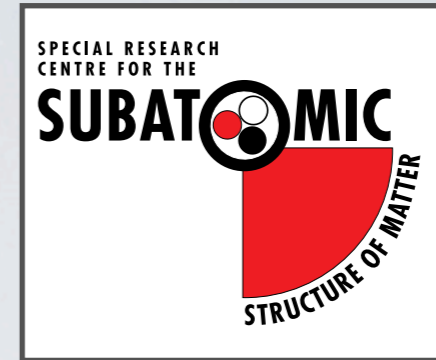




THE UNIVERSITY  
*of* ADELAIDE



# Inelastic scattering parameters from finite volume eigenstates



Ross D. Young  
CSSM & CoEPP  
University of Adelaide

with J. Hall, D. Leinweber, A. Thomas (Adelaide)  
& T.-S. H. Lee, **Jiajun Wu** (ANL)

Sixth Asia-Pacific Conference on  
Few-Body Problems in Physics  
(APFB 2014)  
Hahndorf, SA, Australia; April 7–11, 2014

# Outline

Resolving scattering in lattice QCD is not straightforward

Lattice QCD can determine (finite volume) energy eigenstates

**Problem:**

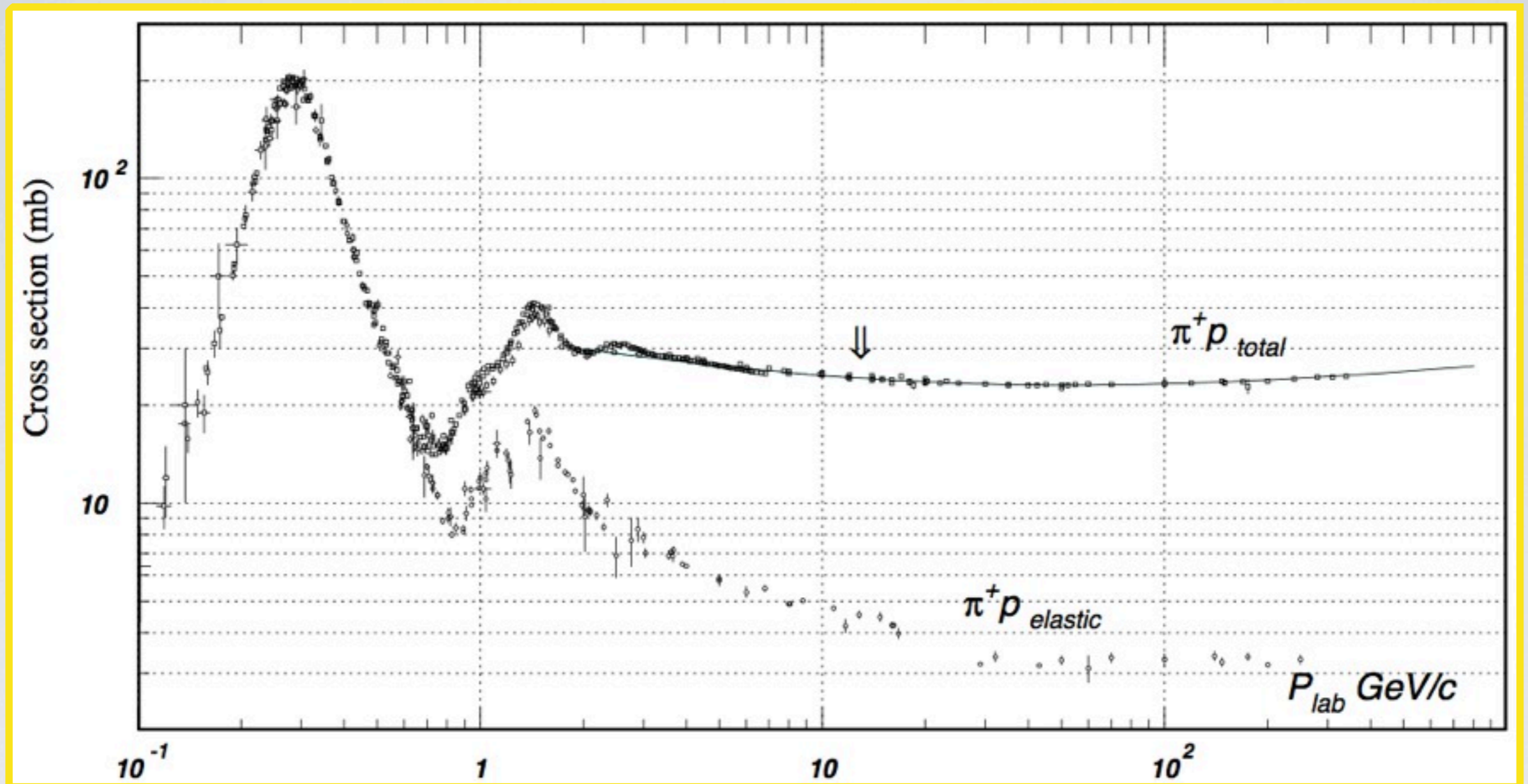
Energy eigenstates  $\rightarrow$  scattering parameters?

Lüscher method

Hamiltonian approach

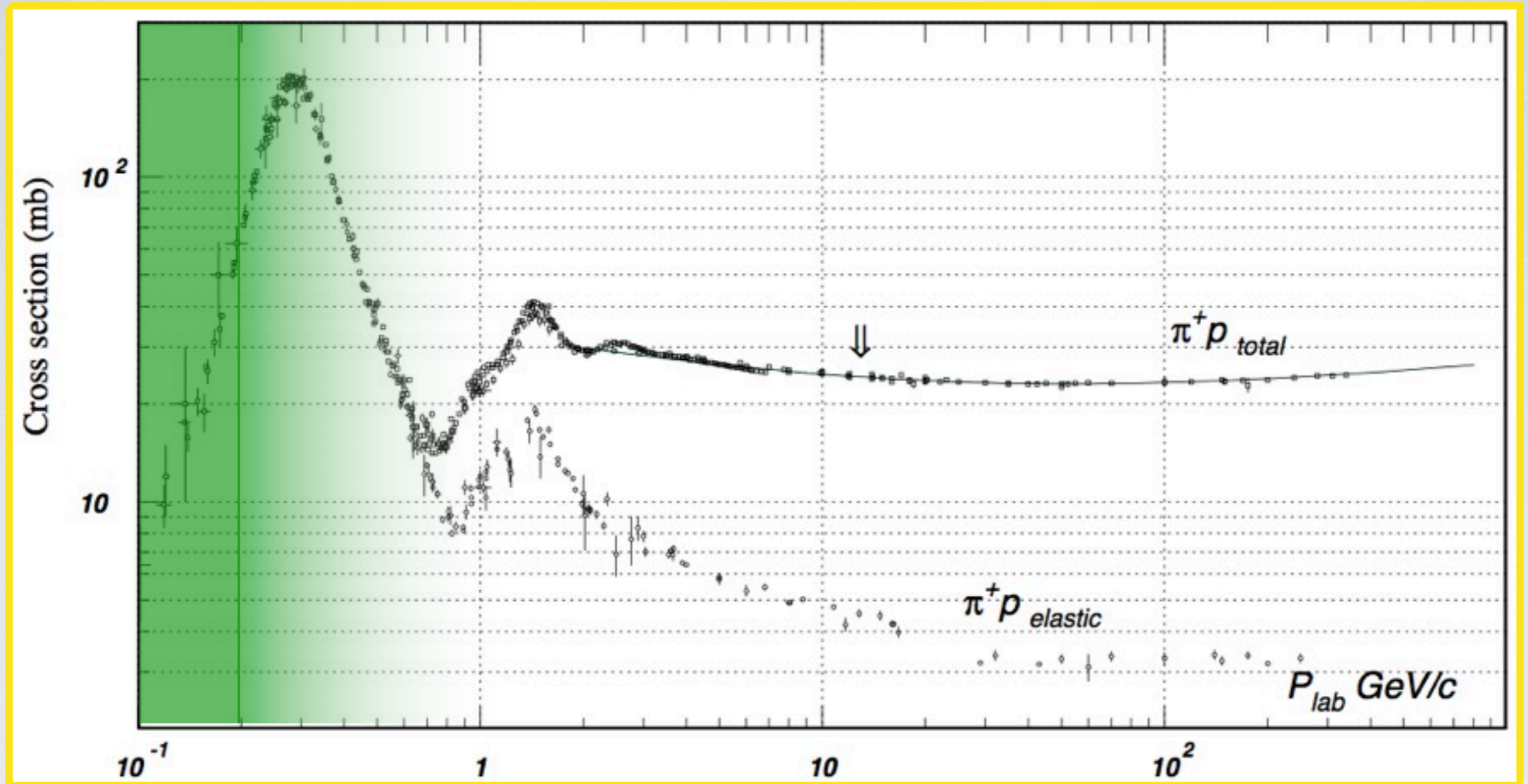
**Multi-channel systems**  $(\pi\pi \leftrightarrow K\bar{K})$





QCD has a rich spectrum

EFT

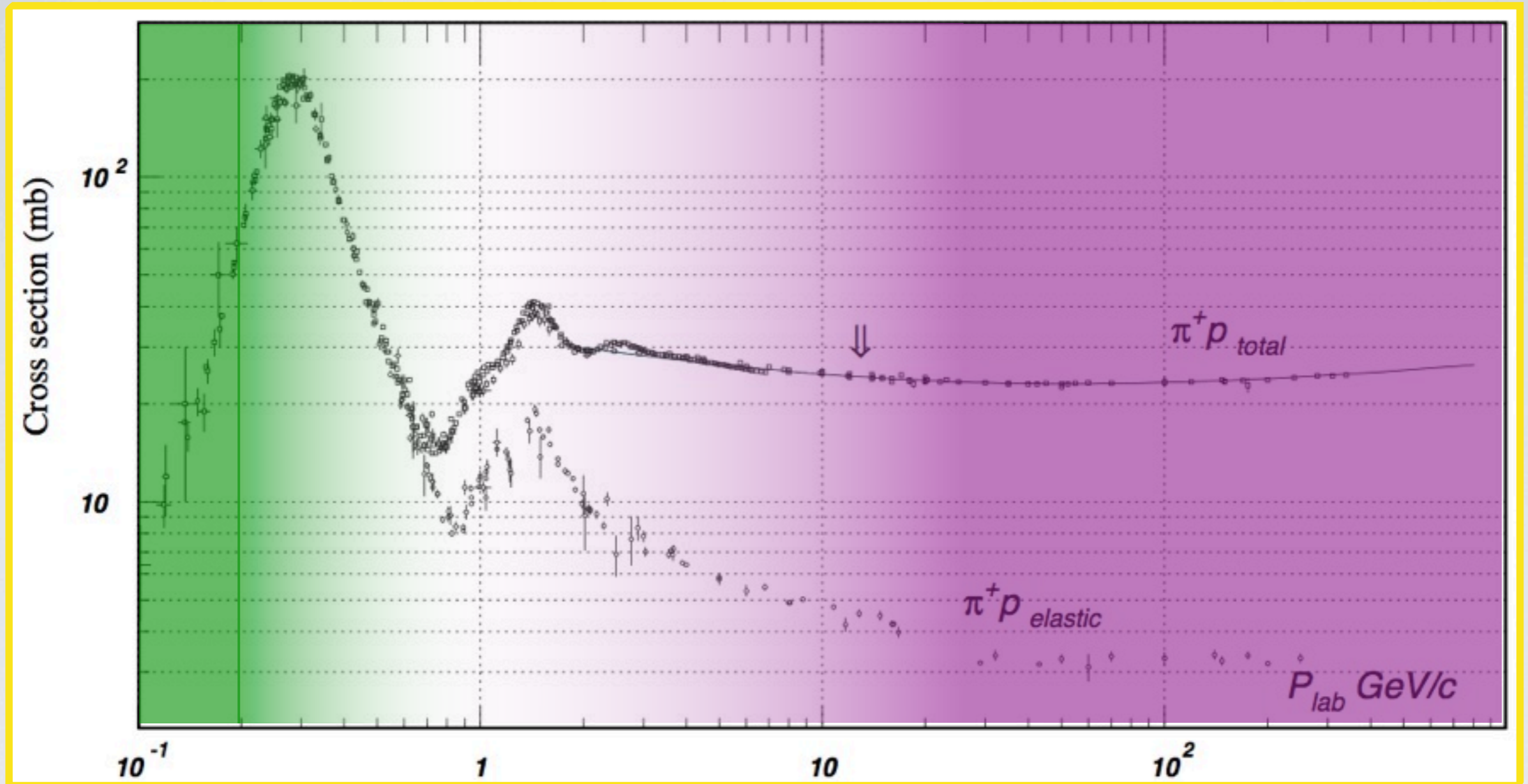


QCD has a rich spectrum



EFT

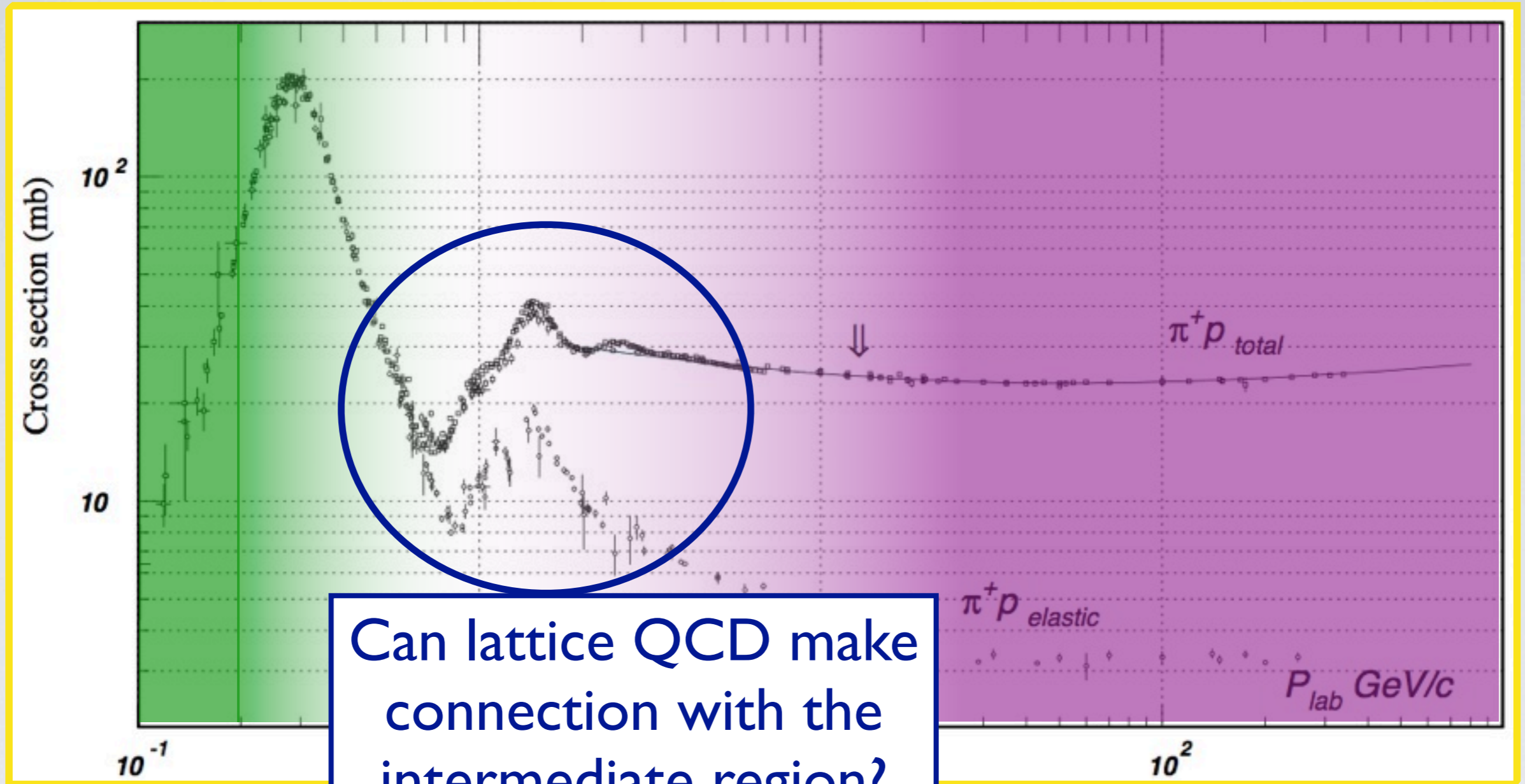
pQCD



QCD has a rich spectrum

EFT

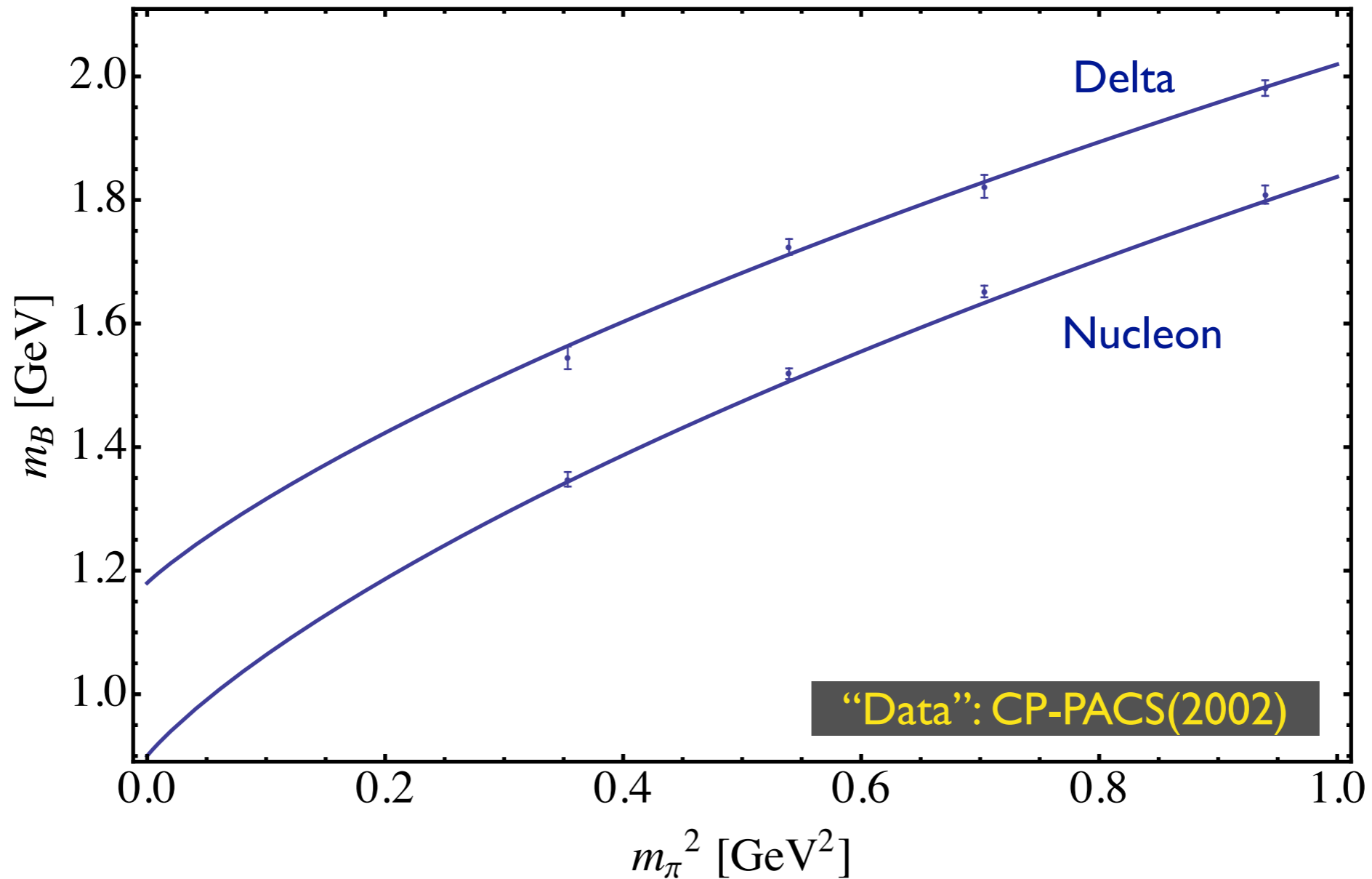
pQCD



QCD has a rich spectrum



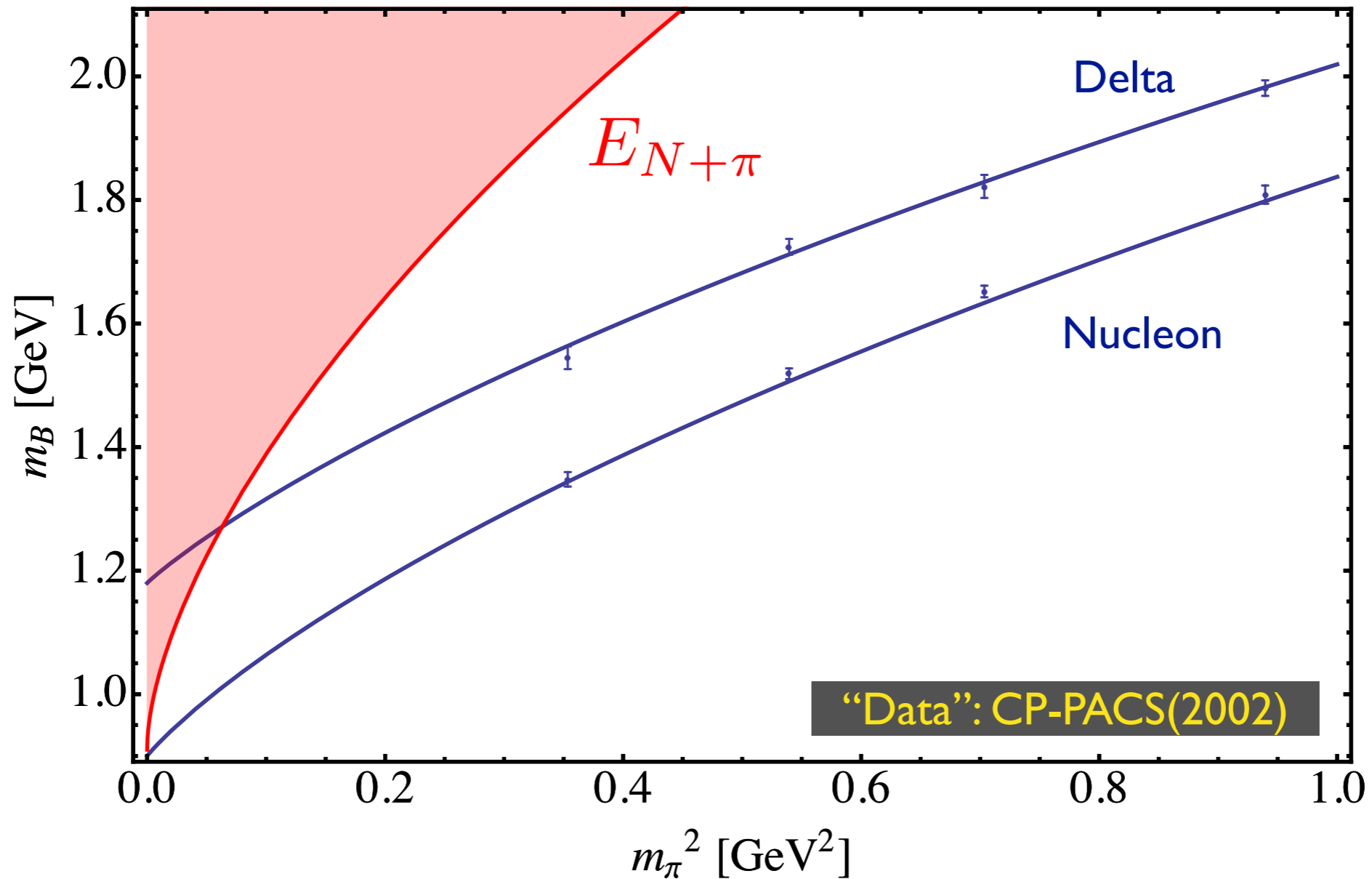
# Spectrum in lattice QCD



## In the good old days

Lattice QCD was forced to use large quark masses  
Low-lying spectrum comprised of stable particles  
[Curves are just a sketch: *not* a fit]



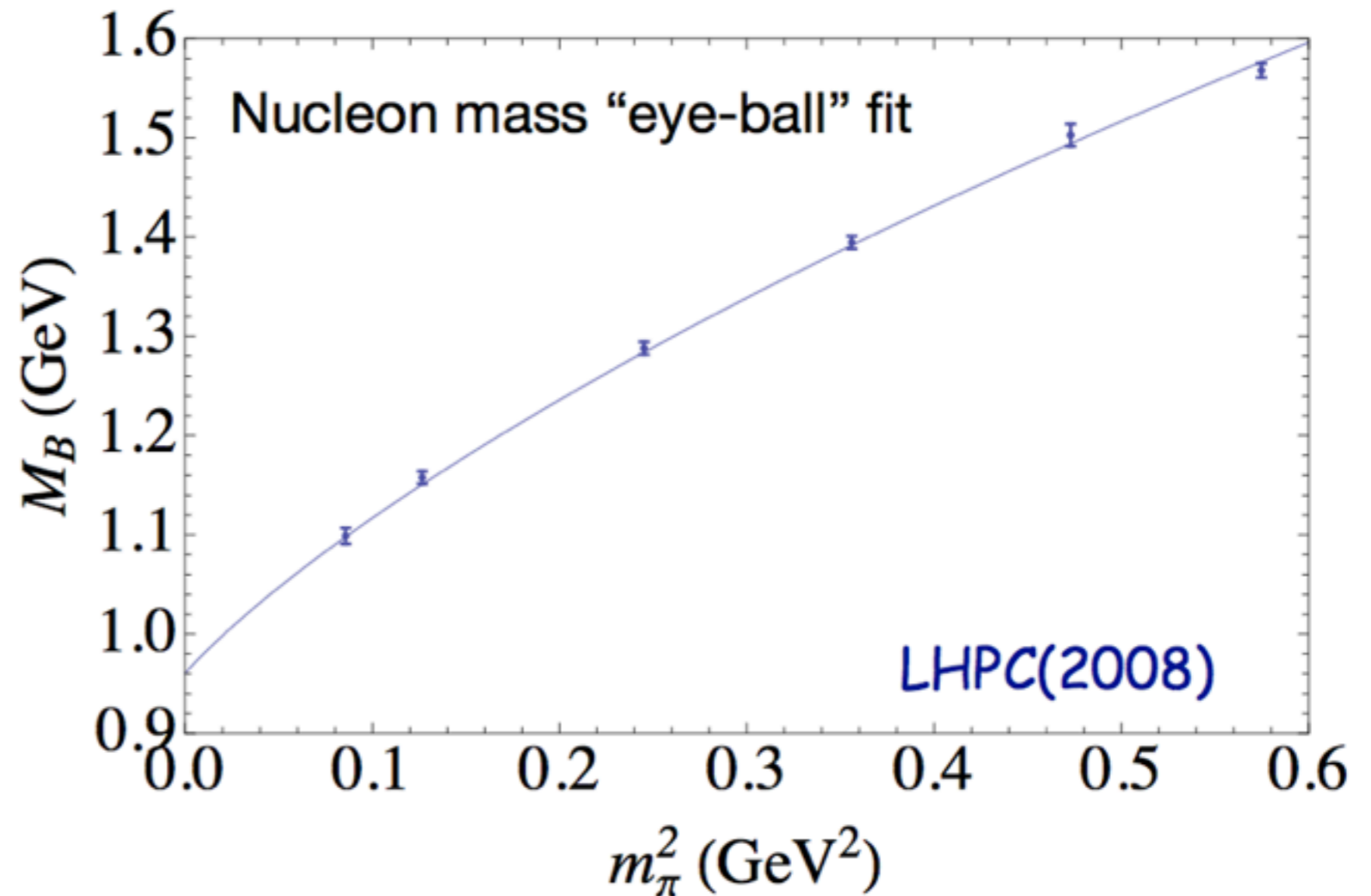


## In the good old days

Lattice QCD was forced to use large quark masses  
 Low-lying spectrum comprised of stable particles  
 [Curves are just a sketch: *not* a fit]

# Baryon mass: “Sketching a curve”

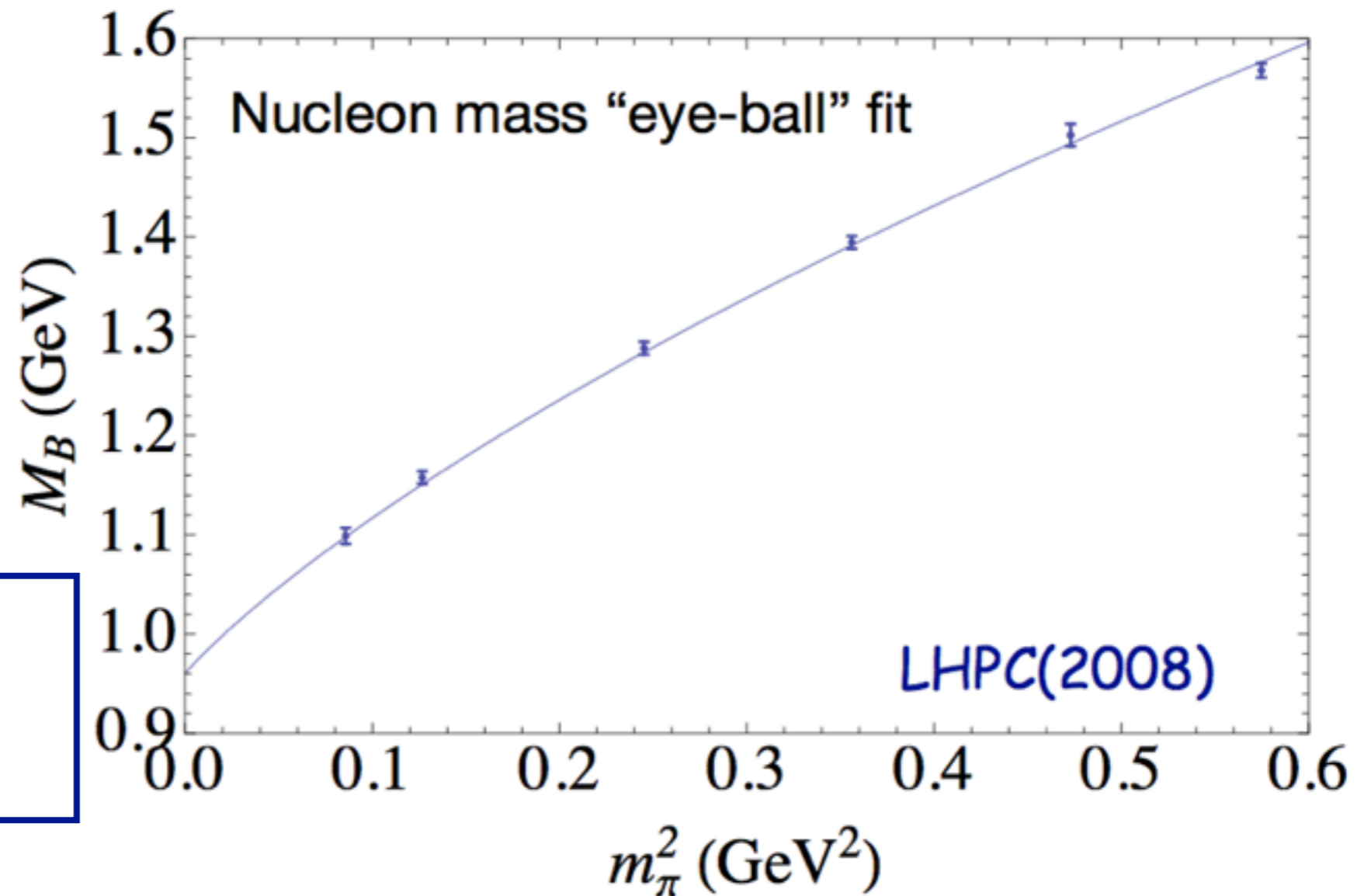
$$M_B = M_B^0 \left( \frac{1}{1 + \frac{3}{2} \frac{m_\pi}{M_B^0}} + \frac{3}{2} \frac{m_\pi}{M_B^0} \right)$$



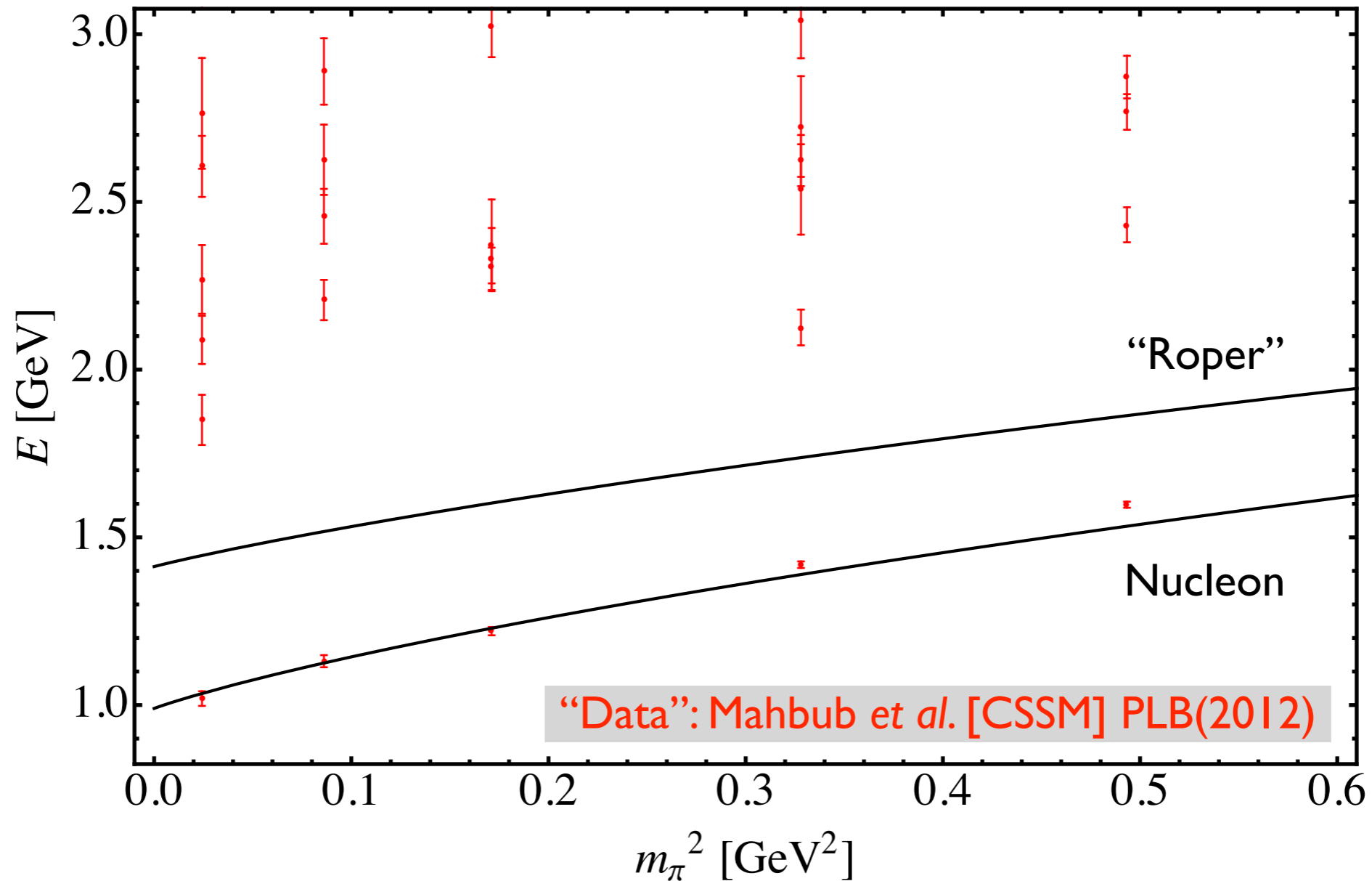


# Baryon mass: “Sketching a curve”

$$M_B = M_B^0 \left( \frac{1}{1 + \frac{3}{2} \frac{m_\pi}{M_B^0}} + \frac{3}{2} \frac{m_\pi}{M_B^0} \right)$$



One parameter =  
magnitude, slope &  
curvature!

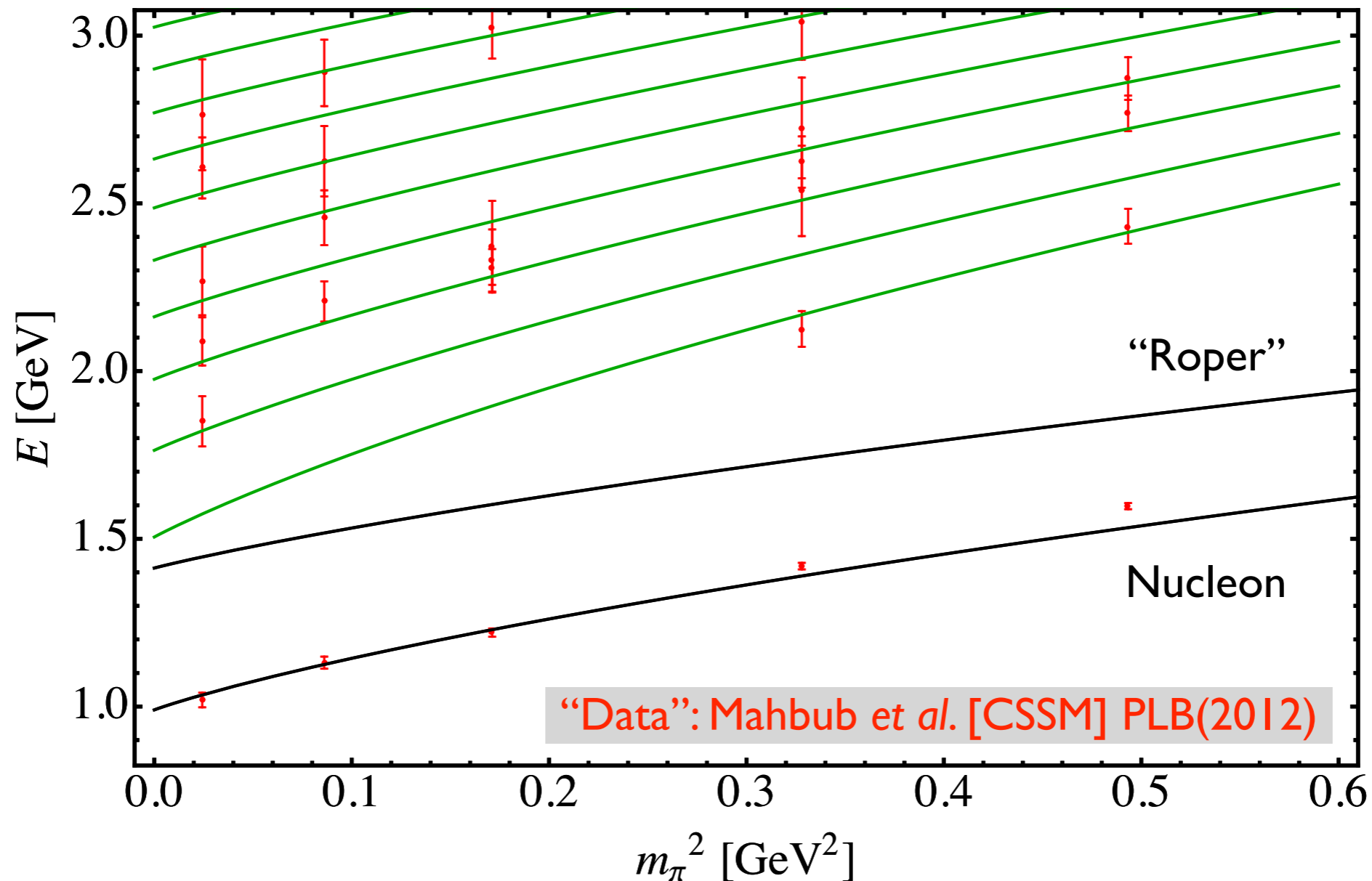


## State-of-the-art LQCD spectrum

$1/2^+$ : "Free-particle" energies on a 2.9 fm box  
 [Interactions will shift energy levels]



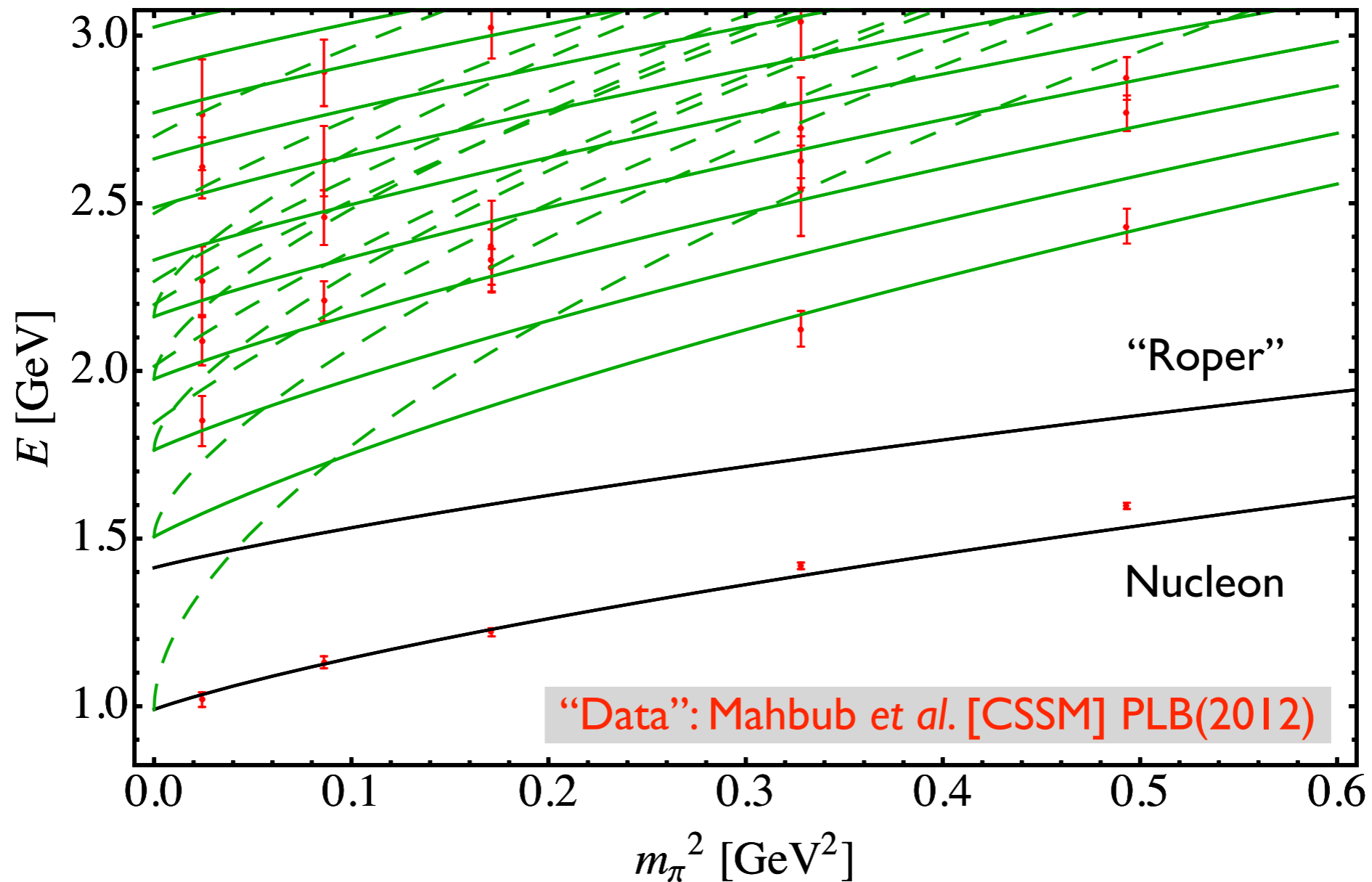
“ $N + \pi$ ”



## State-of-the-art LQCD spectrum

$1/2^+$ : “Free-particle” energies on a 2.9 fm box  
[Interactions will shift energy levels]

“ $N + \pi$ ”  
“ $N + \pi\pi$ ”

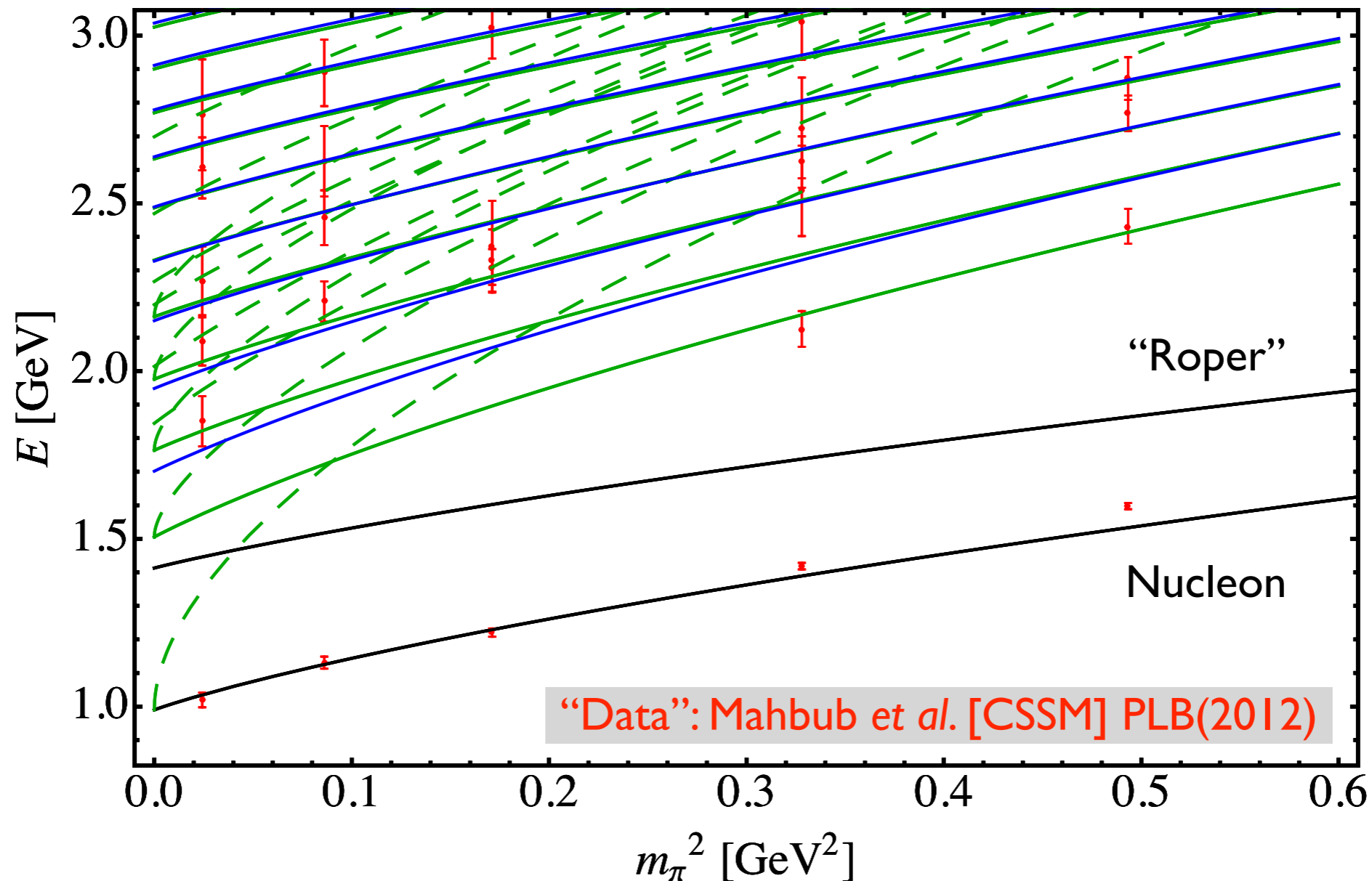


## State-of-the-art LQCD spectrum

$1/2^+$ : “Free-particle” energies on a 2.9 fm box  
[Interactions will shift energy levels]



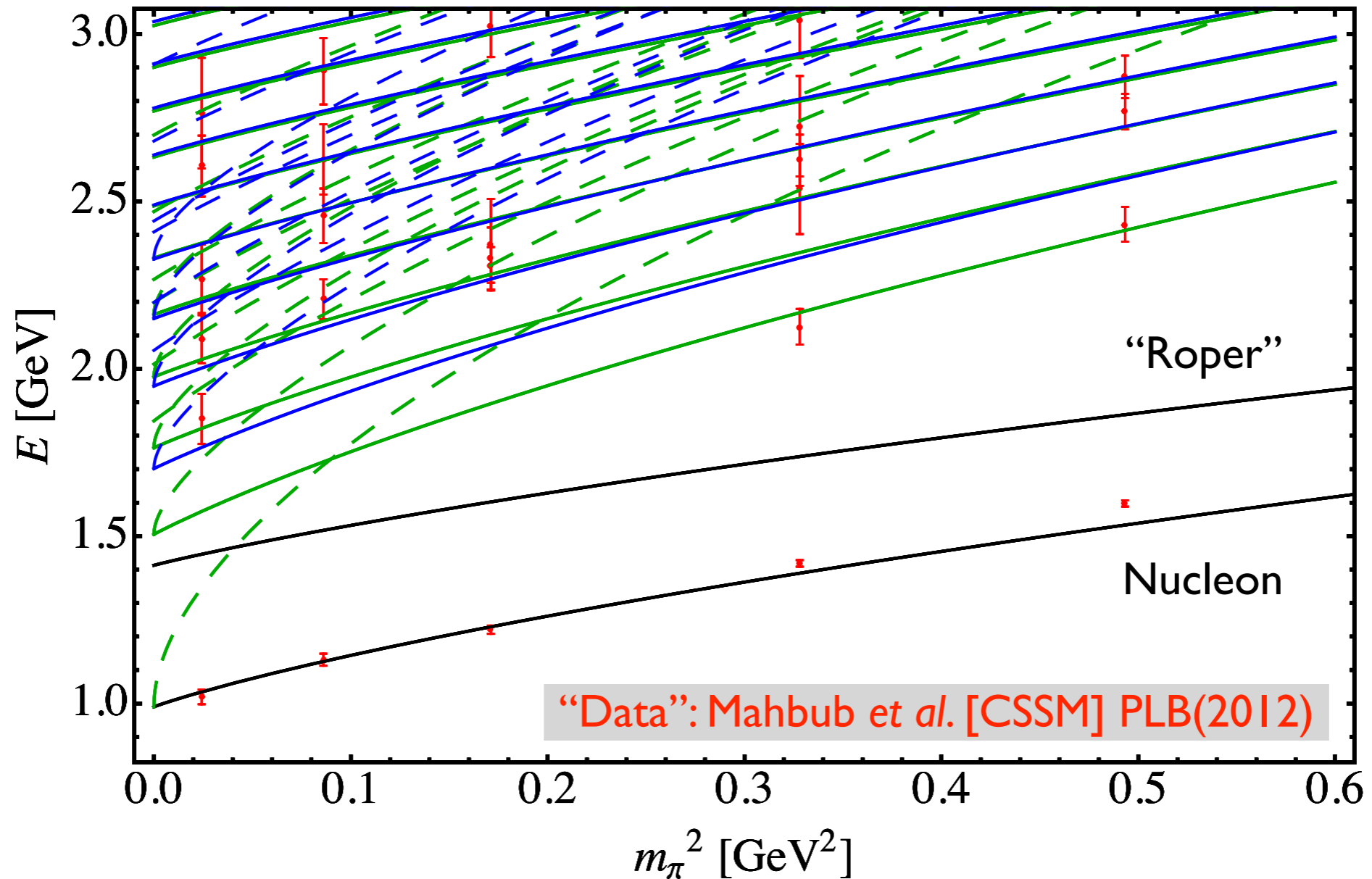
“ $N + \pi$ ”  
“ $N + \pi\pi$ ”  
“ $\Delta + \pi$ ”



## State-of-the-art LQCD spectrum

$1/2^+$ : “Free-particle” energies on a 2.9 fm box  
[Interactions will shift energy levels]

“ $N + \pi$ ”  
“ $N + \pi\pi$ ”  
“ $\Delta + \pi$ ”  
“ $\Delta + \pi\pi$ ”



## State-of-the-art LQCD spectrum

$1/2^+$ : “Free-particle” energies on a 2.9 fm box  
[Interactions will shift energy levels]

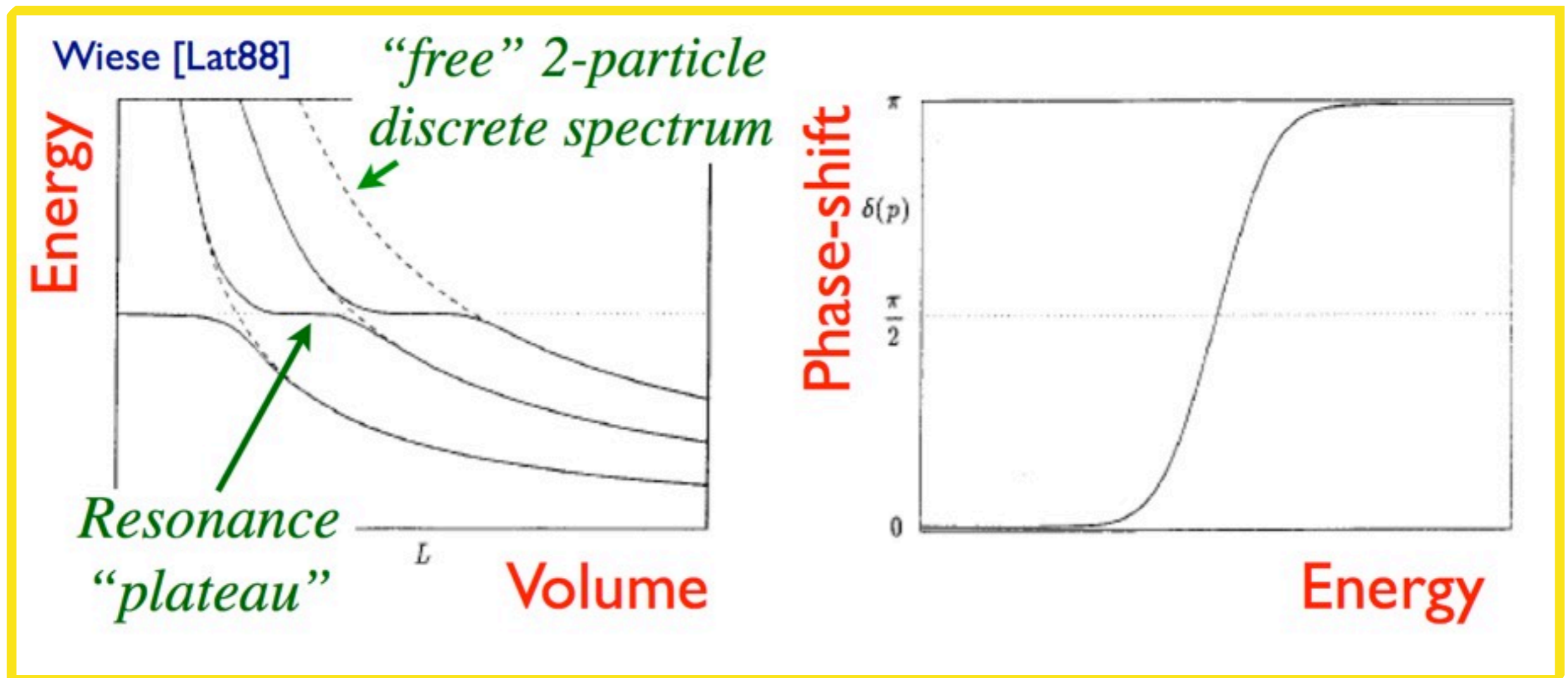
Resonance spectrum is no longer  
separated from multiparticle  
thresholds

[where widths are greater than level spacing]

# Resonance parameters?

Lüscher: 2-body elastic “scattering”, one can map finite-volume energies directly to phase shifts

Map out volume dependence of energy levels





# Scattering parameters?

Finite volume spectra in lattice QCD encode a quantum mechanical admixture of scattering states and resonance-like structures

How can we use finite-volume spectra to identify scattering parameters (eg.  $S$ -matrix)?

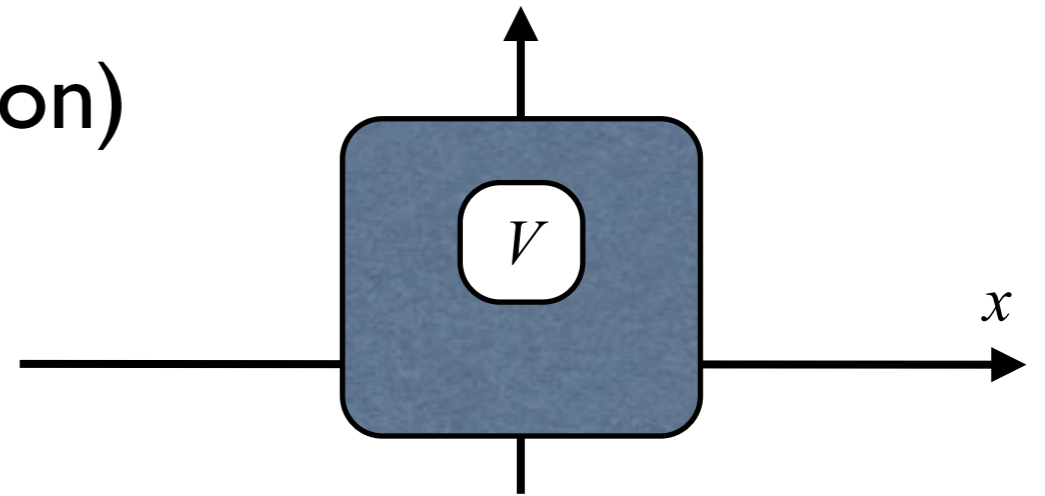
# Lüscher: Qualitative (1-D QM)

1-D Scattering of 2 identical bosons; finite-range potential

General solution (non-interacting region)

Right  $\psi^{(r)}(x) = Ae^{-ikx} + Be^{ikx}$

[Left by Bose symmetry]



Steady-state: conservation of probability  $\Rightarrow |B| = |A|$

Define  $\frac{B}{A} = e^{i2\delta} \Rightarrow \psi^{(r)}(x) = A \left( e^{-ikx} + e^{i2\delta(k)} e^{ikx} \right)$

For any potential, scattering process is entirely determined by the phase shift  $\delta(k)$

# Lüscher: Qualitative (1-D QM)

## Periodic boundary conditions

“Box” length  $L$       $\psi(x + L) = \psi(x)$

**Boundary**      $x = \pm L/2$

Continuous: Bose symmetry      $\Rightarrow \psi^{(r)}(L/2) = \psi^{(l)}(-L/2)$

Smooth      $-ik \left( A e^{ikL/2} - B e^{-iKL/2} \right) = ik \left( A e^{ikL/2} - B e^{-iKL/2} \right),$   
 $\Rightarrow \frac{B}{A} = e^{-ikL}.$

## Potential defines phase shift

Eigenvalue equation

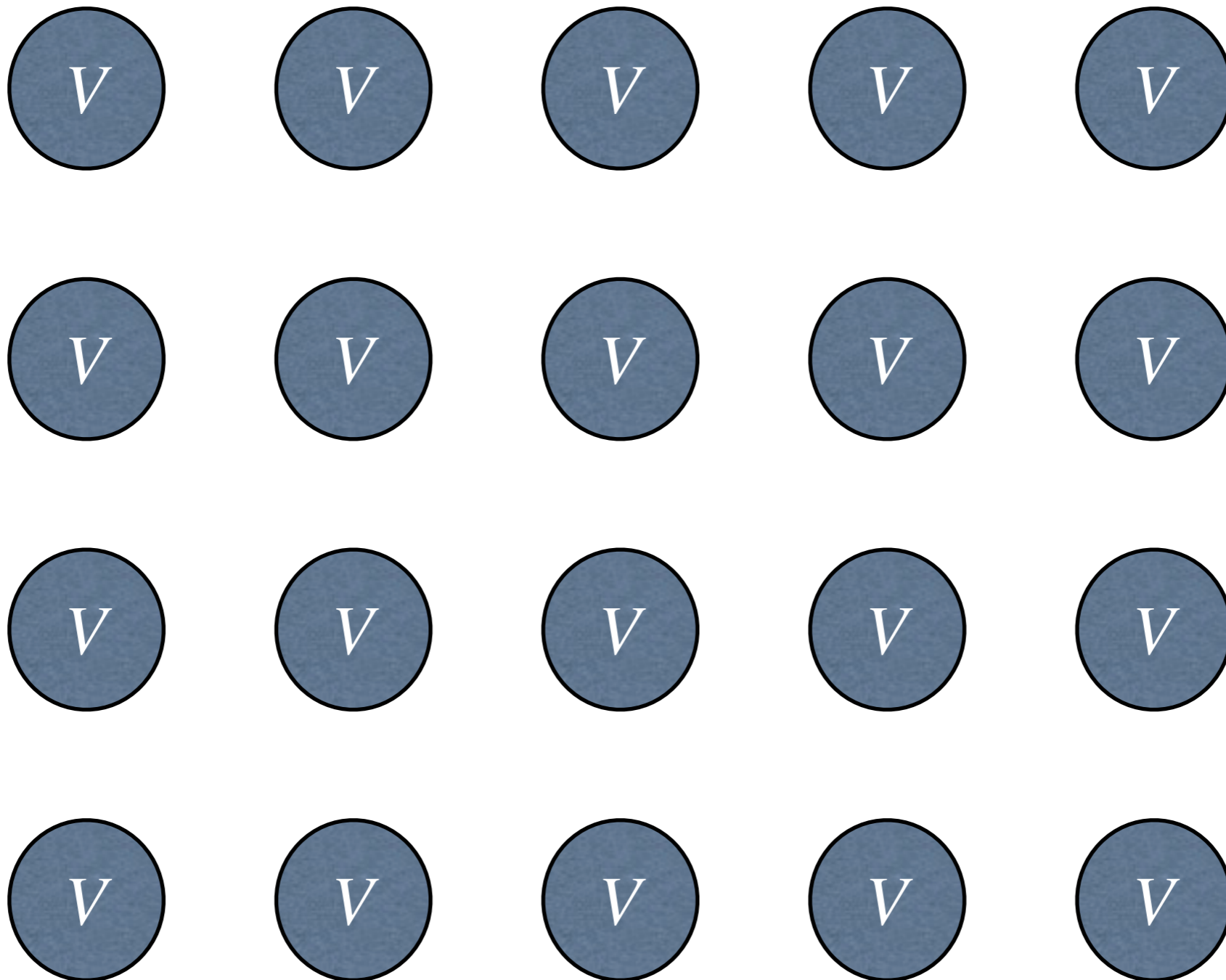
$$\underline{e^{i2\delta(k)}} - \underline{e^{-ikL}} = 0$$

Scattering phase

Lattice geometry

# Lüscher: Extension to 3-D

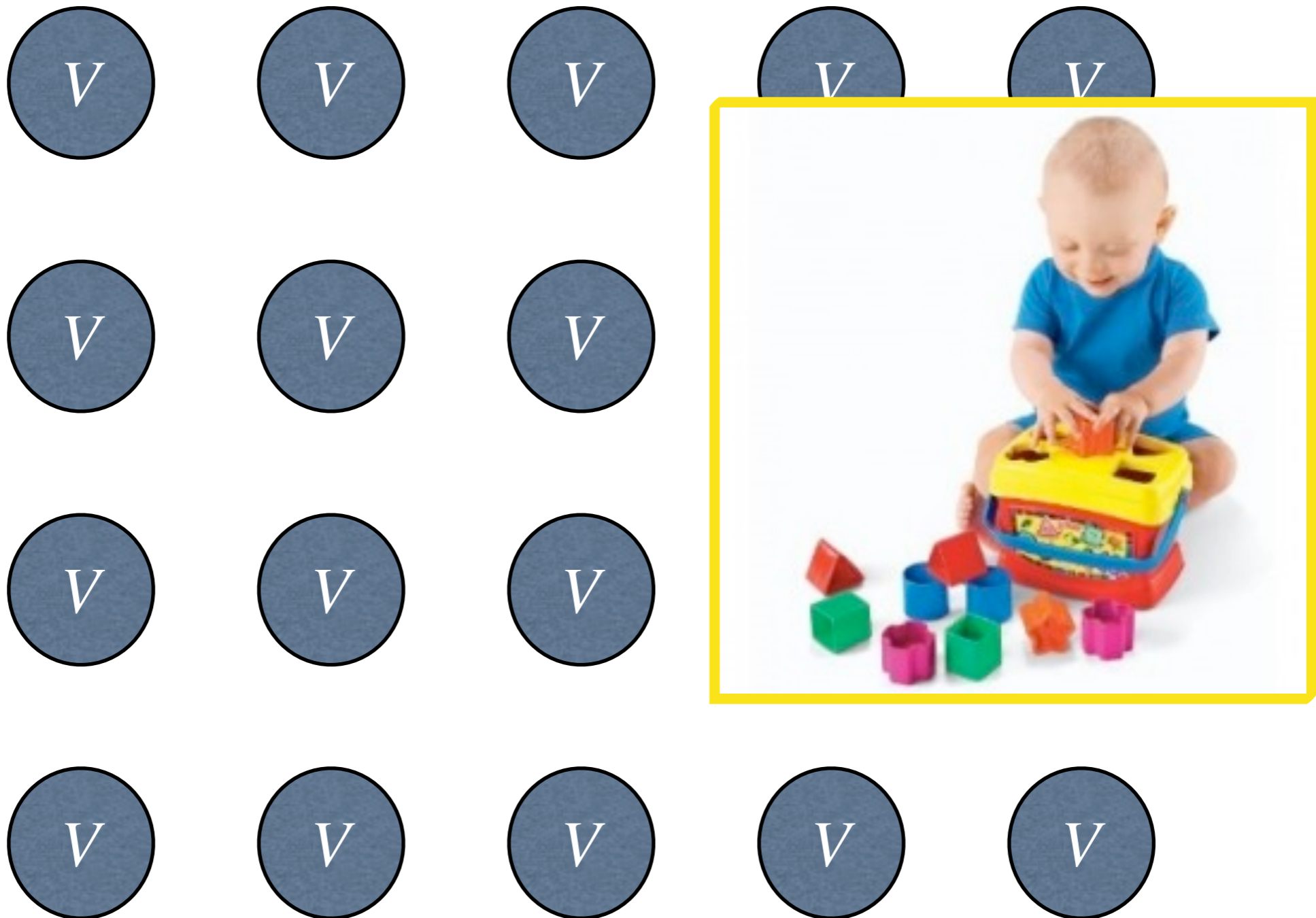
For 3-D, we require general solution for free particles outside spherical potential region with cubic periodicity





# Lüscher: Extension to 3-D

For 3-D, we require general solution for free particles outside spherical potential region with cubic periodicity



# Lüscher (3-D)

Eigenvalue equation

$$\det \left\{ \underbrace{U_{l'm';lm}}_{\text{Lattice geometry}} - \underbrace{e^{i2\delta_{l'}}}_{\text{Scattering phases}} \delta_{ll'} \delta_{mm'} \right\} = 0$$

Lattice geometry

Scattering phases

# Hamiltonian alternative

Construct a basis of non-interacting multi-particle and (bare) resonance states satisfying the periodicity of the lattice

$$\{\Delta, (N\pi)_1, (N\pi)_2, (N\pi)_3, \dots\}$$

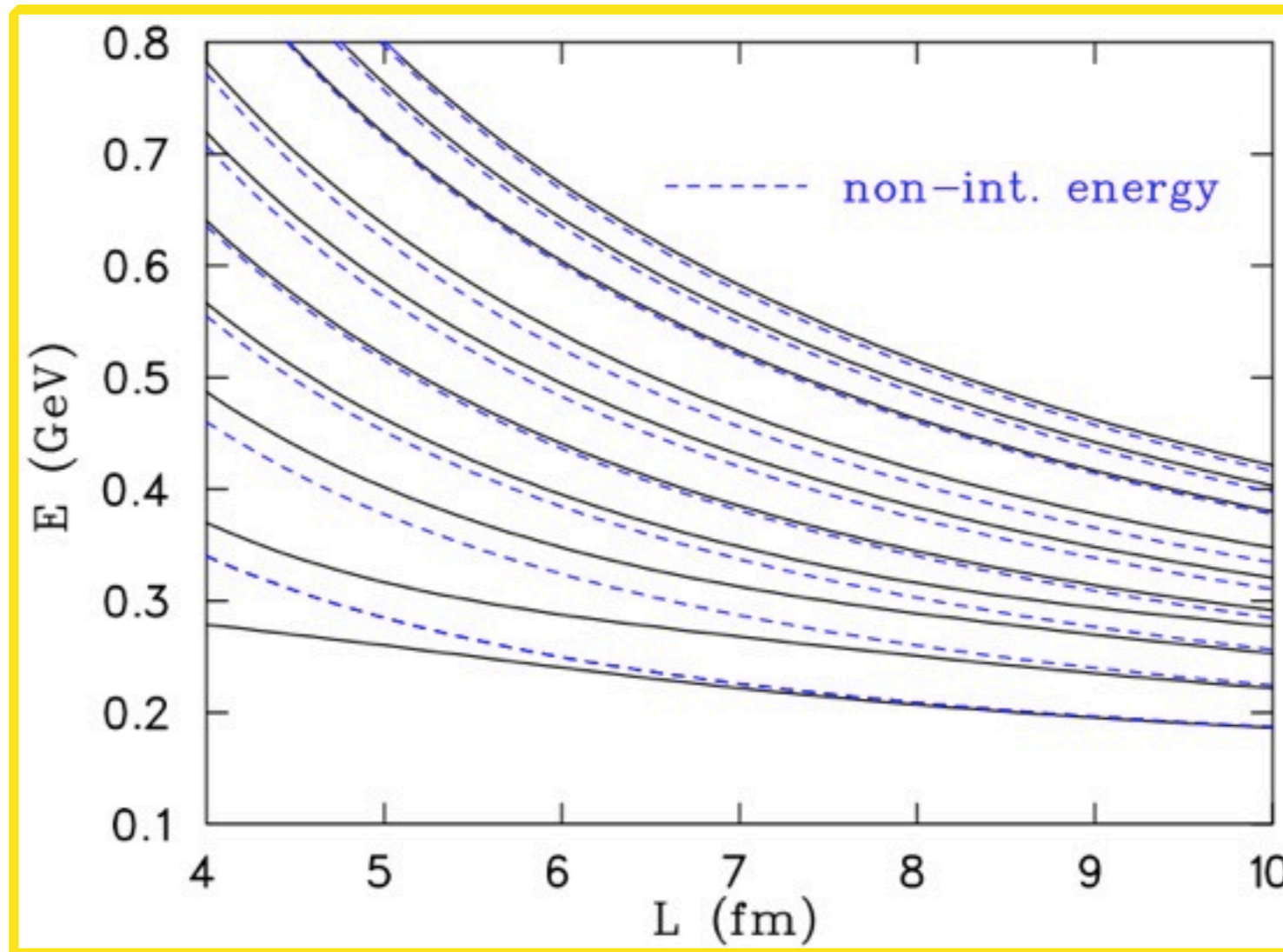
## Bare Hamiltonian

$$H_0|\Delta\rangle = \Delta^0|\Delta\rangle; \quad H_0|(N\pi)_j\rangle = \sqrt{m_\pi^2 + k_j^2}|(N\pi)_j\rangle \quad k_j^2 = j \left(\frac{2\pi}{L}\right)^2$$

## Interaction

$$\langle\Delta|H_I|(N\pi)_j\rangle \sim g(k_j)$$

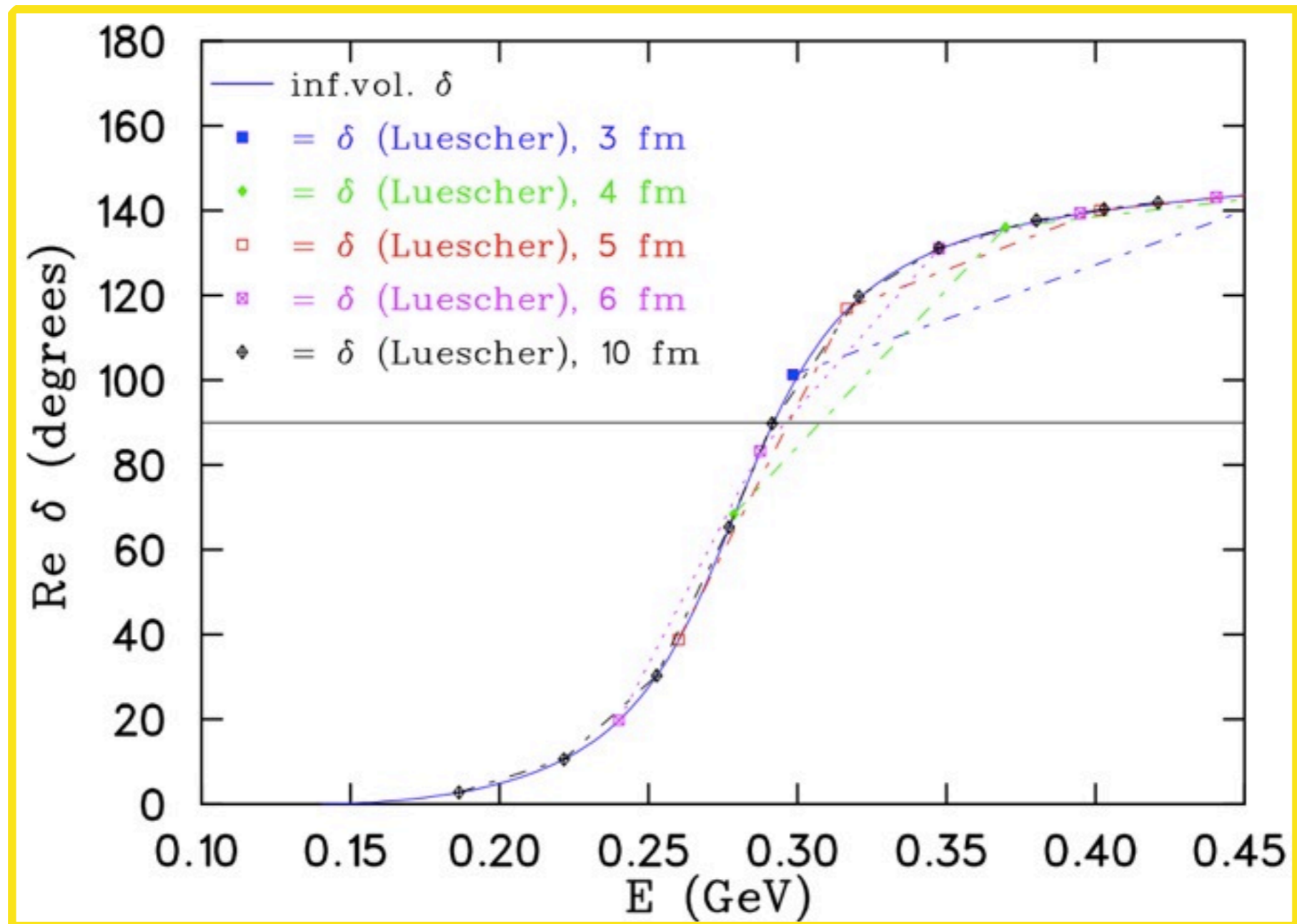
Lattice eigenstates determined by  $\det(H - \lambda I) = 0$



# Hamiltonian spectrum

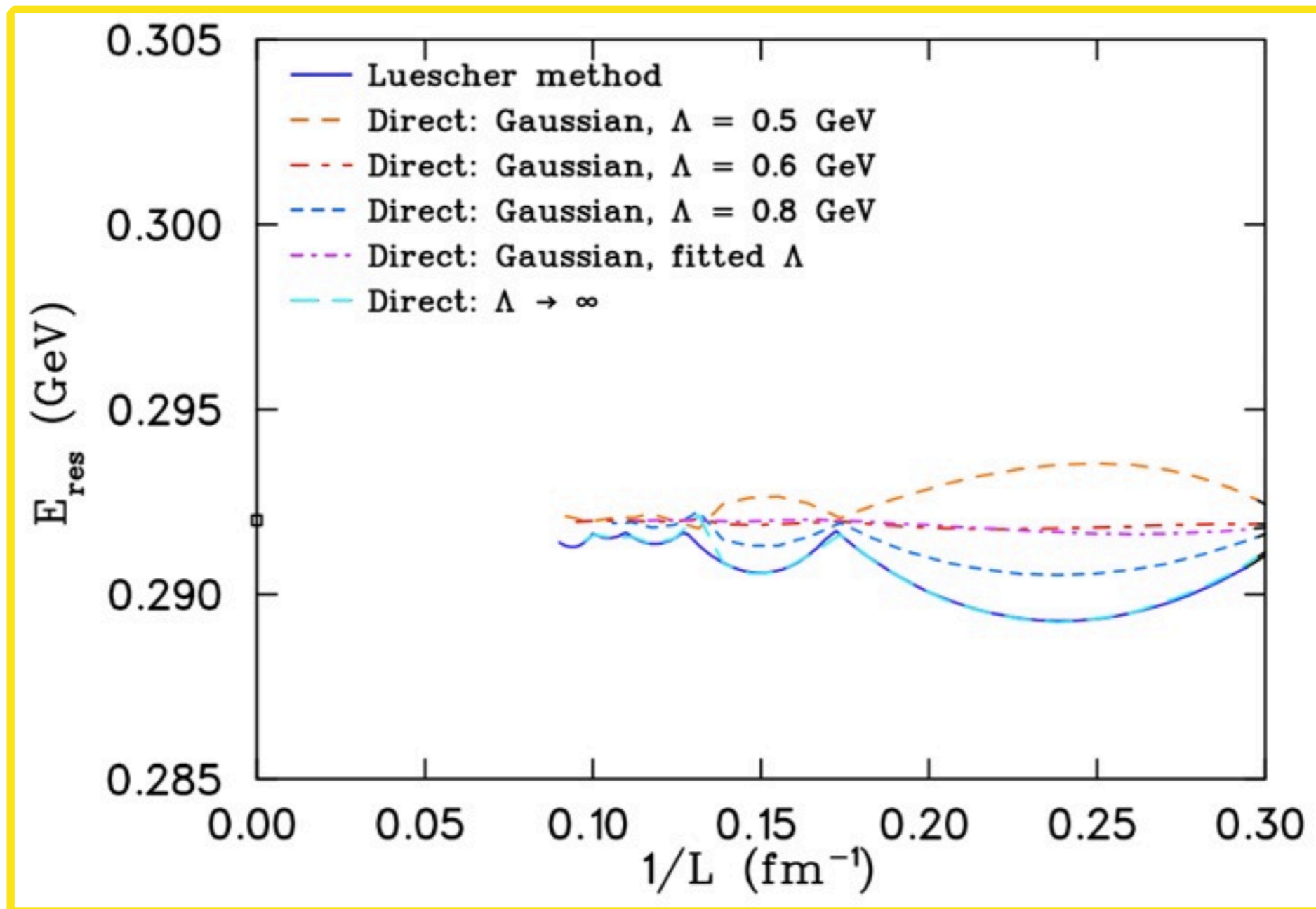
Identical to Lüscher (to the eye)





## Resonance parameter estimation

On each volume determine phase shift from Lüscher equation  
 Use Breit-Wigner to interpolate points and extract mass and width  
**OR** Directly fit Hamiltonian parameters to spectrum



## Resonance mass

Lüscher & Hamiltonian approach both give a reliable estimate of resonance mass at finite  $L$

# Inelastic channels

# Extension to coupled channel

Extension of Lüscher by He, Feng & Liu JHEP(2005)

$$S^{(l)}(E) = \begin{pmatrix} \eta_l e^{2i\delta_1^l} & i\sqrt{1-\eta_l^2} e^{i(\delta_1^l + \delta_2^l)} \\ i\sqrt{1-\eta_l^2} e^{i(\delta_1^l + \delta_2^l)} & \eta_l e^{2i\delta_2^l} \end{pmatrix}$$

Scattering parameters as a function of E

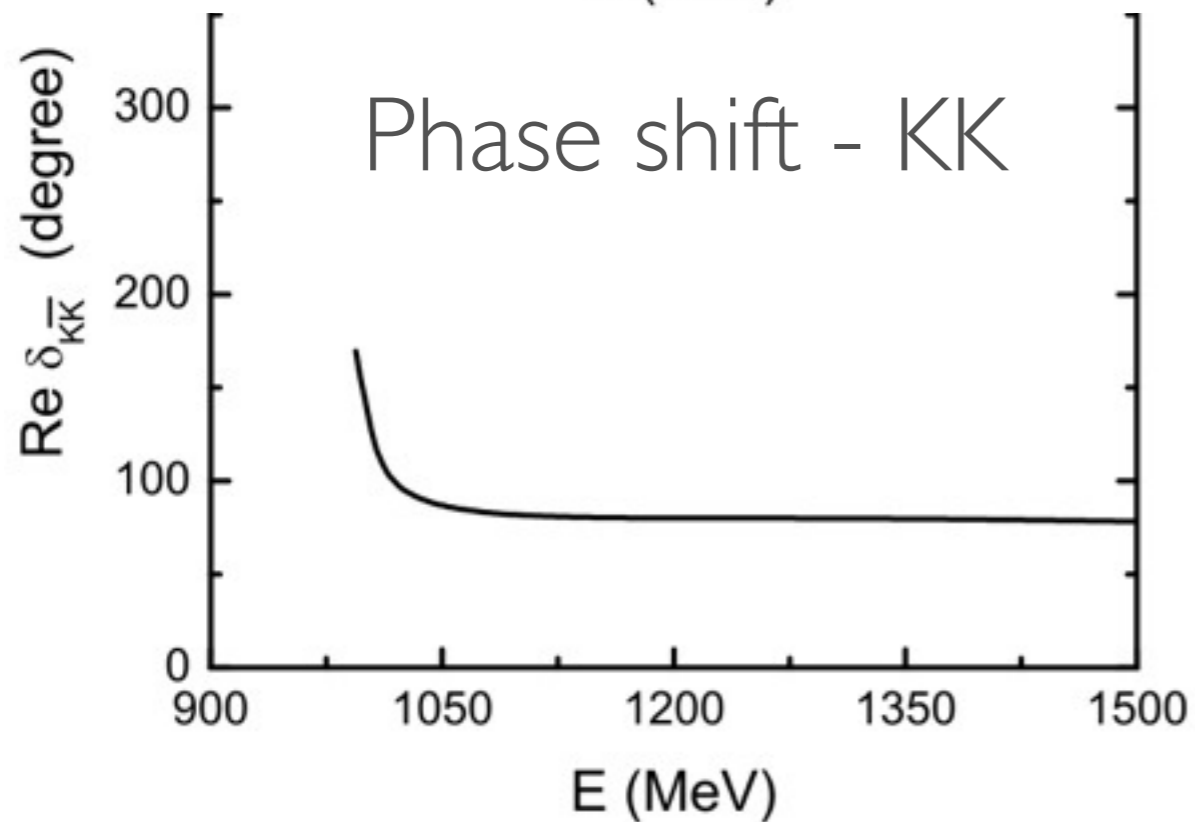
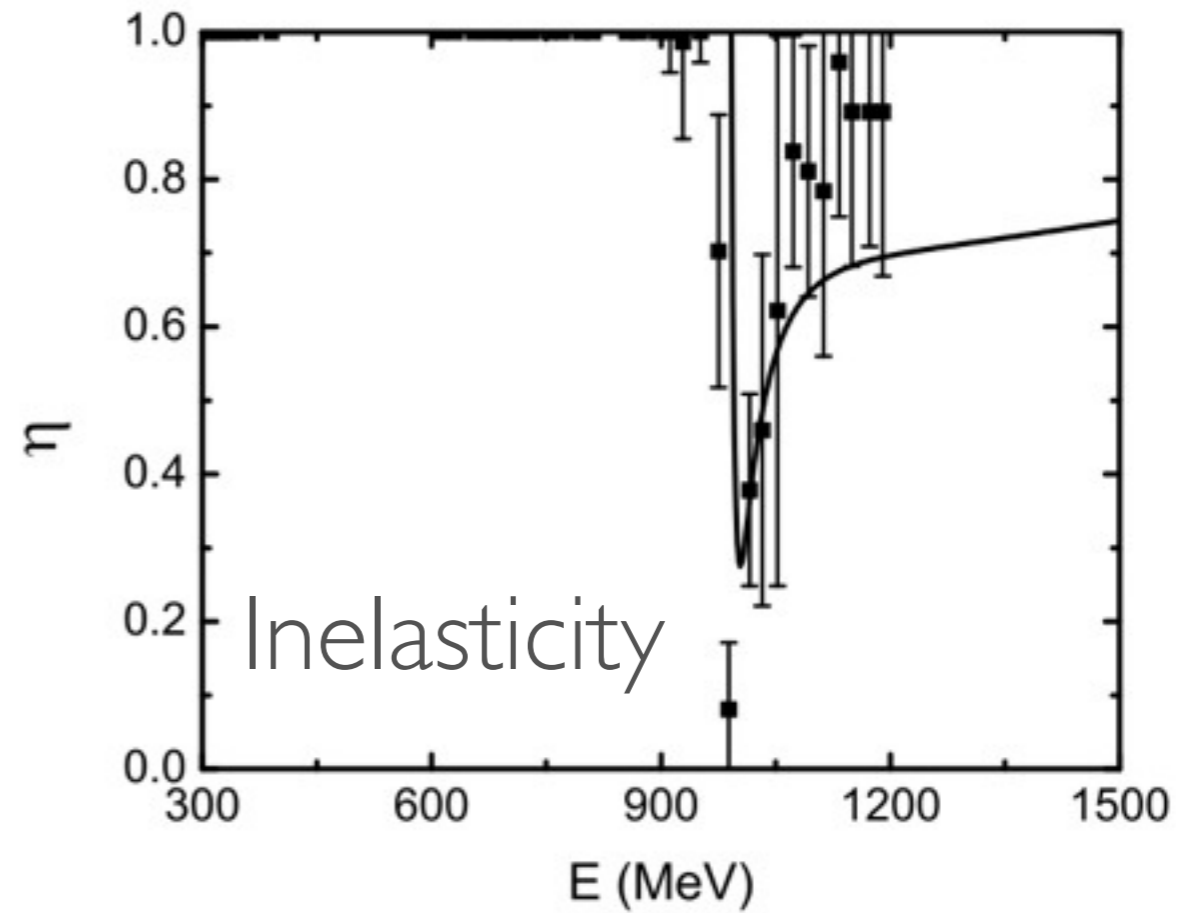
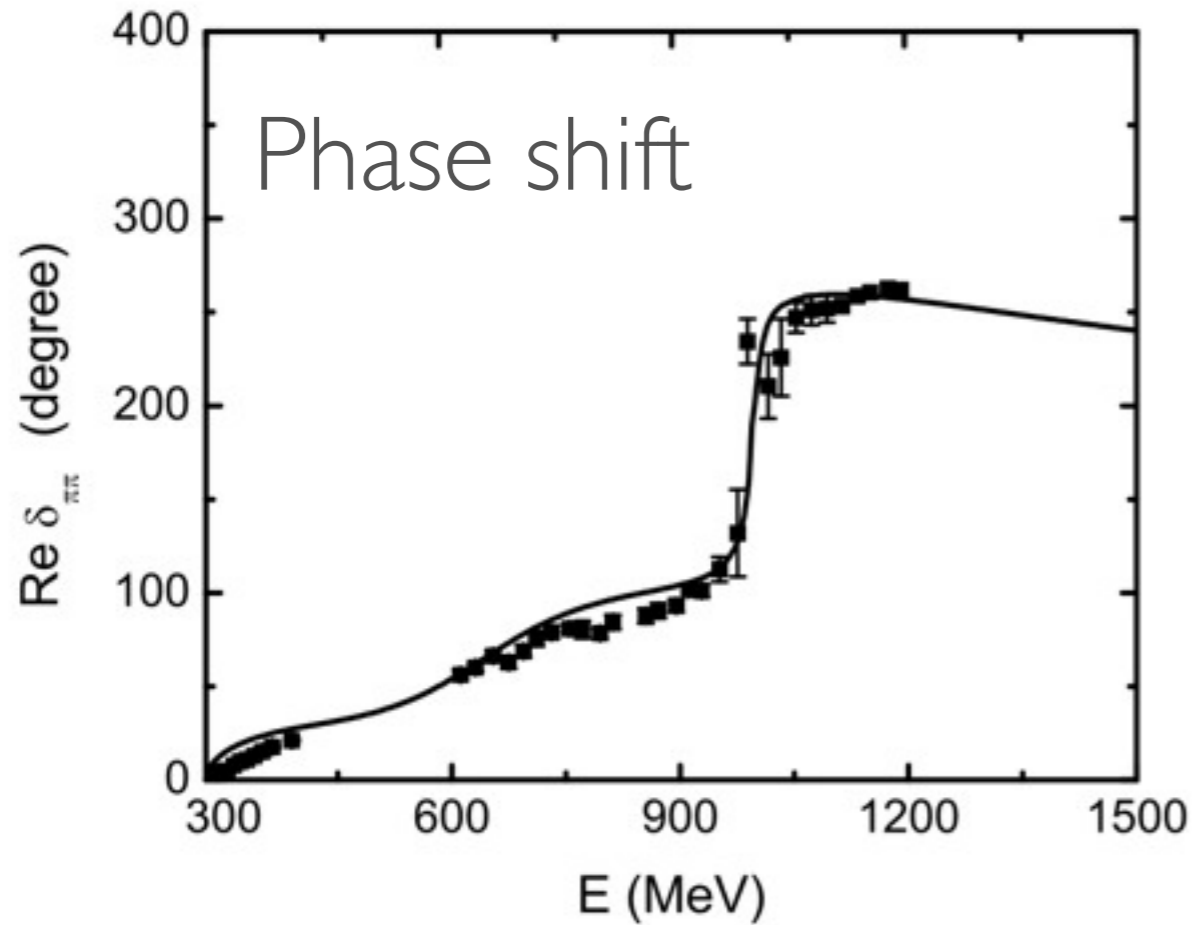
$$\delta_1, \delta_2, \eta$$

Finite-volume eigenvalue equation (s-wave)

$$\cos(\Delta_1 + \Delta_2 - \delta_1^0 - \delta_2^0) = \eta \cos(\Delta_1 - \Delta_2 - \delta_1^0 + \delta_2^0)$$

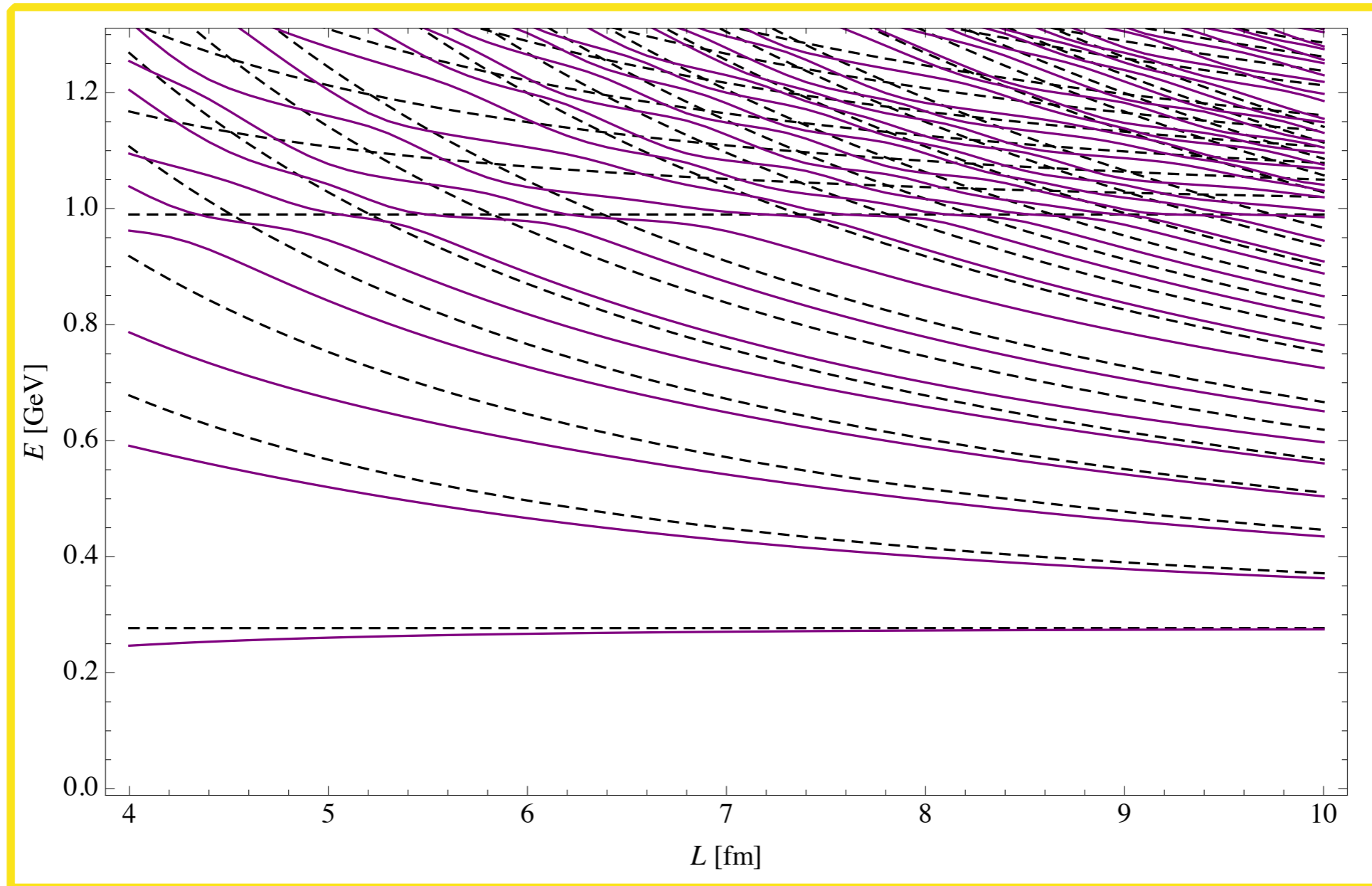
$$\Delta_i = \mathcal{M}_{00;00}(q_i^2) = \frac{\mathcal{Z}_{00}(1; q_i^2)}{\pi^{3/2} q_i}$$

# J=0 Isoscalar $\pi\pi, K\bar{K}$



Wu, Lee, Thomas, RDY, arXiv:1402.4868

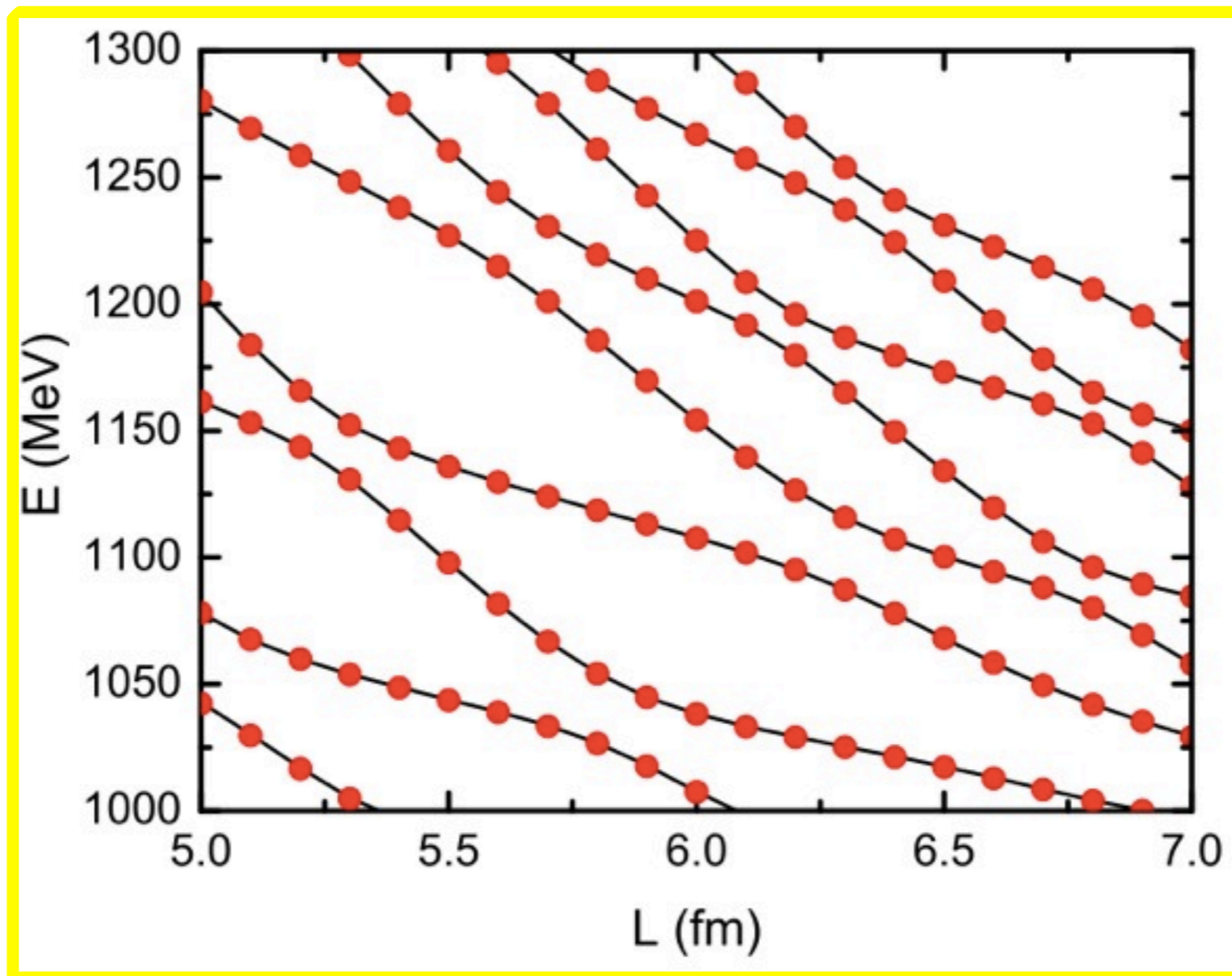




## Hamiltonian spectrum

$$\{\sigma, (\pi\pi)_0, (\pi\pi)_1, (\pi\pi)_2, \dots, (K\bar{K})_0, (K\bar{K})_1, (K\bar{K})_2, \dots\}$$

Dashed lines are non-interacting levels

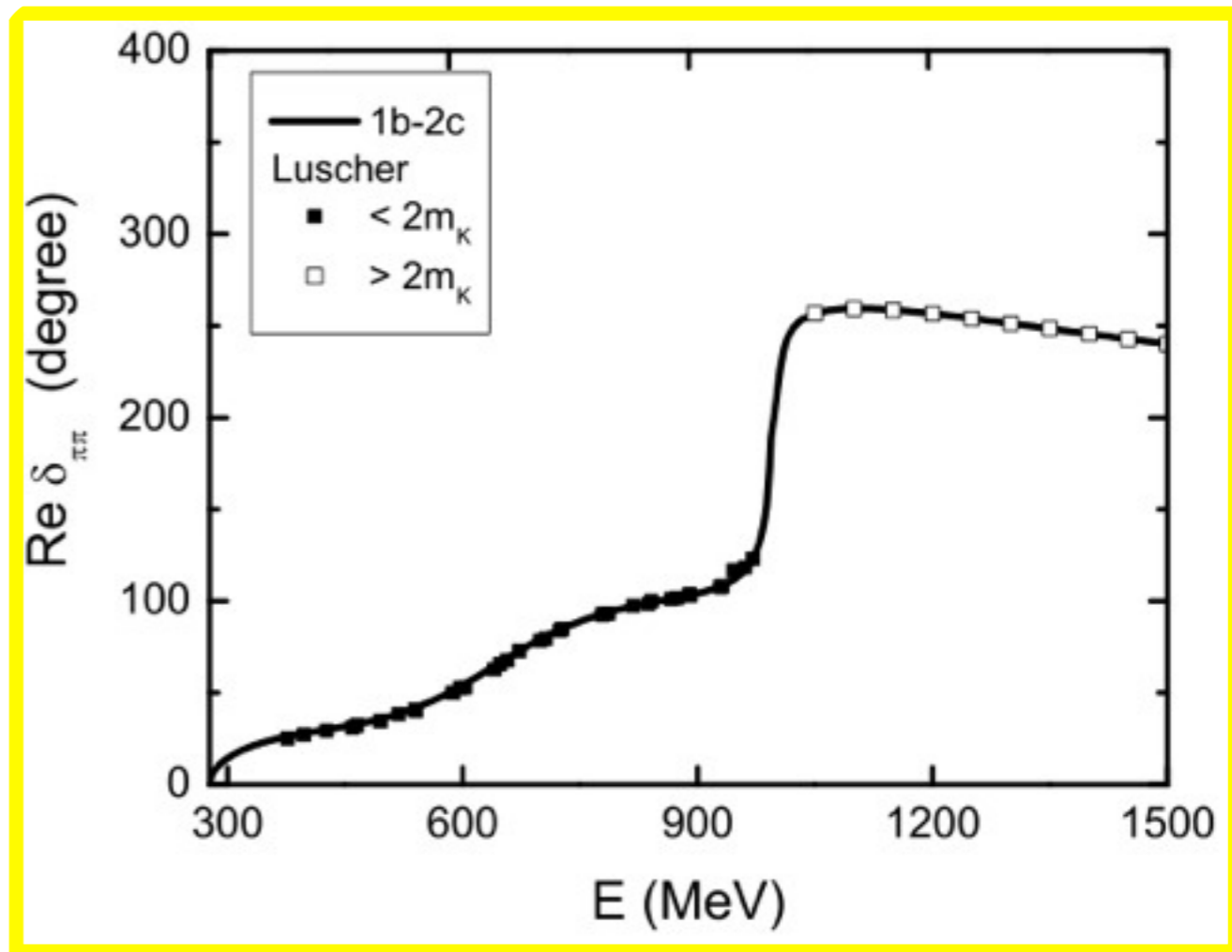
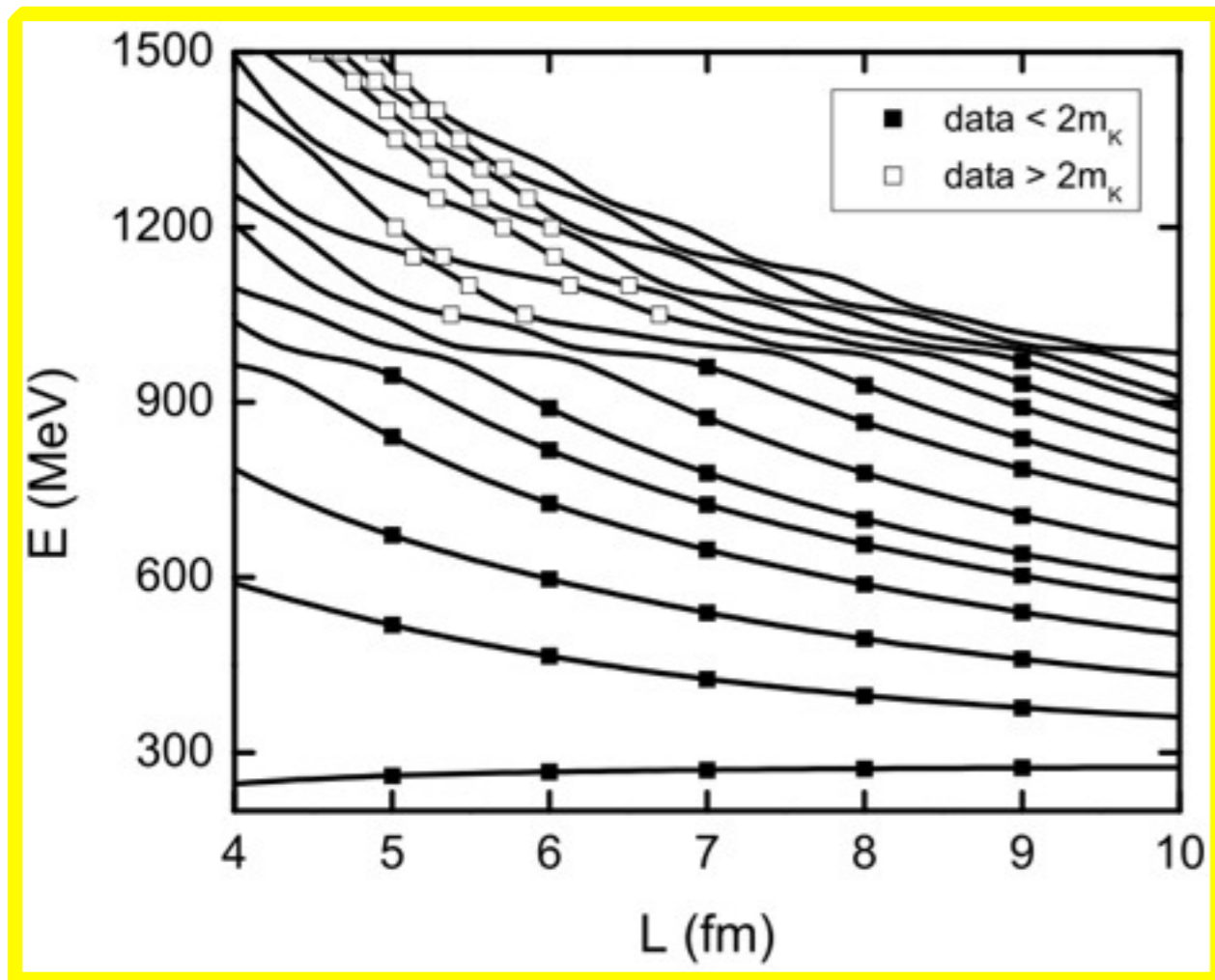


## FV Spectrum “Prediction”

Lüscher: red dots

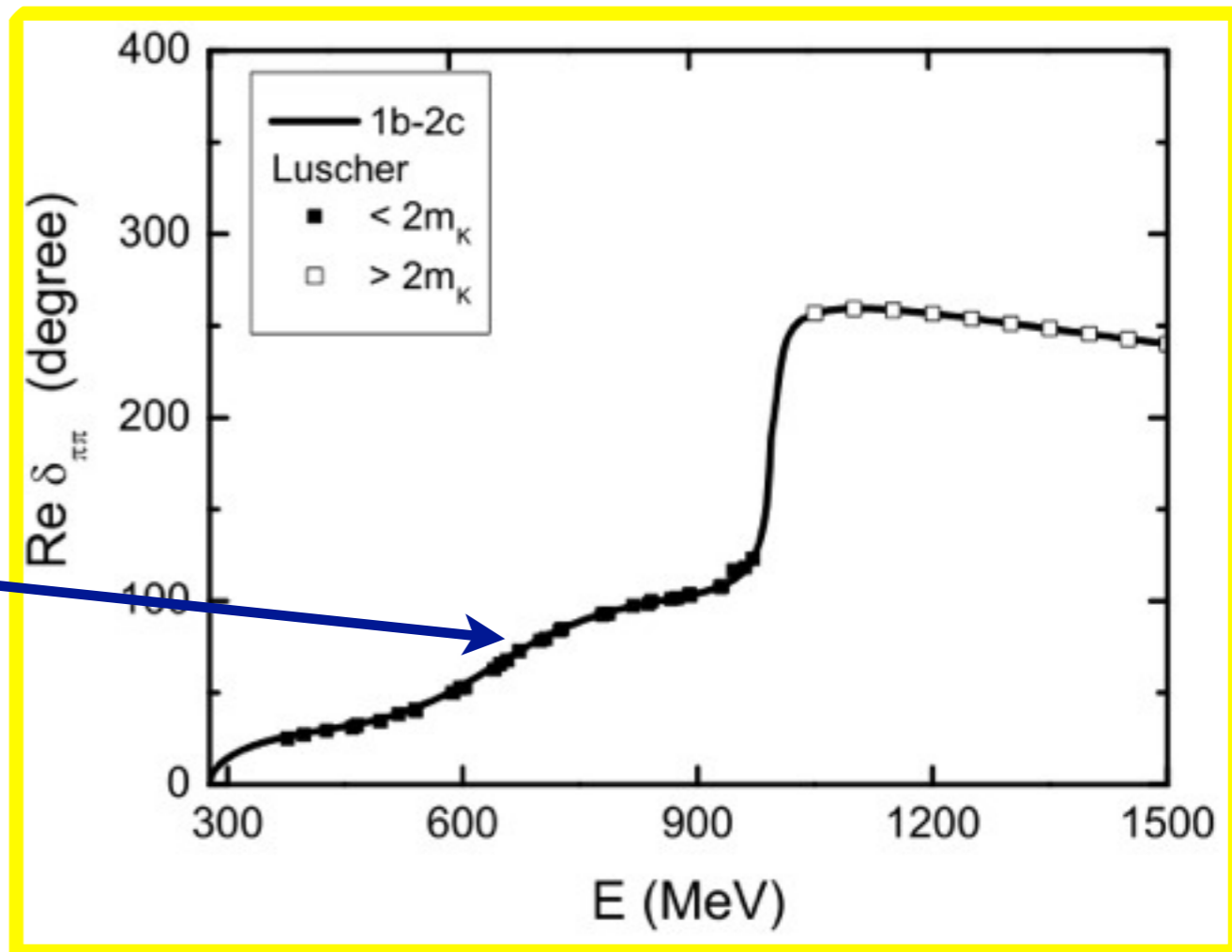
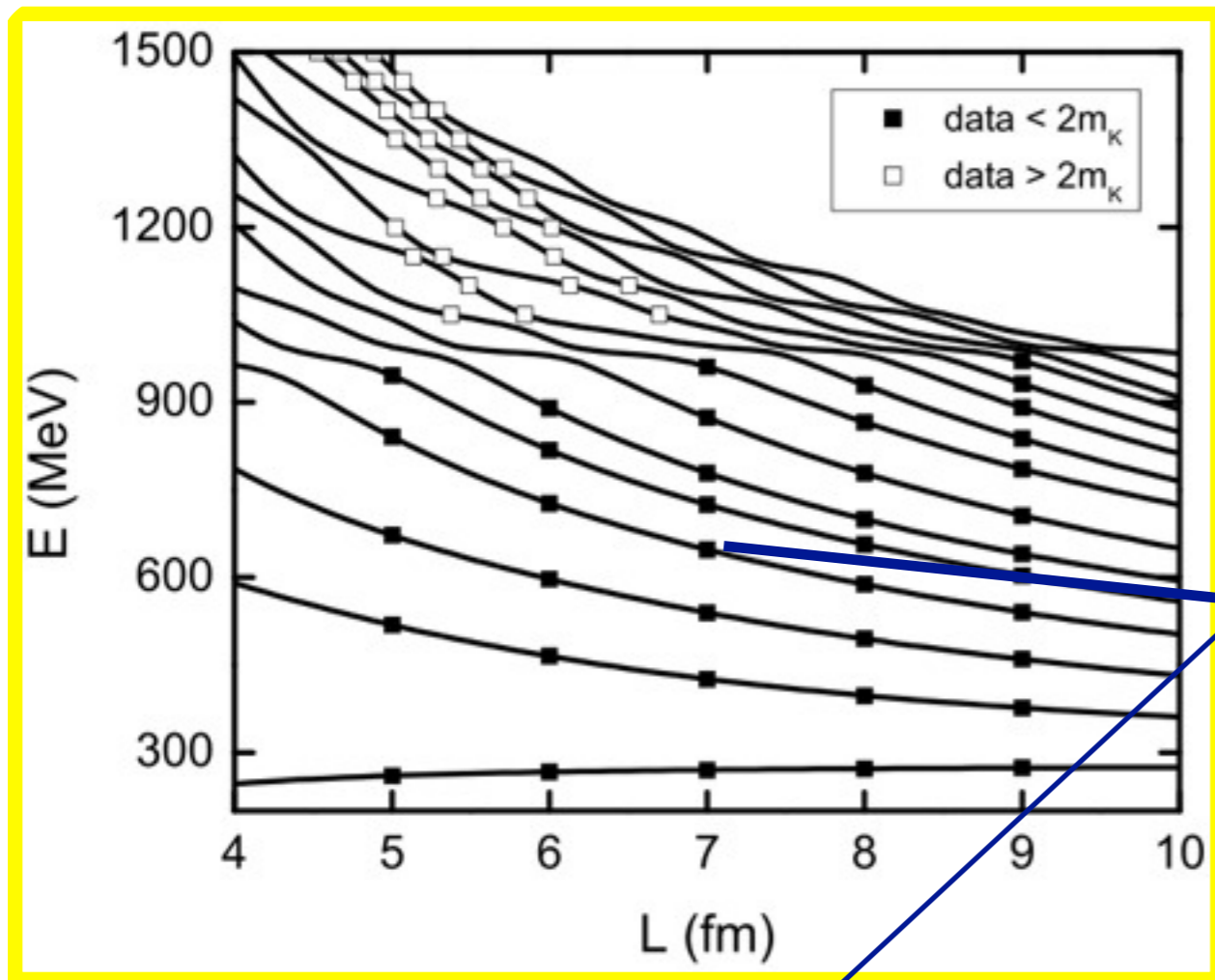
Hamiltonian: black curves

Applications for lattice QCD,  
we want the inverse:  
ie. *Spectrum*  $\rightarrow$  *Phase shifts*



Inversion with multi-channel Lüscher

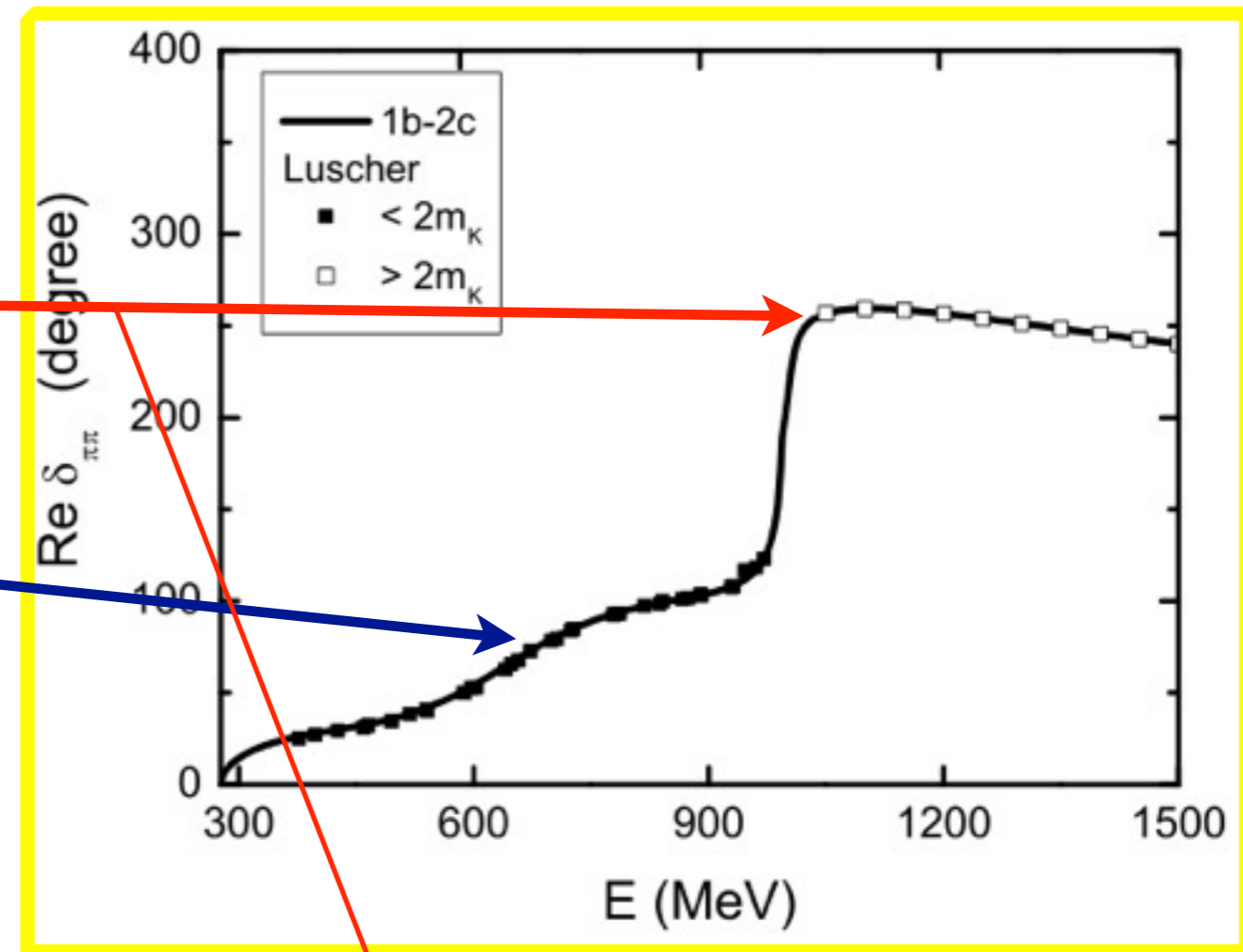
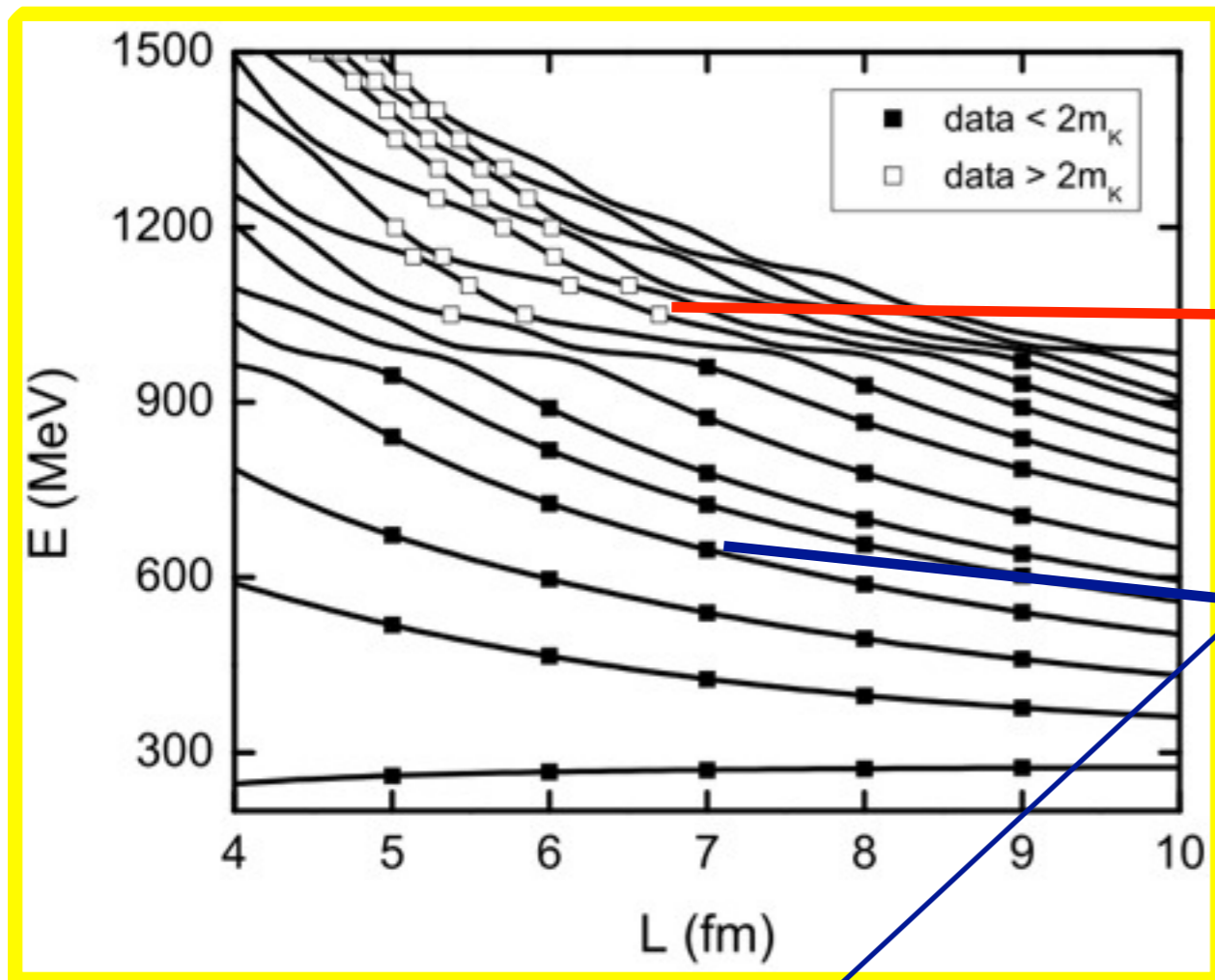




Single channel: straightforward

Inversion with multi-channel Lüscher





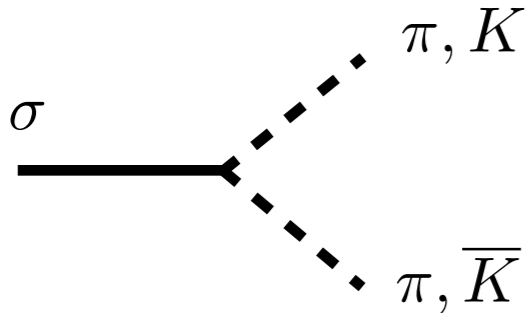
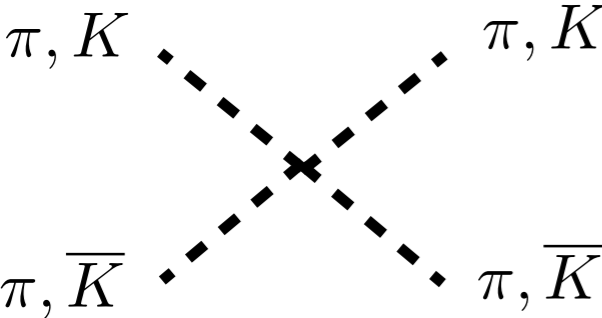
Single channel: straightforward

2 channels: requires degeneracy at independent L

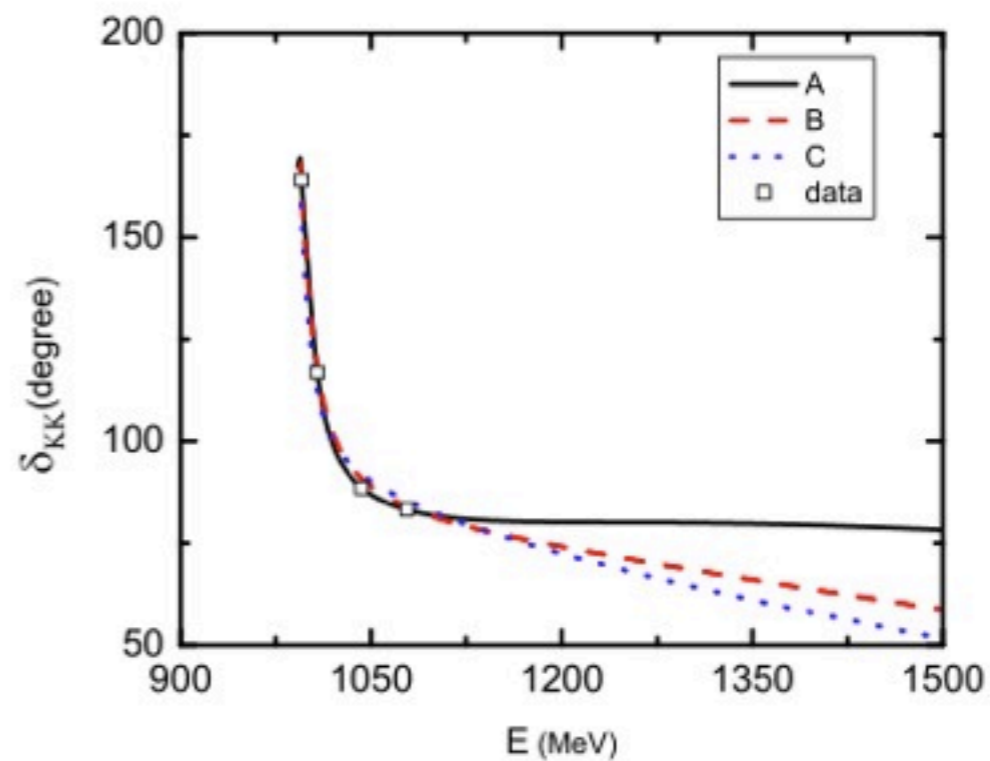
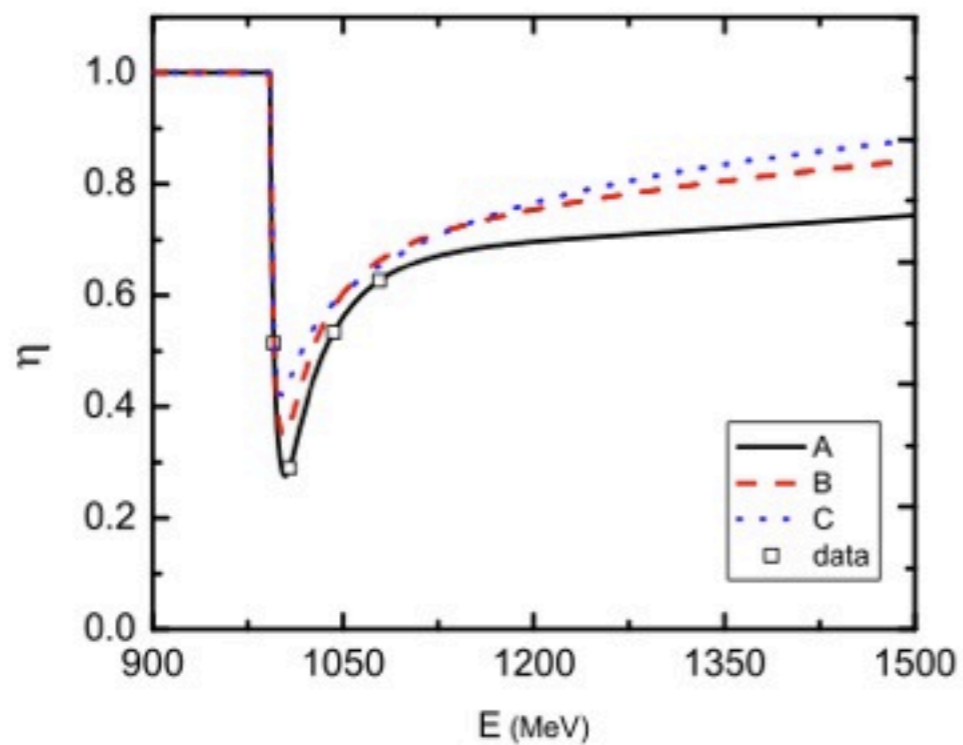
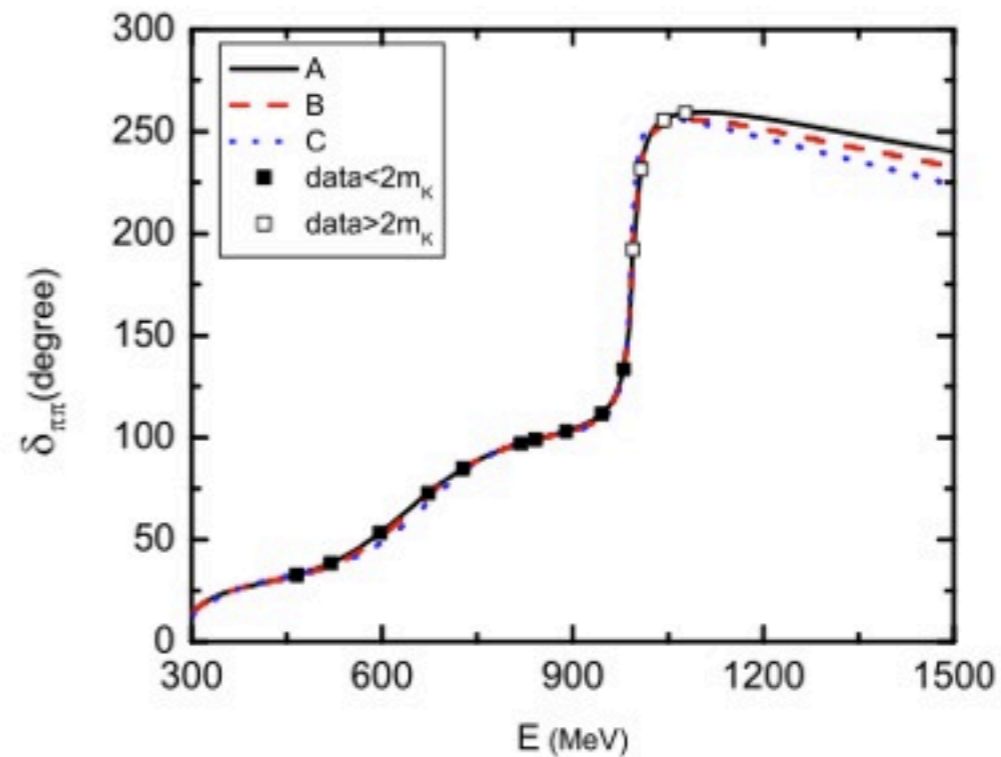
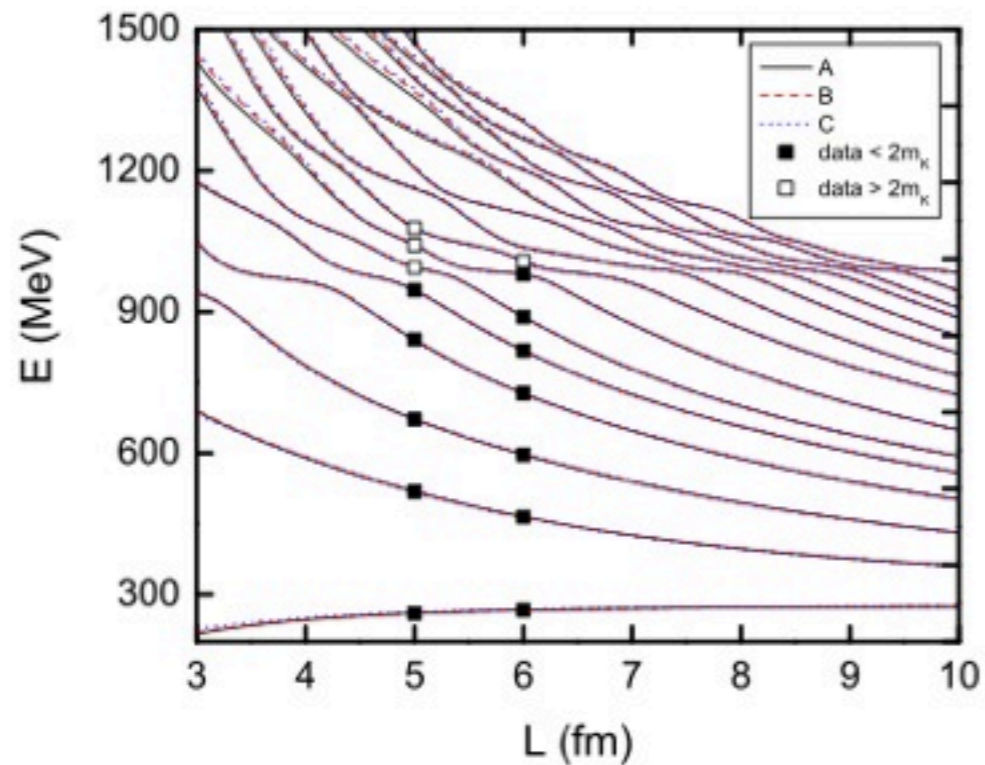
# Inversion with multi-channel Lüscher

# Inversion with Hamiltonian

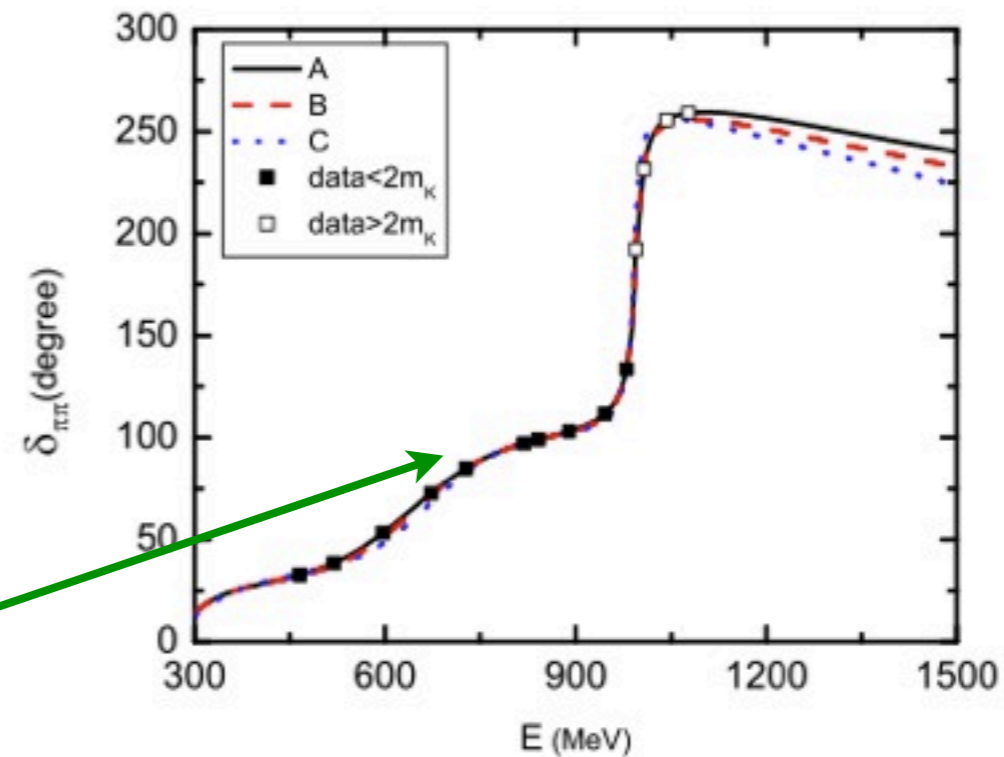
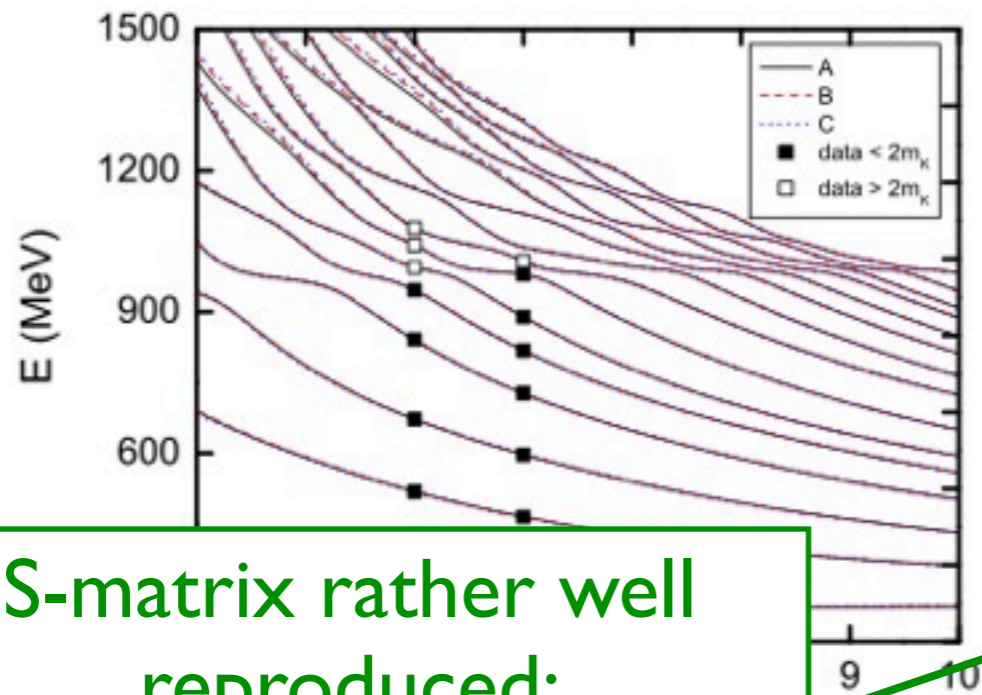
Requires a choice of Hamiltonian

	Model A	Model B	Model C
	$\frac{g_{\sigma,\alpha}}{(1 + (c_{\alpha}k)^2)}$	$\frac{g_{\sigma,\alpha}}{(1 + (c_{\alpha}k)^2)^2}$	$g_{\sigma,\alpha}e^{-(c_{\alpha}k)^2}$
	$\frac{G_{\sigma,\alpha}}{(1 + (d_{\alpha}k)^2)(1 + (d_{\beta}k)^2)}$	$\frac{G_{\sigma,\alpha}}{(1 + (d_{\alpha}k)^2)^2(1 + (d_{\beta}k)^2)^2}$	$G_{\alpha,\beta}e^{-(d_{\alpha}k)^2}e^{-(d_{\beta}k)^2}$

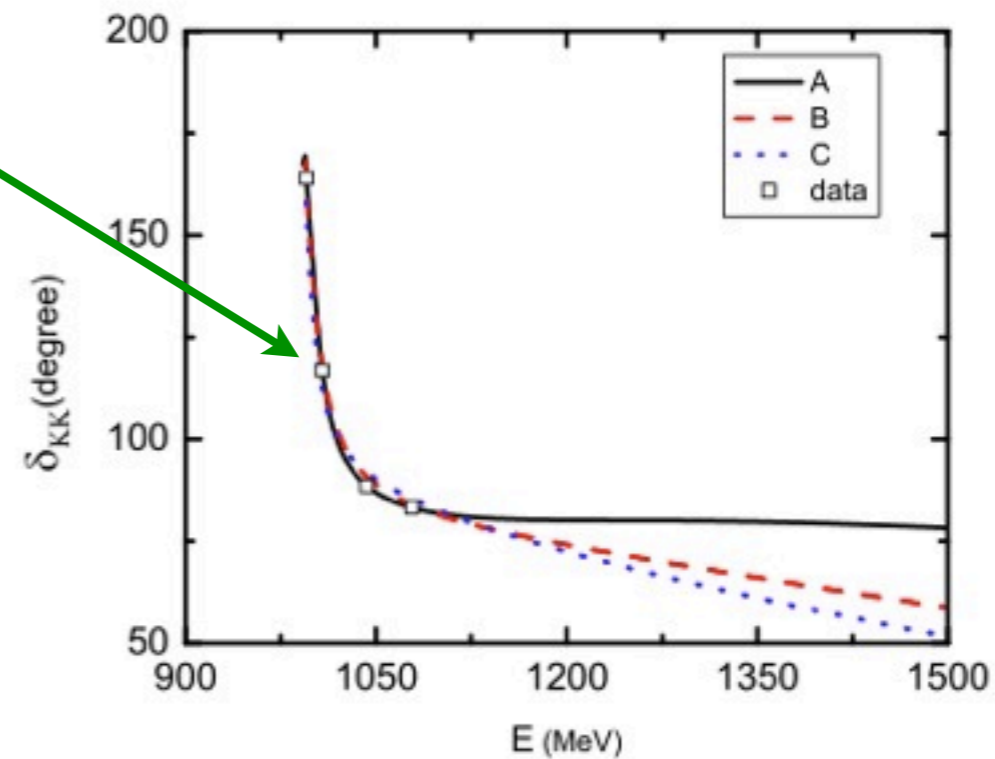
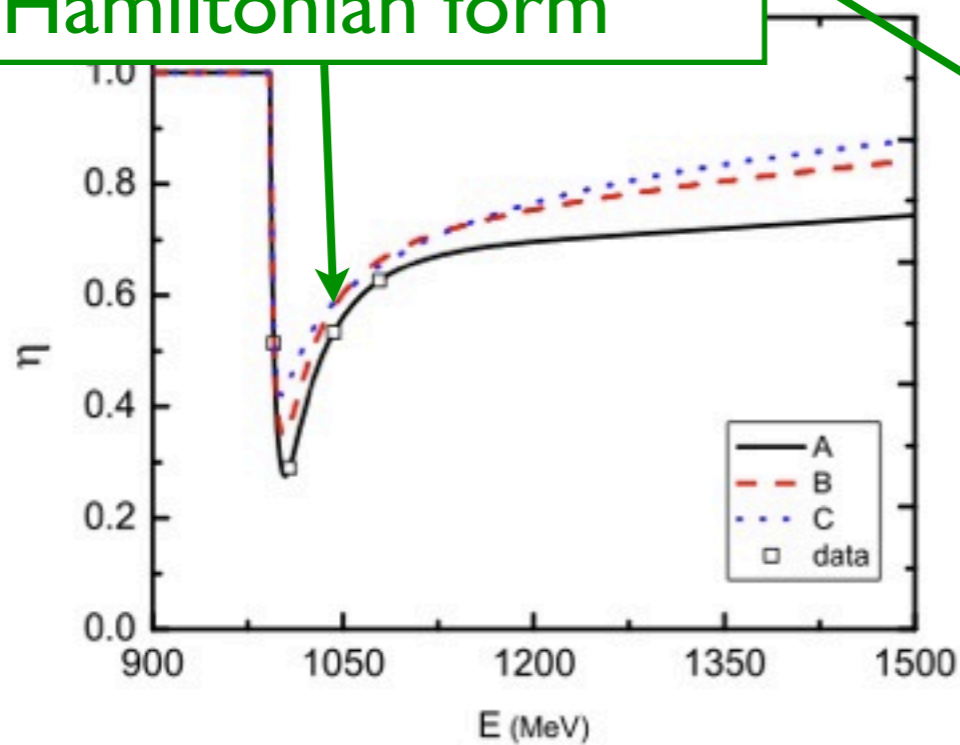
$$\alpha, \beta = (\pi\pi), (K\bar{K})$$



**Inversion with Hamiltonian**  
 Just 2 box sizes (8 eigenstates)

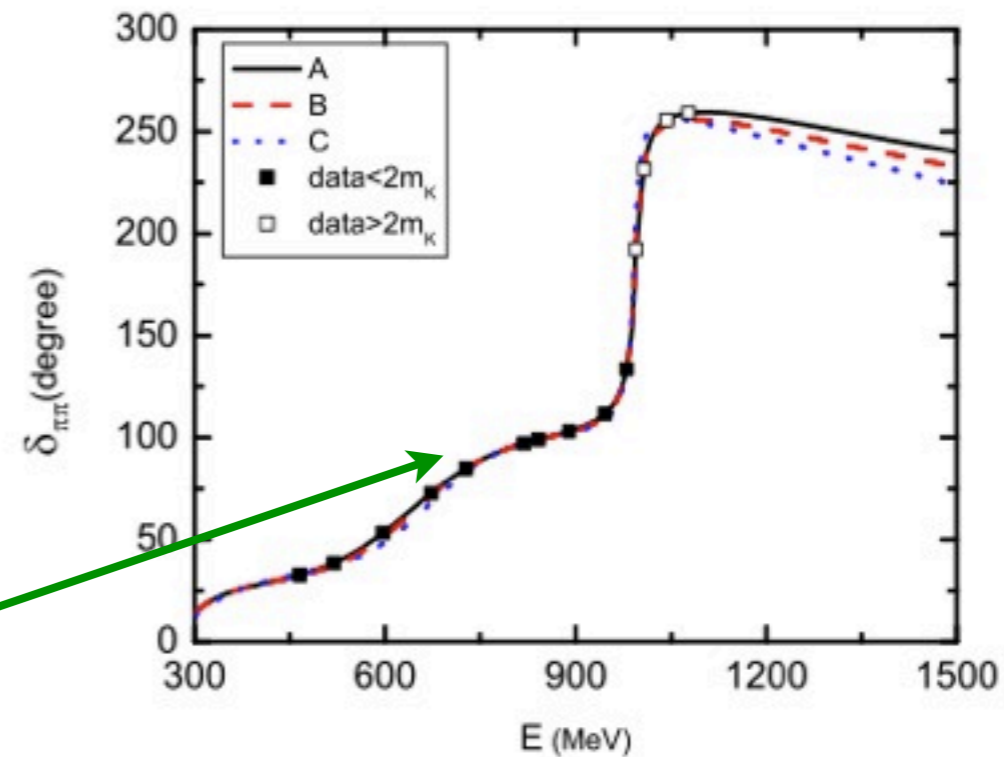
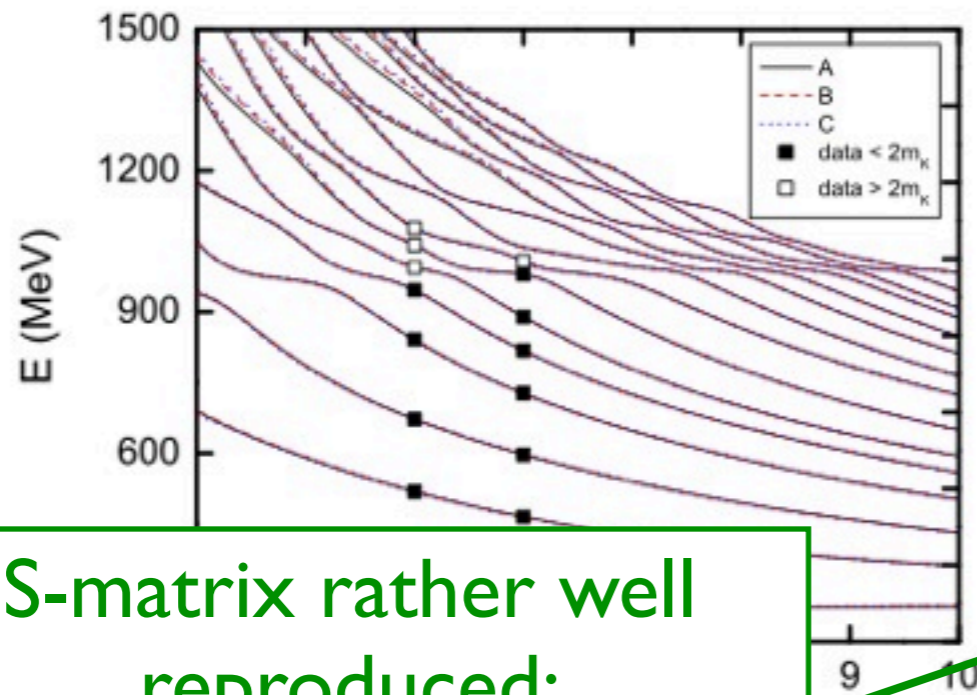


S-matrix rather well reproduced:  
Independent of Hamiltonian form

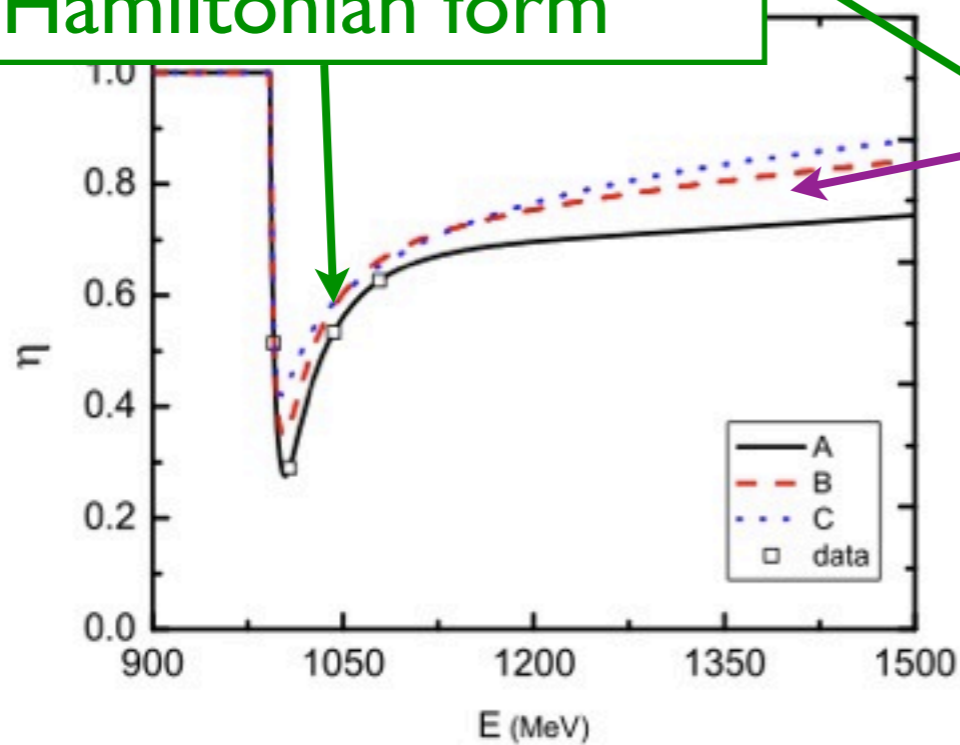


Inversion with Hamiltonian  
Just 2 box sizes (8 eigenstates)

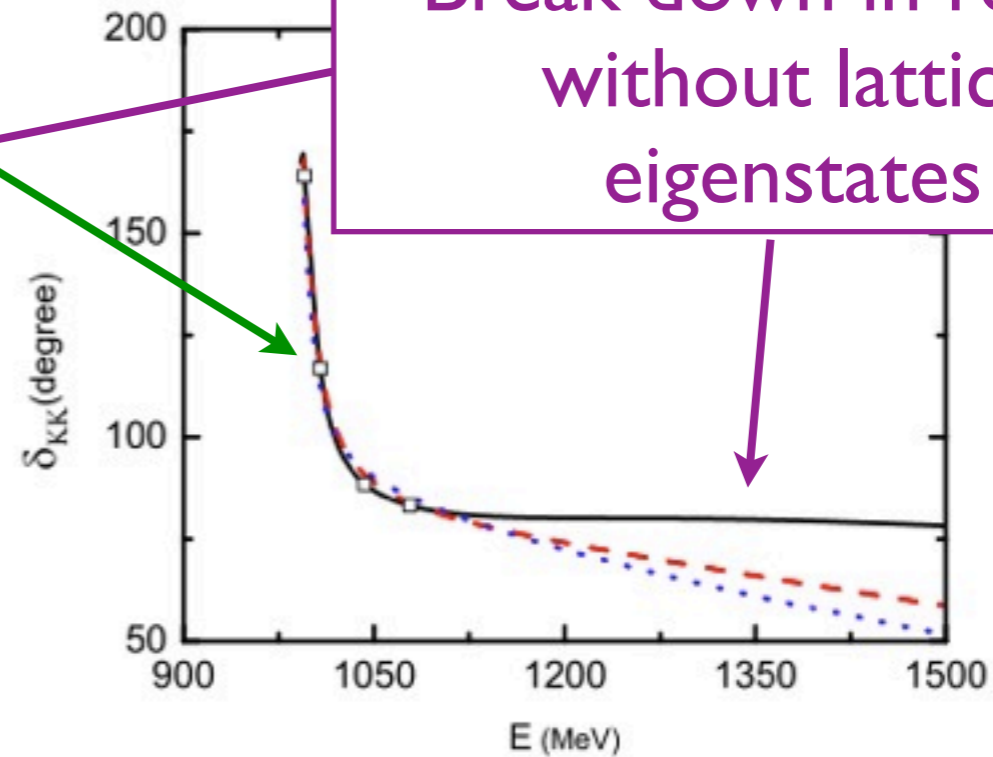




S-matrix rather well reproduced:  
Independent of Hamiltonian form

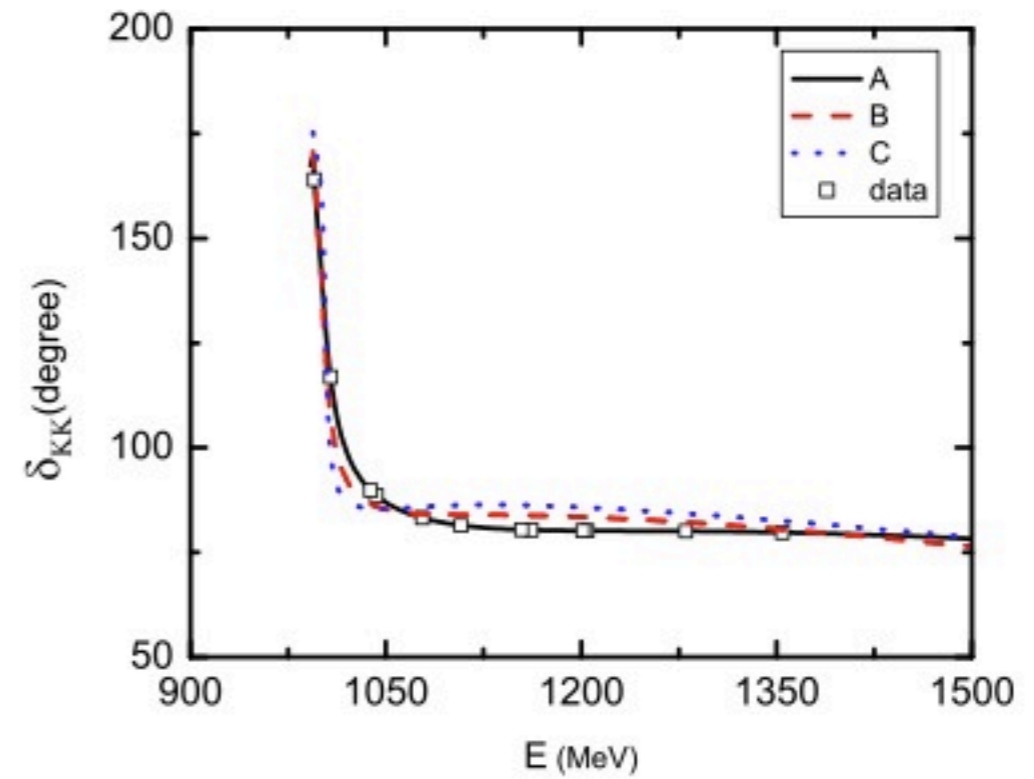
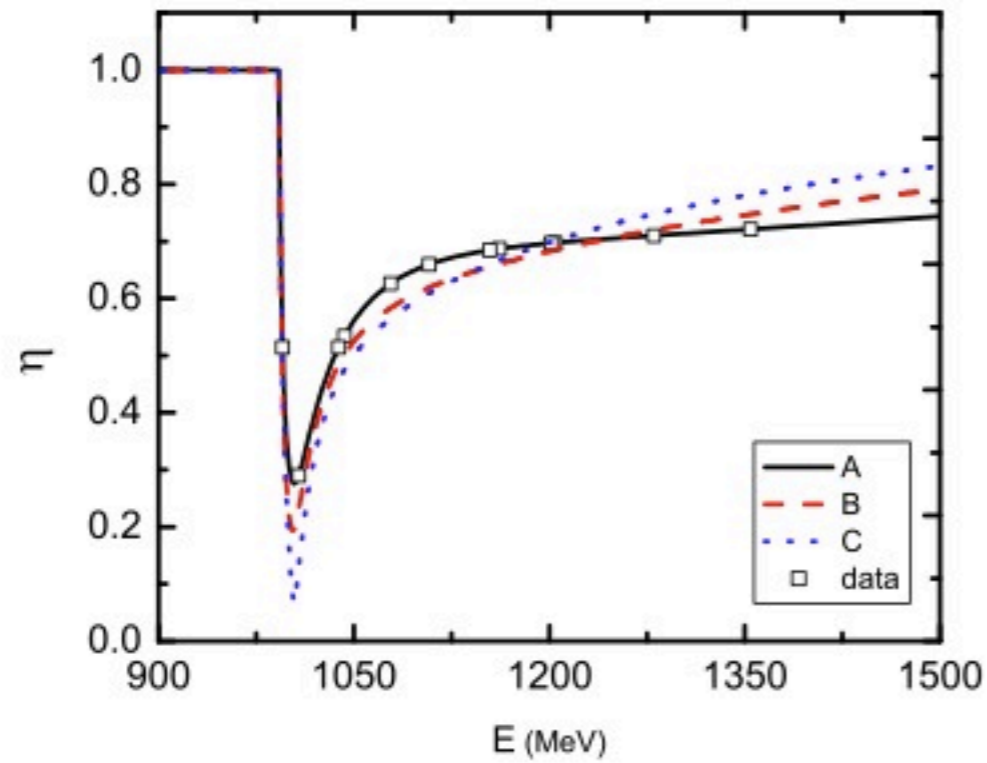
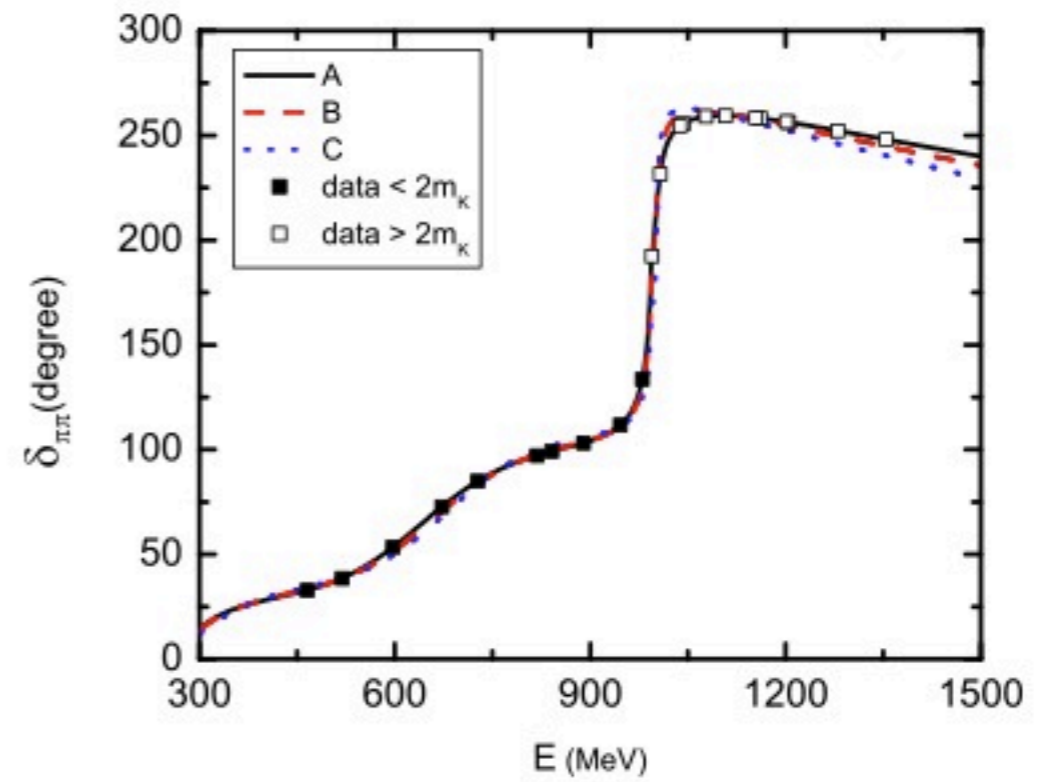
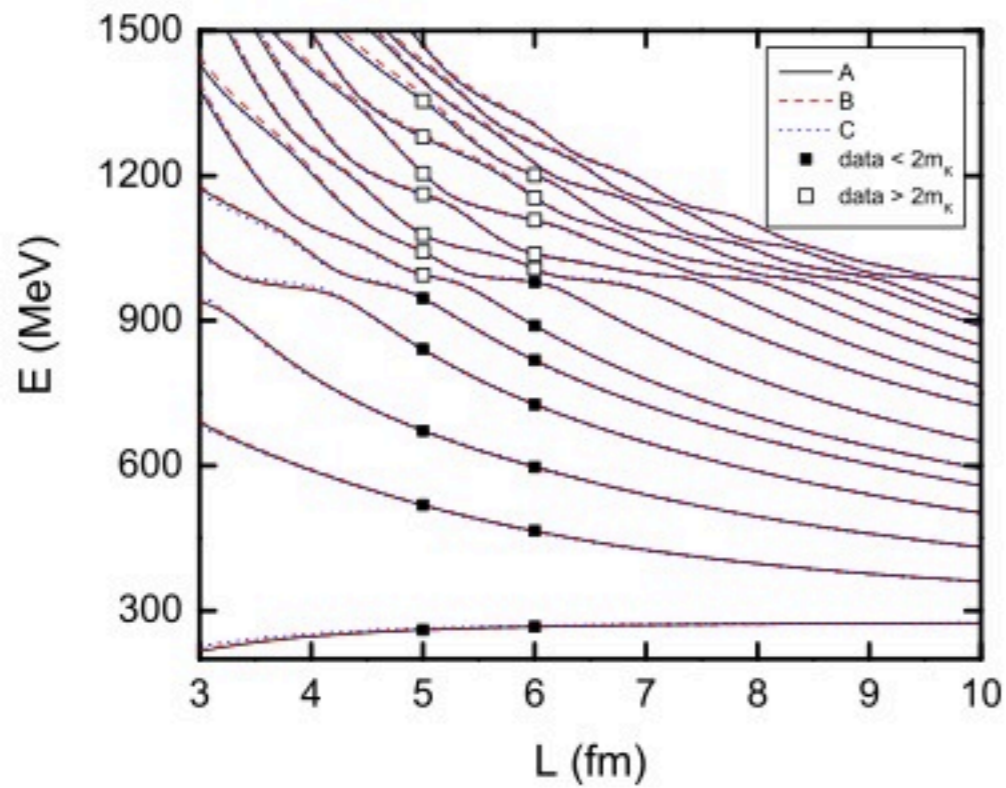


Break down in region without lattice eigenstates

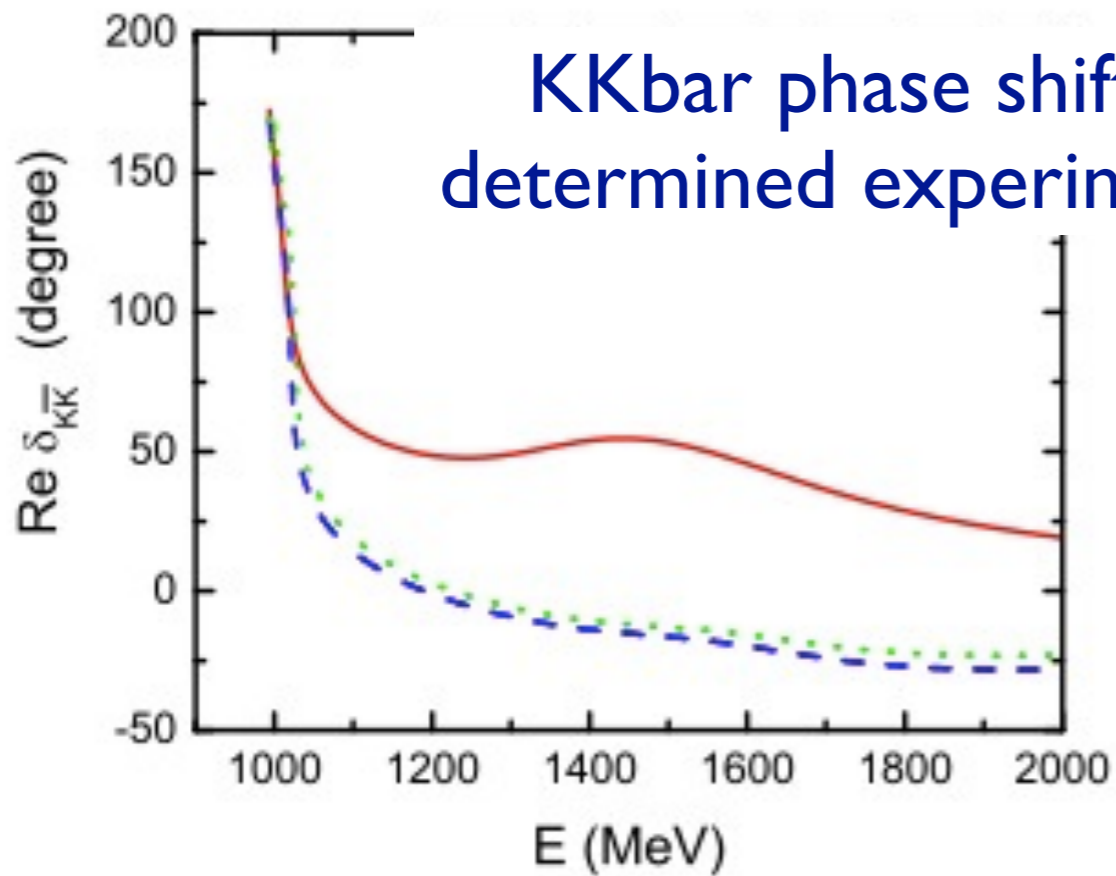
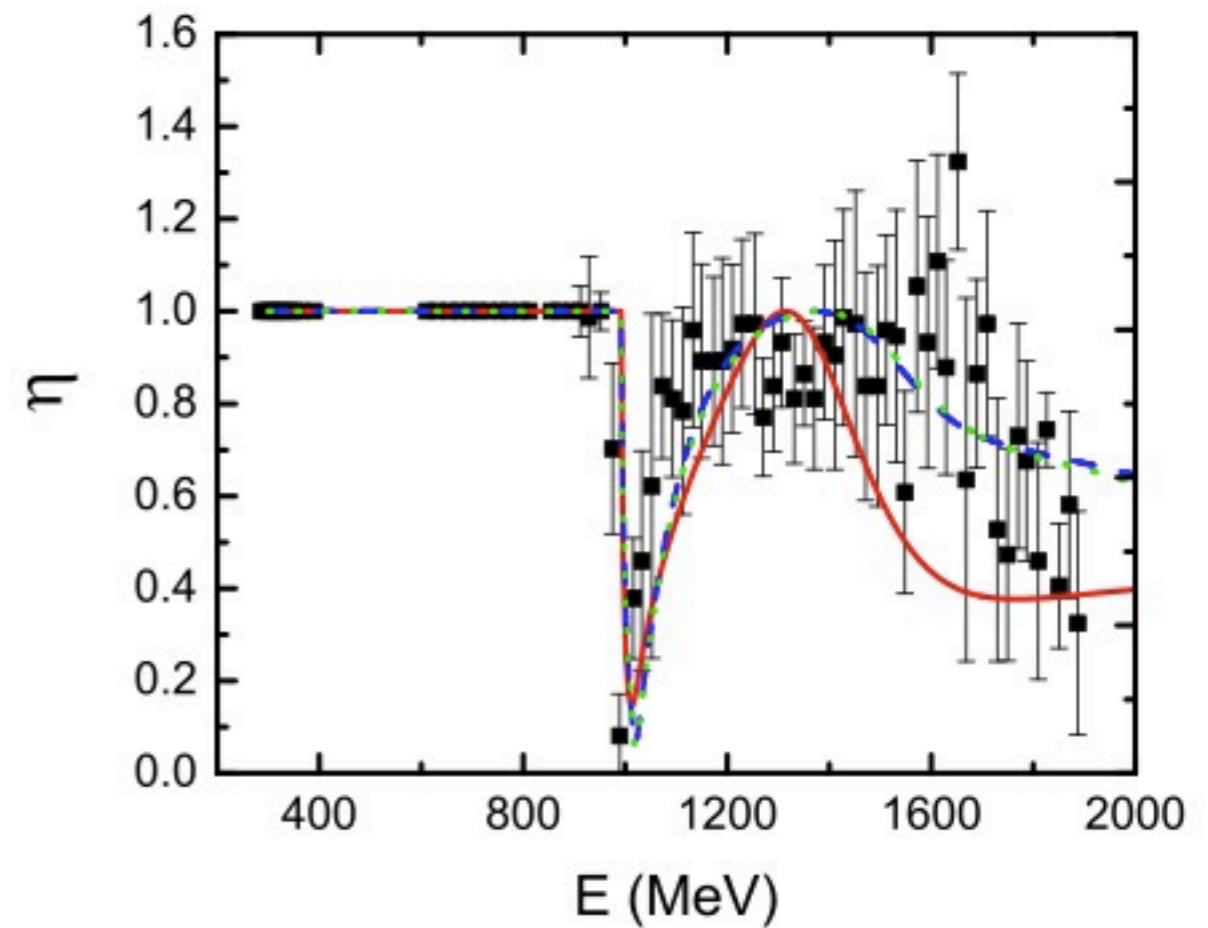
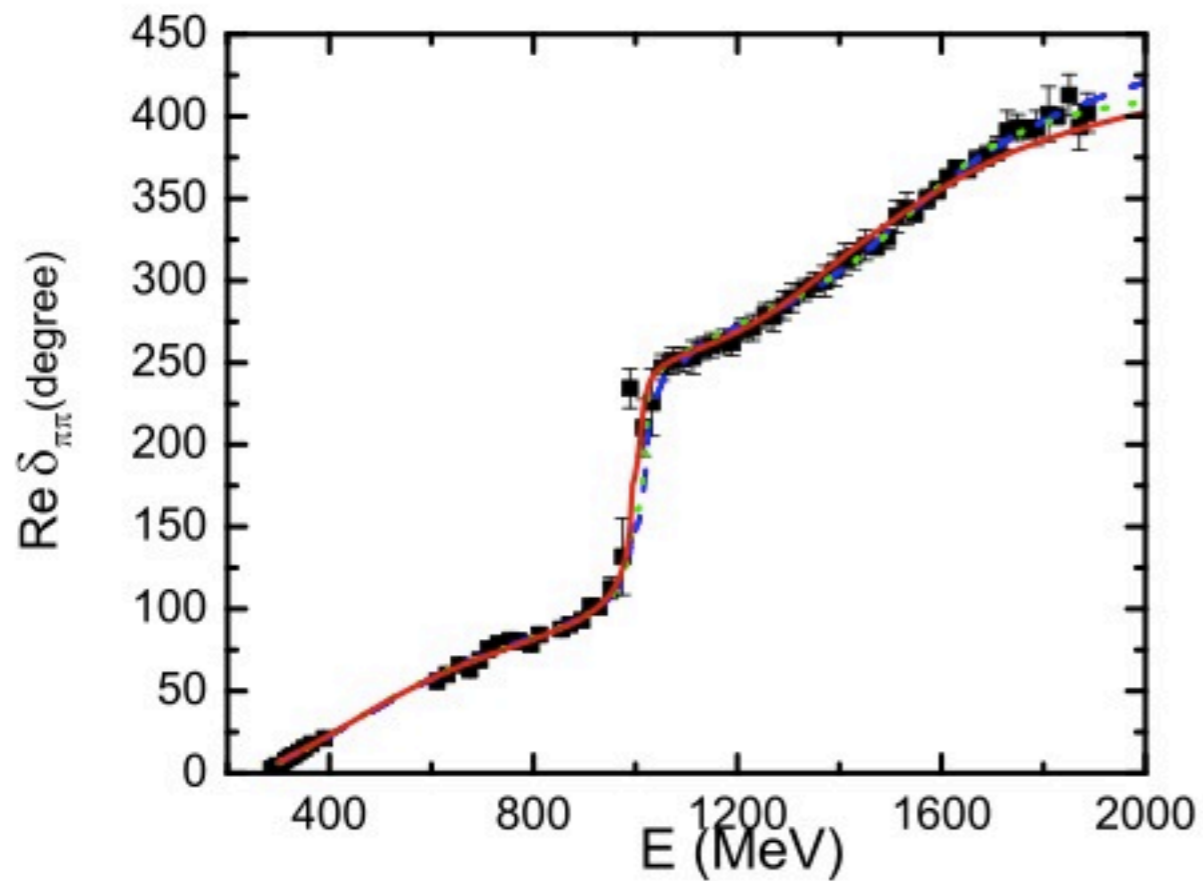


Inversion with Hamiltonian  
Just 2 box sizes (8 eigenstates)



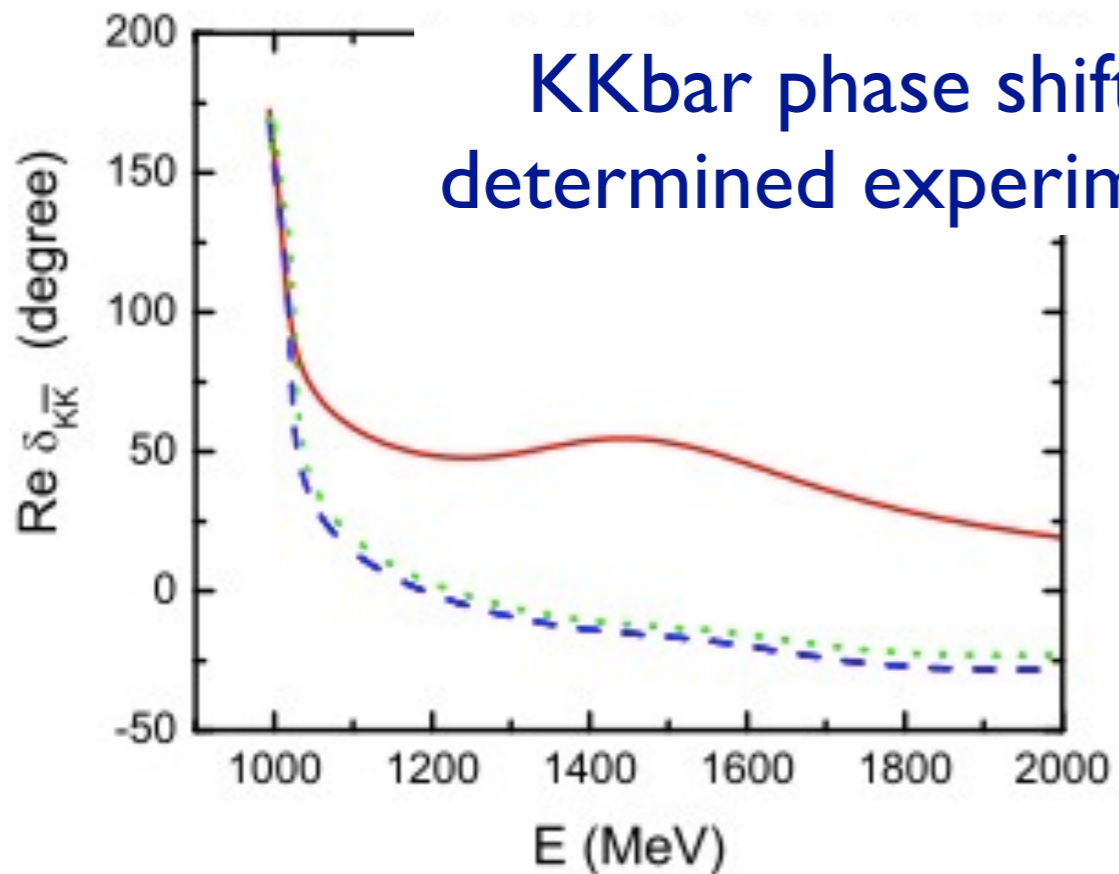
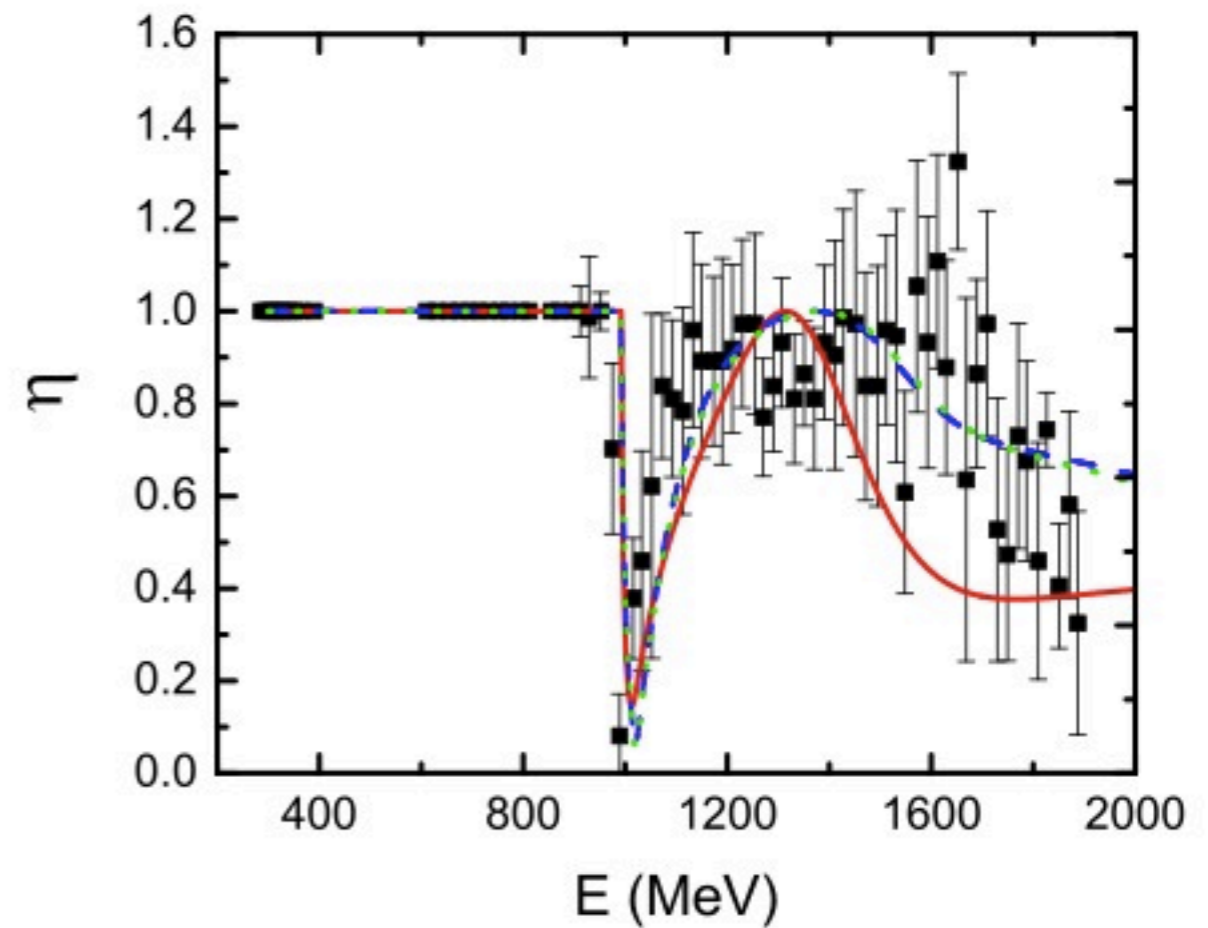
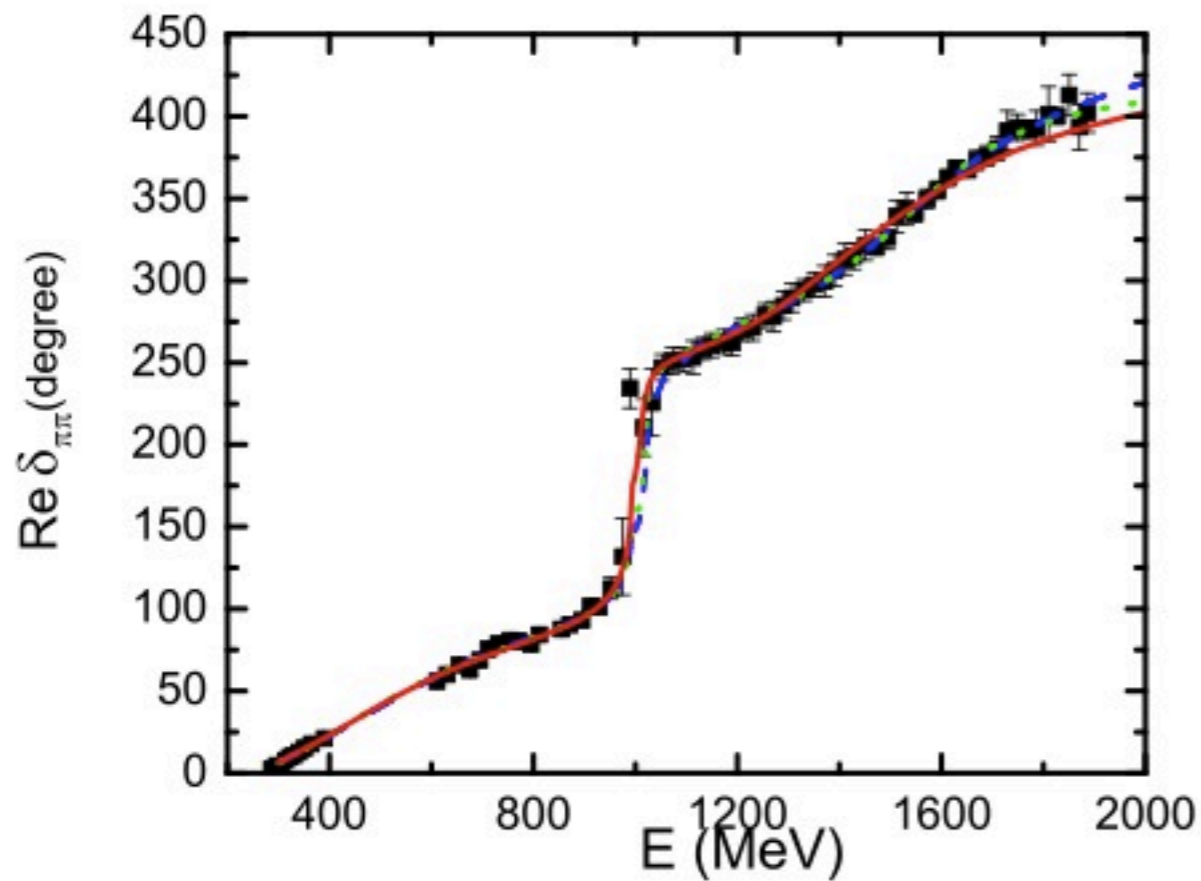


**Inversion with Hamiltonian**  
 Just 2 box sizes (12 eigenstates)



KKbar phase shift not determined experimentally

Application?



KKbar phase shift not determined experimentally

Application?

Lattice eigenstates of  $\sim 50$  MeV precision could discriminate model dependence

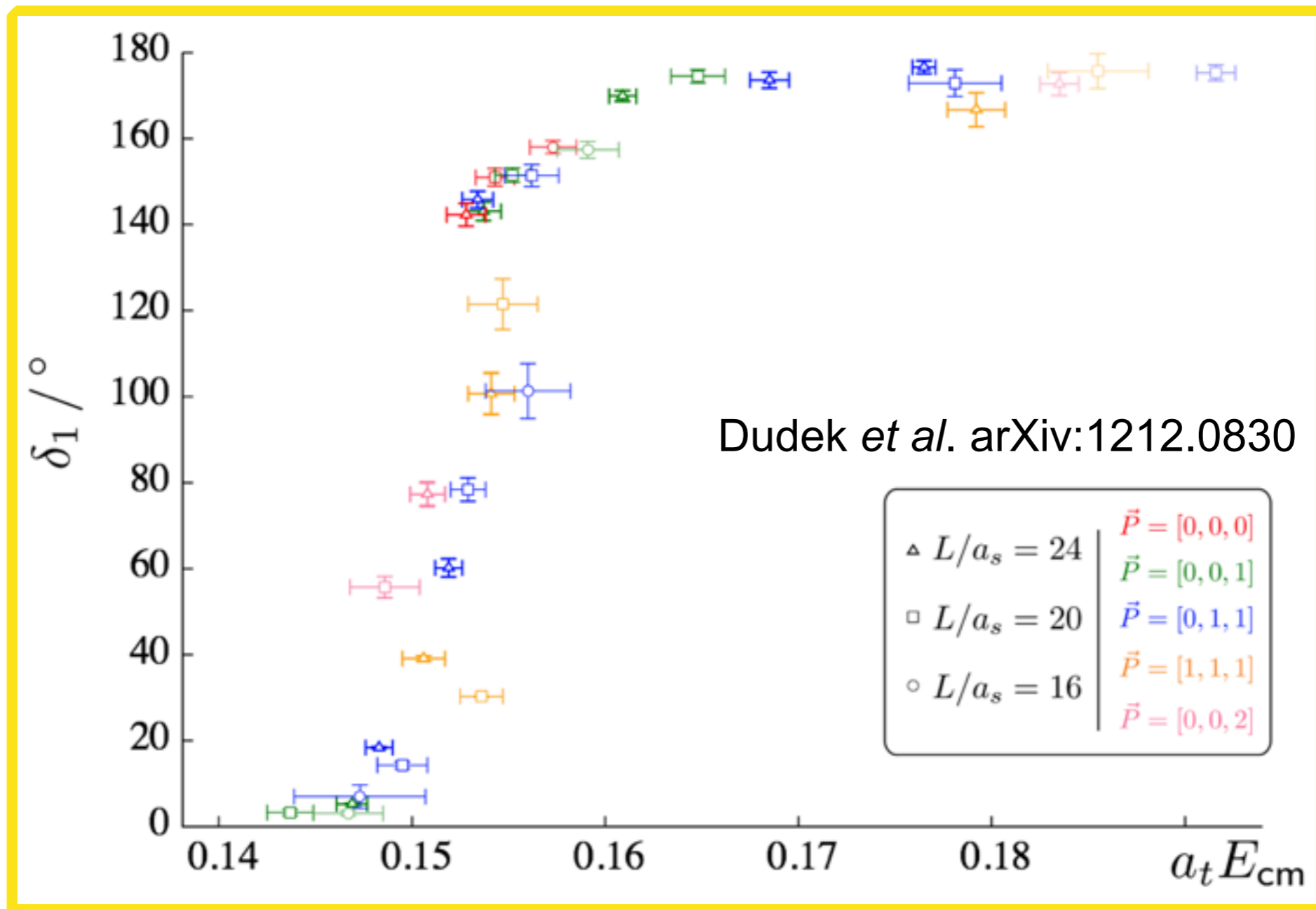
# Summary

For elastic 2-body scattering, we have a straightforward relation that maps finite-volume spectra to scattering parameters (Lüscher and Hamiltonian equivalent)

For coupled-channel systems, Hamiltonian approach appears to give a much more accessible determination of S-matrix

Hamiltonian formulation offers a natural method to extend to more complicated systems

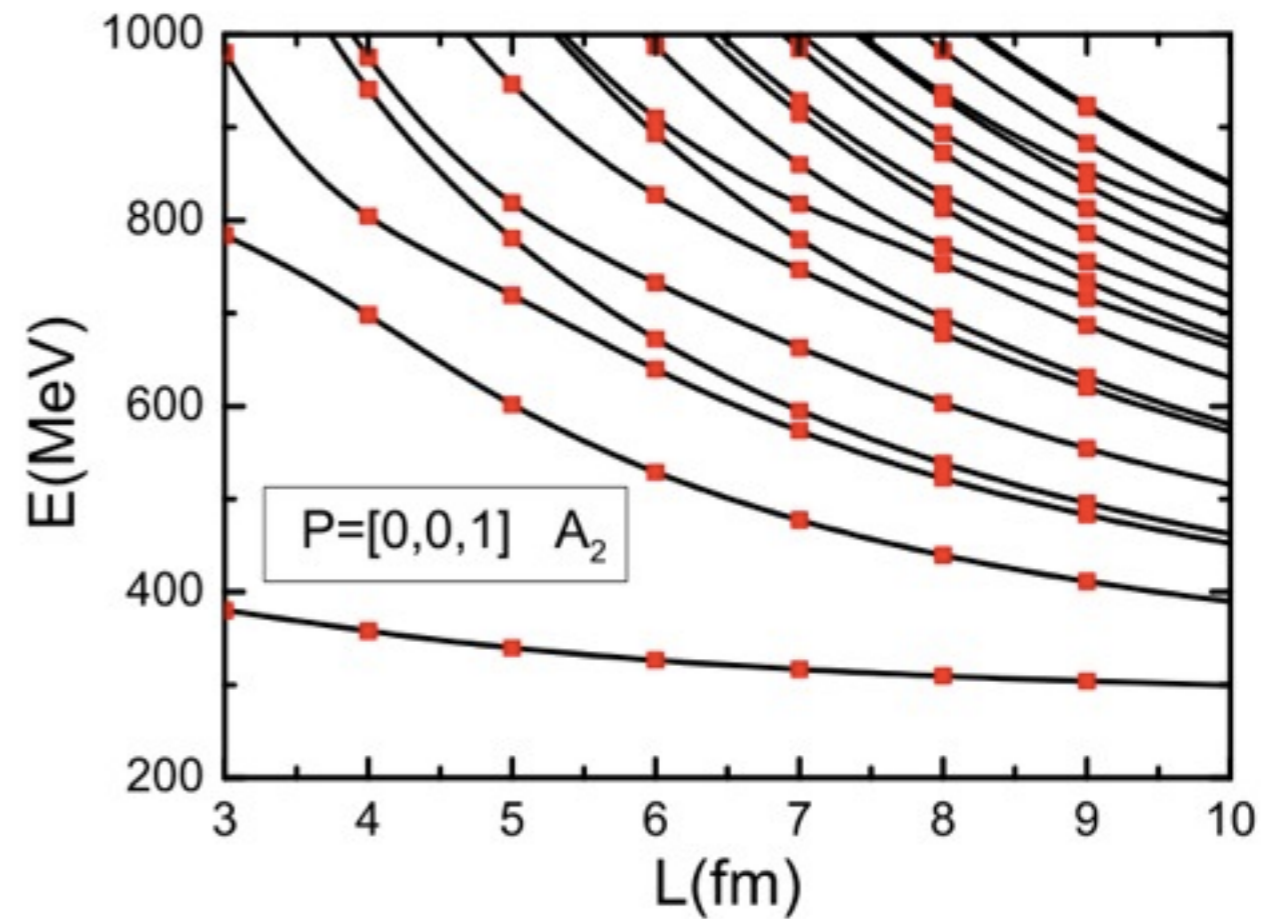
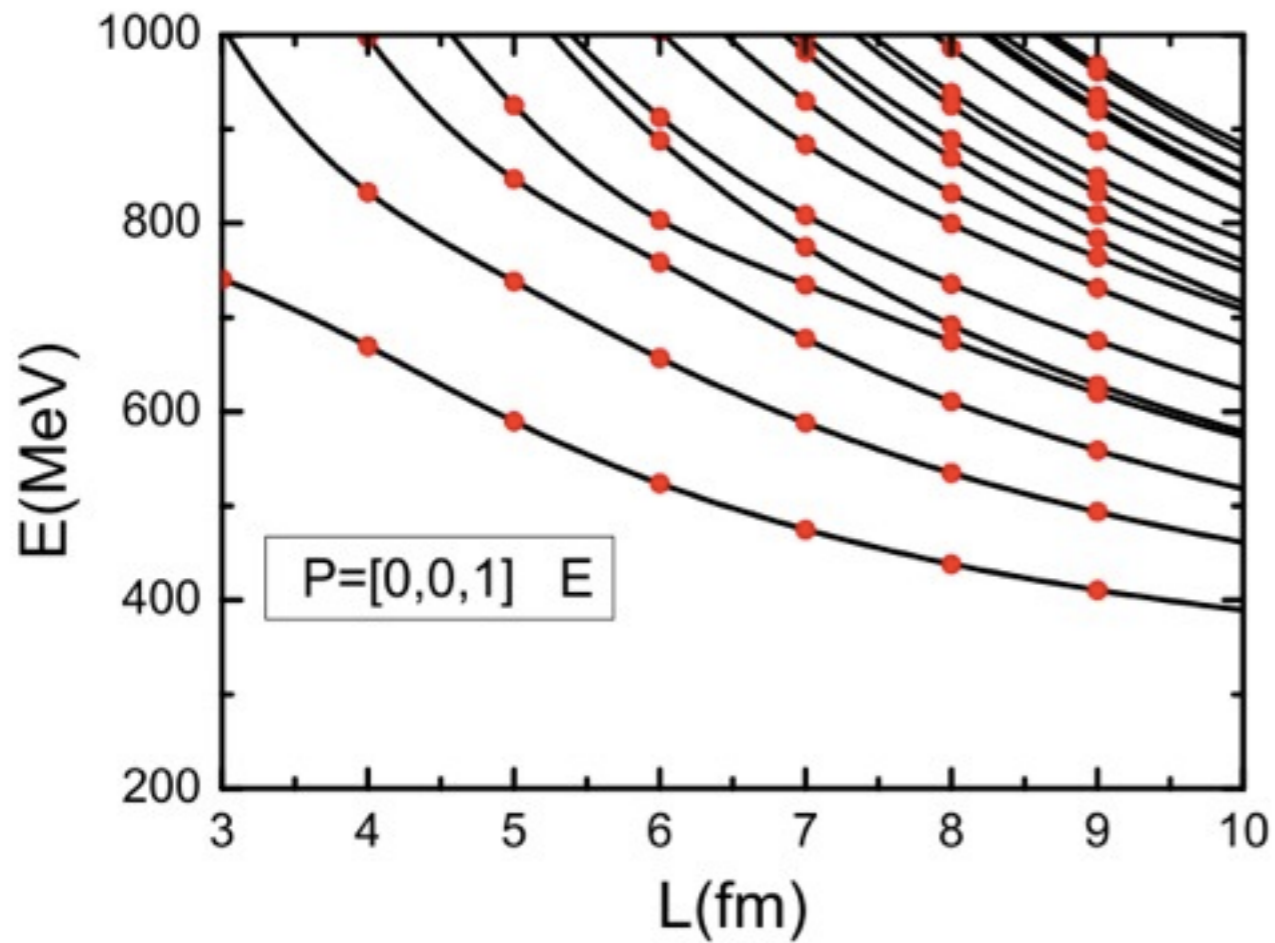
eg. many-body systems; transition matrix elements



## Boosted momentum frame

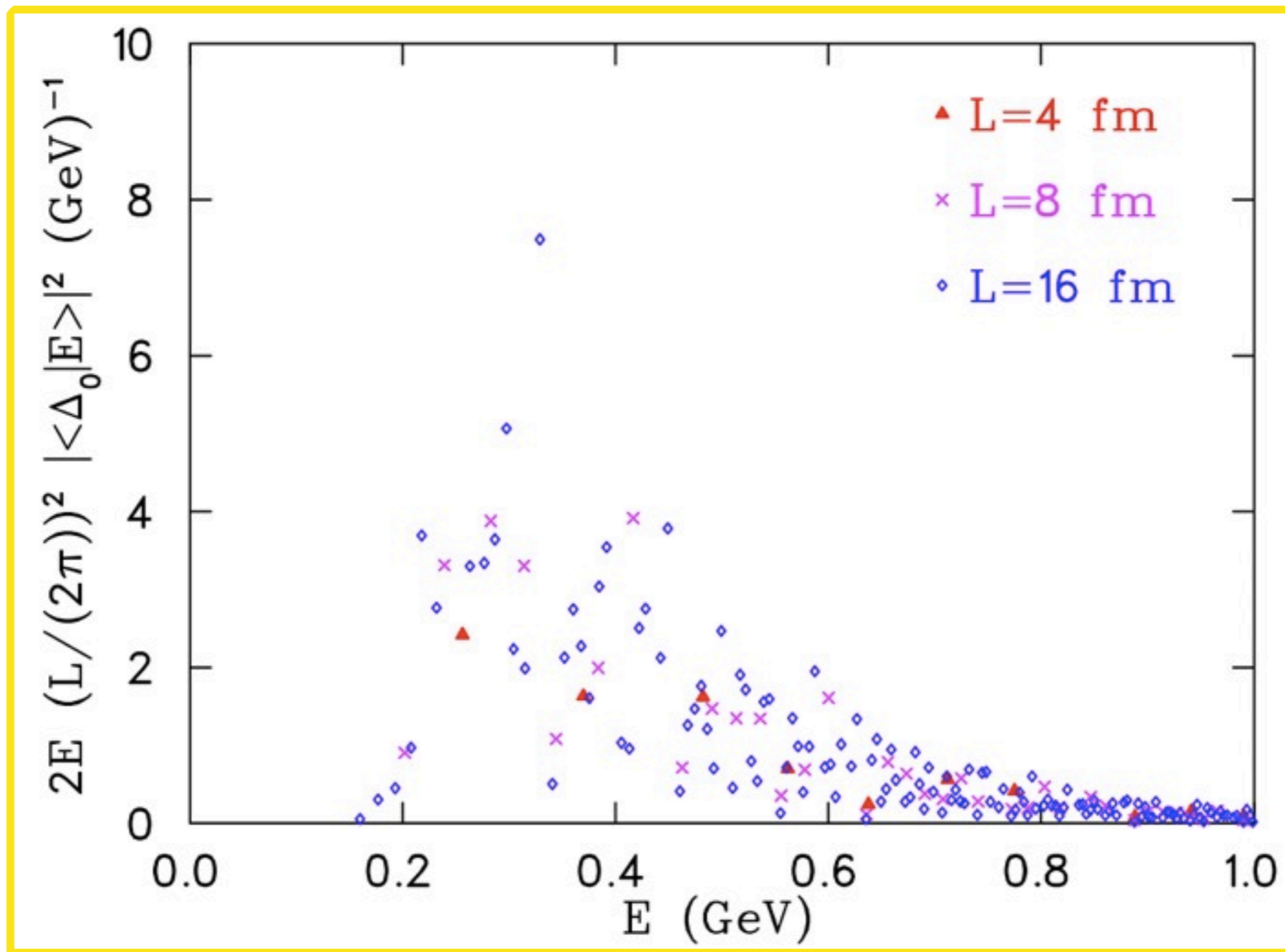
Using boosted frames, can extract the scattering phases at many more eigenenergies (momenta)





## Boosted momentum frame

Using boosted frames, can extract the scattering phases at many more eigenenergies (momenta)



## Eigenvectors

Strength of coupling of bare Delta to energy eigenstates.  
Clearly no distinction between “local” and scattering state (within width of resonance)