# Magnetic properties of the nucleon in a uniform background field

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CSSM

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# Introduction

Introduction

- I The Magnetic Polarisability  $(\beta)$  is a fundamental property of a system of charged particles that describes the systems response to an external magnetic field.
- To calculate these on the Lattice we use,
  - the Background Field Method.

## Outline

Introduction

- 1. How is it done?
  - Background Field Method
- 2. Magnetic Polarisability
  - Correlator Ratios
  - Landau levels
  - Projections & Smearings
- 3. Results
  - Energy Shifts
  - Energy vs. Field Strength fits

Background Field Method

How is the uniform magnetic field put across the lattice?

$$\mathcal{D}'_{\mu} = \partial_{\mu} + g \, G_{\mu} + q e \, A_{\mu} q e \, A_{\mu}, \quad U'_{\mu}(x) = U_{\mu}(x) \, \mathrm{e}^{-i \, q e \, a \, A_{\mu}}$$

Causes a shift in energy (small field limit) of the baryon.

$$E(B) = M - \vec{\mu} \cdot \vec{B} - \frac{|qeB|}{2M} - \frac{4\pi}{2}\beta B^{2} + \mathcal{O}(B^{3})$$

- Magnetic moment  $\mu$  and magnetic polarisability  $\beta$ .
- Use of periodic boundary conditions impose a quantisation condition:

$$qe B a^2 = \frac{2\pi k}{N_x N_y}$$

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Simulation Details

- Through the International Lattice Data grid and PACS-CS Collaboration: S. Aoki, et al., Phys. Rev. D79 (2009) 034503.
  - Lattice Volume:  $32^3 \times 64$
  - Non-perturbative  $\mathcal{O}(a)$ -improved Wilson quark action
  - ♦ Iwasaki gauge action
  - ◆ 2 + 1 flavour dynamical-fermion QCD
  - Physical lattice spacing a = 0.0907 fm
  - $m_{\pi} = 413 \text{ MeV}$
- Standard Interpolating Fields:  $\chi_{p1} = (u^T C \gamma_5 d) u$ ,  $\chi_{n1} = (u^T C \gamma_5 d) d$
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Magnetic Polarisability

Recall the energy of baryon is

$$E(B) = M - \vec{\mu} \cdot \vec{B} - \frac{|qeB|}{2M} - \frac{4\pi}{2}\beta B^{2} + \mathcal{O}(B^{3})$$

Construct ratios of different spin and field direction different 2pt correlation functions.

$$R(B,t) = \left(\frac{G_{\downarrow}(B+,t) + G_{\uparrow}(B-,t)}{G_{\downarrow}(0,t) + G_{\uparrow}(0,t)}\right) \left(\frac{G_{\downarrow}(B-,t) + G_{\uparrow}(B+,t)}{G_{\downarrow}(0,t) + G_{\uparrow}(0,t)}\right)$$

Then extract an effective energy in the standard manner.

$$\delta E(B) = \frac{1}{\delta t} \log \left( \frac{R(B,t)}{R(B,t+\delta t)} \right) = \frac{1}{2} \left( \frac{|qeB|}{2M} - \frac{4\pi}{2}\beta B^2 \right)$$

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$$E^{2} = m^{2} + |qeB| (2n + 1 - \alpha) + p_{z}^{2}$$

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- The Landau levels are closely grouped due to the small fields used.
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## **QED Eigenmodes**

Magnetic Polarisability



Lowest lying eigenmode probability densities of lattice Laplacian operator.

• Origin is centre of the x-y plane illustrated by bottom surface of the grid.

Set source wavefunction to be exactly these modes, i.e.  $\langle x|\psi\rangle = \sum_{i=1}^{n} \langle x|\nu_i\rangle$ 

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September 16, 2016

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Define QED eigenmode projection operator

$$P_{QED}^{n}(x,y) = \sum_{i=1}^{n=|3q_f k_d|} \langle x|\nu_i \rangle \langle \nu_i|y \rangle$$

Also define QCD+QED eigenmode projection operator

$$P_{QCD+QED}^{n}(x,y) = \sum_{i=1}^{n=n_{max}} \langle x|\lambda_i \rangle \langle \lambda_i|y \rangle$$

and project the propagator

$$S(x, y, I, J) = P_I(x, z) S(z, z') P_J^{\dagger}(z', y)$$

I and J describe which projection operator is used, i.e. I=QED or I=QCD+QED.

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- Can apply standard Gaussian smearing in any combination of spatial dimensions.
- Previous results found that a background field can change the wavefunction's spatial extent.
- Can be combined with eigenmode projections.
- Smearing along field axis found to be essential to remove excited state contamination.
  - This is particularly relevant as the projections are 2-dimensional processes only.
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Results

$$\delta E(B) = \left(\frac{|qeB|}{2M} - \frac{4\pi}{2}\beta B^2\right)$$

To choose where to fit and obtain polarisability values, a number of factors are considered.

- The constant fits to the energy shifts as function of time.
- The fits to energy shifts as function of field strength, i.e. to
- We only consider the same fit window across all field strengths.
- The  $\chi^2_{dof}$  of each of these fits must be in an acceptable range.

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## Neutron Energy Shifts for polarisability

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Smeared Source to QED eigenmode projected sink neutron energy shift

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## **Proton Energy Shifts for polarisability**

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Smeared Source to QCD+QED eigenmode projected sink proton energy shift

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#### Proton

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## **Nucleon Polarisability**

Results

From the quadratic term extract the polarisability

$$\delta E(B) \propto -\frac{4\pi}{2} \beta B^2$$

Nucleon polarisabilities are found to be

	Experiment $(m_{\pi} = 138 \text{ MeV})$	This Work $(m_{\pi} = 413 \text{ MeV})$
proton	$2.5(4) \times 10^{-4} \text{ fm}^3$	$1.15(24) \times 10^{-4} \text{ fm}^3$
neutron	$3.7(12) \times 10^{-4} \text{ fm}^3$	$1.31(38) \times 10^{-4} \ { m fm}^3$

Potential exists to make interesting predictions from lattice QCD.

Summary

- Expand to correlation matrix techniques
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  - Fine tune source & sink combinations

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Summary

Effective energy of spin up & down components for neutron in smallest quantised field



#### Chiral Extrapolations based on previous work

**Bonus Slides** 



Magnetic polarisability of Neutron, compared with experimental results.

## **Magnetic Moment**

Bonus Slides

- Considerably easier than magnetic polarisability
- Take a different ratio

$$R(B,t) = \left(\frac{G_{\downarrow}(B-,t) + G_{\uparrow}(B+,t)}{G_{\downarrow}(B+,t) + G_{\uparrow}(B-,t)}\right)$$

■ to get an energy shift of

$$\delta E_{\mu}(B) = -\mu B + \mathcal{O}(B^3)$$

#### **Magnetic Moment**

Bonus Slides



Energy shift for magnetic moment of the neutron.

Ryan Bignell (CSSM)

## **Magnetic Moment**

Bonus Slides

- Extract magnetic moment from linear term
- Background field results are prelimnary only

	$BFM\ (m_{\pi} = 413 \ MeV)$	$3PT~(m_{\pi}=413~MeV)$
proton $(\beta_p)$	$2.184(22)\mu_N$	$2.244(61)\mu_N$
neutron $(\beta_n)$	$-1.371(14)\mu_N$	$-1.36(10)\mu_N$