

Electromagnetic Form Factors of Nucleon Excitations from Lattice QCD

Finn M. Stokes

Waseem Kamleh, Derek B. Leinweber and Benjamin J. Owen

Centre for the Subatomic Structure of Matter



THE UNIVERSITY
*of*ADELAIDE

Introduction

- The isolation of excitations of baryons at nonzero momentum is important for the evaluation of baryon form factors and transition moments

Introduction

- The isolation of excitations of baryons at nonzero momentum is important for the evaluation of baryon form factors and transition moments
- Existing parity projection techniques are vulnerable to opposite parity contaminations at nonzero momentum

Introduction

- The isolation of excitations of baryons at nonzero momentum is important for the evaluation of baryon form factors and transition moments
- Existing parity projection techniques are vulnerable to opposite parity contaminations at nonzero momentum
- We propose the Parity Expanded Variational Analysis (PEVA) technique to resolve this issue

Step 1: Construct correlation matrix

Start with basis of N operators and construct an $N \times N$ correlation matrix $G_{ij}(\mathbf{p}; \Gamma; t)$.

Baryon structure via variational analysis

Step 1: Construct correlation matrix

Start with basis of N operators and construct an $N \times N$ correlation matrix $G_{ij}(\mathbf{p}; \Gamma; t)$.

Step 2: Perform variational analysis

Perform a variational analysis to construct optimised operators.

Baryon structure via variational analysis

Step 1: Construct correlation matrix

Start with basis of N operators and construct an $N \times N$ correlation matrix $G_{ij}(\mathbf{p}; \Gamma; t)$.

Step 2: Perform variational analysis

Perform a variational analysis to construct optimised operators.

Step 3: Compute three point correlation function

Use these optimised operators to construct relevant three point correlation functions

Baryon structure via variational analysis

Step 1: Construct correlation matrix

Start with basis of N operators and construct an $N \times N$ correlation matrix $G_{ij}(\mathbf{p}; \Gamma; t)$.

Step 2: Perform variational analysis

Perform a variational analysis to construct optimised operators.

Step 3: Compute three point correlation function

Use these optimised operators to construct relevant three point correlation functions

Step 4: Extract form factors

Take ratios of three point to to point functions and extract e.g. $G_E(Q^2)$ and $G_M(Q^2)$.

Step 1: Construct correlation matrix

Start with basis of N operators and construct an $N \times N$ correlation matrix $G_{ij}(\mathbf{p}; \Gamma; t)$.

Step 2: Perform variational analysis

Perform a variational analysis to construct optimised operators.

Step 3: Compute three point correlation function

Use these optimised operators to construct relevant three point correlation functions

Step 4: Extract form factors

Take ratios of three point to to point functions and extract e.g. $G_E(Q^2)$ and $G_M(Q^2)$.

Step 1: Construct correlation matrix

- Use local three-quark spin-1/2 nucleon operators

$$\begin{aligned}\chi_1 &= \epsilon^{abc} [u^a{}^\top (C\gamma_5) d^b] u^c \\ \chi_2 &= \epsilon^{abc} [u^a{}^\top (C) d^b] \gamma_5 u^c\end{aligned}$$

Step 1: Construct correlation matrix

- Use local three-quark spin-1/2 nucleon operators

$$\chi_1 = \epsilon^{abc} [u^a{}^\top (C \gamma_5) d^b] u^c$$

$$\chi_2 = \epsilon^{abc} [u^a{}^\top (C) d^b] \gamma_5 u^c$$

- Apply 16, 35, 100 and 200 sweeps of gauge invariant gaussian smearing

Step 1: Construct correlation matrix

- Use local three-quark spin-1/2 nucleon operators

$$\begin{aligned}\chi_1 &= \epsilon^{abc} [u^a{}^\top (C\gamma_5) d^b] u^c \\ \chi_2 &= \epsilon^{abc} [u^a{}^\top (C) d^b] \gamma_5 u^c\end{aligned}$$

- Apply 16, 35, 100 and 200 sweeps of gauge invariant gaussian smearing

Correlation Matrix

$$G_{ij}(\mathbf{p}; t) := \sum_{\mathbf{x}} e^{i\mathbf{p} \cdot \mathbf{x}} \langle \Omega | \chi^i(x) \bar{\chi}^j(0) | \Omega \rangle$$

Step 1: Construct correlation matrix

- Use local three-quark spin-1/2 nucleon operators

$$\chi_1 = \epsilon^{abc} [u^a{}^\top (C\gamma_5) d^b] u^c$$

$$\chi_2 = \epsilon^{abc} [u^a{}^\top (C) d^b] \gamma_5 u^c$$

- Apply 16, 35, 100 and 200 sweeps of gauge invariant gaussian smearing

Correlation Matrix

$$G_{ij}(\mathbf{p}; t) := \sum_{\mathbf{x}} e^{i\mathbf{p}\cdot\mathbf{x}} \langle \Omega | \chi^i(x) \bar{\chi}^j(0) | \Omega \rangle$$

Step 1: Construct correlation matrix

- Use local three-quark spin-1/2 nucleon operators

$$\chi_1 = \epsilon^{abc} [u^a{}^\top (C\gamma_5) d^b] u^c$$

$$\chi_2 = \epsilon^{abc} [u^a{}^\top (C) d^b] \gamma_5 u^c$$

- Apply 16, 35, 100 and 200 sweeps of gauge invariant gaussian smearing

Correlation Matrix

$$G_{ij}(\mathbf{p}; t) := \sum_{\mathbf{x}} e^{i\mathbf{p}\cdot\mathbf{x}} \langle \Omega | \chi^i(x) \bar{\chi}^j(0) | \Omega \rangle$$

Step 1: Construct correlation matrix

- Use local three-quark spin-1/2 nucleon operators

$$\chi_1 = \epsilon^{abc} [u^a{}^\top (C\gamma_5) d^b] u^c$$

$$\chi_2 = \epsilon^{abc} [u^a{}^\top (C) d^b] \gamma_5 u^c$$

- Apply 16, 35, 100 and 200 sweeps of gauge invariant gaussian smearing

Correlation Matrix

$$G_{ij}(\mathbf{p}; t) := \sum_{\mathbf{x}} e^{i\mathbf{p} \cdot \mathbf{x}} \langle \Omega | \chi^i(x) \bar{\chi}^j(0) | \Omega \rangle$$

Step 1: Construct correlation matrix

- Use local three-quark spin-1/2 nucleon operators

$$\begin{aligned}\chi_1 &= \epsilon^{abc} [u^a{}^\top (C\gamma_5) d^b] u^c \\ \chi_2 &= \epsilon^{abc} [u^a{}^\top (C) d^b] \gamma_5 u^c\end{aligned}$$

- Apply 16, 35, 100 and 200 sweeps of gauge invariant gaussian smearing

Correlation Matrix

$$\mathcal{G}_{ij}(\mathbf{p}; t) := \sum_{\mathbf{x}} e^{i\mathbf{p} \cdot \mathbf{x}} \langle \Omega | \chi^i(x) \bar{\chi}^j(0) | \Omega \rangle$$

Projected Correlation Matrix

$$G_{ij}(\Gamma; \mathbf{p}; t) := \text{tr } \Gamma \mathcal{G}_{ij}(\mathbf{p}; t)$$

Step 1: Construct correlation matrix

- Use local three-quark spin-1/2 nucleon operators

$$\begin{aligned}\chi_1 &= \epsilon^{abc} [u^a{}^\top (C\gamma_5) d^b] u^c \\ \chi_2 &= \epsilon^{abc} [u^a{}^\top (C) d^b] \gamma_5 u^c\end{aligned}$$

- Apply 16, 35, 100 and 200 sweeps of gauge invariant gaussian smearing

Correlation Matrix

$$G_{ij}(\mathbf{p}; t) := \sum_{\mathbf{x}} e^{i\mathbf{p} \cdot \mathbf{x}} \langle \Omega | \chi^i(x) \bar{\chi}^j(0) | \Omega \rangle$$

Projected Correlation Matrix

$$G_{ij}(\Gamma; \mathbf{p}; t) := \text{tr } \Gamma G_{ij}(\mathbf{p}; t)$$

Parity-Expanded Variational Analysis (PEVA)

- Expand basis to simultaneously isolate finite momentum energy eigenstates, regardless of parity

Parity-Expanded Variational Analysis (PEVA)

- Expand basis to simultaneously isolate finite momentum energy eigenstates, regardless of parity
- Define PEVA projector

$$\Gamma_{\mathbf{p}} = \frac{1}{4}(\mathbb{I} + \gamma_4)(\mathbb{I} - i\gamma_5\gamma_k\hat{p}_k)$$

Parity-Expanded Variational Analysis (PEVA)

- Expand basis to simultaneously isolate finite momentum energy eigenstates, regardless of parity
- Define PEVA projector

$$\Gamma_{\mathbf{p}} = \frac{1}{4}(\mathbb{I} + \gamma_4)(\mathbb{I} - i\gamma_5\gamma_k\hat{p}_k)$$

- $\chi_{\mathbf{p}}^i := \Gamma_{\mathbf{p}}\chi^i$ couples to positive parity states at zero momentum

Parity-Expanded Variational Analysis (PEVA)

- Expand basis to simultaneously isolate finite momentum energy eigenstates, regardless of parity
- Define PEVA projector

$$\Gamma_{\mathbf{p}} = \frac{1}{4}(\mathbb{I} + \gamma_4)(\mathbb{I} - i\gamma_5\gamma_k\hat{p}_k)$$

- $\chi_{\mathbf{p}}^i := \Gamma_{\mathbf{p}}\chi^i$ couples to positive parity states at zero momentum
- $\chi_{\mathbf{p}}^{i''} := \Gamma_{\mathbf{p}}\gamma_5\chi^i$ couples to negative parity states at zero momentum

Baryon structure via variational analysis

Step 1: Construct correlation matrix

Start with basis of N operators and construct an $N \times N$ correlation matrix $G_{ij}(\mathbf{p}; \Gamma; t)$.

Step 2: Perform variational analysis

Perform a variational analysis to construct optimised operators.

Step 3: Compute three point correlation function

Use these optimised operators to construct relevant three point correlation functions

Step 4: Extract form factors

Take ratios of three point to to point functions and extract e.g. $G_E(Q^2)$ and $G_M(Q^2)$.

Step 2: Perform variational analysis

Seek operators $\{\phi_{\mathbf{p}}^{\alpha}\}$ that couple strongly to a single energy eigenstate

$$\phi_{\mathbf{p}}^{\alpha} = \sum_i v_i^{\alpha}(\mathbf{p}) \chi_{\mathbf{p}}^i + \sum_{i'} v_{i'}^{\alpha}(\mathbf{p}) \chi_{\mathbf{p}}^{i'}$$

$$\bar{\phi}_{\mathbf{p}}^{\alpha} = \sum_i u_i^{\alpha}(\mathbf{p}) \bar{\chi}_{\mathbf{p}}^i + \sum_{i'} u_{i'}^{\alpha}(\mathbf{p}) \bar{\chi}_{\mathbf{p}}^{i'}$$

Step 2: Perform variational analysis

Seek operators $\{\phi_{\mathbf{p}}^{\alpha}\}$ that couple strongly to a single energy eigenstate

$$\phi_{\mathbf{p}}^{\alpha} = \sum_i v_i^{\alpha}(\mathbf{p}) \chi_{\mathbf{p}}^i + \sum_{i'} v_{i'}^{\alpha}(\mathbf{p}) \chi_{\mathbf{p}}^{i'}$$

$$\bar{\phi}_{\mathbf{p}}^{\alpha} = \sum_i u_i^{\alpha}(\mathbf{p}) \bar{\chi}_{\mathbf{p}}^i + \sum_{i'} u_{i'}^{\alpha}(\mathbf{p}) \bar{\chi}_{\mathbf{p}}^{i'}$$

$$G_{\alpha}(\mathbf{p}; t) := v_i^{\alpha}(\mathbf{p}) G_{ij}(\mathbf{p}; t) u_j^{\alpha}(\mathbf{p})$$

Step 2: Perform variational analysis

Seek operators $\{\phi_{\mathbf{p}}^{\alpha}\}$ that couple strongly to a single energy eigenstate

$$\phi_{\mathbf{p}}^{\alpha} = \sum_i v_i^{\alpha}(\mathbf{p}) \chi_{\mathbf{p}}^i + \sum_{i'} v_{i'}^{\alpha}(\mathbf{p}) \chi_{\mathbf{p}}^{i'}$$

$$\bar{\phi}_{\mathbf{p}}^{\alpha} = \sum_i u_i^{\alpha}(\mathbf{p}) \bar{\chi}_{\mathbf{p}}^i + \sum_{i'} u_{i'}^{\alpha}(\mathbf{p}) \bar{\chi}_{\mathbf{p}}^{i'}$$

$$G_{\alpha}(\mathbf{p}; t) := v_i^{\alpha}(\mathbf{p}) G_{ij}(\mathbf{p}; t) u_j^{\alpha}(\mathbf{p})$$

$$E_{\alpha}^{\text{eff}}(\mathbf{p}, t) := \frac{1}{\delta t} \ln \frac{G_{\alpha}(\mathbf{p}; t)}{G_{\alpha}(\mathbf{p}; t + \delta t)}$$

Step 2: Perform variational analysis

Seek operators $\{\phi_{\mathbf{p}}^{\alpha}\}$ that couple strongly to a single energy eigenstate

$$\phi_{\mathbf{p}}^{\alpha} = \sum_i v_i^{\alpha}(\mathbf{p}) \chi_{\mathbf{p}}^i + \sum_{i'} v_{i'}^{\alpha}(\mathbf{p}) \chi_{\mathbf{p}}^{i'}$$

$$\bar{\phi}_{\mathbf{p}}^{\alpha} = \sum_i u_i^{\alpha}(\mathbf{p}) \bar{\chi}_{\mathbf{p}}^i + \sum_{i'} u_{i'}^{\alpha}(\mathbf{p}) \bar{\chi}_{\mathbf{p}}^{i'}$$

$$G_{\alpha}(\mathbf{p}; t) := v_i^{\alpha}(\mathbf{p}) G_{ij}(\mathbf{p}; t) u_j^{\alpha}(\mathbf{p})$$

$$E_{\alpha}^{\text{eff}}(\mathbf{p}, t) := \frac{1}{\delta t} \ln \frac{G_{\alpha}(\mathbf{p}; t)}{G_{\alpha}(\mathbf{p}; t + \delta t)}$$

Expect E_{α}^{eff} to approximately obey dispersion relation

$$E_{\alpha}^{\text{eff}}(\mathbf{p}, t) \approx \sqrt{m_{\alpha}^2 + \mathbf{p}^2}$$

Step 2: Perform variational analysis

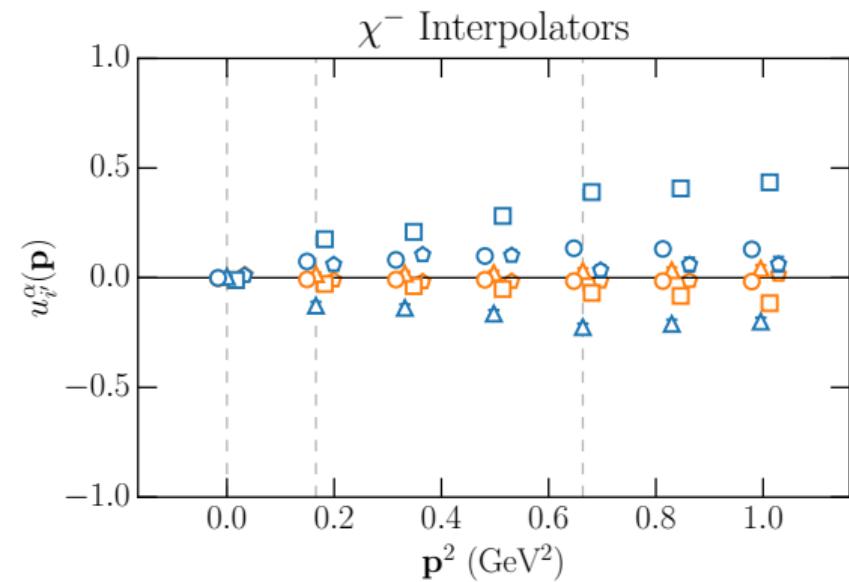
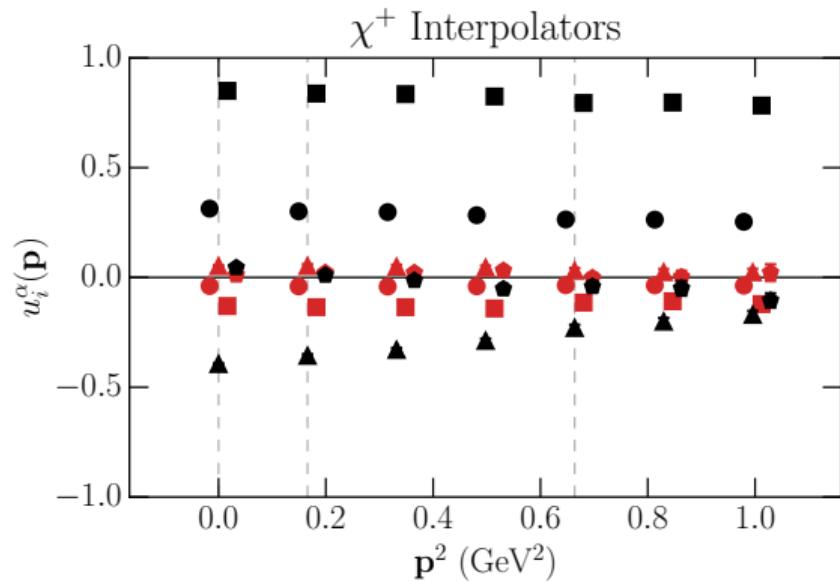
Lattice results

Second lightest PACS-CS (2 + 1)-flavour full-QCD ensemble

- $32^3 \times 64$ lattices
- $a = 0.0951(14)$ fm by Sommer parameter
- $\kappa_{u,d} = 0.1377$, corresponding to $m_\pi = 280(5)$ MeV

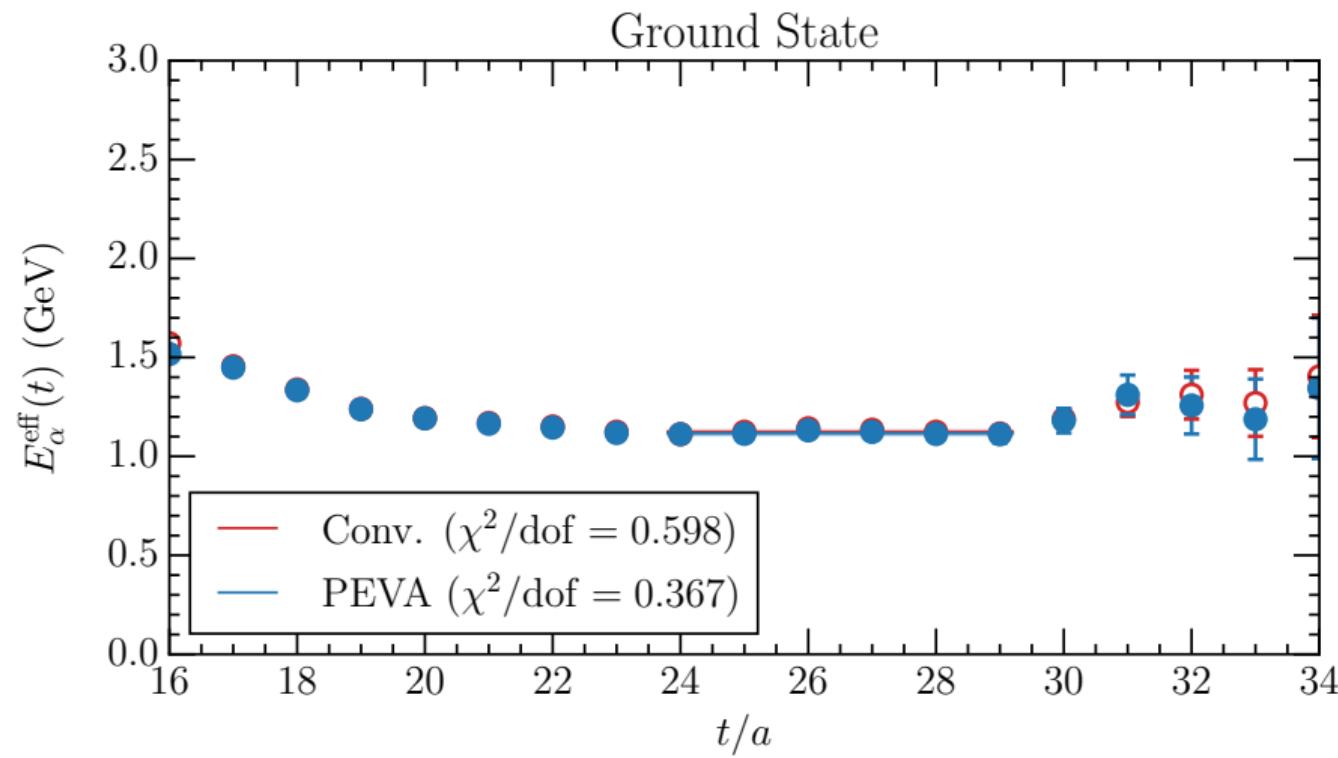
Eigenvector components

Ground state



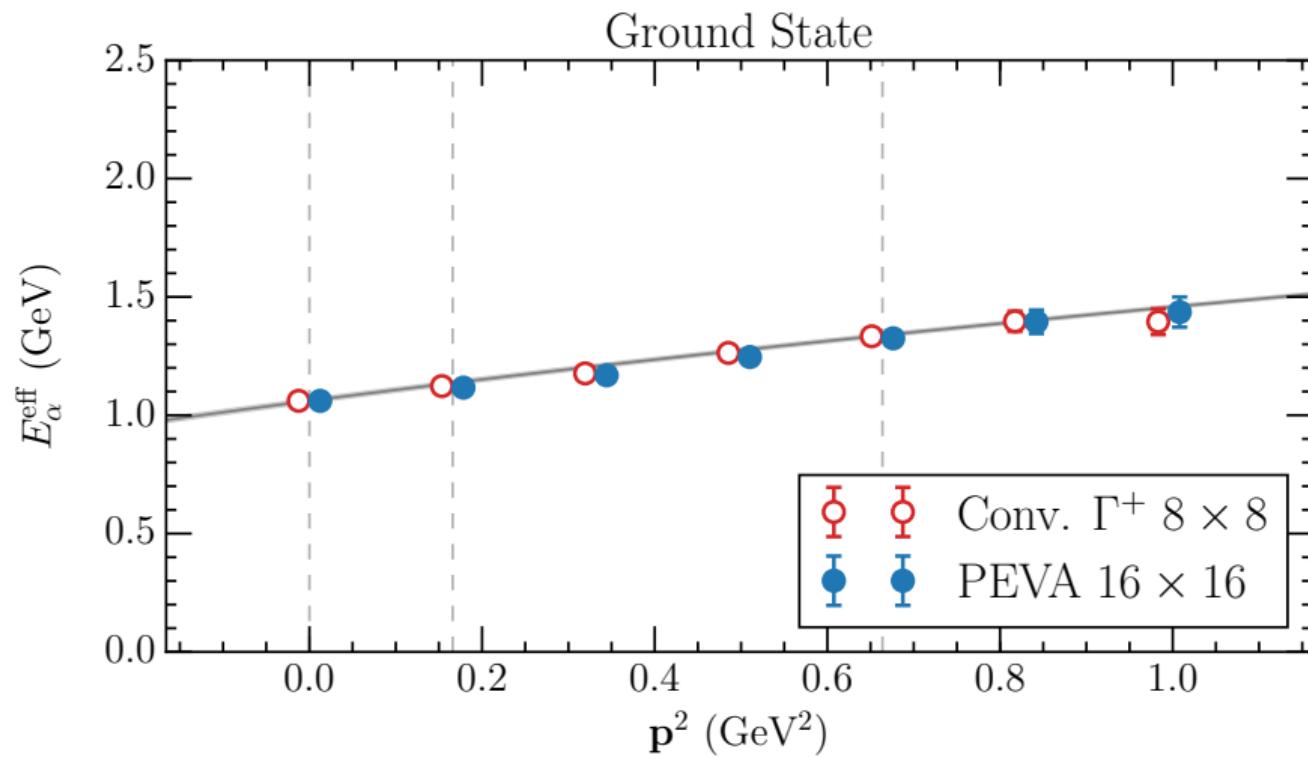
Effective energy

Ground state - $\mathbf{p}^2 \simeq 0.166 \text{ GeV}^2$



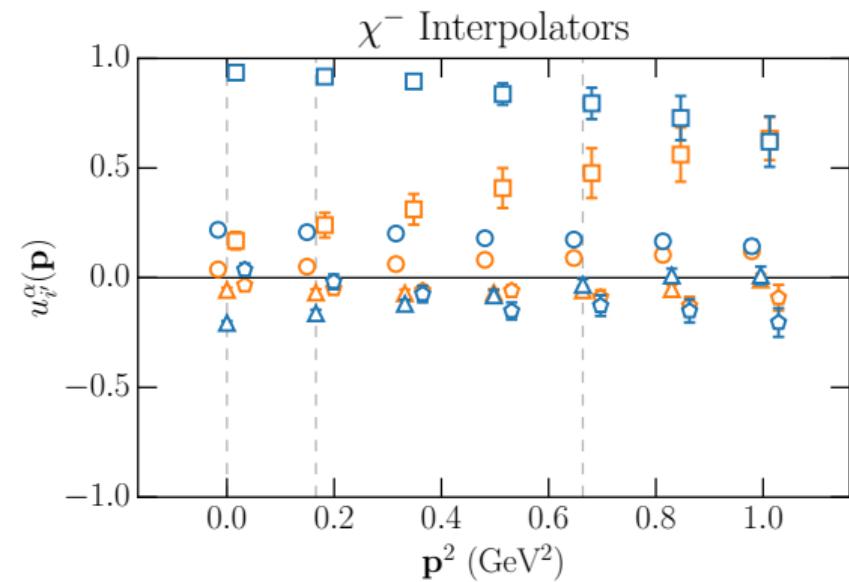
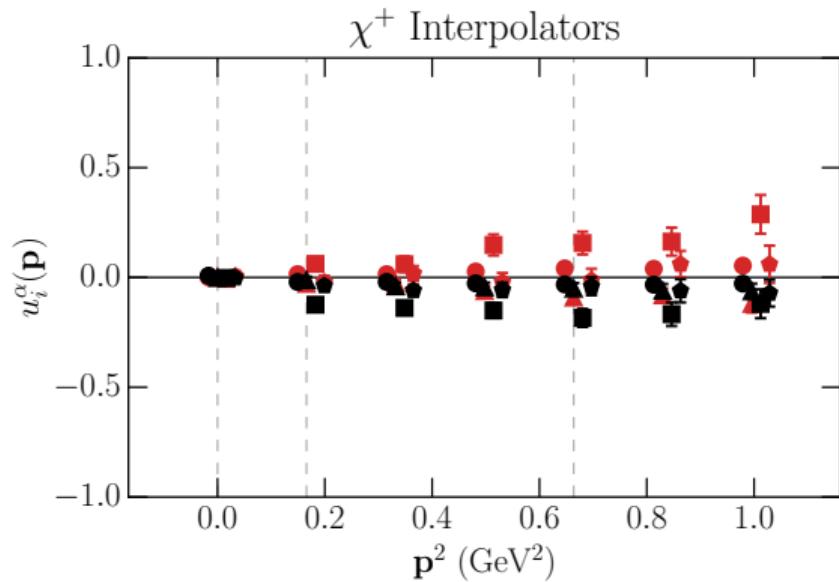
Effective energy

Ground state



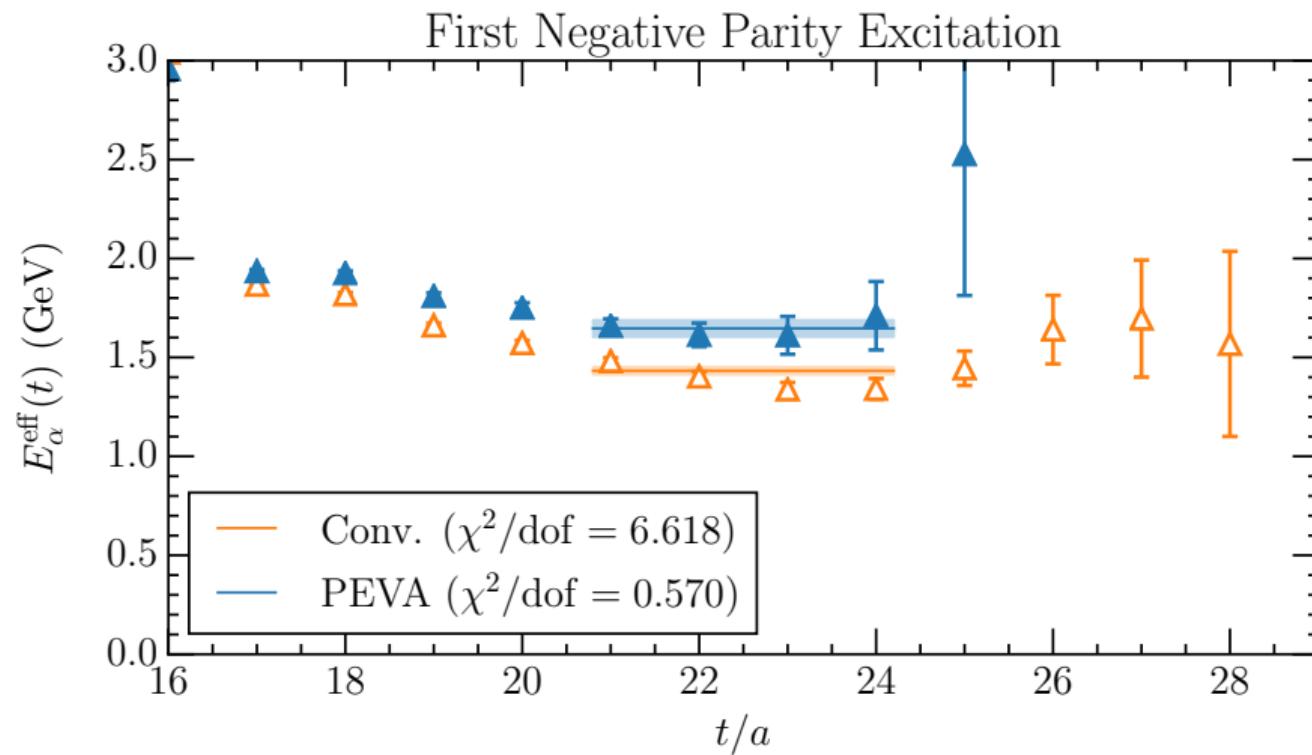
Eigenvector components

First negative parity excitation



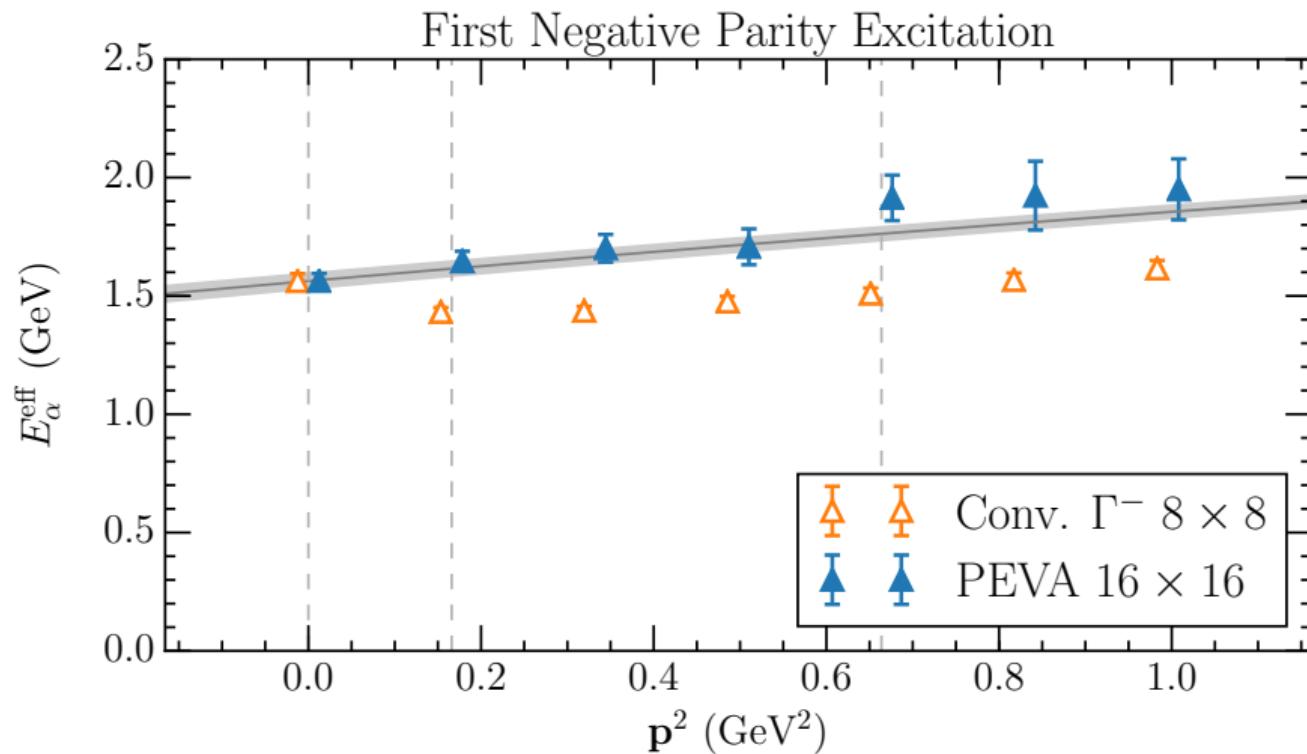
Effective energy

First negative parity excitation - $\mathbf{p}^2 \simeq 0.166 \text{ GeV}^2$



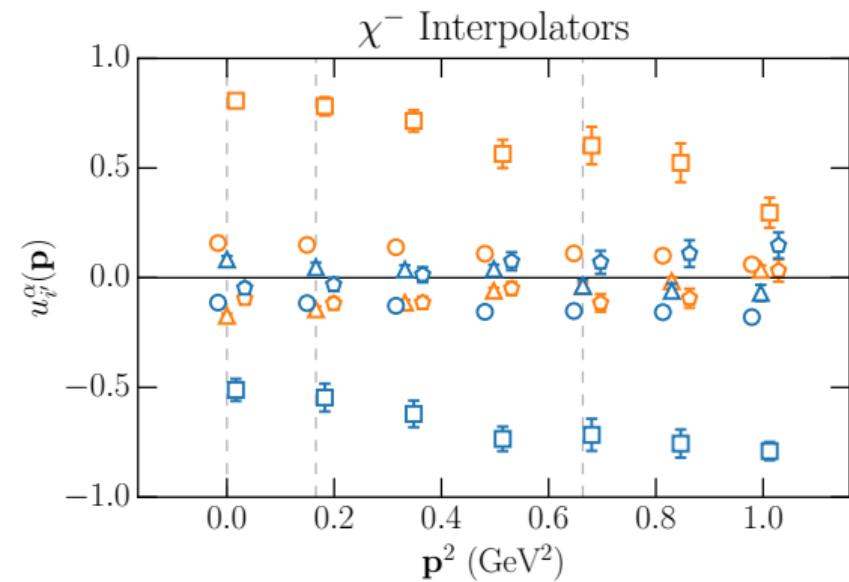
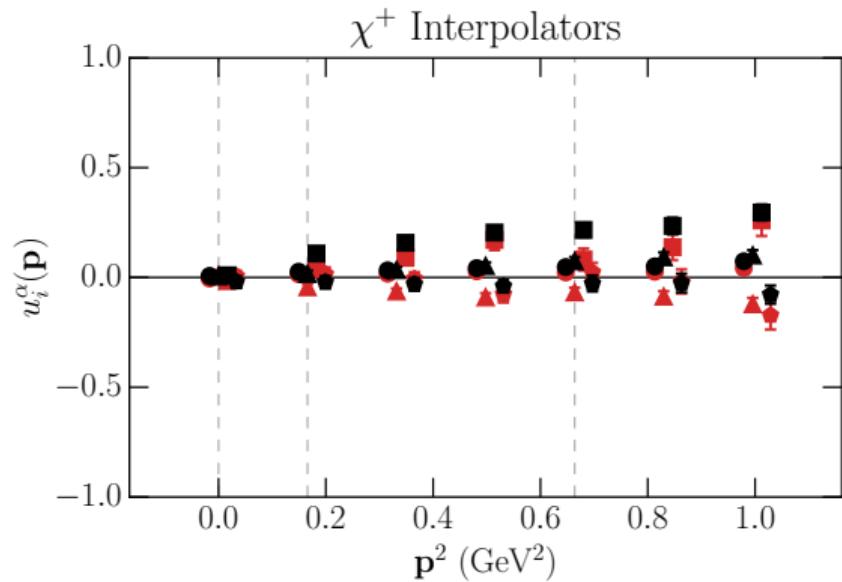
Effective energy

First negative parity excitation



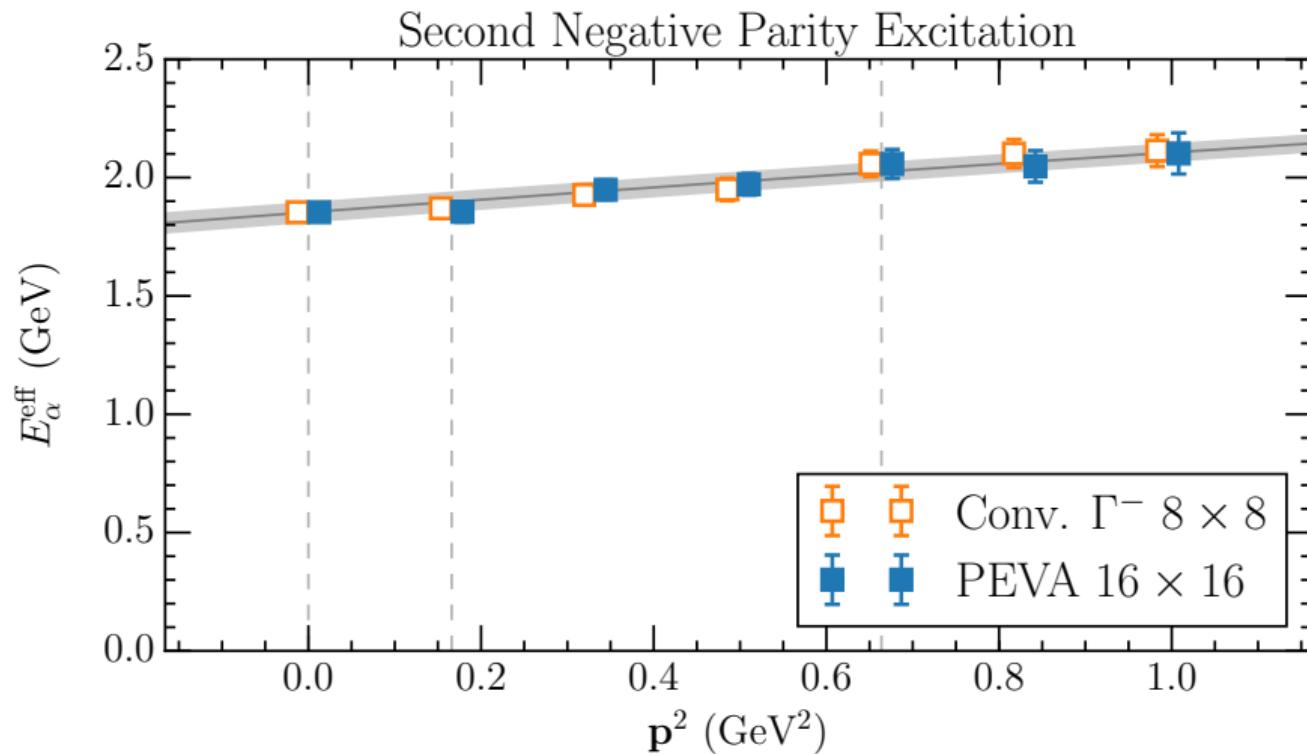
Eigenvector components

Second negative parity excitation



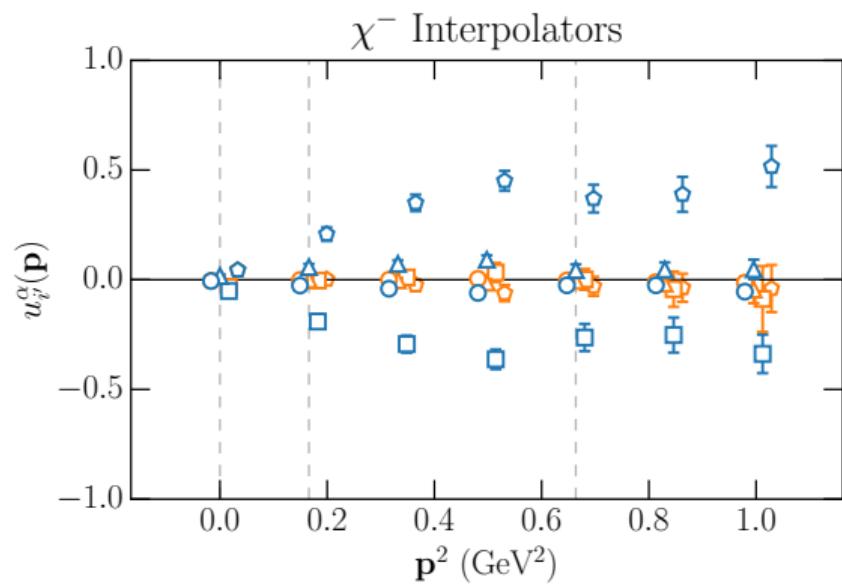
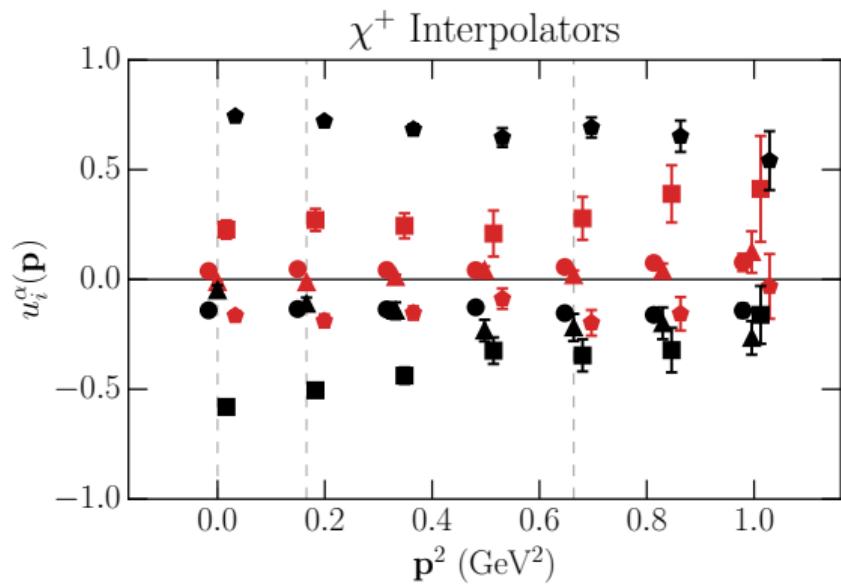
Effective energy

Second negative parity excitation



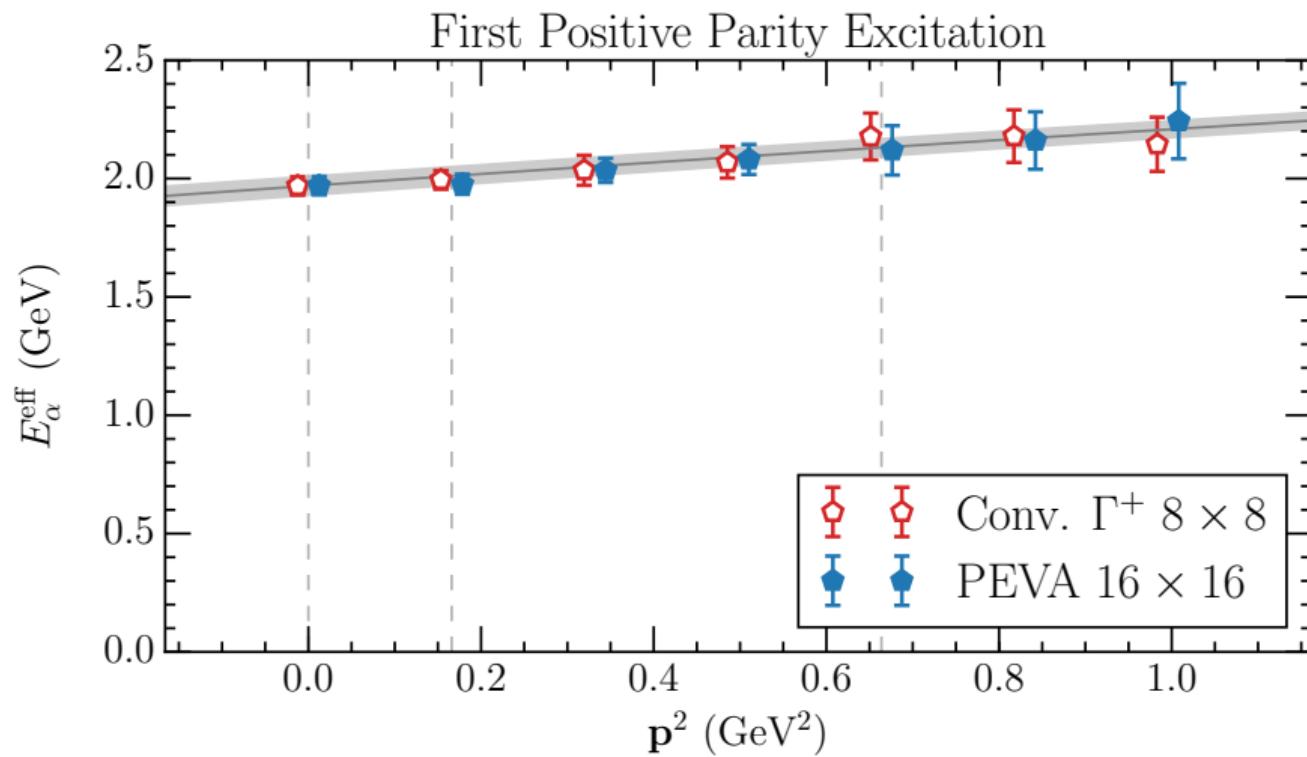
Eigenvector components

First positive parity excitation



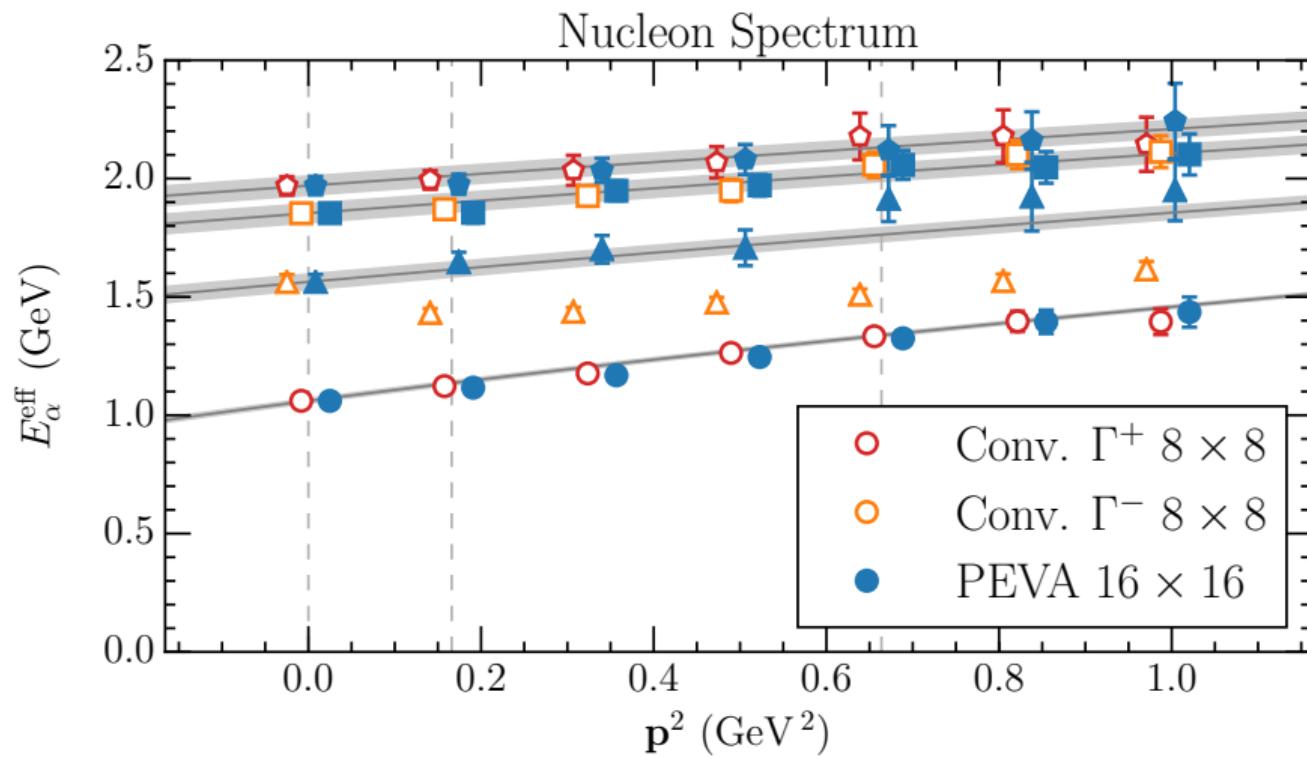
Effective energy

First positive parity excitation



Effective energy

Nucleon spectrum



Baryon structure via variational analysis

Step 1: Construct correlation matrix

Start with basis of N operators and construct an $N \times N$ correlation matrix $G_{ij}(\mathbf{p}; \Gamma; t)$.

Step 2: Perform variational analysis

Perform a variational analysis to construct optimised operators.

Step 3: Compute three point correlation function

Use these optimised operators to construct relevant three point correlation functions

Step 4: Extract form factors

Take ratios of three point to to point functions and extract e.g. $G_E(Q^2)$ and $G_M(Q^2)$.

Step 3: Compute three point correlation function

Eigenstate-projected three point correlation function

$$G_{\alpha \rightarrow \beta}^{\mu}(\mathbf{p}', \mathbf{p}; \Gamma; t_2, t_1) := \text{tr} \left[\Gamma \sum_{\mathbf{x}, \mathbf{y}} e^{i\mathbf{p}' \cdot \mathbf{x}} e^{i(\mathbf{p}' - \mathbf{p}) \cdot \mathbf{y}} \langle \Omega | \phi_{\mathbf{p}'}^{\beta}(x) J^{\mu}(y) \bar{\phi}_{\mathbf{p}}^{\alpha}(0) | \Omega \rangle \right]$$

Step 3: Compute three point correlation function

Eigenstate-projected three point correlation function

$$G_{\alpha \rightarrow \beta}^{\mu}(\mathbf{p}', \mathbf{p}; \Gamma; t_2, t_1) := \text{tr} \left[\Gamma \sum_{\mathbf{x}, \mathbf{y}} e^{i\mathbf{p}' \cdot \mathbf{x}} e^{i(\mathbf{p}' - \mathbf{p}) \cdot \mathbf{y}} \langle \Omega | \phi_{\mathbf{p}'}^{\beta}(x) J^{\mu}(y) \bar{\phi}_{\mathbf{p}}^{\alpha}(0) | \Omega \rangle \right]$$

- Define ratio

$$R_{\alpha \rightarrow \beta}^{\mu}(\mathbf{p}', \mathbf{p}; \Gamma; t_1) := \sqrt{\frac{G_{\alpha \rightarrow \beta}^{\mu}(\mathbf{p}', \mathbf{p}; \Gamma; t_2, t_1) G_{\beta \rightarrow \alpha}^{\mu}(\mathbf{p}, \mathbf{p}'; \Gamma; t_2, t_1)}{G_{\beta}(\mathbf{p}'; t_2) G_{\alpha}(\mathbf{p}; t_2)}}$$

Step 3: Compute three point correlation function

Eigenstate-projected three point correlation function

$$G_{\alpha \rightarrow \beta}^{\mu}(\mathbf{p}', \mathbf{p}; \Gamma; t_2, t_1) := \text{tr} \left[\Gamma \sum_{\mathbf{x}, \mathbf{y}} e^{i\mathbf{p}' \cdot \mathbf{x}} e^{i(\mathbf{p}' - \mathbf{p}) \cdot \mathbf{y}} \langle \Omega | \phi_{\mathbf{p}'}^{\beta}(x) J^{\mu}(y) \bar{\phi}_{\mathbf{p}}^{\alpha}(0) | \Omega \rangle \right]$$

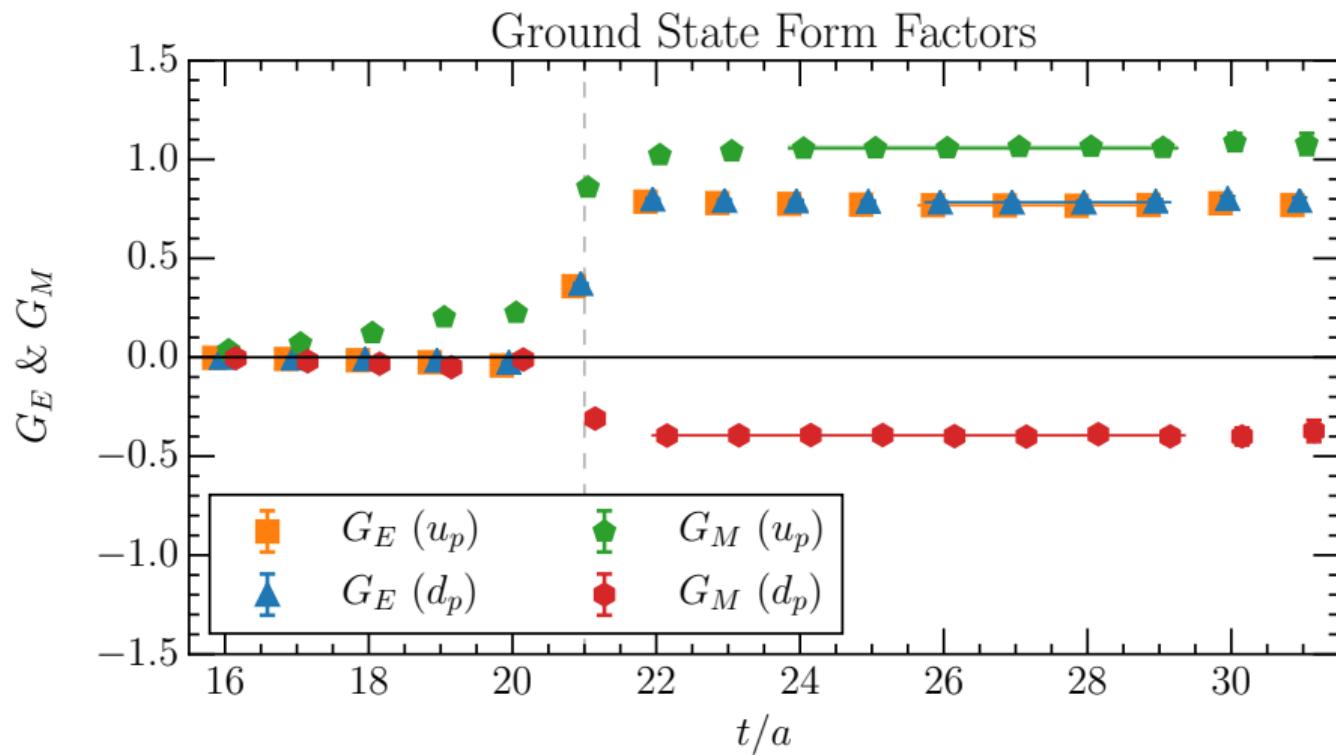
- Define ratio

$$R_{\alpha \rightarrow \beta}^{\mu}(\mathbf{p}', \mathbf{p}; \Gamma; t_1) := \sqrt{\frac{G_{\alpha \rightarrow \beta}^{\mu}(\mathbf{p}', \mathbf{p}; \Gamma; t_2, t_1) G_{\beta \rightarrow \alpha}^{\mu}(\mathbf{p}, \mathbf{p}'; \Gamma; t_2, t_1)}{G_{\beta}(\mathbf{p}'; t_2) G_{\alpha}(\mathbf{p}; t_2)}}$$

- Extract $G_E(Q^2)$ and $G_M(Q^2)$ from ratios

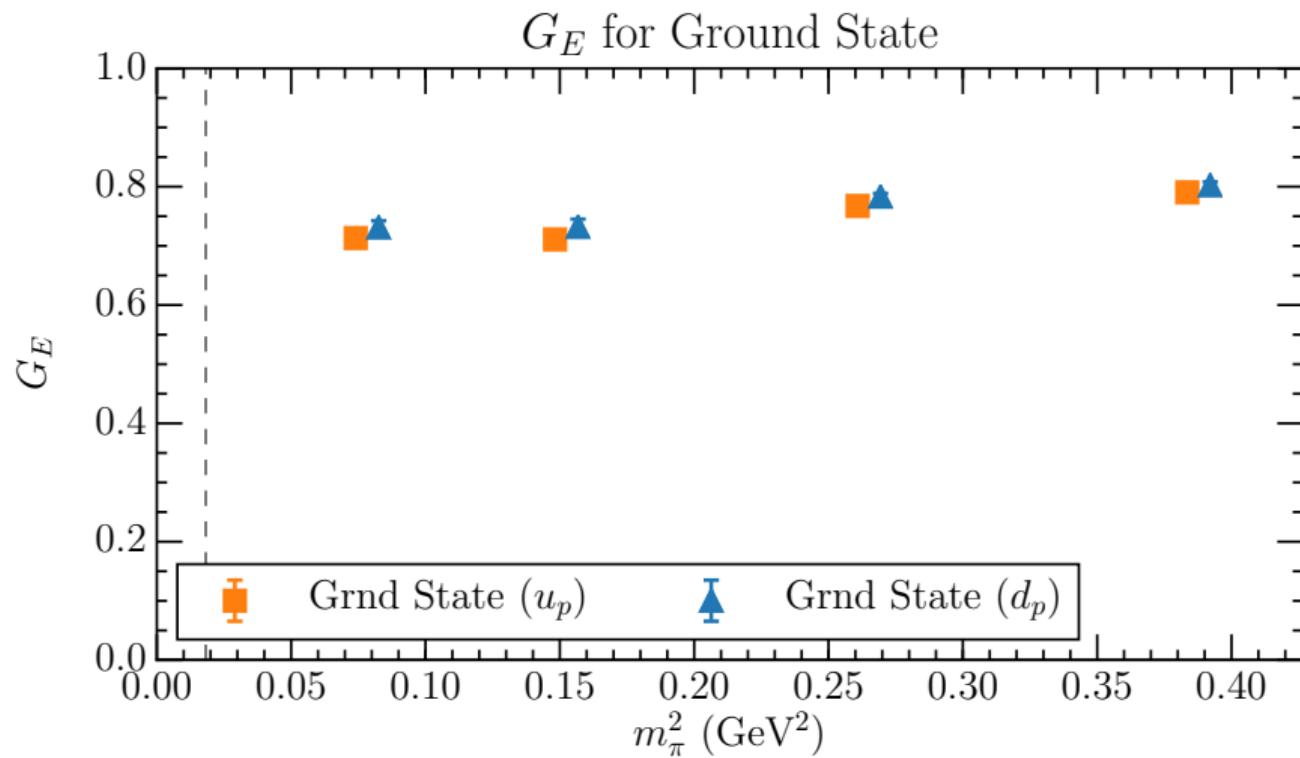
Step 4: Extract form factors

Fits to ground state form factors



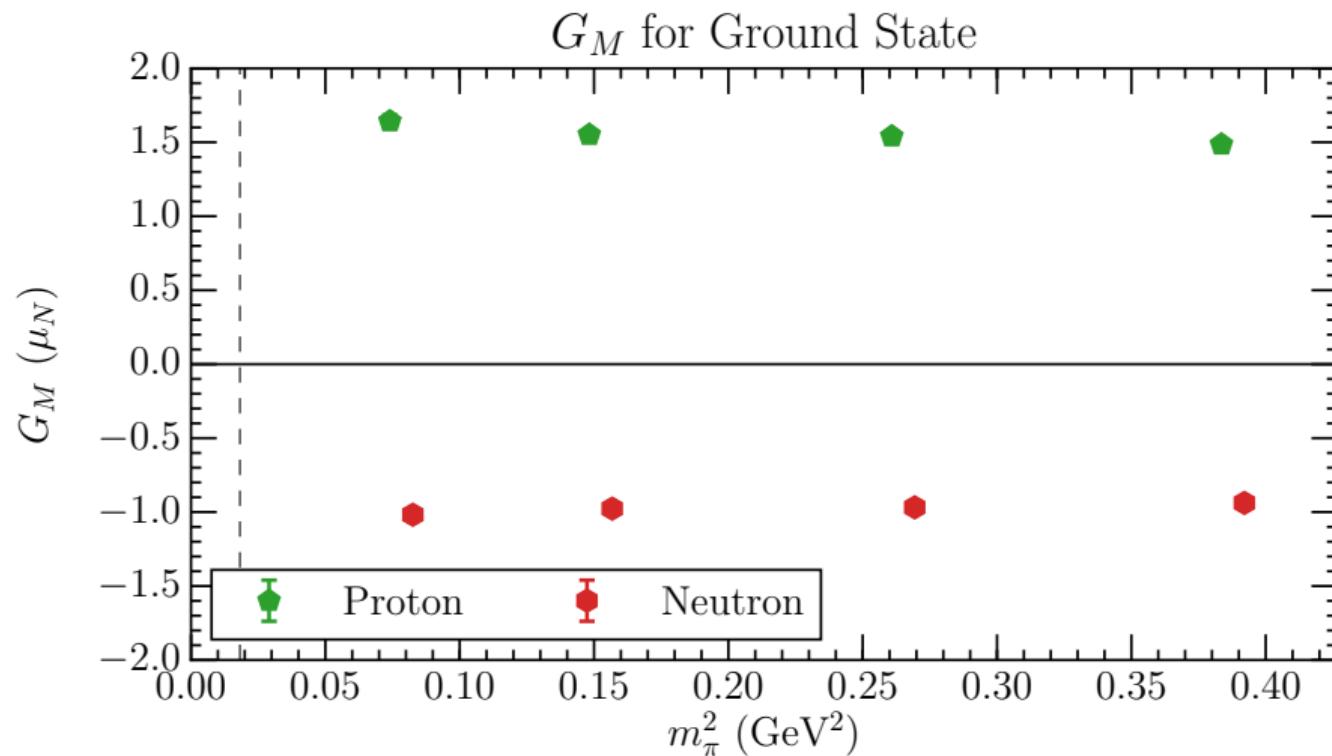
Step 4: Extract form factors

$G_E(Q^2 = 0.15(1) \text{ GeV}^2)$ for ground state



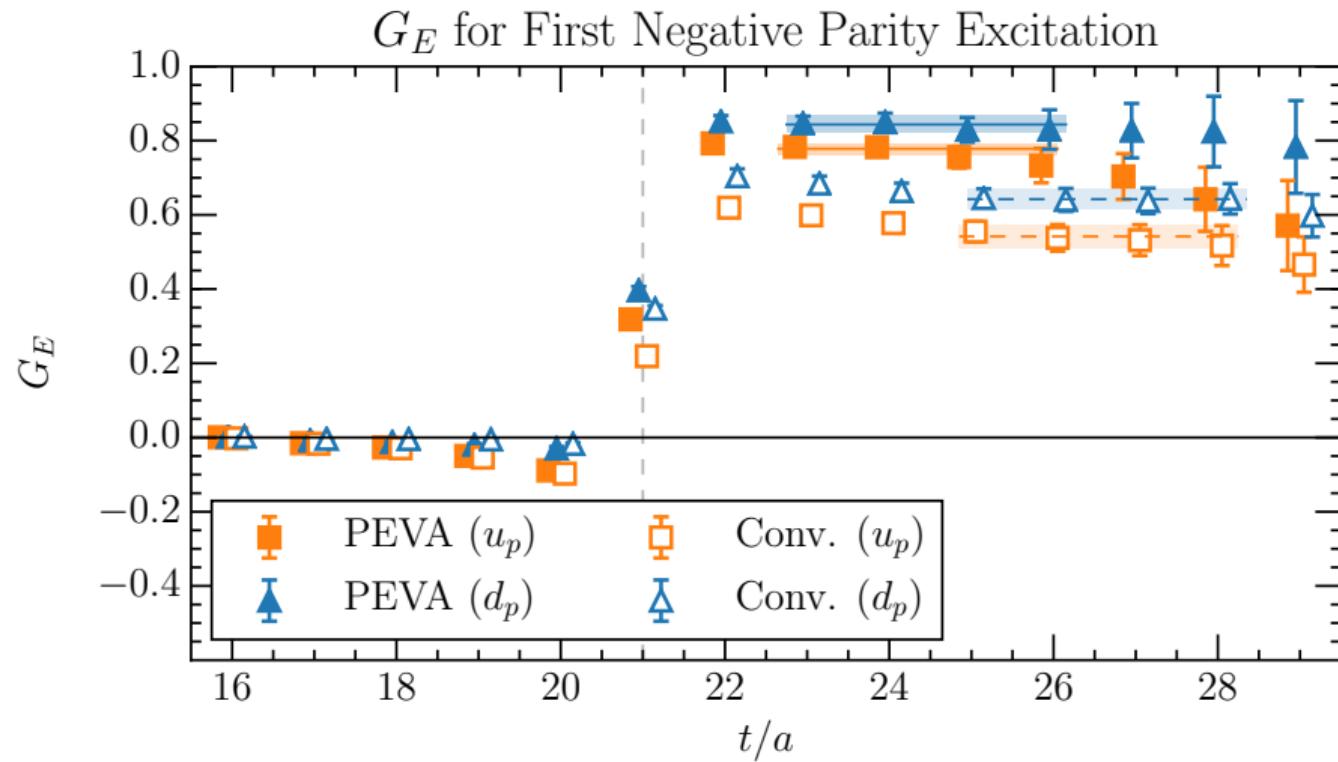
Step 4: Extract form factors

$G_M(Q^2 = 0.15(1) \text{ GeV}^2)$ for ground state



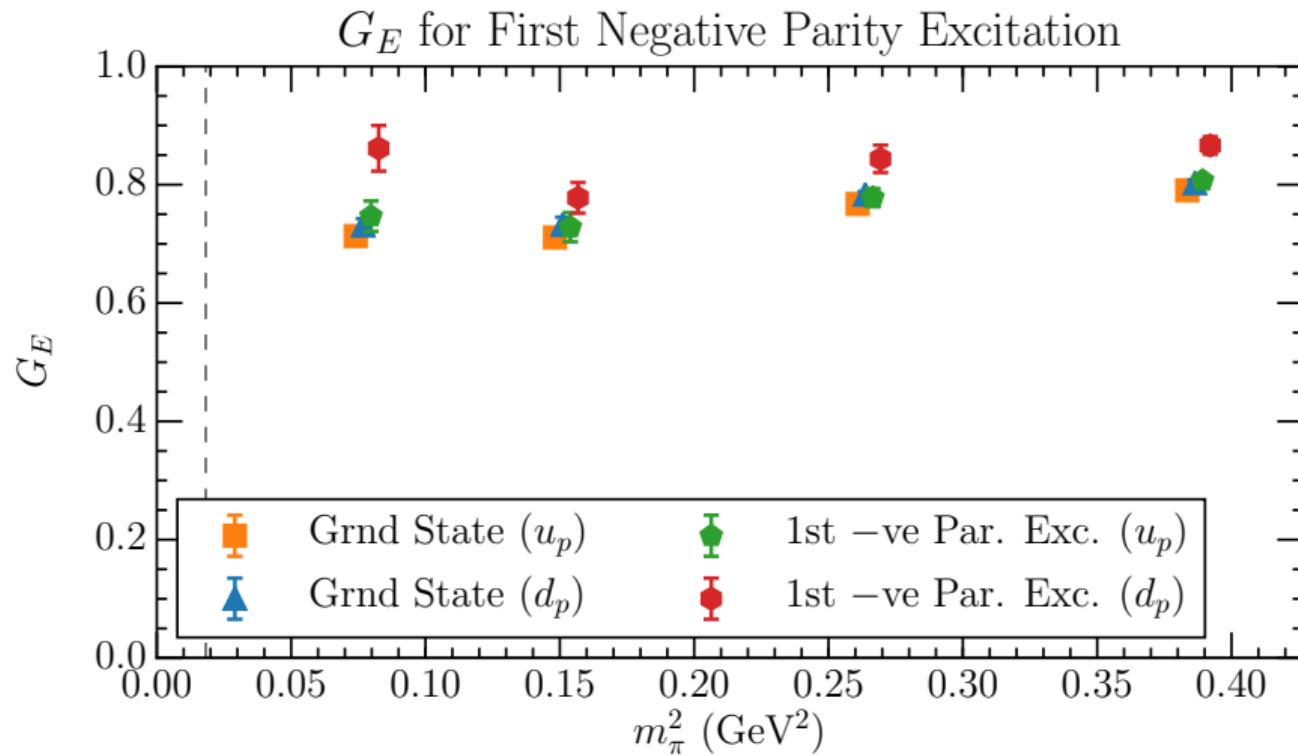
Step 4: Extract form factors

Fits to G_E for first negative parity excitation



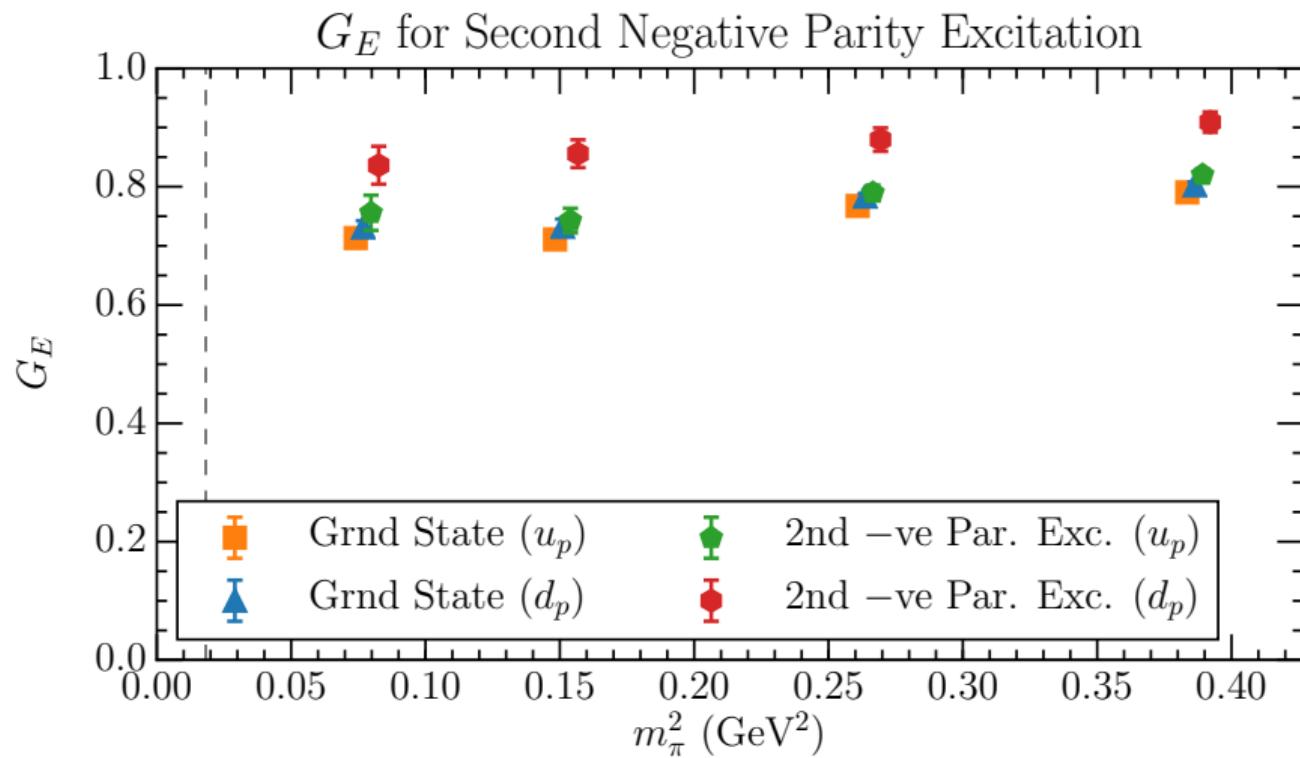
Step 4: Extract form factors

$G_E(Q^2 = 0.15(1) \text{ GeV}^2)$ for first negative parity excitation



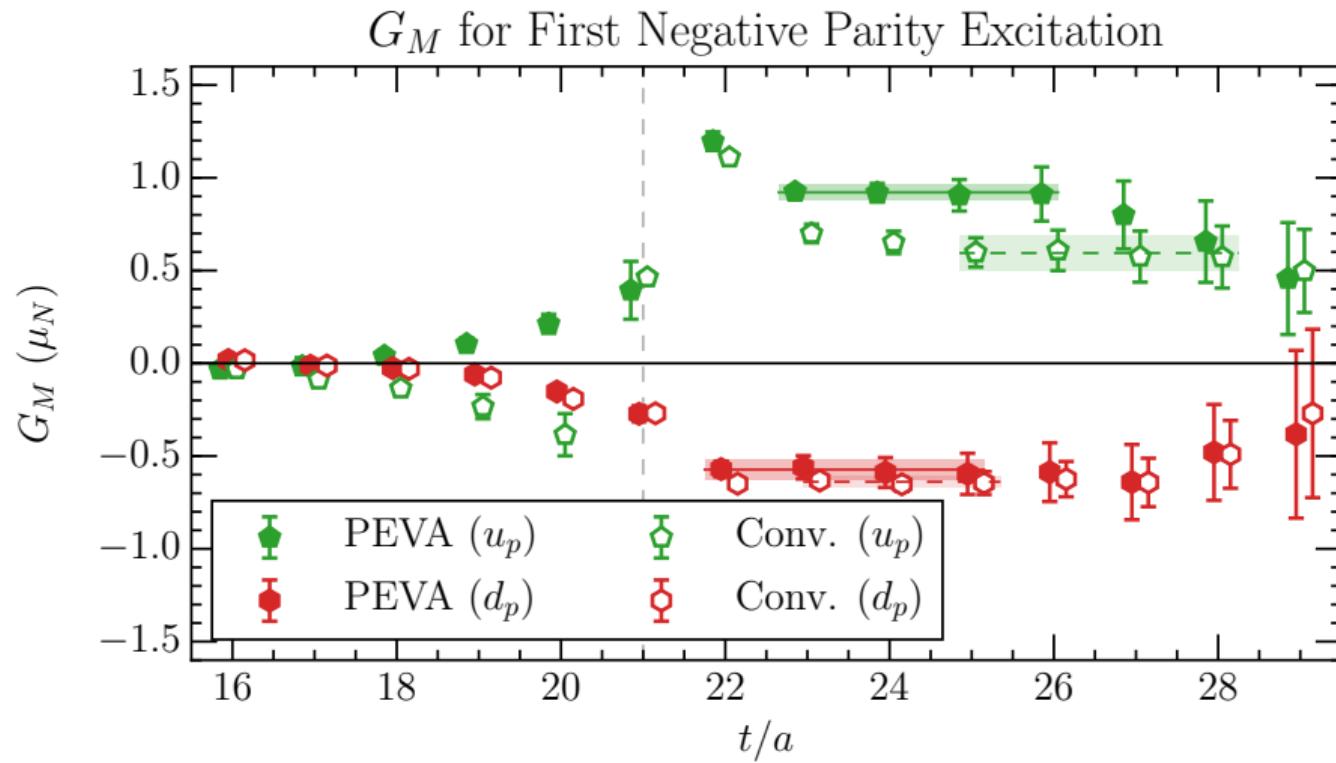
Step 4: Extract form factors

$G_E(Q^2 = 0.15(1) \text{ GeV}^2)$ for second negative parity excitation



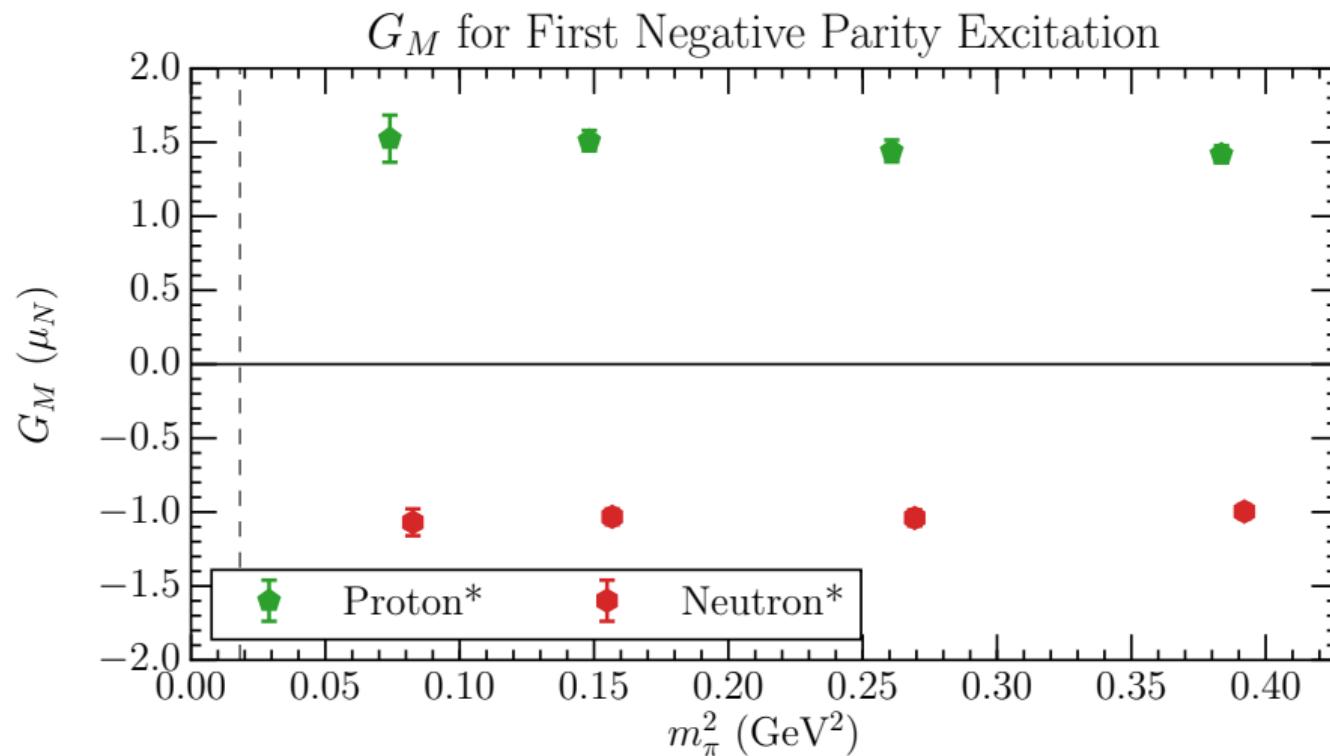
Step 4: Extract form factors

Fits to G_M for first negative parity excitation



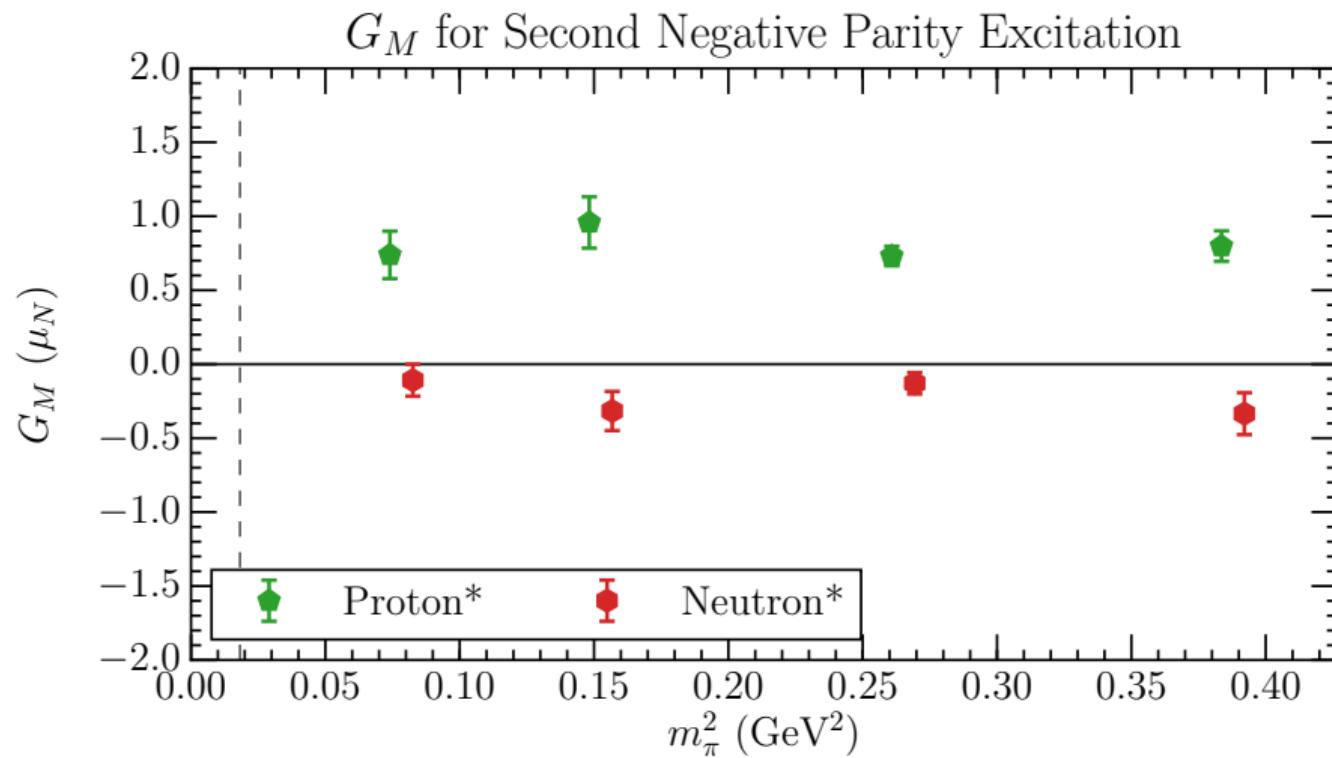
Step 4: Extract form factors

$G_M(Q^2 = 0.15(1) \text{ GeV}^2)$ for first negative parity excitation



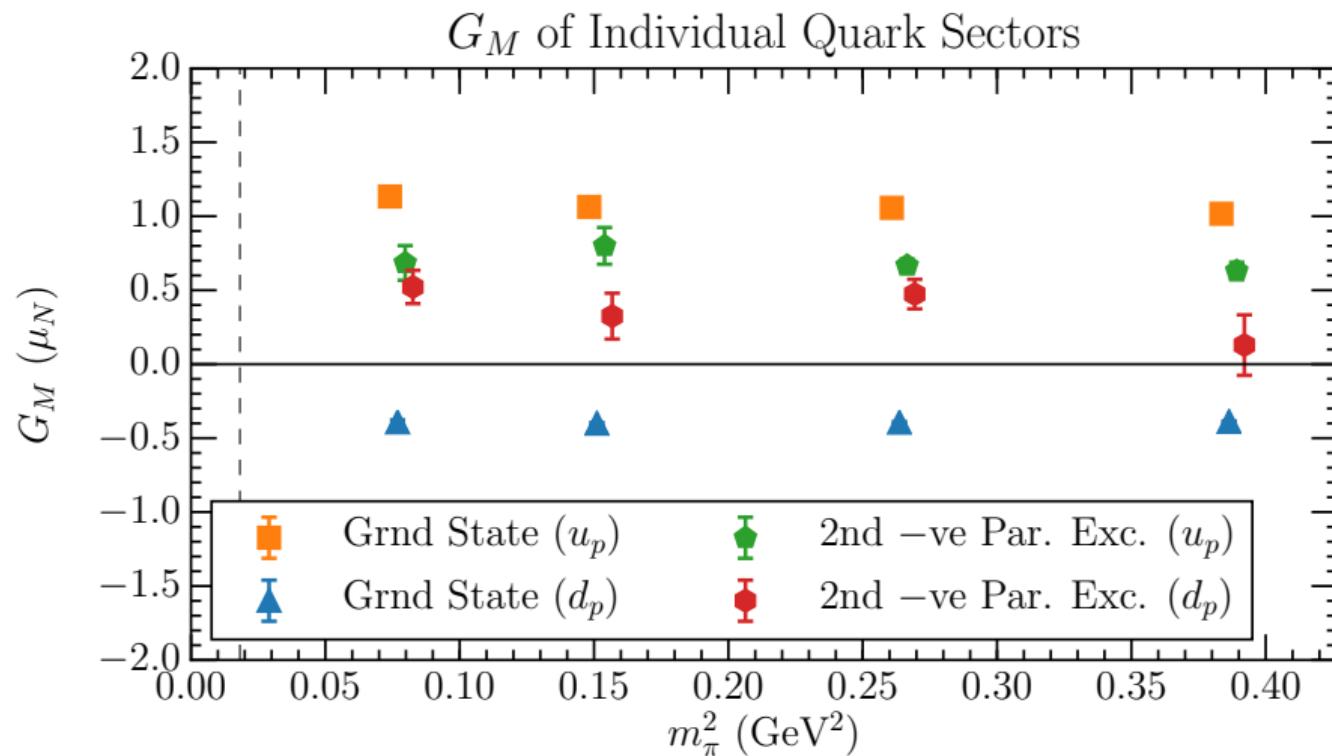
Step 4: Extract form factors

$G_M(Q^2 = 0.15(1) \text{ GeV}^2)$ for second negative parity excitation



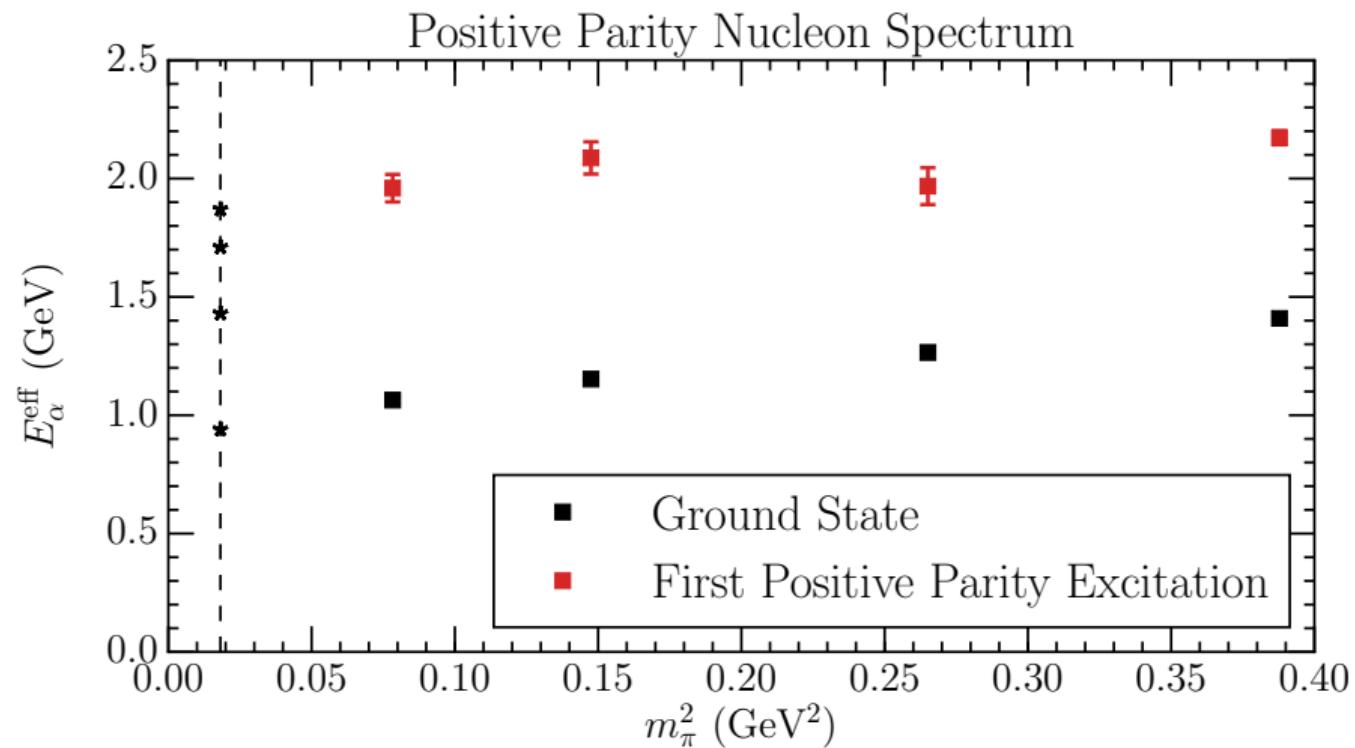
Step 4: Extract form factors

$G_M(Q^2 = 0.15(1) \text{ GeV}^2)$ for second negative parity excitation



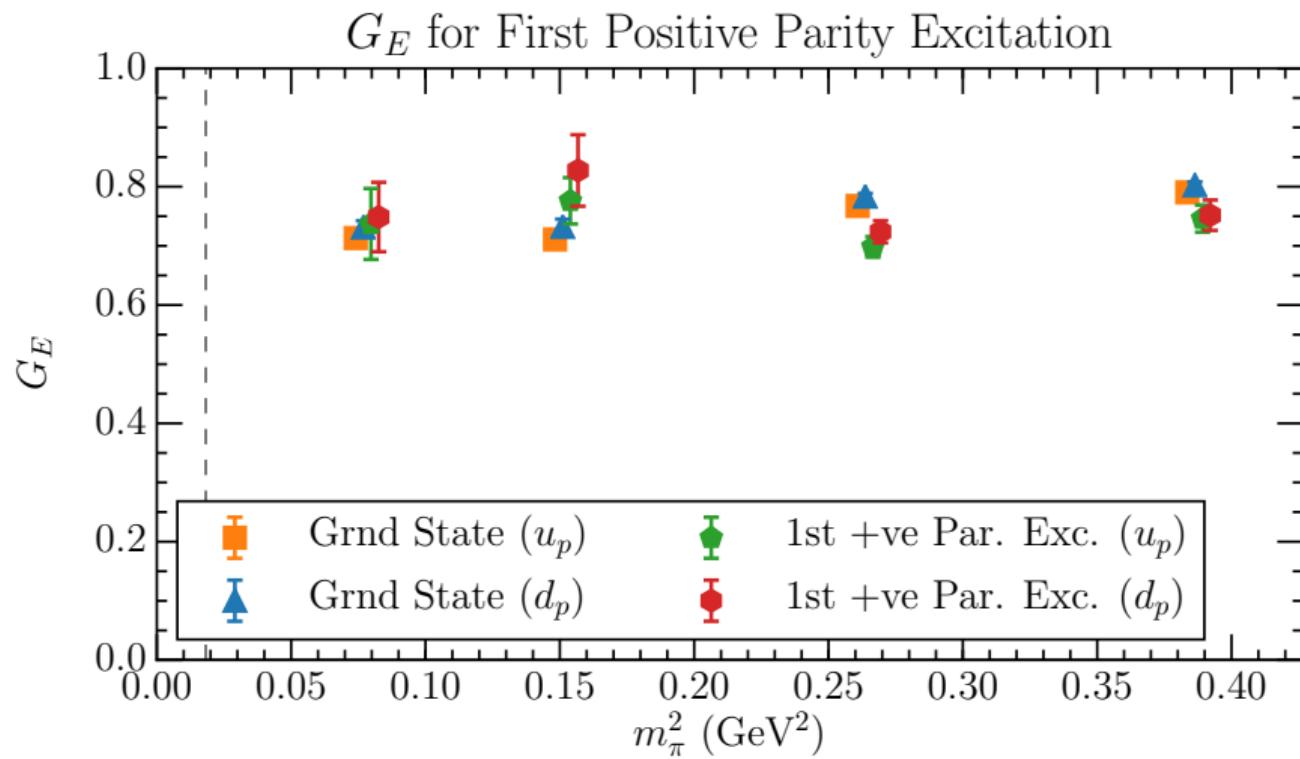
Nucleon mass spectrum

Positive Parity



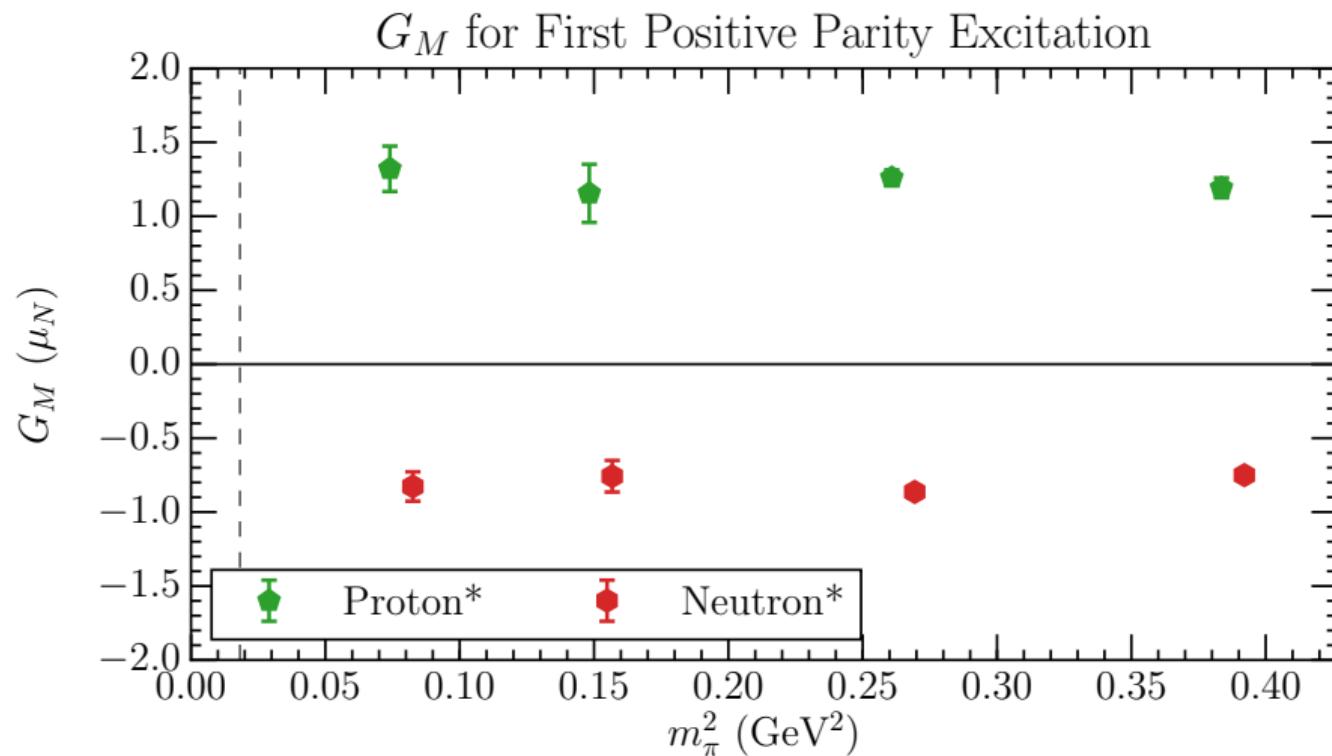
Step 4: Extract form factors

$G_E(Q^2 = 0.15(1) \text{ GeV}^2)$ for first positive parity excitation



Step 4: Extract form factors

$G_M(Q^2 = 0.15(1) \text{ GeV}^2)$ for first positive parity excitation



Conclusion

- Conventional baryon spectroscopy at nonzero momentum is contaminated by opposite parity states

Conclusion

- Conventional baryon spectroscopy at nonzero momentum is contaminated by opposite parity states
- The PEVA technique can effectively remove these opposite parity contaminations

Conclusion

- Conventional baryon spectroscopy at nonzero momentum is contaminated by opposite parity states
- The PEVA technique can effectively remove these opposite parity contaminations
- Clear effect on two point function for lowest lying negative parity excitation

Conclusion

- Conventional baryon spectroscopy at nonzero momentum is contaminated by opposite parity states
- The PEVA technique can effectively remove these opposite parity contaminations
- Clear effect on two point function for lowest lying negative parity excitation
- Has significant effects on three point functions for excited states

Conclusion

- Conventional baryon spectroscopy at nonzero momentum is contaminated by opposite parity states
- The PEVA technique can effectively remove these opposite parity contaminations
- Clear effect on two point function for lowest lying negative parity excitation
- Has significant effects on three point functions for excited states
- Calculations of finite volume baryon transition moments are almost complete

Conclusion

- Conventional baryon spectroscopy at nonzero momentum is contaminated by opposite parity states
- The PEVA technique can effectively remove these opposite parity contaminations
- Clear effect on two point function for lowest lying negative parity excitation
- Has significant effects on three point functions for excited states
- Calculations of finite volume baryon transition moments are almost complete
- This is an important step towards making contact with experiment

More information

“Parity-expanded variational analysis for nonzero momentum”

F. M. Stokes, W. Kamleh, D. B. Leinweber, M. S. Mahbub, B. J. Menadue, B. J. Owen

Phys. Rev. D **92** (2015) 11, 114506

doi:10.1103/PhysRevD.92.114506

arXiv:1302.4152 (hep-lat).