

The 'Exotic Glue' Structure Function

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Motivation

Electron Ion Collider: The Next QCD Frontier

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Understanding the glue that binds us all

'Exotic' Glue in the Nucleus



'Exotic' Glue in the Nucleus



'Exotic' Glue

Contributions to gluon observables that are not from nucleon degrees of freedom.

Exotic glue operator:

operator in nucleon = 0 operator in nucleus $\neq 0$

Jaffe and Manohar (1989)

Leading-twist, double-helicity-flipping structure function $\Delta(x,Q^2)$

- Clear signature for exotic glue in nuclei with spin ≥ 1: NO analogous twist-2 quark PDF → unambiguous
- In single hadrons: gluon transversity structure function
- Experimentally measurable (JLab LOI 2016)
- Moments are calculable using lattice QCD

First Lattice Study: arXiv:1606.04505 (PRD) • First moment of $\Delta(x, Q^2)$ in spin-1 ϕ meson



Double helicity flip amplitude:

$$\Delta(x, Q^2) = A_{+-,-+} = A_{-+,+-}$$

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Double helicity flip amplitude:

Photon helicity

$$\Delta(x, Q^2) = A_{+-, -+} = A_{-+, +-}$$

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Exotic Glue in the Nucleus



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Photon helicity Target helicity

Optical theorem, dispersion relation for hadronic forward scatt. amplitude, analytic continuation give **moments**:



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Operator Product Expansion to relate to matrix elements of operator

Gluonic Operator

$$\langle pE' | S[G_{\mu\mu_1} \overleftrightarrow{D}_{\mu_3} \dots \overleftrightarrow{D}_{\mu_n} G_{\nu\mu_2}] | pE \rangle$$

$$= (-2i)^{n-2} S[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu})(p_{\nu}E_{\mu_2} - p_{\mu_2}E_{\nu})$$

$$+ (\mu \leftrightarrow \nu)] p_{\mu_3} \dots p_{\mu_n} \underline{A_n(Q^2)} \dots,$$
Reduced Matrix Element

Optical theorem, dispersion relation for hadronic forward scatt. amplitude, analytic continuation give **moments**:



Operator Product Expansion to relate to matrix elements of operator

$$\begin{array}{c} & \left\langle pE' \left[\overbrace{S[G_{\mu\mu_1} \overleftarrow{D}_{\mu_3} \dots \overrightarrow{D}_{\mu_n} G_{\nu\mu_2}} \right] | pE \right\rangle \\ & \left\langle \overbrace{P} = (-2i)^{n-2} \underbrace{S[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu})(p_{\nu}E_{\mu_2} - p_{\mu_2}E_{\nu})} \right. \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu_1})(p_{\nu}E_{\mu_2} - p_{\mu_2}E_{\nu}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu_1})(p_{\nu}E_{\mu_2} - p_{\mu_2}E_{\nu}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu_1})(p_{\nu}E_{\mu_2} - p_{\mu_2}E_{\nu}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu_1})(p_{\nu}E_{\mu_2} - p_{\mu_2}E_{\nu}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu_1})(p_{\nu}E_{\mu_2} - p_{\mu_2}E_{\nu}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu_1})(p_{\nu}E_{\mu_2} - p_{\mu_2}E_{\nu}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu_1})(p_{\nu}E_{\mu_2} - p_{\mu_2}E_{\nu}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\nu}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\nu}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\mu_2}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_1}E'^*_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\mu_2}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_2}E_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\mu_2}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_1} - p_{\mu_2}E_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\mu_2}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\mu_2}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\mu_2}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\mu_2}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu}E_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2} - p_{\mu_2}E_{\mu_2}) \right] \\ & \left[(p_{\mu}E'^*_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2} - p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_2})(p_{\mu_2}E_{\mu_$$

Lattice QCD

Numerical first-principles approach

Discretise space-time (4D box)

Lattice spacing *a*, volume $L^3 \times T$ order $32^3 \times 64 \approx 2 \times 10^6$ lattice sites



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Lattice simulation in spin-1 ϕ meson

Lattice Details

Luscher-Weisz gauge action with a clover-improved quark action

L/a	T/a	eta	am_l	am_s
24	64	6.1	-0.2800	-0.2450
a (fm)	L (fm)	T (fm)	m_π (MeV)	m_K (MeV)
0.1167(16)	2.801(29)	7.469(77)	450(5)	596(6)
m_{ϕ} (MeV)	$m_{\pi}L$	$m_{\pi}T$	$N_{ m cfg}$	$N_{ m src}$
1040(3)	6.390	17.04	1042	10^{5}

- All ϕ polarization states ({1,2,3} or {+,-,0})
 - on-diagonal
 - off-diagonal
- $\bullet\,$ Momenta up to (1,1,1) in lattice units (1 unit \sim 0.4GeV)
- Different discretisations of the operator (different irreps.)

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L/a	SYSTE	am_s		
24	Quark m	-0.2450		
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Extraction of A_2



We calculate on the lattice:

 $\begin{bmatrix} \frac{C_{3\text{pt}}^{EE'}}{C_{2\text{pt}}^{EE'}} \end{bmatrix} (t_{\text{sink}}, \tau) \propto A_2, \qquad 0 \ll \tau \ll t_{\text{sink}}$

Extraction of A_2 : 3pt/2pt ratio



Extraction of A_2 : 3pt/2pt ratio



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Explore gluon structure of ϕ meson more generally

Soffer bound for transversity quark distributions:

$$|\delta q(x)| \leq \frac{1}{2} \left(q(x) + \Delta q(x) \right)$$

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Direct analogue for leading moments of gluon distributions:

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Direct analogue for leading moments of gluon distributions:

$$\begin{array}{ccc}
G_{\mu\mu_1}G_{\nu\mu_2} & G_{\mu_1\alpha}G_{\mu_2}^{\alpha} \\
& & & \\
\hline A_2 & \leq \frac{1}{2}B_2 & \widetilde{G}_{\mu_1\alpha}G_{\mu_2}^{\alpha} \to 0
\end{array}$$

UNRENORMALISED reduced matrix element: ϕ meson



If we assume approx. the same renormalisation for A_2 and B_2 :



First two moments of quark distributions: Soffer bound saturated to 80% (lattice QCD, Diehl *et al.* 2005)

Summary

ROBUST NON-ZERO signal for 'exotic glue' operator in the ϕ meson

Proof of principle: similar signal in a nucleus \Leftrightarrow exotic glue

Explore gluon structure of hadrons more generally e.g., Soffer bound analogue

BUT: SYSTEMATICS IGNORED ⇒ no physically meaningful number (yet)