

^6Li structure information from $^2\text{H}(\alpha, \alpha)^2\text{H}$ scattering

P. R. Fraser¹ K. Massen-Hane¹ K. Amos^{2,3}
A. S. Kadyrov¹ I. Bray¹ L. Canton⁴

¹Curtin University, Australia

²University of Melbourne, Australia

³University of Johannesburg, South Africa

⁴Istituto Nazionale di Fisica Nucleare, Sezione di Padova, Italy



Curtin University

Multi-channel algebraic scattering (MCAS)

Low-energy, two-body, light-mass, Lippmann-Schwinger scattering

$$\begin{aligned}
 T_{cc'}^{J\pi}(p, q; E) &= V_{cc'}^{J\pi}(p, q) \\
 &+ \mu \left[\sum_{c''=1}^{\text{open}} \int_0^\infty V_{cc''}^{J\pi}(p, x) \frac{x^2}{k_{c''}^2 - x^2 + i\epsilon} T_{c''c'}^{J\pi}(x, q; E) dx \right. \\
 &\quad \left. - \sum_{c''=1}^{\text{closed}} \int_0^\infty V_{cc''}^{J\pi}(p, x) \frac{x^2}{h_{c''}^2 + x^2} T_{c''c'}^{J\pi}(x, q; E) dx \right]
 \end{aligned}$$

If we assume the T-matrix can be split, viz.

$$T_{cc'} = - \sum_p |\chi_{cp}\rangle \frac{1}{[1 - \eta_p] \eta_p} \langle \chi_{c'p}| ,$$

we are led to a separable potential

$$V_{cc'} = - \sum_p |\chi_{cp}\rangle \frac{1}{\eta_p} \langle \chi_{c'p}| ,$$

and we do this in reverse.

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MCAS includes of the Pauli exclusion principle in $V_{cc'}$, even for targets defined by collective models.

This is done with **orthogonalising pseudo potentials**:

$$\mathcal{V}_{cc'}(r, r') = V_{cc'}(r)\delta(r - r') + \lambda A_c(r)A_{c'}(r')\delta_{c,c'}$$

Many low-energy reactions of light-mass nuclei with α -particles are of interest for nuclear structure and astrophysics.

So far, the MCAS formalism has only considered nucleon projectiles.

Changing the angular momentum algebra and interaction potential, we investigate $X(\alpha, \alpha)X$ reactions at energies below where excitation of the α -particle is relevant.

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The basic potential ingredients:

$$V_{cc}(r) = \left[(V_0 \delta_{cc'} + V_{\ell\ell}[\ell \cdot \ell] + V_{II}[I \cdot I]) f(r) + V_{\ell I} \frac{1}{r} \frac{df(r)}{dr} \right]_{cc'} \\ + V_{\text{mono}} \delta_{cc'} \delta_{J\pi=0} f(r)$$

where

$$f(r) = \frac{1}{1 + \exp\left(\frac{r-R}{a}\right)}$$

${}^2\text{H}(\alpha, \alpha){}^2\text{H}$

Consider an **effective** two-body model of ${}^2\text{H}(\alpha, \alpha){}^2\text{H}$ scattering.

Deuterium treated as a rigid drop of nuclear matter, with axial, permanent deformation as

$$R(\theta, \phi) = R_0 \left[1 + \sum_{L(\geq 2)} \sqrt{\frac{4\pi}{2L+1}} \beta_L [\mathbf{Y}_L(\hat{r}) \cdot \mathbf{Y}_L(\hat{\Gamma})] \right] = R_0 [1 + \epsilon] .$$

$f(r)$ is expanded to second order in ϵ .

The spectrum of deuterium is restricted to the 3S_1 1^+ ground state and the hypothetical 1S_0 0^+ state, bound by 67 keV*.

*S. B. Borzakov, N. A. Gundorin, Y. N. Pokotilovski, Phys. Part. Nuclei Lett. 12, 536 (2015).

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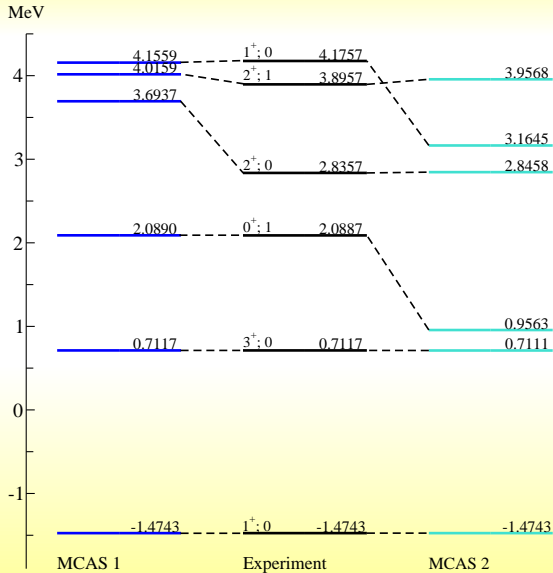
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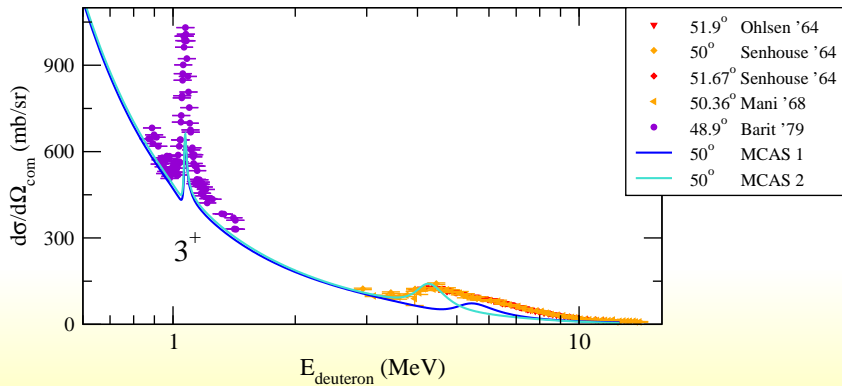
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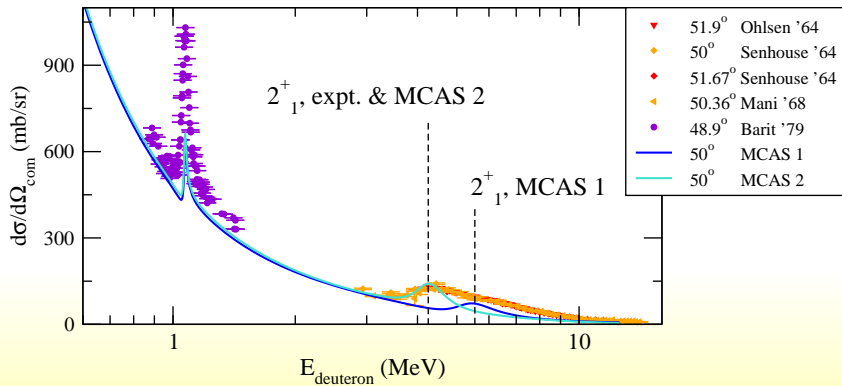
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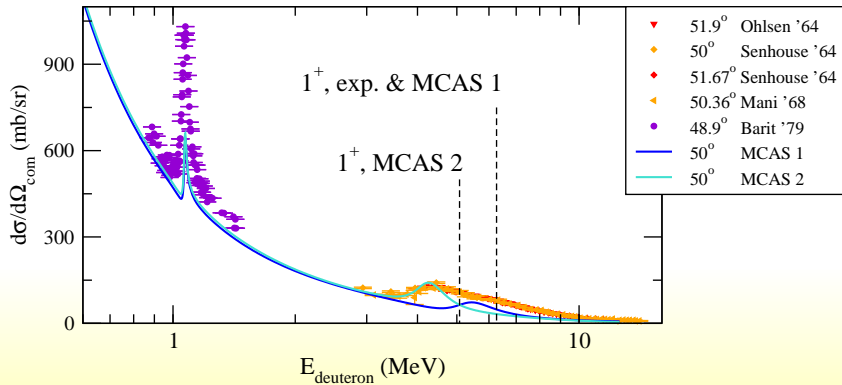
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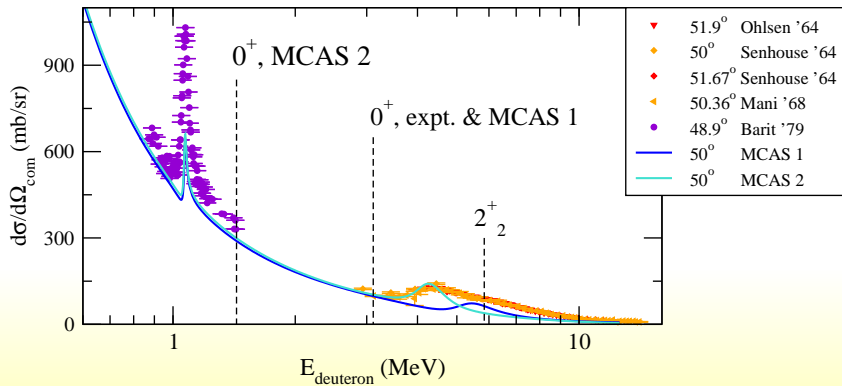
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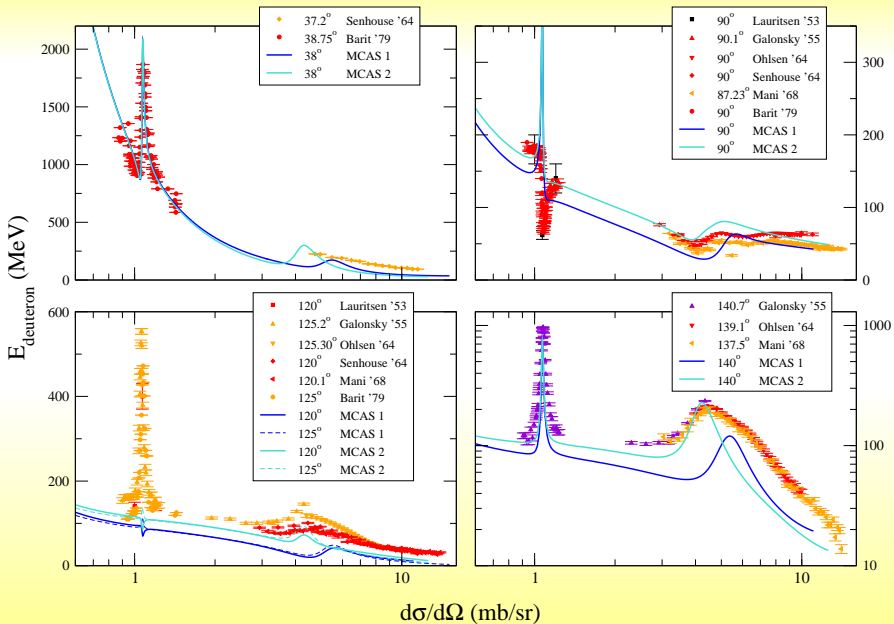


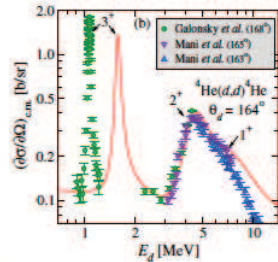
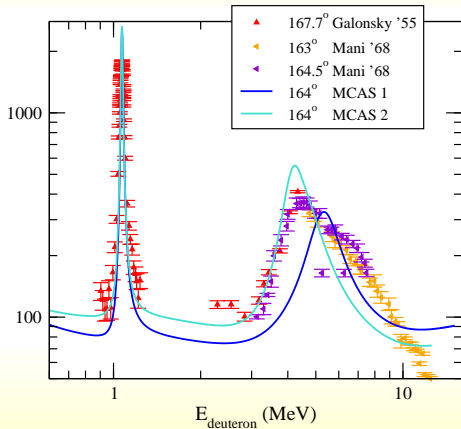




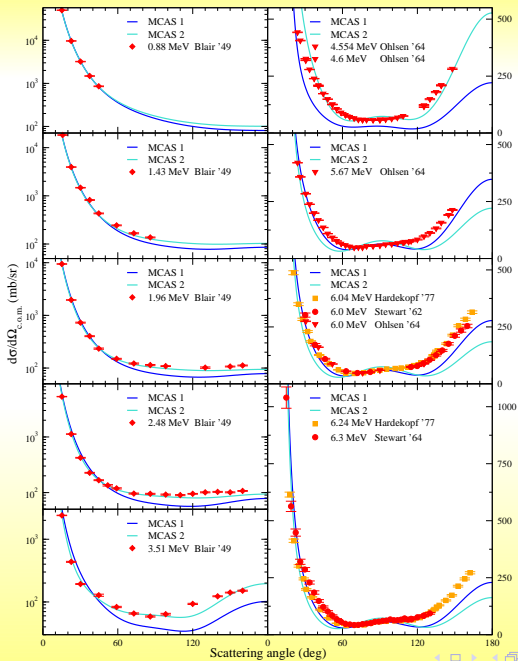


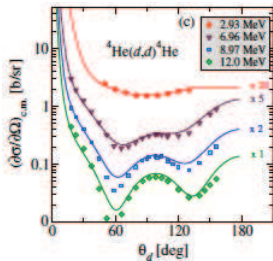
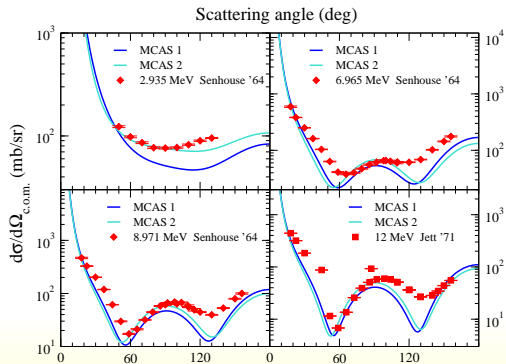






Right panel: G. Hupin, S. Quaglioni, P. Navrátil, PRL 114, 212502 (2015)





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Conclusions

Much data exists for ${}^2\text{H}(\alpha, \alpha){}^2\text{H}$ elastic scattering.

Most theoretical analyses have been fits with no structure input.

In light of a recent 6-nucleon study of the data, we have used an effective 2-body potential to analyse the data.

The correct list of states is recreated, with energies in reasonable agreement with the ${}^6\text{Li}$ data.

The non-resonant part of the cross section, the 3^+ resonance and the 2^+ resonance are well recreated at most angles (and energies). The 1^+ resonance is not.

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Physicist for hire: have laptop,
will travel.

Outline

1 Background

2 ${}^6\text{Li}$

3 Conclusions