Effects of meson-nucleon dynamics in a relativistic approach to medium-mass nuclei

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• Relativistic Nuclear Field Theory (RNFT) solutions to the nuclear many-body problem

Outline

- Recent developments: pion degrees of freedom
- Spin-isospin nuclear response:
 - Gamow-Teller resonance (GTR)
 - Spin-Dipole resonance and higher multipoles: the onset of pion condensation in medium-mass nuclei
- Pion-exchange contributions to the nucleonic self-energy:
 - Shell structure of tin isotopes
- Conclusions and perspectives

Relativistic Nuclear Field Theory (RNFT):

- RNFT as a solution: microscopic, universal, connecting scales from Quantum Hadrodynamics to emergent collective phenomena
 - Lagrangian for mesons and nucleons constrained by QCD symmetries and sum rules
- Lorentz covariance: ~5-10% accuracy at the excitation energy of interest (grows with energy)
- Spin-orbit and tensor "forces" are naturally included
 - Fewer parameters; hidden correlations minimized (4-10 universal parameters)
- Natural extension to the inclusion of the delta isobar, to higher excitation energy ~200-300 MeV and to hypernuclei
 - Non-perturbative self-consistent response theory with high-order NN correlations





Systematic expansion in the RNFT: single-nucleon self-energy



Systematic expansion in the RNFT: nuclear response function

Vibration (phonon) vertex:



Relativistic (Quasiparticle) Random Phase Approximation (QRPA)

2q+phonon coupling amplitude:



2q+Nphonon coupling amplitude:





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Isospin transfer response function: proton-neutron relativistic quasiparticle time blocking approximation (pn-RQTBA)

Response
$$R(\omega) = \tilde{R}^0(\omega) + \tilde{R}^0(\omega)\overline{W}(\omega)R(\omega)$$

Interaction

$$W(\omega) = \underbrace{V_{\rho} + V_{\pi} + V_{\delta\pi}}_{\checkmark} + \underbrace{\Phi(\omega)}_{\checkmark}$$

Static:

$$R(Q)RPA \begin{cases} V_{\rho}(1,2) = g_{\rho}^{2}\vec{\tau_{1}}\vec{\tau_{2}}(\beta\gamma^{\mu})_{1}(\beta\gamma_{\mu})_{2}D_{\rho}(\mathbf{r_{1}},\mathbf{r_{2}}) \\ V_{\pi}(1,2) = -\left(\frac{f_{\pi}}{m_{\pi}}\right)^{2}\vec{\tau_{1}}\vec{\tau_{2}}(\boldsymbol{\Sigma_{1}}\nabla_{1})(\boldsymbol{\Sigma_{2}}\nabla_{2})D_{\pi}(\mathbf{r_{1}},\mathbf{r_{2}}) \\ V_{\delta\pi}(1,2) = g'\left(\frac{f_{\pi}}{m_{\pi}}\right)^{2}\vec{\tau_{1}}\vec{\tau_{2}}\boldsymbol{\Sigma_{1}}\boldsymbol{\Sigma_{2}}\delta(\mathbf{r_{1}}-\mathbf{r_{2}}) \\ fixed strength \end{cases}$$

Dynamic (retardation), 2-nd order:

quasiparticlevibration coupling

in time blocking approximation

$$\Phi_{k_{1}k_{4},k_{2}k_{3}}^{\eta}(\omega) = \\
= \sum_{\mu\xi} \delta_{\eta\xi} \left[\delta_{k_{1}k_{3}} \sum_{k_{6}} \frac{\gamma_{\mu;k_{6}k_{2}}^{\eta;-\xi} \gamma_{\mu;k_{6}k_{4}}^{\eta;-\xi}}{\eta\omega - E_{k_{1}} - E_{k_{6}} - \Omega_{\mu}} + \delta_{k_{2}k_{4}} \sum_{k_{5}} \frac{\gamma_{\mu;k_{1}k_{5}}^{\eta;\xi} \gamma_{\mu;k_{3}k_{5}}^{\eta;\xi*}}{\eta\omega - E_{k_{5}} - E_{k_{2}} - \Omega_{\mu}} \\
- \left(\frac{\gamma_{\mu;k_{1}k_{3}}^{\eta;\xi} \gamma_{\mu;k_{2}k_{4}}^{\eta;-\xi*}}{\eta\omega - E_{k_{3}} - E_{k_{2}} - \Omega_{\mu}} + \frac{\gamma_{\mu;k_{3}k_{1}}^{\eta;\xi*} \gamma_{\mu;k_{4}k_{2}}^{\eta;-\xi}}{\eta\omega - E_{k_{1}} - E_{k_{4}} - \Omega_{\mu}} \right) \right]$$

Spin-isospin excitations in proton-neutron relativistic time blocking approximation (pn-RTBA)



E.L., B.A. Brown, D.-L. Fang, T. Marketin, R.G.T. Zegers, PLB 730, 307 (2014)



T. Marketin, E.L., D. Vretenar, P. Ring, PLB 706, 477 (2012).

Spin-isospin response in open-shell nuclei: see talk of Caroline Robin tomorrow afternoon at Nuclear Structure A

Spin-dipole resonance in neutron-rich nuclei



$\Delta L = 1$
$\Delta T = 1$
∆S = 1
$\lambda = 0, 1, 2$

Earlier studies:

 $P_{\pm}^{\lambda} = \sum r(i) \left[\boldsymbol{\sigma}(i) \otimes Y_{1}(i) \right]_{\lambda} t_{\pm}(i)$

W.H. Dickhoff et al., PRC 23, 1154 (1981) J. Meyer-Ter-Vehn, Phys. Rep. 74, 323 (1981) A.B. Migdal et al., Phys. Rep. 192, 179 (1990)

p↓

n1

p↑

n↓

Existence of low-lying unnatural parity states indicates that nuclei are close to the pion condensation point. However, it is not clear which observables are sensitive to this phenomenon.





Low-lying states in *ΔT=1* channel and nucleonic self-energy

(N,Z) (N+1,Z-1)

In spectra of neighboring odd-odd nuclei we see low-lying (collective) states with natural and unnatural parities: 0+, 0-, 1+, 1-, 2+, 2-, 3+, 3-,... Their contribution to the nucleonic self-energy is expected to affect single-particle states:



Single-particle states in 100-Sn: effects of pion dynamics

Truncation scheme: phonons below 20 MeV Phonon basis: $\Delta T=0$ phonons: 2+, 3-, 4+, 5-, 6+ $\Delta T=1$ phonons: 0±, 1±, 2±, 3±, 4±, 5±, 6±

Single-particle states:



E.L., Phys. Lett. B 755, 138 (2016)

Next step: isovector ground state correlations (backward going diagrams), in progress. Influence on GTR: see talk of Caroline Robin tomorrow afternoon Relativistic NFT offers a powerful framework for the treatment of meson-nucleon correlations beyond the Hartree-Fock approximation.

Outlook

- Non-perturbative self-consistent response theory based on QHD and including highorder correlations is available for a large class of nuclear excited states.
- ✤ Effects of isospin dynamics are introduced into RNFT. Pion exchange is included non-perturbatively in both shell structure and response calculations.
- Perspectives: dynamical proton-neutron pairing (deuteron transfer channel), higherorder ground-state correlations; deformations etc.

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