Variational Study of Hyperon Effects on the Nuclear Equation of State at Finite Temperature

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Outline

1: Introduction

2: Variational method for hyperon matter

3: Application to supernova matter

4: Summary

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The nuclear equation of state (EOS) plays important roles for astrophysical studies.

1. Lattimer-Swesty EOS : <u>*The Skyrme-type interaction*</u> (NPA 535 (1991) 331)

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$\begin{array}{c c c c c c c c c c c c c c c c c c c $	Nuclear	$n_{\rm sat}$	BE/A	K	Q	J	L	type of int.	used in
$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	Interaction	(fm^{-3})	(MeV)	(MeV)	$\left(\frac{\text{MeV}}{\text{fm}^3}\right)$	(MeV)	(MeV)		SN-EOS list by M. Hempel
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	\mathbf{SKa}	0.155	16.0	263	-300	32.9	74.6	Skyrme	H&W
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	LS180	0.155	16.0	180	-451	28.6	73.8	Skyrme	LS180
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	LS220	0.155	16.0	220	-411	28.6	73.8	Skyrme	LS220, LS220A, LS220 π
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	LS375	0.155	16.0	375	176	28.6	73.8	Skyrme	LS375
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	\mathbf{TMA}	0.147	16.0	318	-572	30.7	90.1	$\mathbf{R}\mathbf{MF}$	HS(TMA)
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	NL3	0.148	16.2	272	203	37.3	118.2	$\mathbf{R}\mathbf{MF}$	SHT, HS(NL3)
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	FSUgold	0.148	16.3	230	-524	32.6	60.5	$\mathbf{R}\mathbf{MF}$	SHO(FSU1.7), HS(FSUgold)
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	FSUgold2.1	0.148	16.3	230	-524	32.6	60.5	$\mathbf{R}\mathbf{MF}$	SHO(FSU2.1)
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	IUFSU	0.155	16.4	231	-290	31.3	47.2	$\mathbf{R}\mathbf{MF}$	HS(IUFSU)
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	DD2	0.149	16.0	243	169	31.7	55.0	\mathbf{RMF}	HS(DD2), BHBA, BHBA ϕ
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	SFHo	0.158	16.2	245	-468	31.6	47.1	$\mathbf{R}\mathbf{MF}$	SFHo
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	\mathbf{SFHx}	0.160	16.2	239	-457	28.7	23.2	\mathbf{RMF}	SFHx
STOSY, STOSY π , STOS π , STOS π Q, STOSQ, STOSB139, STOSB145, STOSB155, STOSB162, STOSB165	TM1	0.145	16.3	281	-285	36.9	110.8	$\mathbf{R}\mathbf{MF}$	STOS, FYSS, $HS(TM1)$, $STOS\Lambda$,
STOSQ, STOSB139, STOSB145, STOSB155, STOSB162, STOSB165									STOSY, STOSY π , STOS π , STOS π Q,
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There were no supernova EOSs based on the microscopic many-body theory.

SN-EOS with the microscopic many-body theory

We constructed a new nuclear EOS for core-collapse simulations.

<u>Uniform matter : Cluster variational method (AV18 + UIX)</u> <u>Non-uniform matter : Thomas-Fermi approximation</u>

Collaboration with M. Takano (Waseda Univ.), K. Sumiyoshi (Numazu College of Tech.), K. Nakazato (Kyushu Univ.), Y. Takehara, S. Yamamuro, H. Suzuki (Tokyo Univ. of Science)



We aim to extend the microscopic nuclear EOS table to consider A hyperon mixing.

2: Variational method for hyperon matter



- NN potential: AV18 two-body potential (PRC 51 (1995) 38)
- YN and YY potentials: Central three-range Gaussian potentials

 ΛN interaction
 (E. Hiyama et al., PRC 74 (2006) 054312)

 ΛΛ interaction
 (E. Hiyama et al., PRC 66 (2002) 024007)

Hyperon interactions are determined so as to reproduce the experimental data on single-and double- Λ hypernuclei.



Energy of hyperonic nuclear matter

Free energy of hyperon matter at finite temperature

We follow the prescription proposed by Schmidt and Pandharipande.

(K. E. Schmidt et al., Phys. Lett. 87B(1979) 11) (A. Mukherjee et al., PRC 75(2007) 035802)

Free energy F is expressed by <u>the average occupation probabilities</u>.



Free energy of hyperonic nuclear matter

3: Application to supernova matter

Using the obtained EOS of hot hyperon matter,

We calculate the thermodynamic quantities of supernova (SN) matter under the equilibrium condition $\mu_n = \mu_{\Lambda}$

at fixed baryon number density $n_{\rm B}$, proton fraction $Y_{\rm p}$, and temperature T.



The effect of Λ hyperons on SN matter becomes larger at higher temperatures.

Composition of supernova matter





The effect of Λ hyperons on SN matter becomes larger at lower proton fraction.

Composition of supernova matter



- X_{Λ} with our EOS is considerably smaller than that of the Shen EOS for $Y_{\rm p} = 0$.
- X_{Λ} with our EOS is close to that of the Shen EOS for $Y_{\rm p} = 0.5$.
 - → Symmetry energy of our EOS ($E_{sym} = 30.0 \text{ MeV}$) is smaller than that of the Shen EOS ($E_{sym} = 36.9 \text{ MeV}$).

Free energy and entropy of hot hyperon matter



Thermodynamic quantities of our EOS at smaller Y_p are less affected by the Λ hyperon mixing than in the case of the Shen EOS.



We construct the EOS of nuclear matter including Λ hyperons at zero and finite temperatures by <u>the variational method</u>.

Variational method for hyperon matter

• The obtained thermodynamic quantities are reasonable.

Application of the EOS to supernova matter

• The effect of Λ hyperons on supernova matter becomes larger at higher temperatures and lower proton fractions.

Comparison with the Shen EOS

• Our EOS at smaller Y_p are less affected by the Λ hyperon mixing due to the smaller symmetry energy.

Future Plans

- Construction of the hyperon EOS table for core-collapse simulations
- Contribution of the hyperon three-baryon force

Hyperon Matter with Three-Baryon Force

We consider a phenomenological three-baryon repulsive force (TBF) as a density dependent two-body effective potential.



Mass-radius relations of neutron stars

Phase diagram of hot hyperon matter