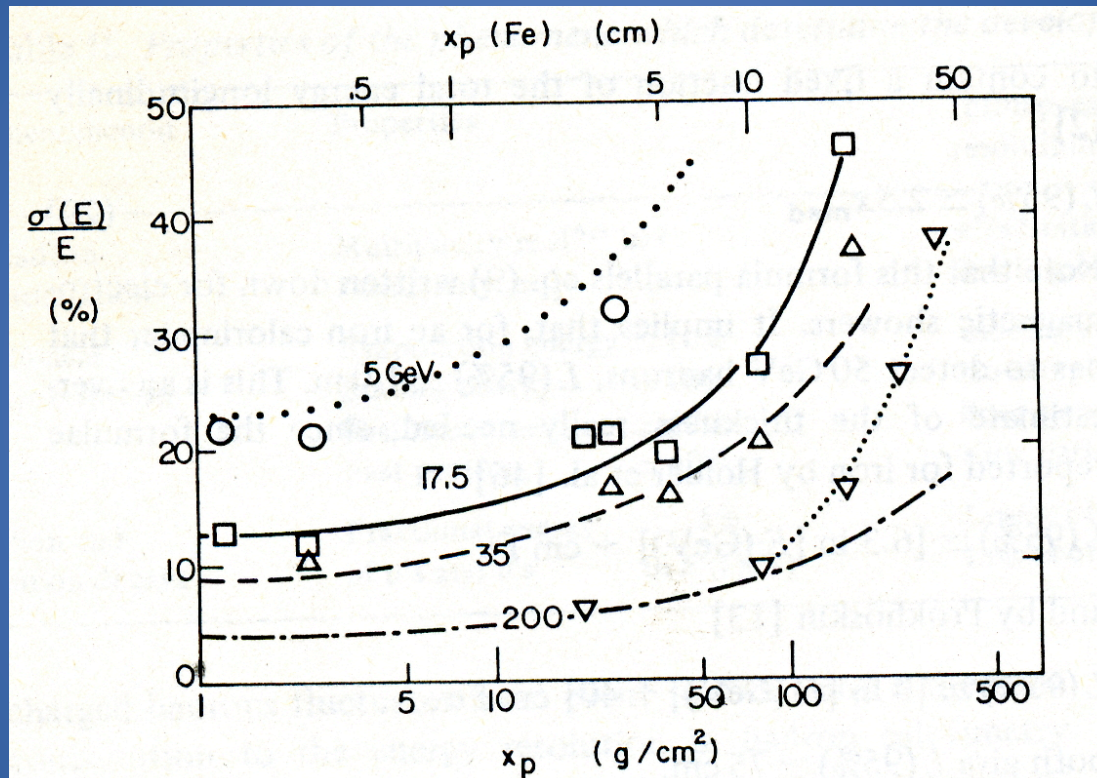


Fig. 19. The compilation by Iwata [2] on the thickness dependence of measured hadronic resolutions at 5, 17.5, 35 and 200 GeV is compared with the predictions of eq. (31) with $R' = 40\%$. The agreement is good below $\sim 100 \text{ g cm}^{-2}$. The dotted curve above 100 g cm^{-2} is only meant to guide the eye through the 200 GeV points.

EMULSION CHAMBERS ON BALLOONS

JACEE

RUNJOB



Nuclear Emulsions

high energy
hadronic interaction
production of secondaries
fragmentation regions

==> elementary particle
and high energy physics

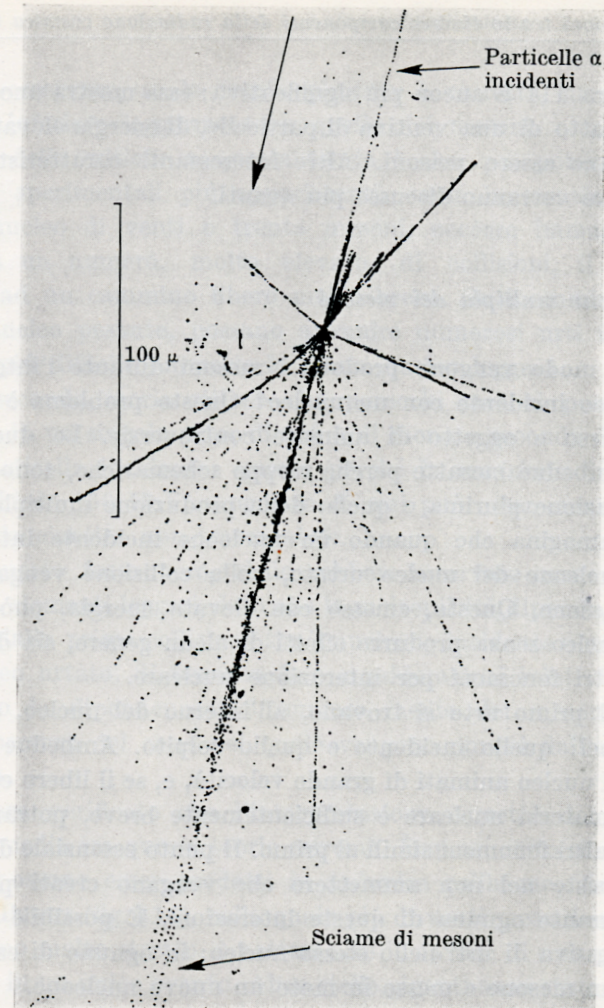


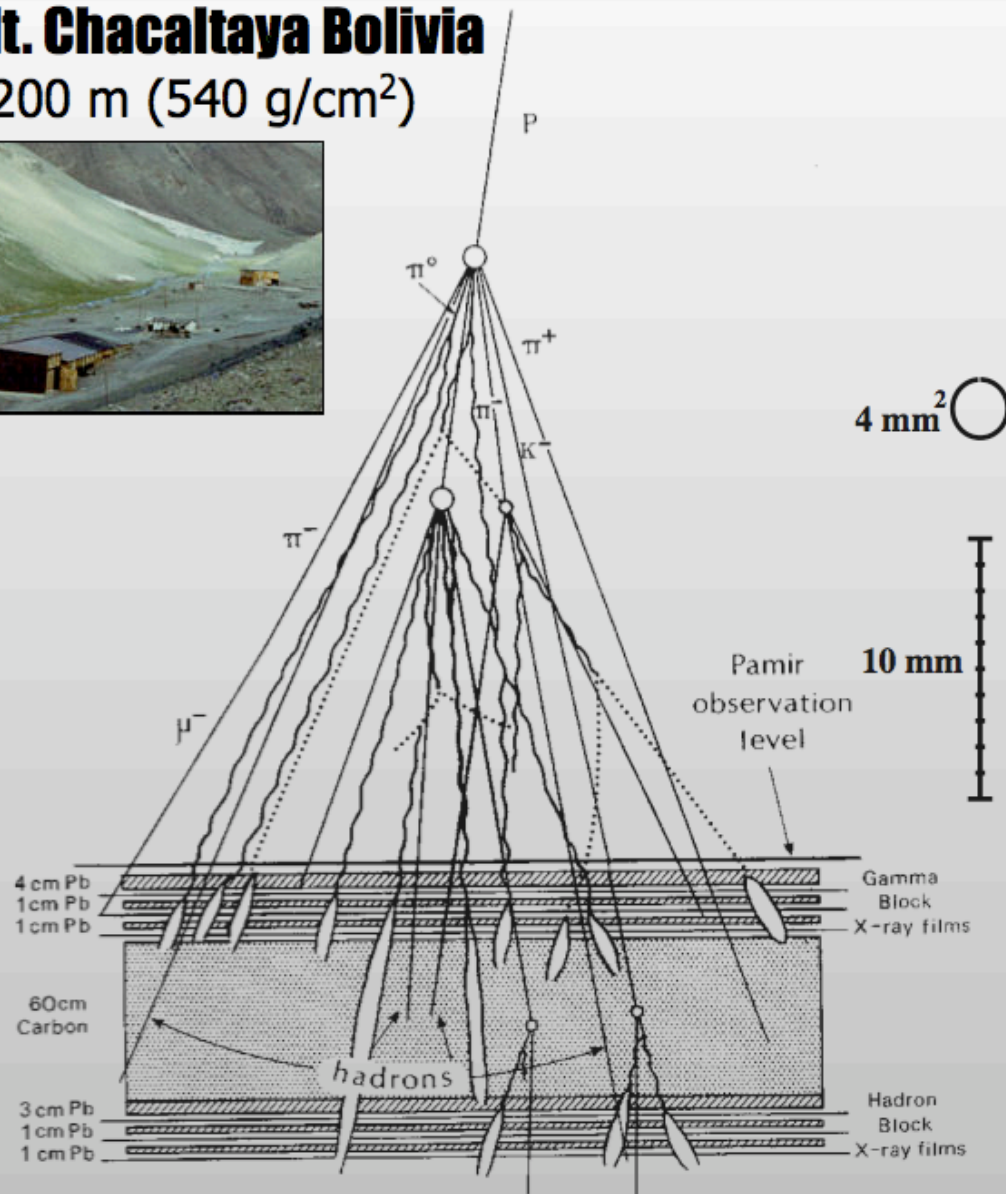
Fig. 114. Ecco uno dei fenomeni nucleari più straordinari osservati fino ad ora. Una particella α dotata di energia superiore a 1000 GeV produce uno sciame formato da due gruppi di particelle relativistiche: il primo di 33 particelle più allargato, e il secondo, di 23, emesso in un cono di piccolissima apertura. Le particelle di questo gruppo sono così collimate che non è possibile distinguerne le traiettorie nell'immediata vicinanza della stella. A una certa distanza dal centro di questa si osservano delle coppie di elettroni all'interno del cono di emissione: si tratta di coppie prodotte dai fotoni di disintegrazione dei mesoni neutri emessi insieme a quelli carichi nello sciame collimato.

KAPLON, PETERS e BRADT, « Phys. Rev. », **76**, 1735 (1949).

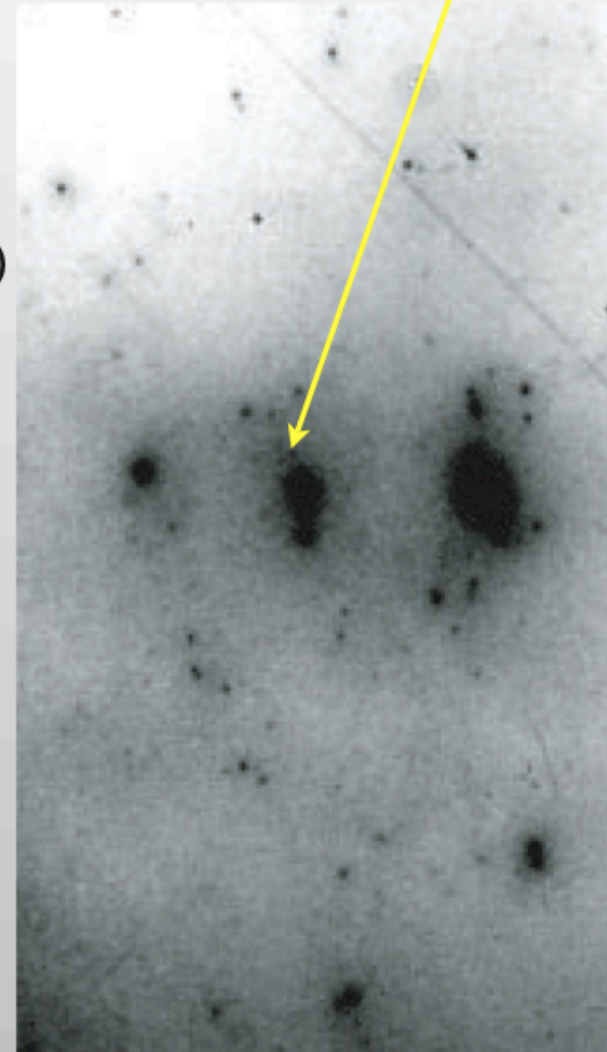
Emulsions at high Altitude

Mt. Chacaltaya Bolivia

5200 m (540 g/cm²)

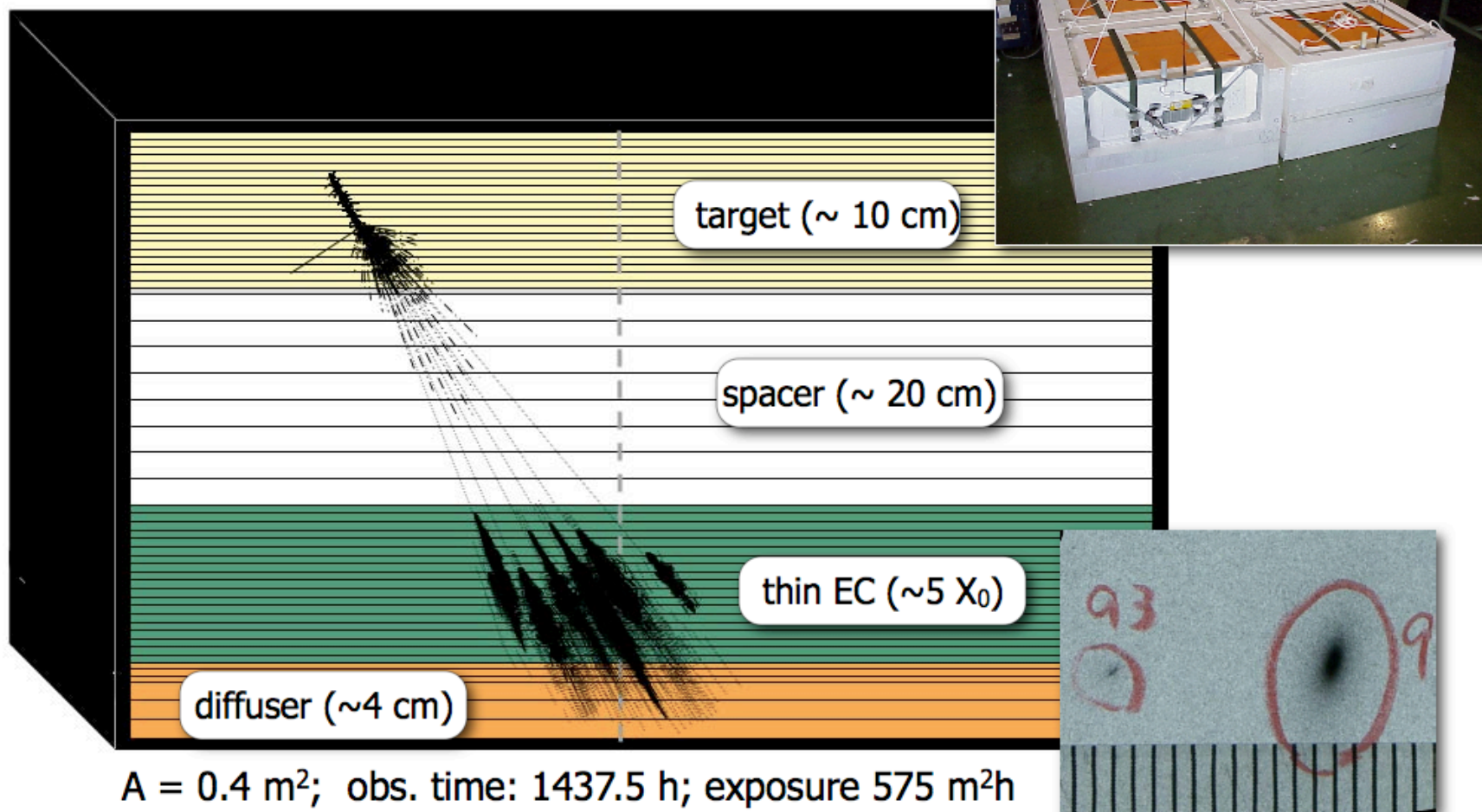


TeV γ -rays

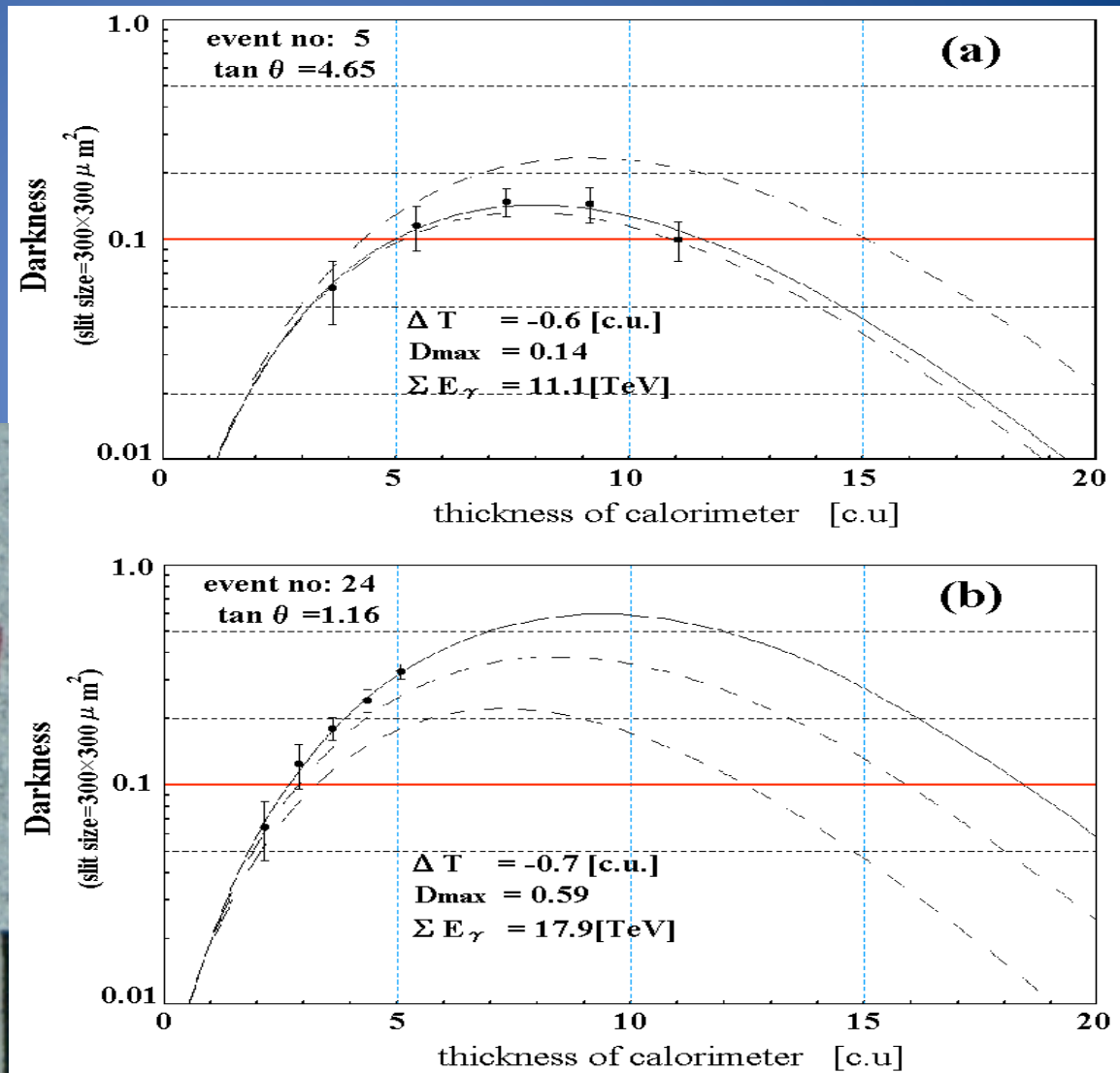
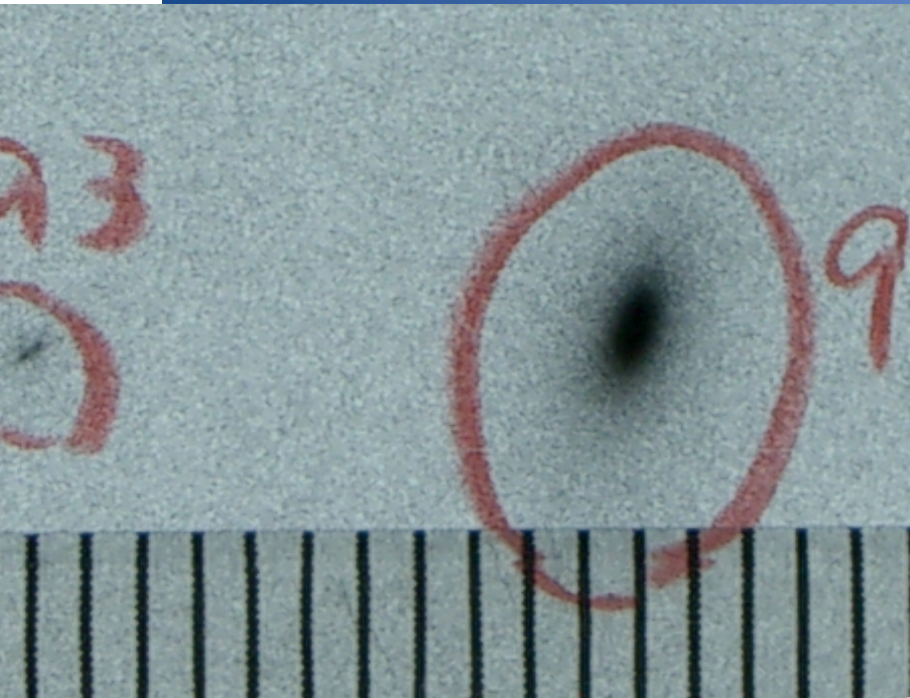


RUNJOB – Emulsion Chamber on Balloon

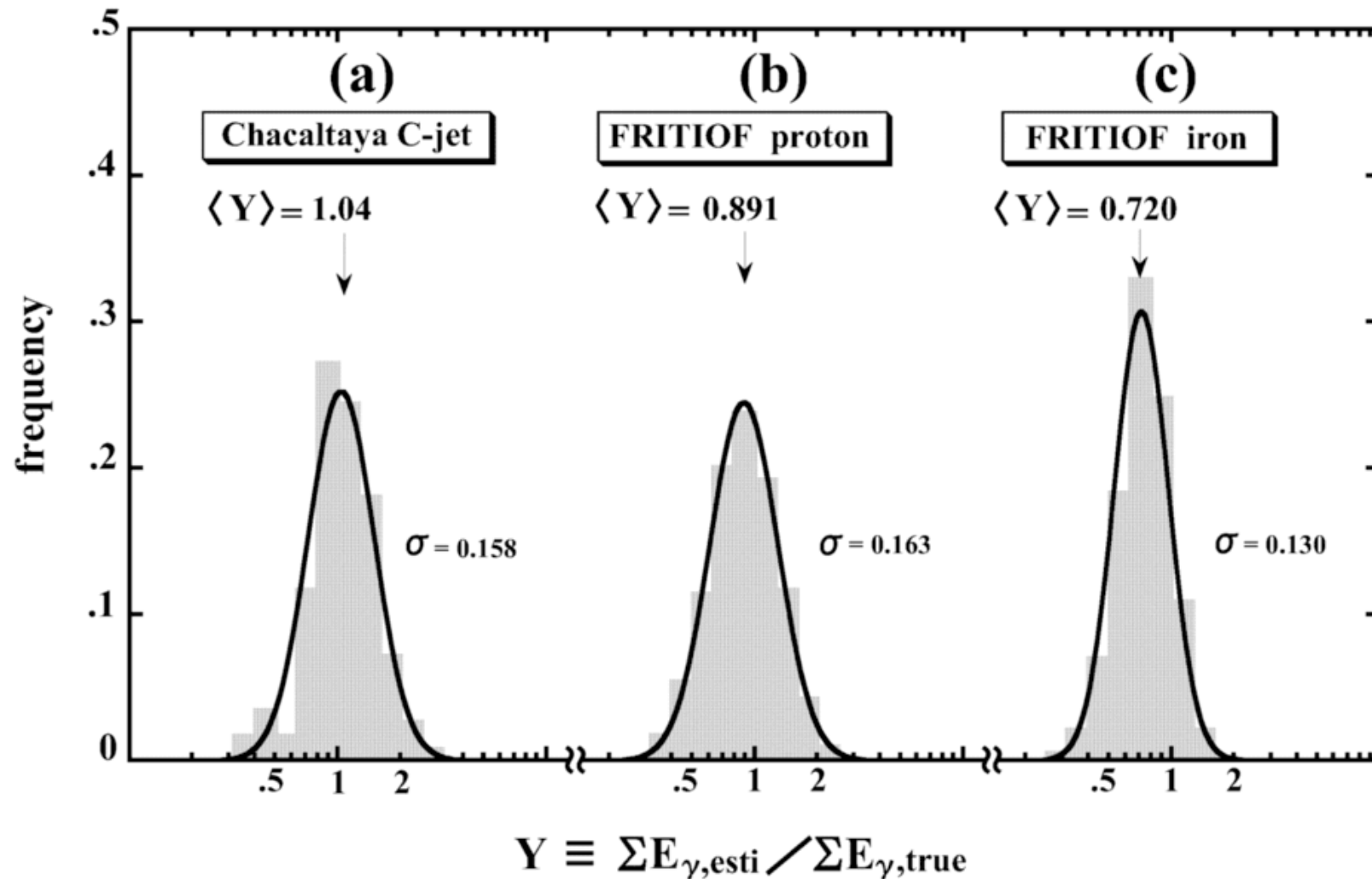
Gamma rays from first interaction detected



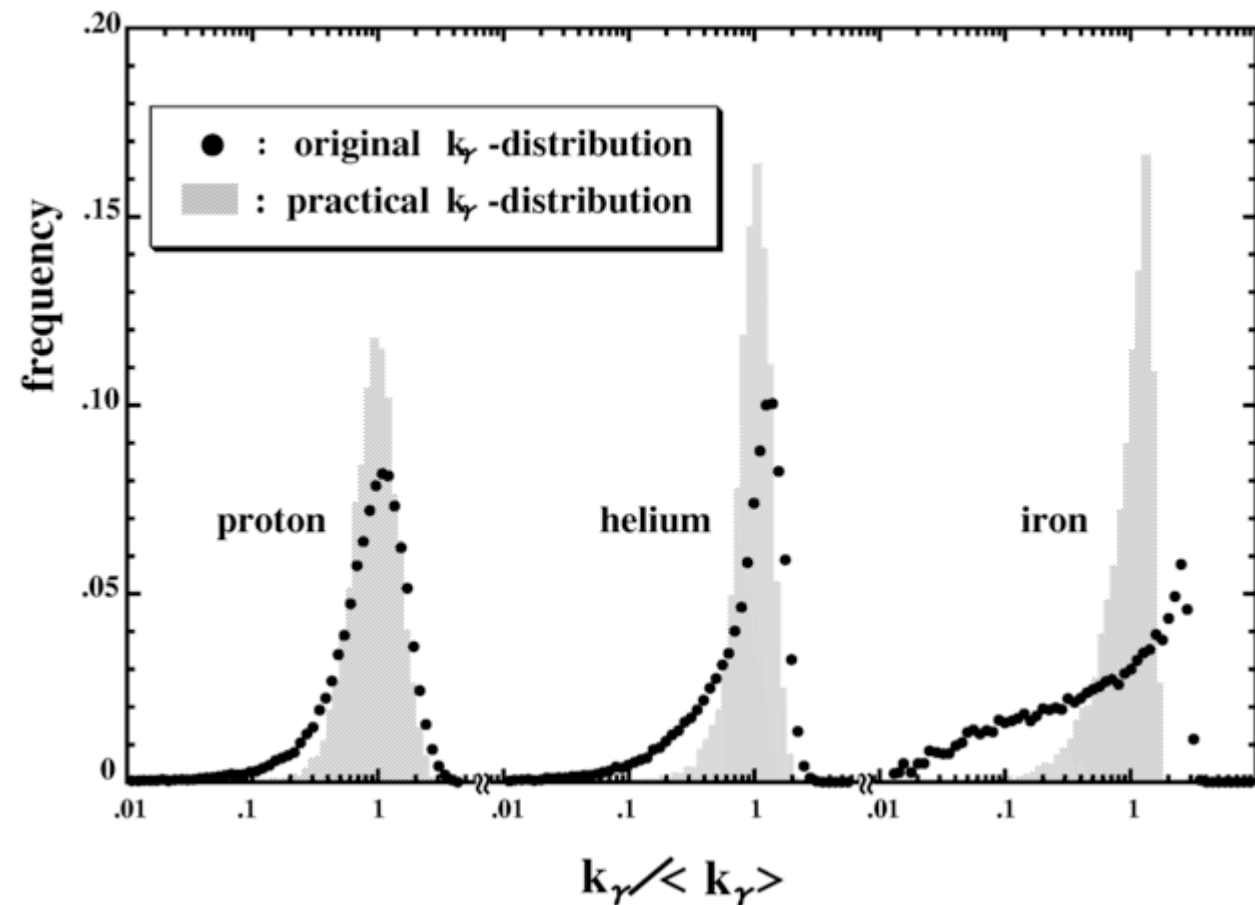
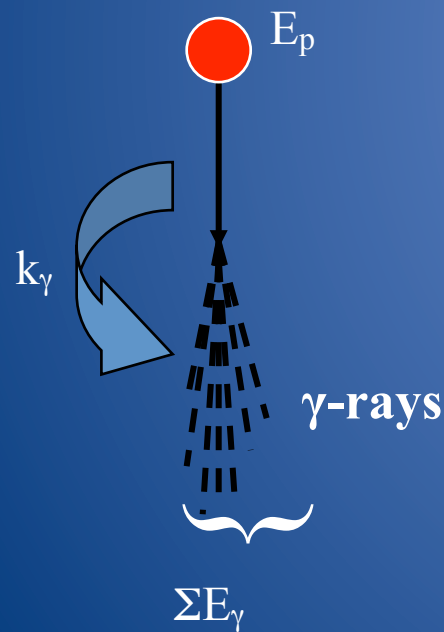
RUNJOB e.m. calorimeter



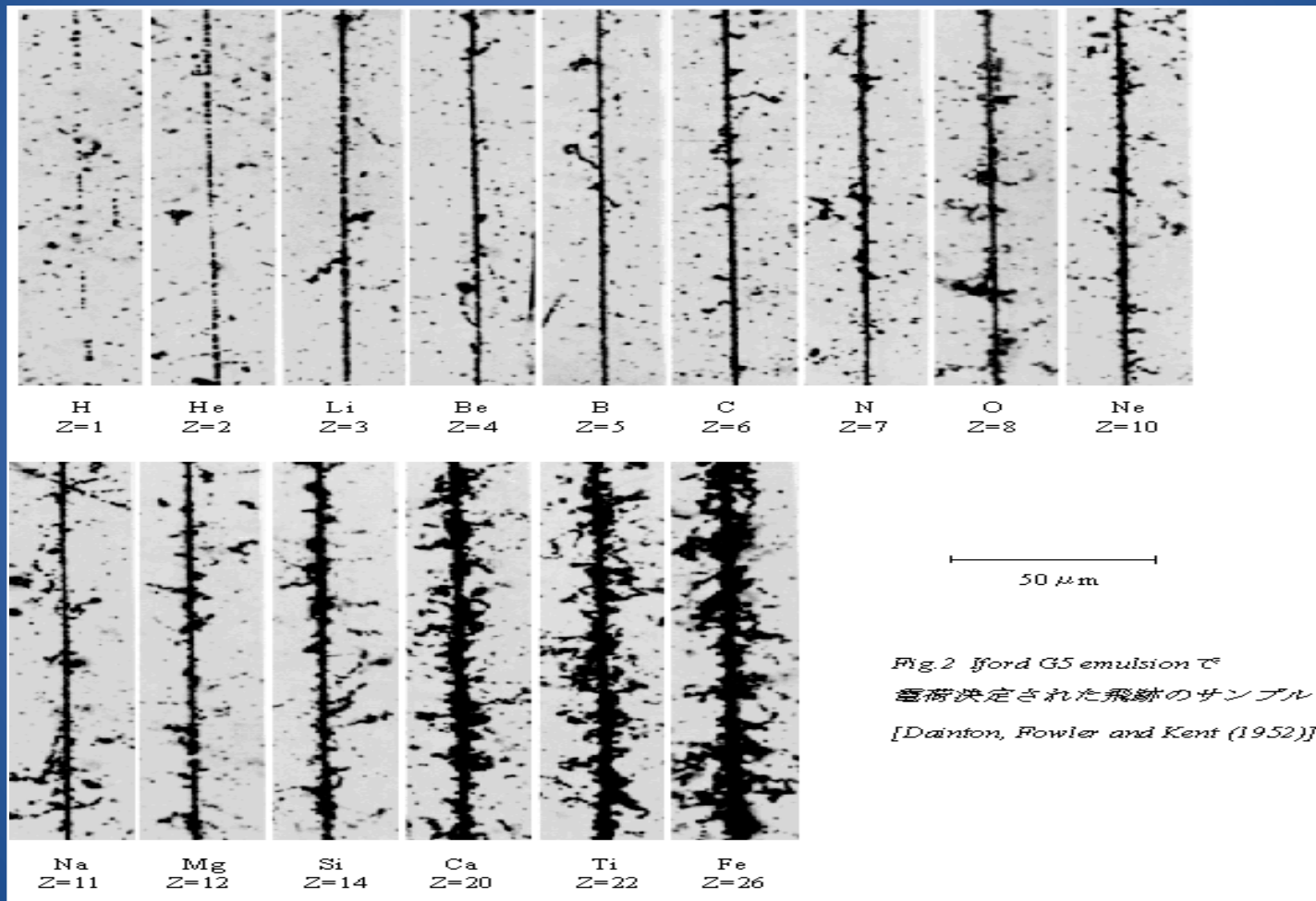
Energy resolution of γ -core method



*Conversion from shower energy to
primary energy k_γ : fraction of energy
transferred to photons in first interaction*



Charge determination in RUNJOB



Charge resolution

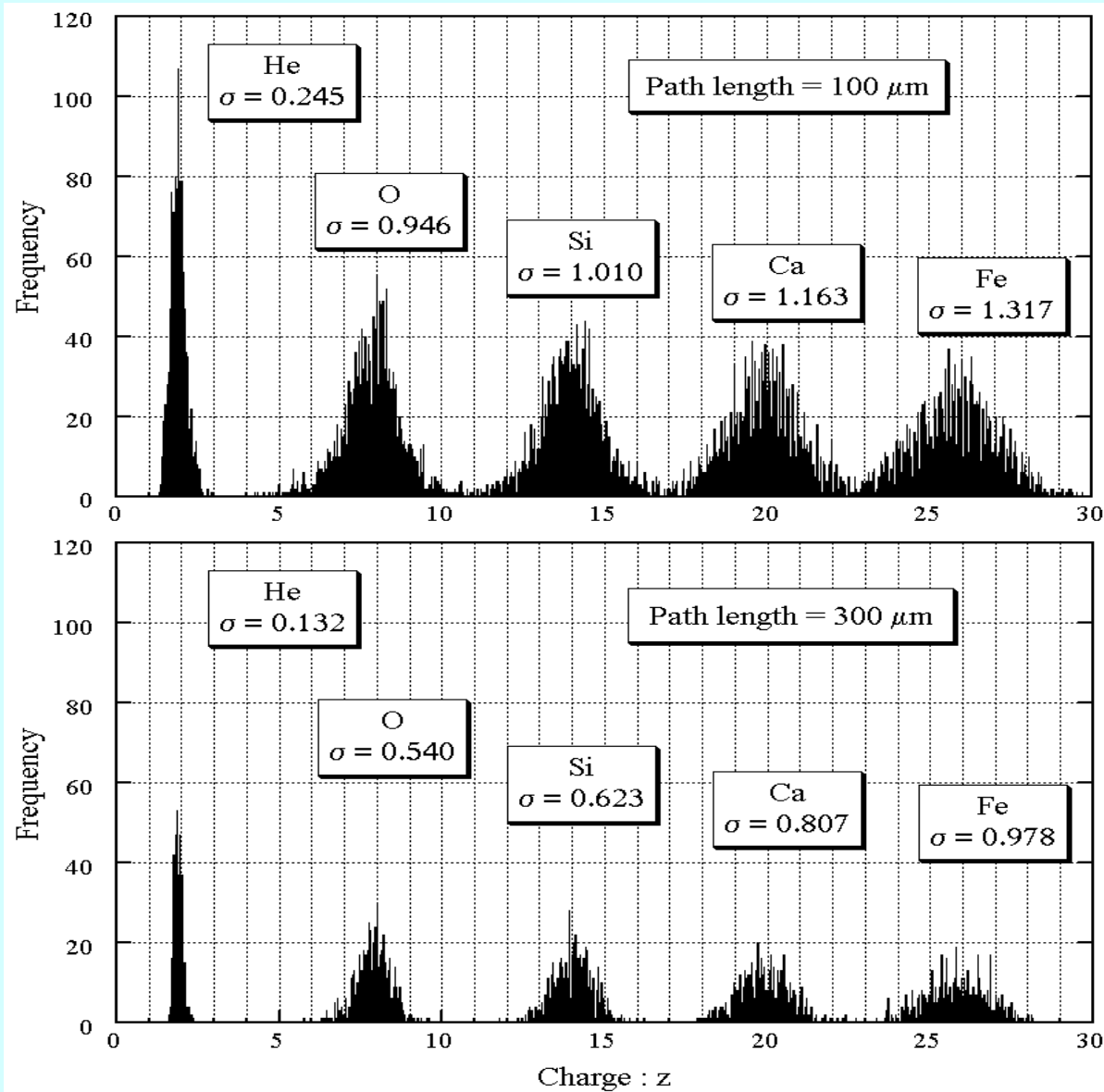
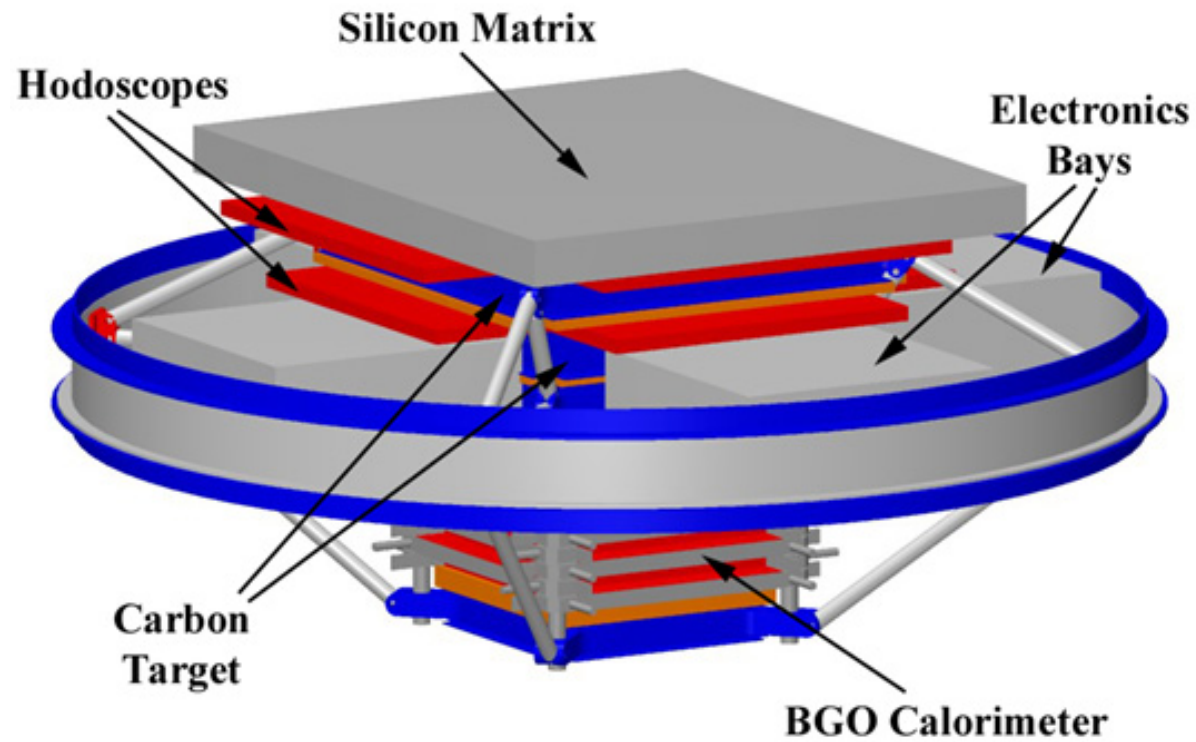
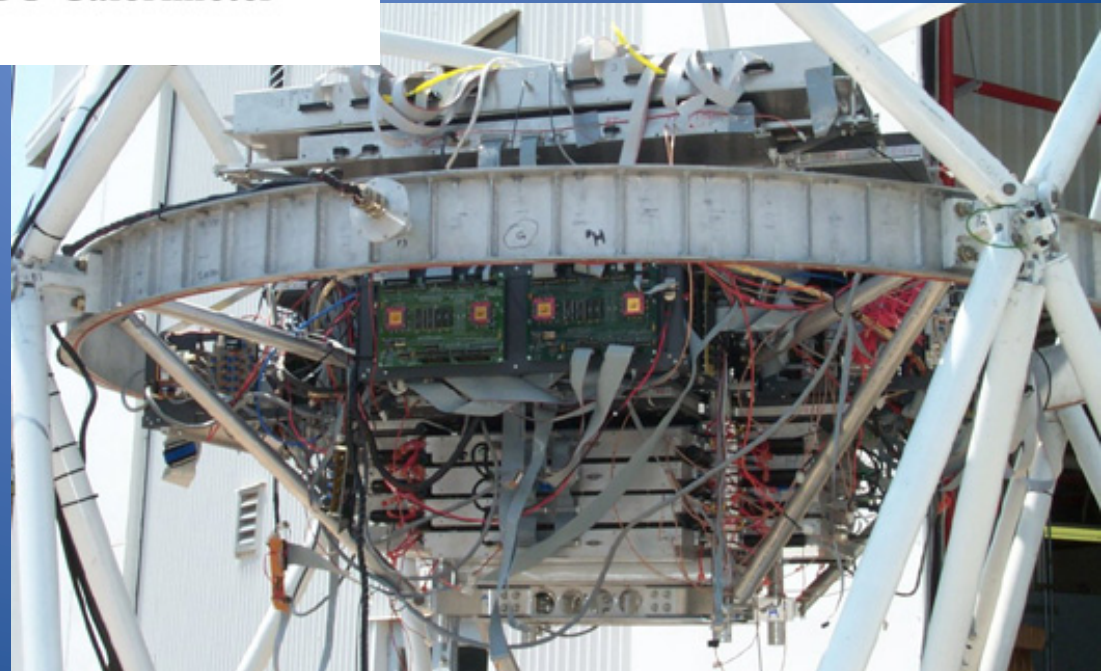


Fig. 5.5b シミュレーション計算による電荷分解能 (2)

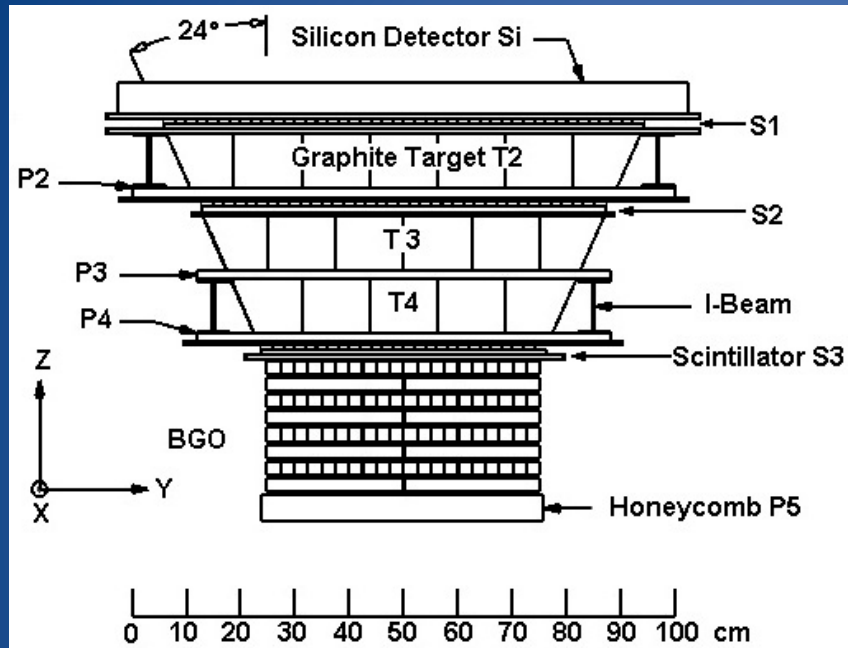
New space detectors ATIC



Long duration flights
In Antarctica



ATIC Instrument Details



- **Si-Matrix:** 4480 pixels each 2 cm x 1.5 cm mounted on offset ladders; 0.95 m x 1.05 m area; 16 bit ADC; CR-1 ASIC's; sparsified readout.

- **Scintillators:** 3 x-y layers; 2 cm x 1 cm cross section; Bicron BC-408; Hamamatsu R5611 pmts both ends; two gain ranges; ACE ASIC. S1 – 336 channels; S2 – 280 channels; S3 – 192 channels; First level trigger: S1-S3

- **Calorimeter:** 8 layers (10 for ATIC-3); 2.5 cm x 2.5 cm x 25 cm BGO crystals, 40 per layer, each crystal viewed by R5611 pmt; three gain ranges; ACE ASIC; 960 channels (1200 for ATIC-3).

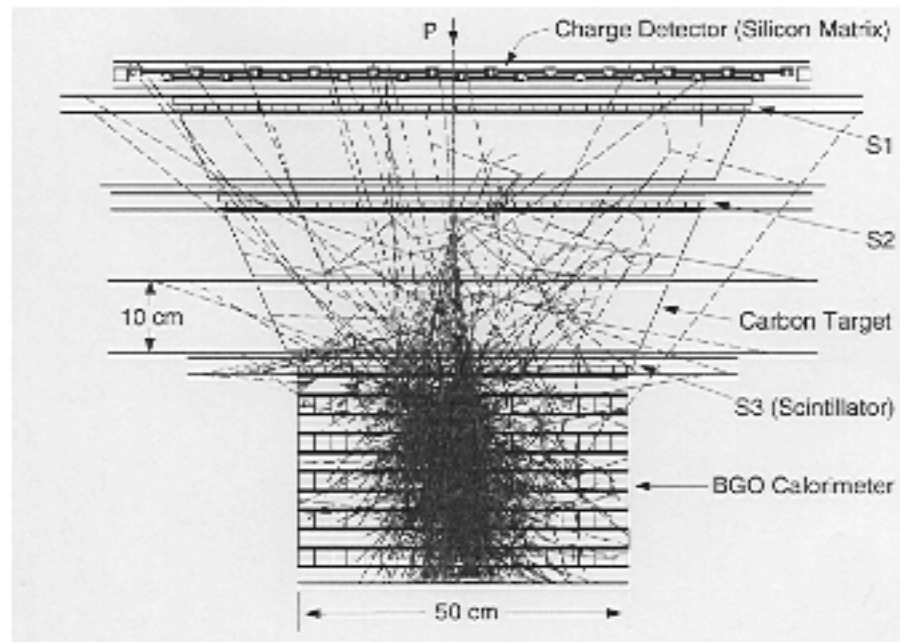
Data System: All data recorded on-board; 70 Gbyte disk (150 Gbyte for ATIC-3); LOS data rate – 330 kbps; TDRSS data rate – 4 kbps (6+ kbps for ATIC-3); Underflight capability (not used).

Housekeeping: Temperature, Pressure, Voltage, Current, Rates, Software Status, Disk status

Command Capability: Power on / off; Trigger type; Thresholds; Pre-scaler; Housekeeping frequency; LOS data rate, Reboot nodes; High Volt settings; Data collection on / off

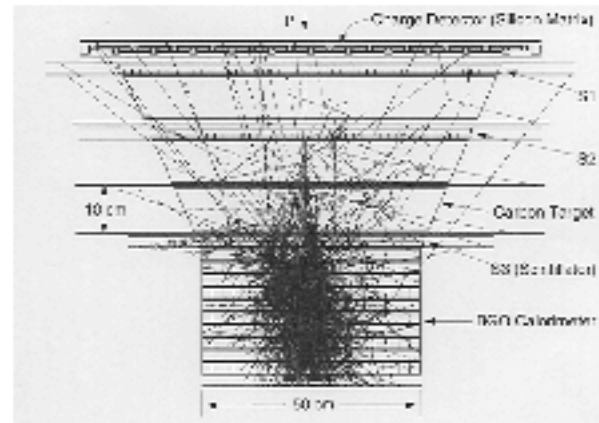
Geometry Factors: S1-S3: 0.42 m²sr; S1-S3-BGO 6: 0.24 m²sr; S1-S3-BGO 8: 0.21 m²sr

ATIC Instrument Summary

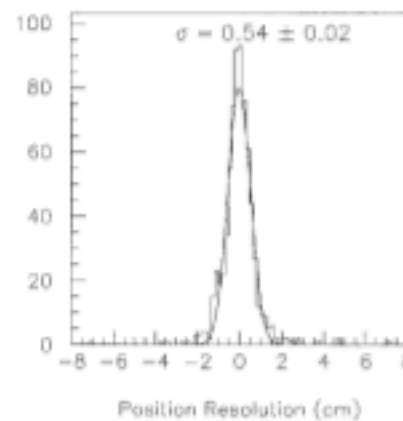


- Measure charge, energy and number
- Ionization Calorimetry only practical method to measure high energy light elements
- Silicon Matrix (Si) has 4,480 pixels to measure GCR charge in presence of shower backscatter
- Graphite Target to interact the primary particle and generate fragments that, in turn, will start an electromagnetic cascade. Also provides some backscatter shielding
- Plastic scintillator hodoscopes (S1, S2, S3) , embedded in Carbon target, provides event trigger plus charge & trajectory information
- Fully active calorimeter includes 400 Bismuth Germinate (BGO) crystals to foster and measure the nuclear - electromagnetic cascade showers
- Geometrical factor: $0.24 \text{ m}^2\text{sr}$ (S1 – S3 – BGO6)

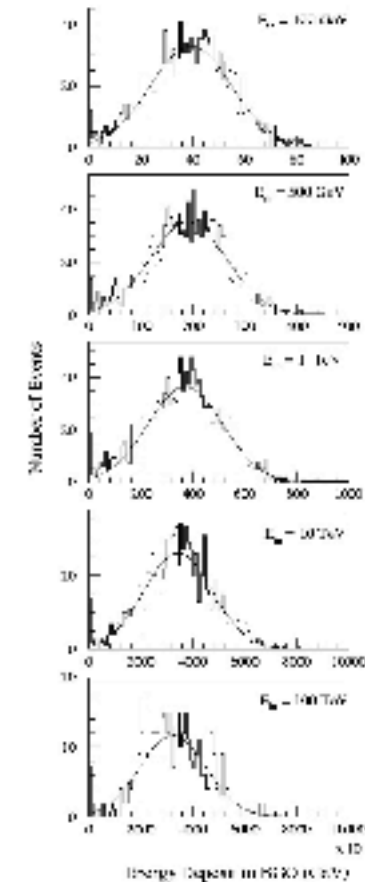
ATIC has been extensively simulated



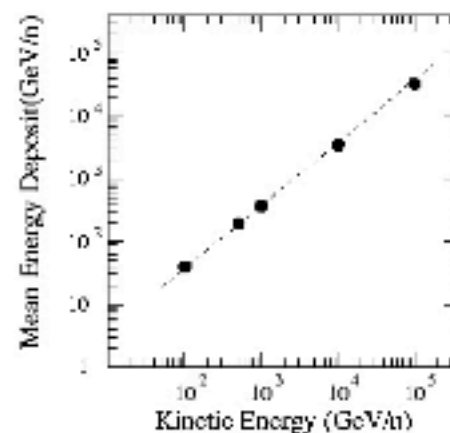
An example of a proton shower in the McMurdo flight configuration



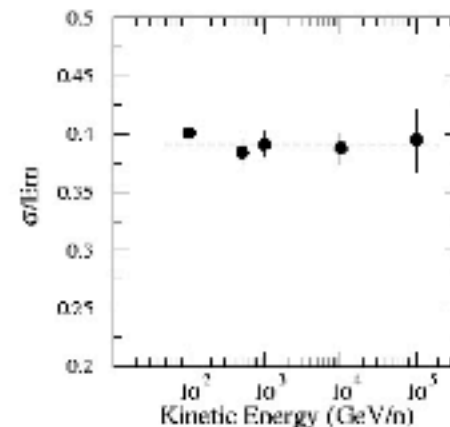
Trajectory resolution in S1 calculated from the BGO shower profile for 1 TeV protons.



Energy deposition in 20 cm depth BGO calorimeter for 10^2 , 5×10^2 , 10^3 , 10^4 , and 10^5 GeV protons.



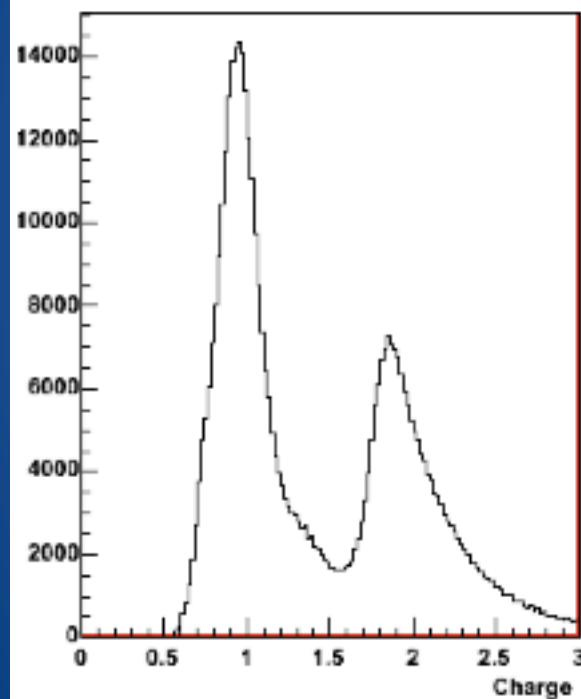
Energy dependence of the proton mean energy deposit and the energy resolution.



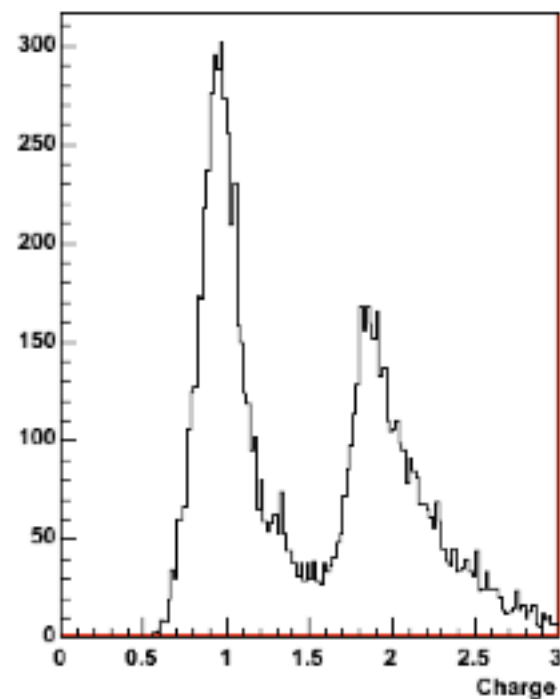
40%

Charge resolution in the p-He group

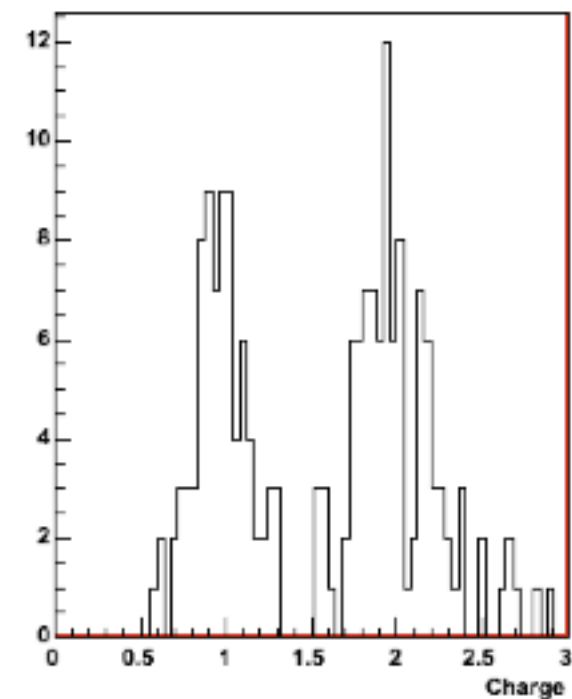
EBGO > 50 GeV



EBGO > 500 GeV



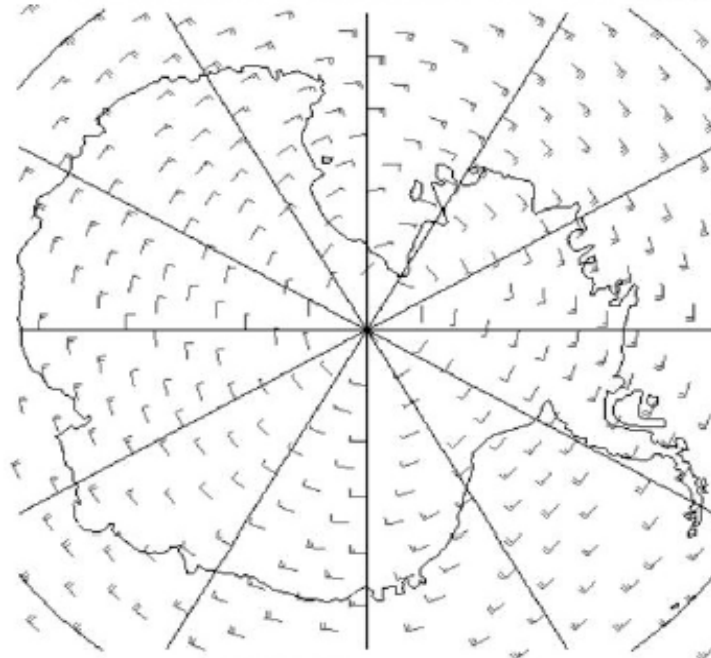
EBGO > 5 TeV



Antarctica long duration flights

High altitude winds are circumpolar during summer

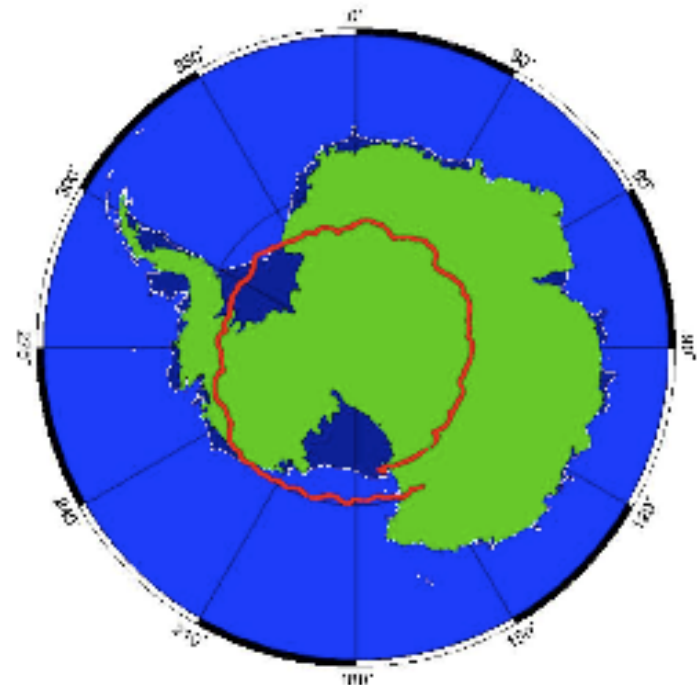
18 December, 2000 on the 3.0 mb surface



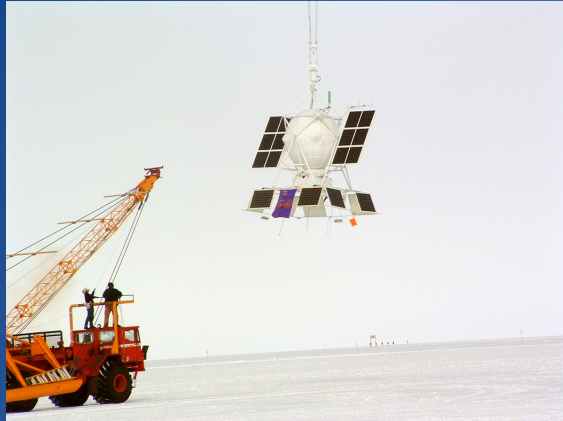
NASA GSFC

ATIC Science - McMurdo

LSU 11/14/05



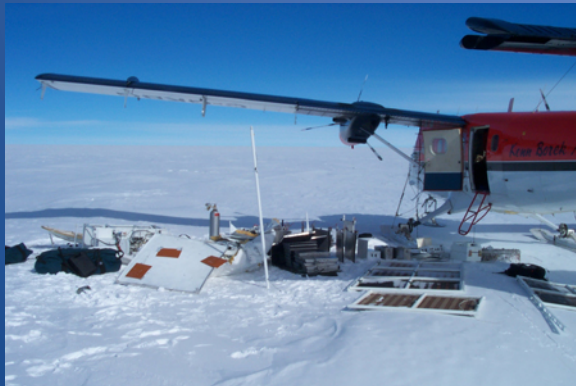
Flight and Recovery



Launch of ATIC-2 in Dec. 2002



The good ATIC-1 landing on 1/13/01 (left) and the not so good landing of ATIC-2 on 1/18/03 (right)



ATIC is designed to be disassembled in the field and recovered with Twin Otters. Two recovery flights are necessary to return all the ATIC components. Pictures show 1st recovery flight of ATIC-1



CREAM

(long balloon flights: 42/28 days)

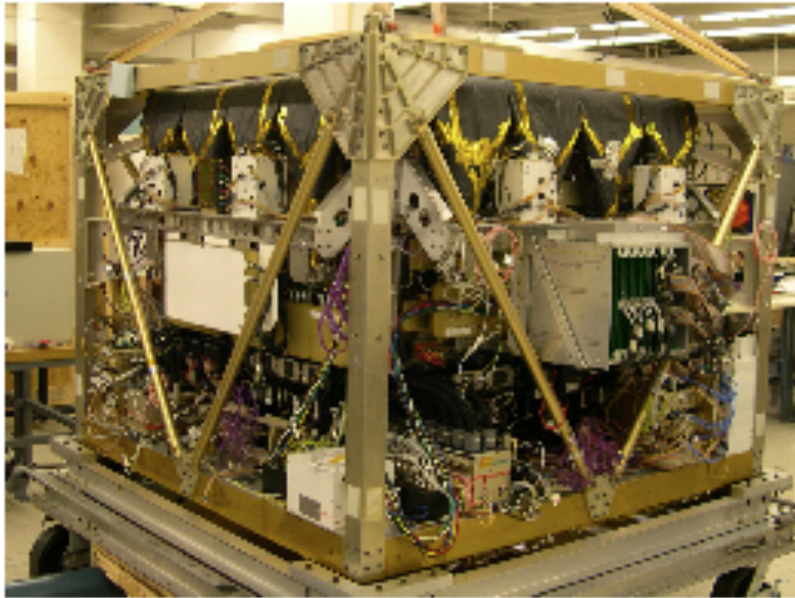


Figure 1. Photo of the CREAM Instrument during integration.

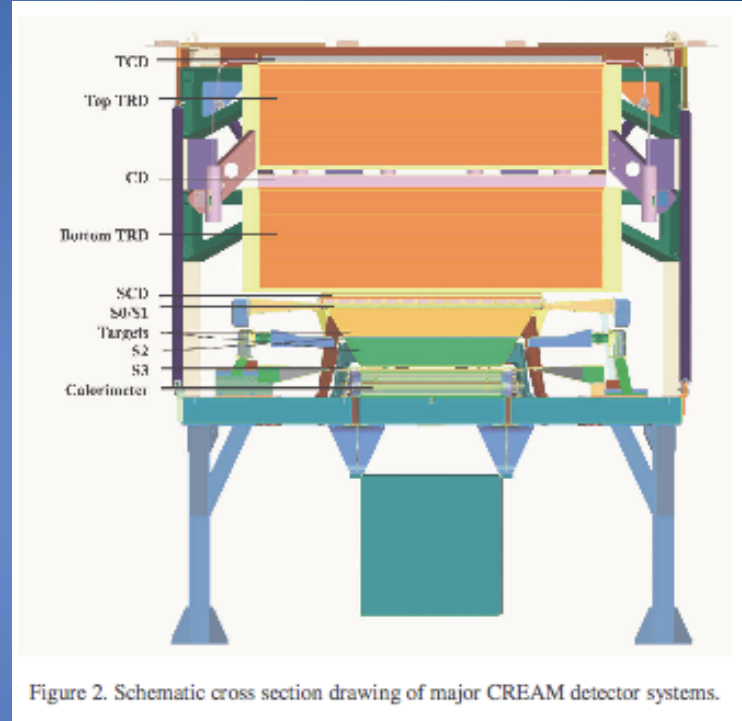


Figure 2. Schematic cross section drawing of major CREAM detector systems.

W = 1143 kg P = 379 W dim = 180 x 180 x 128 cm³

Acc = 0.3 m² sr (2xATIC)

Calorimeter: 20 Xo tungsten sampled every Xo;
sampling fraction 0.3% (scintillating fibers)

SEE YOU TOMORROW....

..... GROUND BASED DETECTORS