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Microscopic studies on nuclear spin-isospin properties --- a personal perspective on covariant density functional theory

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Acknowledgments



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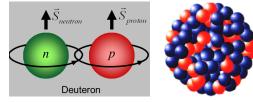
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- C. Stoyanov, D. Tarpanov
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- P. Ring
- G. Colò, V. De Donno, Y.F. Niu, X. Roca-Maza
- Y. Kim, Y. Lim
 - M. Anguiano, M. Moreno-Torres
 - Y. Tanizaki, P.W. Zhao

Research interests and tools

Research interests: Spin and Isospin properties in atomic nuclei

Spin and **Isospin** are essential degrees of freedom in nuclear physics



Relevant studies in nuclear physics × *nuclear astrophysics* × *particle physics*

Research tools: Covariant density functional theory (CDFT)

- Fundamental: Kohn-Sham Density Functional Theory
- Scheme: Yukawa meson-exchange nuclear interactions

$$\mathscr{L} = \bar{\psi} \left[i\gamma^{\mu}\partial_{\mu} - M - g_{\sigma}\sigma - \gamma^{\mu} \left(g_{\omega}\omega_{\mu} + g_{\rho}\vec{\tau} \cdot \vec{\rho}_{\mu} + e\frac{1 - \tau_{3}}{2}A_{\mu} \right) - \frac{f_{\pi}}{m_{\pi}}\gamma_{5}\gamma^{\mu}\partial_{\mu}\vec{\pi} \cdot \vec{\tau} \right] \psi$$
$$+ \frac{1}{2}\partial^{\mu}\sigma\partial_{\mu}\sigma - \frac{1}{2}m_{\sigma}^{2}\sigma^{2} - \frac{1}{4}\Omega^{\mu\nu}\Omega_{\mu\nu} + \frac{1}{2}m_{\omega}^{2}\omega_{\mu}\omega^{\mu} - \frac{1}{4}\vec{R}_{\mu\nu} \cdot \vec{R}^{\mu\nu} + \frac{1}{2}m_{\rho}^{2}\vec{\rho}^{\mu} \cdot \vec{\rho}_{\mu}$$
$$+ \frac{1}{2}\partial_{\mu}\vec{\pi} \cdot \partial^{\mu}\vec{\pi} - \frac{1}{2}m_{\pi}^{2}\vec{\pi} \cdot \vec{\pi} - \frac{1}{4}F^{\mu\nu}F_{\mu\nu}$$
$$\rho_{\mu}(\mathbf{T}=0)$$



Covariant density functional theory

Why DFT?

- applicable to almost whole nuclear chart (ground states × excited states)
- no other method achieves comparable accuracy at the similar computational costs



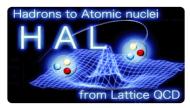
http://www.unedf.org/

Why covariant (relativistic)?

Dirac equation consistent treatment of spin d.o.f. & nuclear saturation properties (3-body effect)

- Lorentz covariant symmetry unification of time-even and time-odd components
- Effective Lagrangian

connections to underlying theories, QCD at low energy



cf. HAL QCD Collaboration

Covariant density functional theory

From Walecka model in 1974 to now

in Jan 2016

International Review of Nuclear Physics - Vol. 10

Relativistic Density Functional Im Nuclear Structure



- □ Concept of Covariant Density Functional Theory (P Ring)
- □ Relativistic Mean-Field Theory (J Meng, P Ring and P W Zhao)

cf. S.-G. Zhou's plenary talk

- Relativistic Mean Field Description of Exotic Nuclei (J Meng, P Ring, P W Zhao and S G Zhou)
- Relativistic Hartree–Fock–Bogoliubov Theory: Ground States and Excitations (W H Long, J Meng and N Van Giai)
- Superheavy Nuclei and Fission Barriers (B N Lu, J Zhao, E G Zhao and S G Zhou)
- Relativistic Symmetries in Nuclear Single-Particle Spectra (J Y Guo, H Z Liang, J Meng and S G Zhou)
- Structure of Hypernuclei in Relativistic Approaches (K Hagino and J M Yao)
- Rotating Nuclei: From Ground State to the Extremes of Spin and Deformation (A V Afanasjev)
- □ Novel Rotational Excitations (J Meng, S Q Zhang and P W Zhao)
- □ Small Amplitude Motion (N Paar and Y Niu)
- Nuclear Shell Structure and Response with Quasiparticle-Vibration Coupling (E Litvinova and P Ring)
- Beyond the Relativistic Mean-Field Approximation Collective Correlations (Z P Li, T Nikšić, D Vretenar and J M Yao)
- Heavy Element in Astrophysical Nucleosynthesis (B H Sun and Z M Niu)
- Relativistic Density Functional Theory for Finite Nuclei and Neutron Stars (J Piekarewicz)
- □ Relativistic Versus Non-Relativistic Mean Field (P-G Reinhard)

Our studies on pseudospin symmetry

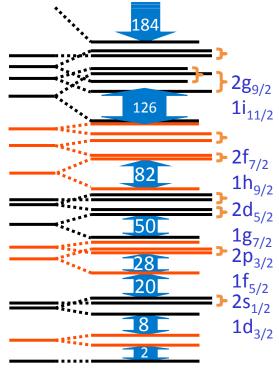


Pseudospin symmetry near degeneracy between ----- a relativistic symmetry

$$[(n-1, l+2, j = l+3/2)]$$

(n, l, j = l+1/2)

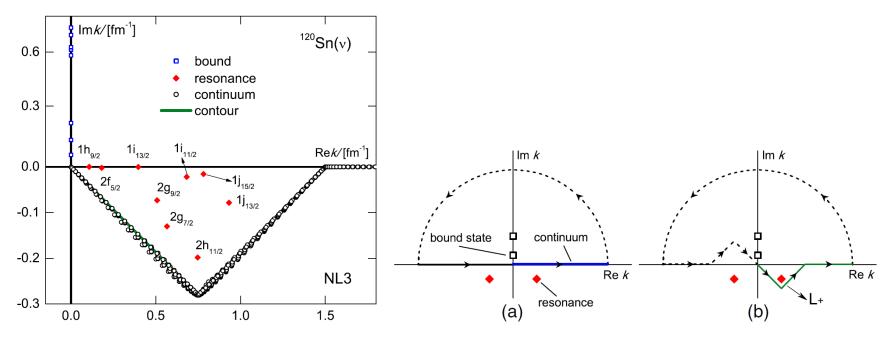
- □ The origin of PSS deeply hidden in original Hamiltonian H_1 can be traced in its SUSY partner Hamiltonian H_2 .
- $\square \Delta E_{PSO}$ can be understood in an explicit and quantitative way.



HZL, Meng, Zhou, Phys. Rep. 570, 1-84 (2015)

Our studies on single-particle resonances

Resonances in relativistic scheme



Li, Shi, Guo, Niu, HZL, Phys. Rev. Lett. 117, 062502 (2016)

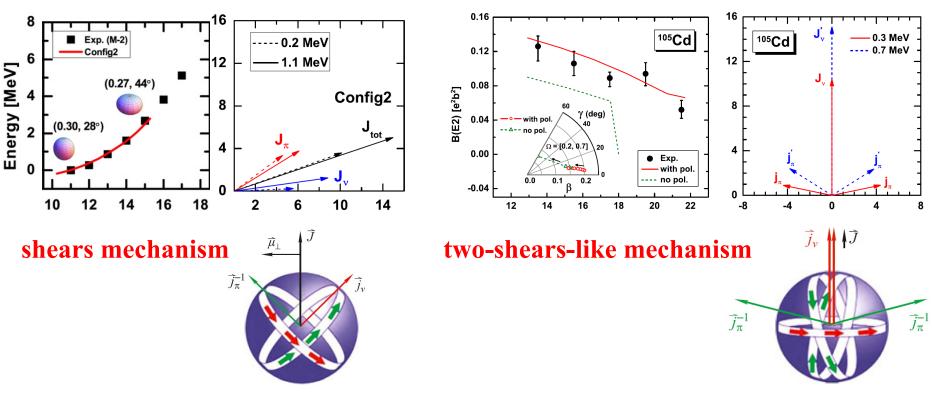
- □ solving nucleon equation of motion in complex momentum space \rightarrow singleparticle resonances in relativistic scheme
- This method is not only very effective for narrow resonances, but also can be reliably applied to broad resonances.

→ halo nuclei

Our studies on (anti-)magnetic rotations



Anti-magnetic rotation in ¹⁰⁵Cd



Zhao et al., Phys. Lett. B 699, 181 (2011); Phys. Rev. Lett. 107, 122501 (2011)

- □ 2D tilted axis cranking model with CDFT \rightarrow nuclear rotations
- □ Shears and two-shears-like mechanisms are described and understood self-consistently and microscopically.

Covariant density functional theory

From Walecka model in 1974 to now

CDFT achieves a great success in nuclear ground-state and excited-state properties

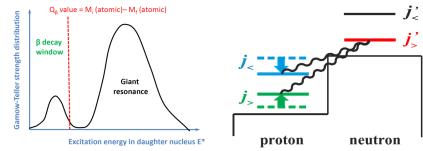
For simplicity, in most of versions \leftarrow relativistic mean-field (**RMF**) theory

- \succ with local Hartree terms \checkmark
- without non-local Fock terms *



However, "Hartree terms only" show limitations in

- properties in spin-isospin channel
- effects of tensor interaction



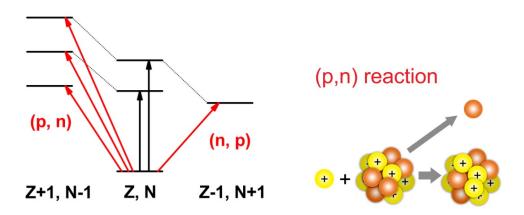
Our works: both Hartree and Fock terms

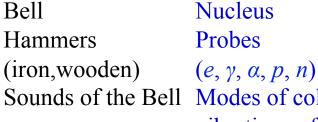
Nuclear spin-isospin properties

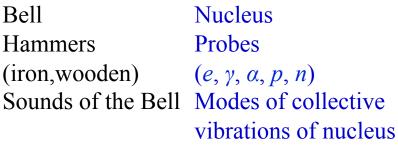
Nuclear spin-isospin excitations

 $\succ \beta$ -decays in nature T β⁻ decay electron antineutrino Z, N Z+1, N-1

charge-exchange reactions in lab









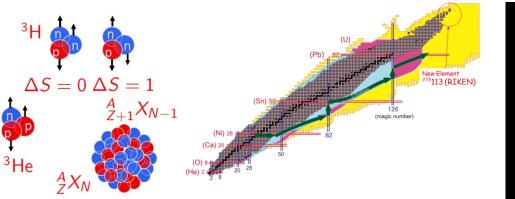
Spin-isospin excitations

These excitations are important to understand

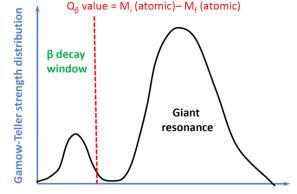
"What are the spin and isospin properties of nuclear force and nuclei?" (nuclear physics)

"Where and how does the rapid neutroncapture process (*r*-process) happen?" (nuclear astrophysics)

"Does Cabibbo-Kobayashi-Maskawa matrix satisfy the unitary condition?" (particle physics)



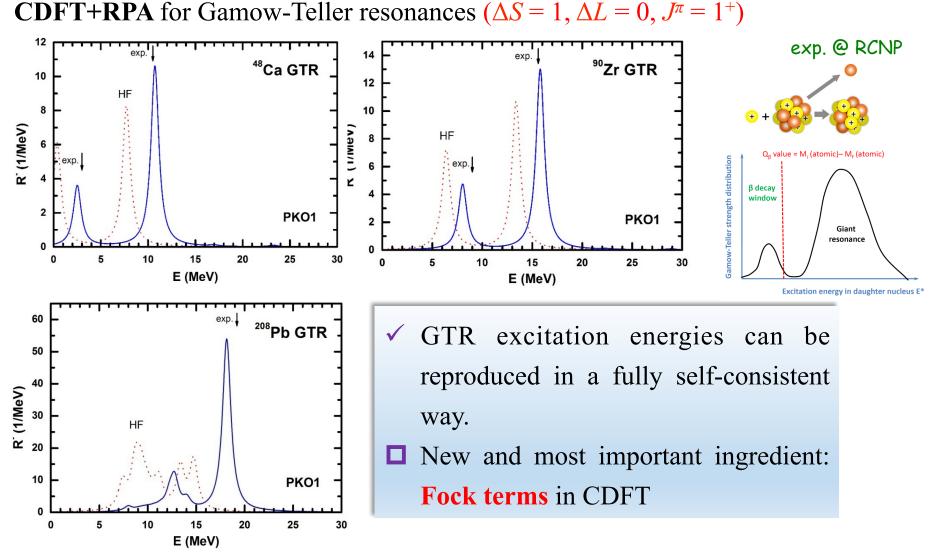




Excitation energy in daughter nucleus E*

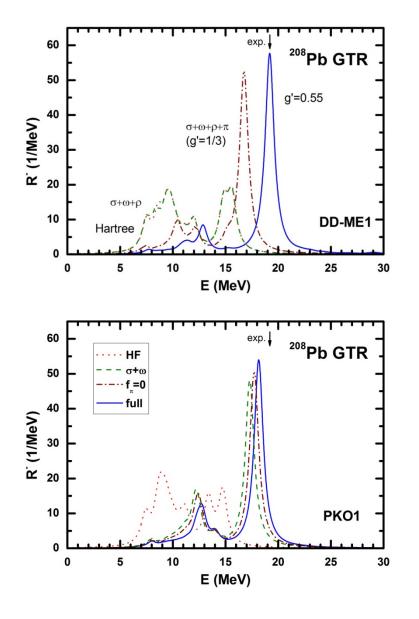
Key exp. @ RIKEN RCNP MSU GSI TRIUMF CERN

Gamow-Teller resonances



HZL, Giai, Meng, Phys. Rev. Lett. 101, 122502 (2008)

Physical mechanisms of GTR



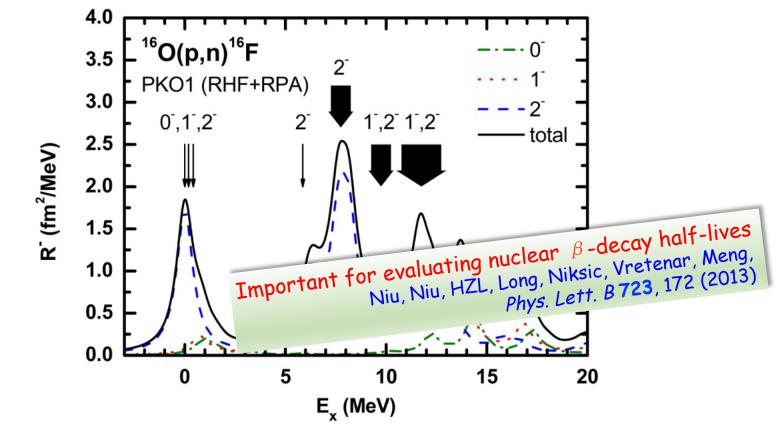
With only Hartree terms

- No contribution from isoscalar σ and ω mesons, because exchange terms are missing.
- \square *π*-meson is dominant in this resonance.
- □ g' has to be re-fitted to reproduce the experimental data.

With both Hartree & Fock terms

- Isoscalar σ and ω mesons play an essential role via the exchange terms.
- $\square \pi$ -meson plays a minor role.
- **\Box** g' = 1/3 is kept for self-consistency.

HZL, Giai, Meng, Phys. Rev. Lett. 101, 122502 (2008) HZL, Zhao, Ring, Roca-Maza, Meng, Phys. Rev. C 86, 021302(R) (2012) **CDFT+RPA** for spin-dipole resonances ($\Delta S = 1, \Delta L = 1, J^{\pi} = 0^{-}, 1^{-}, 2^{-}$)



(Exp.) Wakasa et al., PRC 84, 014614 (2011); (Theory) HZL, Zhao, Meng, Phys. Rev. C 85, 064302

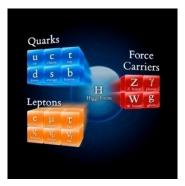
a crucial test for the theoretical predictive power

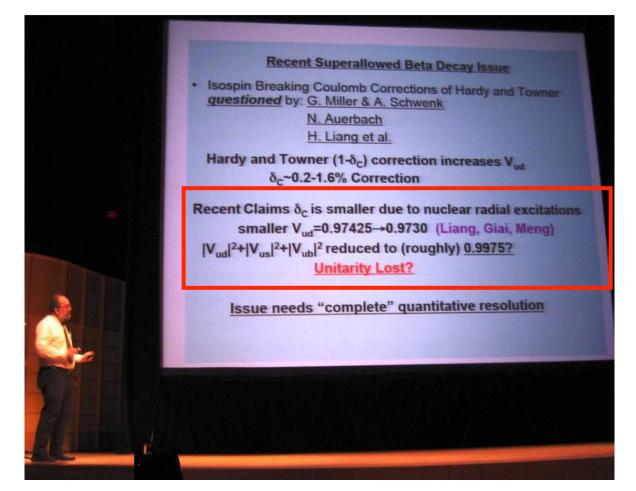
CKM matrix and its unitarity test

Cabibbo-Kobayashi-Maskawa matrix



Nobel Prize 2008 "There exist at least three families of quarks in nature." "Only three?"





Plenary talk in INPC2010 "Precision Electroweak Tests of the Standard Model" by Professor William Marciano

CKM matrix and its unitarity test

Cabibbo-Kobayashi-Maskawa matrix

- > quark eigenstates of weak interaction (>> quark mass eigenstates
- unitarity of CKM matrix (test of Standard Model

$\left(\left V_{ud} \right \right)$	$ V_{us} $	$ V_{ub} $	(0.97425 ± 0.00022)	0.2252 ± 0.0009	0.00415 ± 0.00049
$ V_{cd} $	$ V_{cs} $	$ V_{cb} =$	0.230 ± 0.011	1.006 ± 0.023	0.0409 ± 0.0011
			0.0084 ± 0.0006		•

Unitarity test Particle Data Group 2014

- > the most precise test comes from $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2$
- → the most precise $|V_{ud}|$ comes from nuclear $0^+ \rightarrow 0^+$ superallowed β transitions

Nuclear superallowed β transitions

> experimental measurements

- $|M_F|^2 = |\langle f| T_+ |i\rangle|^2 = |M_0|^2(1-\delta_c)$
- theoretical corrections (isospin symmetry-breaking corrections)

"Only three families of quarks in nature?"

Isospin symmetry-breaking corrections δ_c

Isospin symmetry-breaking corrections δ_c . All values are expressed in %. with Fock terms w/o Fock terms

		rock u		w/o rock terms				
	PKO1	PKO2	PKO3	DD-ME1	DD-ME2	NL3	TM1	
$^{10}C \rightarrow {}^{10}B$	0.082	0.083	0.088	0.149	0.150	0.124	0.133	
$^{14}\text{O} \rightarrow ^{14}\text{N}$	0.114	0.134	0.110	0.189	0.197	0.181	0.159	
$^{18}{ m Ne} ightarrow {}^{18}{ m F}$	0.270	0.277	0.288	0.424	0.430	0.344	0.373	
26 Si $ ightarrow$ 26 Al	0.176	0.176	0.184	0.252	0.252	0.213	0.226	
$^{30}S \rightarrow {}^{30}P$	0.497	0.550	0.507	0.612	0.633	0.551	0.648	
$^{34}\mathrm{Ar} ightarrow ^{34}\mathrm{Cl}$	0.268	0.281	0.267	0.368	0.376	0.438	0.320	
38 Ca $ ightarrow$ 38 K	0.313	0.330	0.313	0.431	0.441	0.390	0.572	
$^{42}\mathrm{Ti} ightarrow ^{42}\mathrm{Sc}$	0.384	0.387	0.390	0.515	0.523	0.436	0.443	
$^{26}AI \rightarrow {}^{26}Mg$	0.139	0.138	0.144	0.198	0.198	0.172	0.179	
$^{34}\text{CI} \rightarrow ^{34}\text{S}$	0.234	0.242	0.231	0.302	0.307	0.289	0.267	
$^{38}\text{K} ightarrow ^{38}\text{Ar}$	0.278	0.290	0.276	0.363	0.371	0.334	0.484	
$^{42}\text{Sc} \rightarrow {}^{42}\text{Ca}$	0.333	0.334	0.336	0.442	0.448	0.377	0.383	
$^{54}Co \rightarrow {}^{54}Fe$	0.319	0.317	0.321	0.395	0.393	0.355	0.368	
$^{66}As \rightarrow {}^{66}Ge$	0.475	0.475	0.469	0.568	0.572	0.560	0.524	
$^{70}\mathrm{Br} ightarrow {}^{70}\mathrm{Se}$	1.140	1.118	1.107	1.232	1.268	1.230	1.226	
$^{74}\text{Rb} \rightarrow ^{74}\text{Kr}$	1.088	1.091	1.071	1.233	1.258	1.191	1.234	

HZL, Giai, Meng, Phys. Rev. C 79, 064316 (2009)

Isospin corrections & V_{ud}

PHYSICAL REVIEW C 79, 064316 (2009)

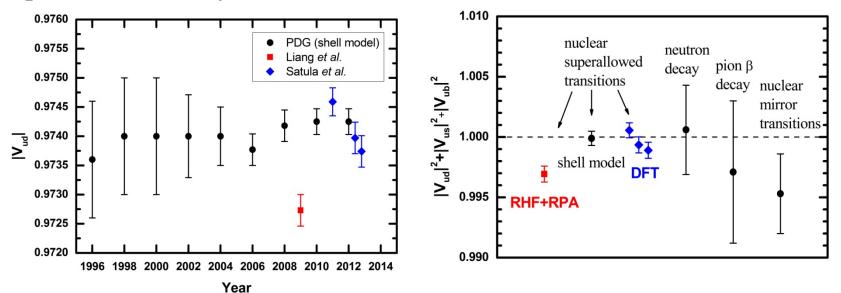
Isospin corrections for superallowed Fermi β decay in self-consistent relativistic random-phase approximation approaches

Haozhao Liang (梁豪兆),^{1,2} Nguyen Van Giai,² and Jie Meng (孟杰)^{1,3}



cited by PDG 2010, 2012, 2014, ...

Isospin corrections by self-consistent CDFT

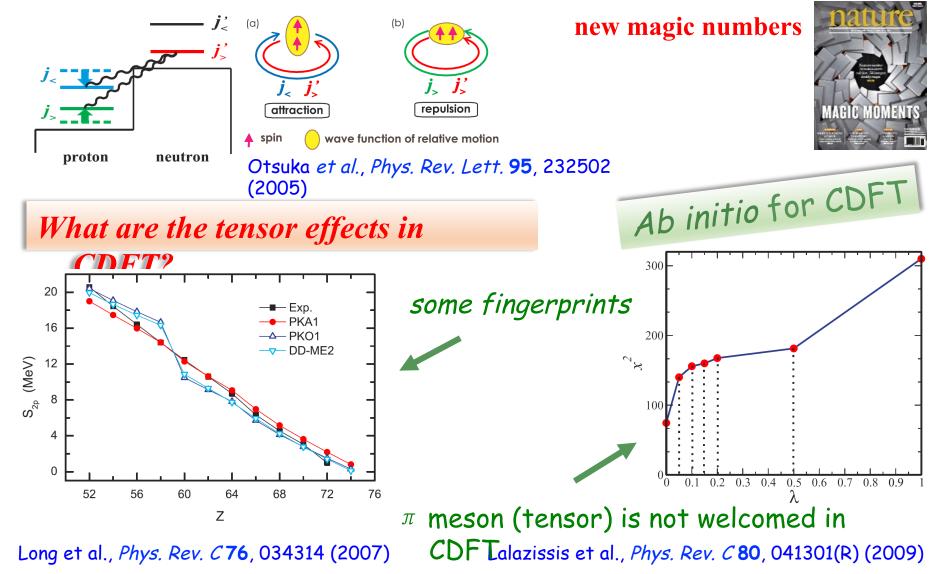


HZL, Giai, Meng, PRC 79, 064316 (2009); Satula et al., PRL 106, 132502 (2011); PRC 86, 054316 (2012)

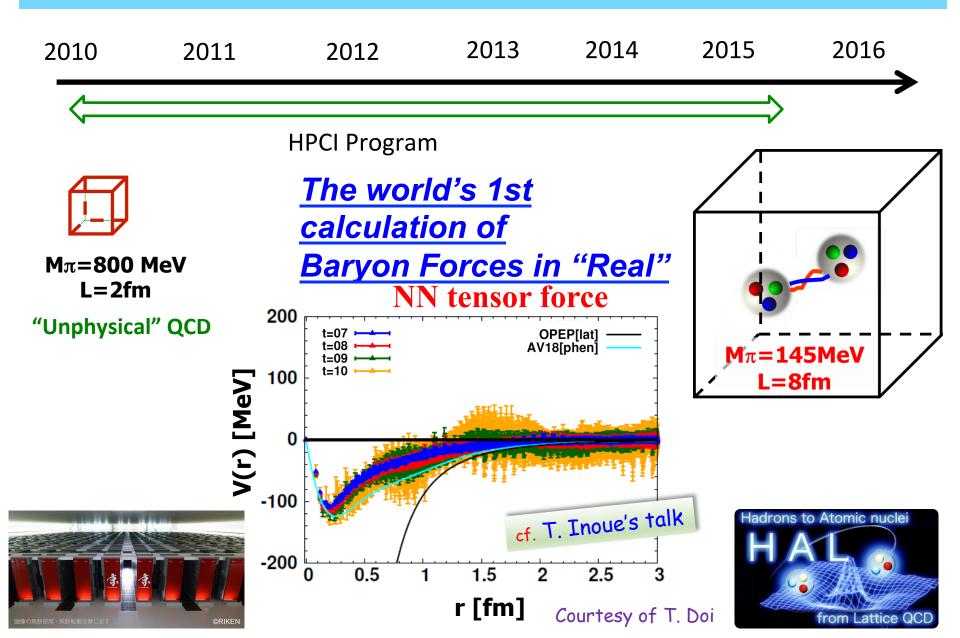
To our best knowledge: |V_{ud}|² + |V_{us}|² + |V_{ub}|²: 0.997 ~ 1.000 (the 4th family?)
 ongoing studies

Tensor effects in CDFT?

Tensor effects are crucial, in particular, for properties of exotic nuclei



CDFT from lattice QCD?



Ab initio for CDFT

Brueckner-Hartree-Fock theory (ladder diagrams to all order)

(with exchange terms as well)

$$c \bigcirc \overset{\mathsf{HF}}{\longrightarrow} \bigcirc d + c \bigcirc \overset{\mathsf{m}}{\xrightarrow{n}} \overset{\mathsf{m}}{\longrightarrow} \overset{\mathsf{m}}{\xrightarrow{n}} \overset{\mathsf{m}}{\longrightarrow} \overset{\mathsf{m}}{\xrightarrow{n}} \overset{\mathsf{m}}{\longrightarrow} \overset{\mathsf{m}}{\xrightarrow{n}} \overset{\mathsf{m}}{\xrightarrow{n}} \overset{\mathsf{m}}{\longrightarrow} d + \dots$$

$$= c \bigcirc \overset{\mathsf{BHF}}{\xrightarrow{\mathsf{o}}} \overset{\mathsf{d}}{\xrightarrow{\mathsf{o}}} \overset{\mathsf{d}}{\xrightarrow{\mathsf{o}}} d$$

➢ Bethe-Goldstone equation → finite nuclei

 $\langle ab|G(W)|cd \rangle = \langle ab|V|cd \rangle + \sum_{mn} \langle ab|V|mn \rangle \frac{Q(m,n)}{W - \varepsilon_m - \varepsilon_n} \langle mn|G(W)|cd \rangle$ where the starting energy $W = \varepsilon_c + \varepsilon_d$

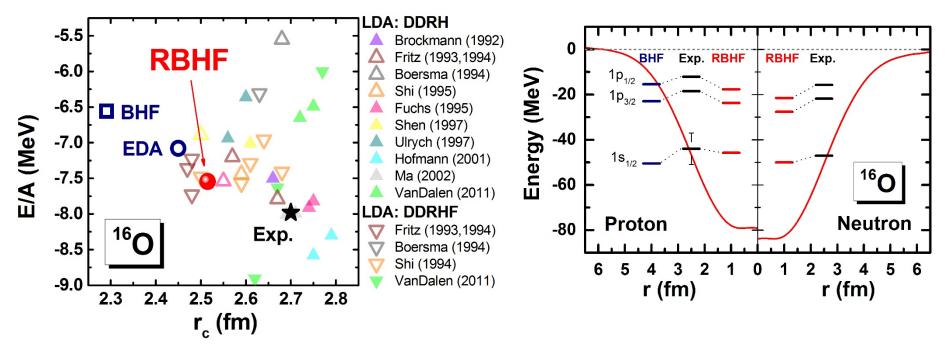
For the first time, **Bethe-Goldstone equation** is solved **self-consistently**

- > in the **two-body** frame
- > in pure **relativistic** scheme

RIKEN IPA project Shen, Hu, HZL, Meng, Ring, Zhang, arXiv:1609.01866

Relativistic BHF for finite nuclei

➢ Relativistic BHF calculations for ¹⁶O with Bonn A interaction



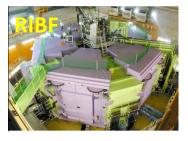
EDA = effective density approximation [Müther(1990)] LDA = local density approximation Shen, Hu, HZL, Meng, Ring, Zhang, arXiv:1609.01866

a first *ab initio* calculations for finite nuclei in relativistic scheme
 Spin-orbit splitting is reproduced well from the bare interaction.
 benchmark for various LDA calculations

A dream for another 10 years

quantum-field-theory oriented DFT

Exp. facilities:







non-perturbative nature by renormalization group $\partial_k \Gamma_k[\rho] =$ $Tr\{...\}$

(flow eq.)

EDF from effective action $E_{\rm HK}[\rho] \sim \Gamma[\rho]/\beta$ (Legendre transform)

> theoretical uncertainties from **EFT** $\Gamma^{(2)}, \Gamma^{(3)}, \Gamma^{(4)} \dots$ (power counting)

Interdisciplinary: (lattice) QCD hadron cold atom condensed matter quantum chemistry

Supercomputers





also cf. Nazarewicz' talk Schwenk & Polonyi, arXiv:0403011 [nucl-th] Kutzelnigg, JMS 768, 163 (2006) Drut, Furnstahl, Platter, PPNP 64, 120 (2010) Braun, JPG 39, 033001 (2012) Metzner et al., RMP 84, 299 (2012)





Acknowledgments

To Jun & our first baby



(~in a month)

Acknowledgments

Thank you very much for your encouragement and supports!