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**NUCLEONIC SYSTEMS
AND IN-MEDIUM NUCLEONS**

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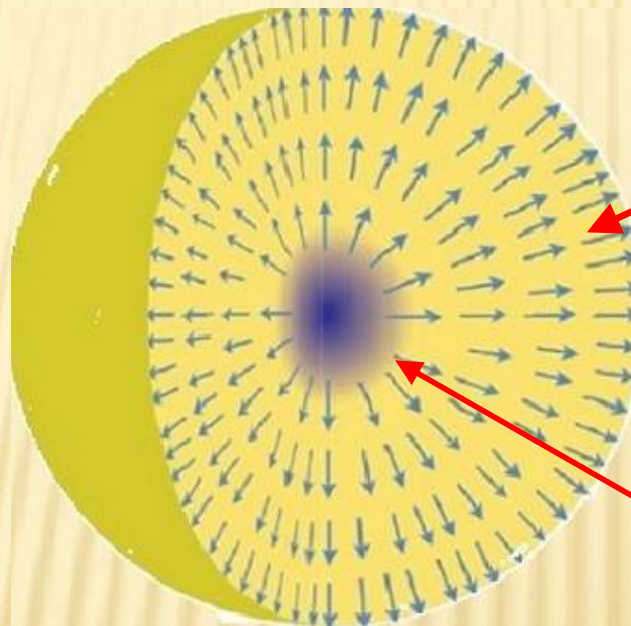
CONTENT

- ❑ Topological models and soliton
- ❑ Medium modifications
 - ❑ “Outer shell” modifications
 - ❑ “Inner core” modifications
- ❑ Nuclear matter
 - ❑ Symmetric matter
 - ❑ Asymmetric matter
- ❑ Hadron properties in nuclear matter
- ❑ Finite nuclei
- ❑ Summary
- ❑ Outlook

TOPOLOGICAL MODELS AND SOLITON

Structure

- ❑ What is a nucleon and, in particular, its core?
- ❑ Core treatment depends from the energy scale
- ❑ At large number of colors it still has the mesonic content



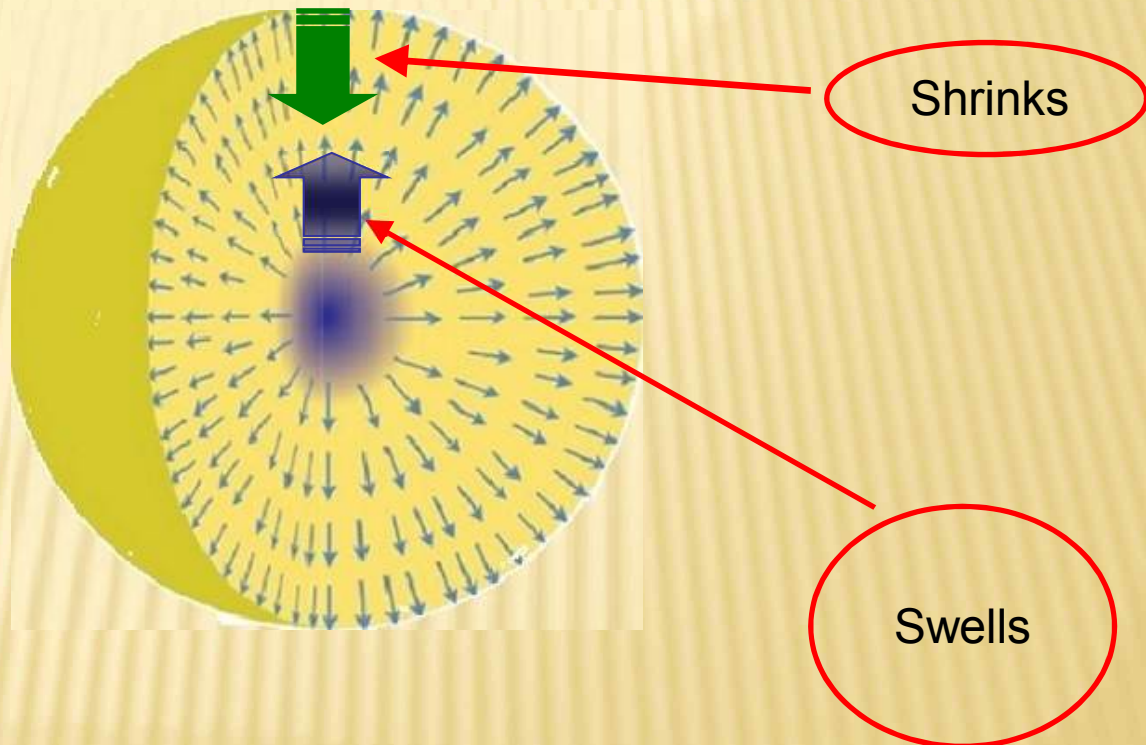
Meson cloud

Core..
Made from
what?

TOPOLOGICAL MODELS AND SOLITON

Stabilization

- ❑ Soliton has a finite size and finite energy
- ❑ One needs at least two contrrterms in the effective Lagrangian



TOPOLOGICAL MODELS AND SOLITON

Skyrme model

[T.H.R. Skyrme, Proc.Roy.Soc.Lond. A260(1961)]

- Nonlinear chiral effective meson (pionic) theory)

$$\mathcal{L} = \frac{F_\pi^2}{16} \text{Tr}(\partial_\alpha U)(\partial^\alpha U^\dagger) + \frac{1}{32e^2} \text{Tr}[U^\dagger \partial_\alpha U, U^\dagger \partial_\beta U]^2$$

Shrinks

Swells



- Hedgehog soliton (nontrivial mapping)

$$U = \exp\left\{\frac{i\bar{\tau}(\bar{\pi})}{2F_\pi}\right\} = \exp\{i\bar{\tau}(\bar{n}F(r))\}$$

TOPOLOGICAL MODELS AND SOLITON

Original Lagrangian in use

[G.S. Adkins *et al.* Nucl. Phys. B228 (1983)]

$$\mathcal{L}_{\text{free}} = \frac{F_\pi^2}{16} \text{Tr}(\partial^\alpha U)(\partial_\alpha U^+) + \frac{1}{32e^2} \text{Tr}[U^+ \partial_\alpha U, U^+ \partial_\beta U]^2 + \frac{F_\pi^2 m_\pi^2}{16} \text{Tr}(U + U^+ - 2)$$

Nontrivial mapping

- It has topologically nontrivial solitonic solutions (separated in the different topological sectors) with the corresponding conserved topological number A
- Nucleon is quantized state of the classical soliton-skyrmion

$$U = \exp\{i\bar{\tau} \bar{\pi} / 2F_\pi\} = \exp\{i\bar{\tau} \bar{n} F(r)\}$$

$$B^\mu = \frac{1}{24\pi^2} \varepsilon^{\mu\nu\alpha\beta} \text{Tr}(L_\nu L_\alpha L_\beta) \quad L_\alpha = U^+ \partial_\alpha U$$

$$A = \int d^3r B^0$$

$$H = M_{cl} + \frac{\bar{S}^2}{2I} = M_{cl} + \frac{\bar{T}^2}{2I},$$

$$|S = T, s, t \rangle = (-1)^{t+T} \sqrt{2T+1} D_{-t,s}^{S=T}(A)$$

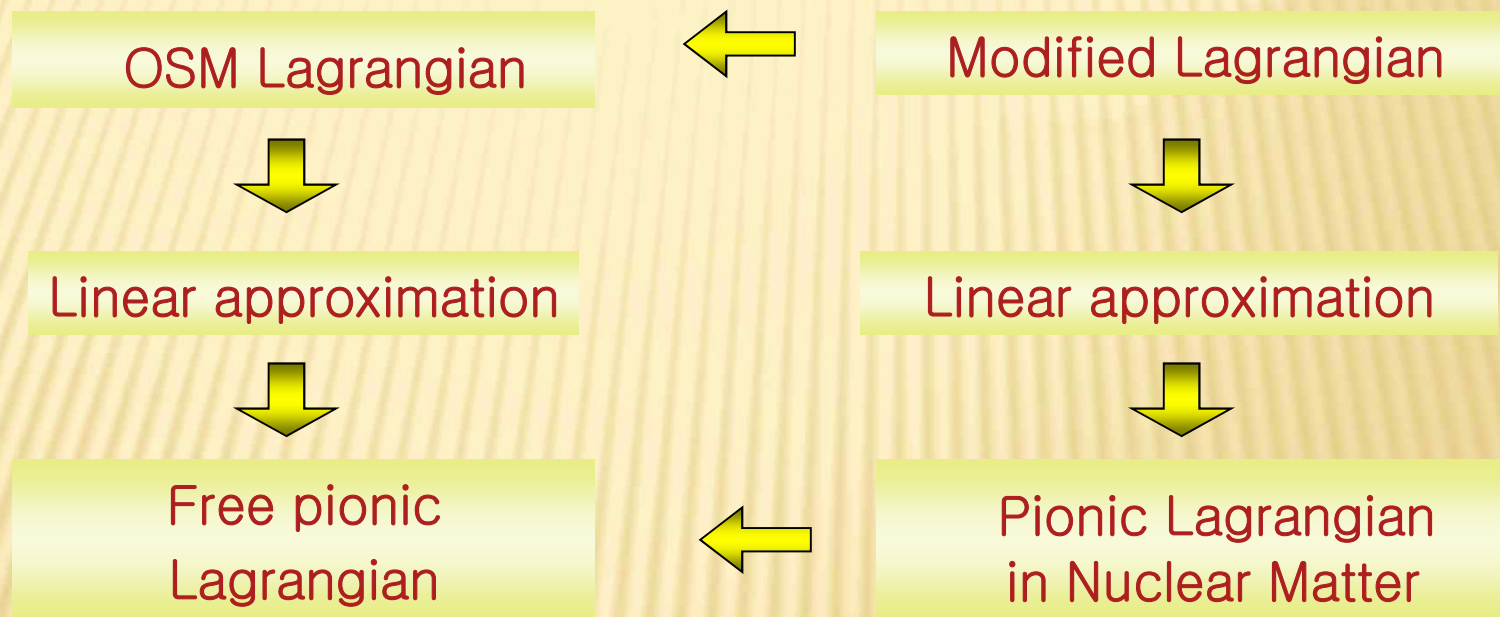
MEDIUM MODIFICATIONS

What happens in a nuclear medium?

- ❑ Medium effect
- ❑ One should be able to describe the possible phenomena
 - Deformations
 - Mass change
 - Swelling or shrinking
 - Change of NN interactions
 - Etc.

MEDIUM MODIFICATIONS

- Modification of the mesonic sector modifies the baryonic sector

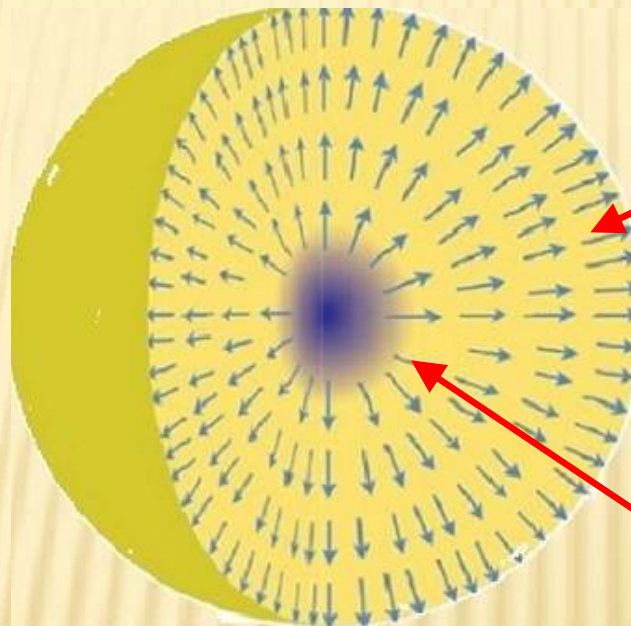


- Question arises: How to modify the mesonic sector?

MEDIUM MODIFICATIONS

Soliton in Nuclear Medium (structure changes)

- ❑ Outer shell modifications
- ❑ Inner core modifications (in particular, at higher densities)



Meson cloud modifications in nuclear medium

Inner core modifications in nuclear medium

MEDIUM MODIFICATIONS

“Outer shell” modifications

- Three types of pions can be treated separately
- In nuclear matter, one considers three types of polarization operators
- There will be some **parameters** which **correspond to the isospin breaking effects** in the surrounding environment

$$\left(\partial^\mu \partial_\mu + m_{\pi^{(\pm,0)}}^2 \right) \vec{\pi}^{(\pm,0)} = 0$$

$$\left(\partial^\mu \partial_\mu + m_{\pi^{(\pm,0)}}^2 + \hat{\Pi}^{(\pm,0)} \right) \vec{\pi}^{(\pm,0)} = 0$$

	π -atom	$T_\pi = 50$ MeV
$b_0 [m_\pi^{-1}]$	- 0.03	- 0.04
$b_1 [m_\pi^{-1}]$	- 0.09	- 0.09
$c_0 [m_\pi^{-3}]$	0.23	0.25
$c_1 [m_\pi^{-3}]$	0.15	0.16
g'	0.47	0.47

MEDIUM MODIFICATIONS

“Outer shell” modifications [U.Meissner *et al.*, EPJ A36 (2008)]

$$\mathcal{L}_2^* = -\frac{F_\pi^2}{16} \left\{ \alpha_s^{02} \text{Tr}(\partial_0 U \partial_0 U^+) - \alpha_p^0 \text{Tr}(\partial_i U \partial_i U^+) \right\}$$

$$\mathcal{L}_{\chi SB}^* = \frac{F_\pi^2 m_\pi^2}{8} \alpha_s^{00} \text{Tr}(U - 1)$$

$$\hat{\Pi} = 2\omega U_{opt} = \chi_s + \vec{\nabla} \cdot \chi_p \vec{\nabla}$$

- Due to the nonlocality of optic potential the kinetic term is also modified
- Due to energy and momentum dependence of the optic potential parameters following parts of the kinetic term is modified in different form:
 - Temporal part
 - Space part

	π -atom	$T_\pi = 50$ MeV
$b_0 [m_\pi^{-1}]$	- 0.03	- 0.04
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MEDIUM MODIFICATIONS

“Inner core” modifications

[UY & H.Ch. Kim, PRC83 (2011); UY, JKPS62 (2013); UY, PRC88 (2013)]

- May be related to
 - vector meson properties in nuclear matter
 - nuclear matter properties

$$\mathcal{L}_4^* = -\frac{1}{16e_\tau^{*2}} \text{Tr}[L_0, L_i]^2 + \frac{1}{32e_s^{*2}} \text{Tr}[L_i, L_j]^2$$

$$e \rightarrow e^* = e \gamma^{1/2}(\rho)$$

MEDIUM MODIFICATIONS

Final Lagrangian

[UY, JKPS62 (2013), UY, PRC88 (2013)]

- Separated into two parts

$$\mathcal{L}^* = \mathcal{L}_{\text{sym}}^* + \mathcal{L}_{\text{asym}}^*$$

- Isoscalar part

$$\mathcal{L}_{\text{sym}}^* = \mathcal{L}_2^* + \mathcal{L}_4^* + \mathcal{L}_{\chi SB}^*$$

- Isovector part

$$\mathcal{L}_{\text{asym}}^* = \Delta\mathcal{L}_{\text{mes}}^* + \Delta\mathcal{L}_{\text{env}}^*$$

- Nuclear matter stabilization
- Asymmetric matter properties

$$\mathcal{L}_2^* = -\frac{F_\pi^2}{16} \left\{ \alpha_s^{02} \text{Tr}(\partial_0 U \partial_0 U^+) - \alpha_p^0 \text{Tr}(\partial_i U \partial_i U^+) \right\}$$

$$\mathcal{L}_4^* = -\frac{1}{16e_\tau^{*2}} \text{Tr}[L_0, L_i]^2 + \frac{1}{32e_s^{*2}} \text{Tr}[L_i, L_j]^2$$

$$\mathcal{L}_{\chi SB}^* = \frac{F_\pi^2 m_\pi^2}{8} \alpha_s^{00} \text{Tr}(U - 1)$$

$$\Delta\mathcal{L}_{\text{mes}}^* = -\frac{F_\pi^2}{32} \sum_{a=1}^2 (m_{\pi_\pm}^2 - m_{\pi_0}^2) \text{Tr}(\tau_a U) \text{Tr}(\tau_a U^+)$$

$$\Delta\mathcal{L}_{\text{env}}^* = -\frac{F_\pi^2}{16} m_\pi \alpha_e \varepsilon_{ab3} \text{Tr}(\tau_a U) \text{Tr}(\tau_b \partial_0 U^+)$$

MEDIUM MODIFICATIONS

Reparametrization [UY, PRC88 (2013)]

- Five medium parameters

$$\begin{array}{ll}
 F_{\pi,\tau} \rightarrow F_{\pi,\tau}^* & e_{\tau} \rightarrow e_{\tau}^*, \quad m_{\pi} \rightarrow m_{\pi}^*, \\
 F_{\pi,s} \rightarrow F_{\pi,s}^* & e_s \rightarrow e_s^* \\
 \text{Shell} & \text{Core} \\
 \text{modifications} & \text{modifications}
 \end{array}$$

- Rearranging

$$\begin{aligned}
 1 + C_1 \frac{\rho}{\rho_0} &= f_1 \left(\frac{\rho}{\rho_0} \right) \equiv \sqrt{\frac{\alpha_p^0}{\gamma_s}} \\
 1 + C_2 \frac{\rho}{\rho_0} &= f_2 \left(\frac{\rho}{\rho_0} \right) \equiv \frac{\alpha_s^{00}}{(\alpha_p^0)^2 \gamma_s} & \frac{\alpha_e}{\gamma_s} &= f_4 \left(\frac{\rho}{\rho_0} \right) \frac{\rho_n - \rho_p}{\rho_0} = \frac{C_4 \frac{\rho}{\rho_0}}{1 + C_5 \frac{\rho}{\rho_0}} \frac{\rho_n - \rho_p}{\rho_0} \\
 1 + C_3 \frac{\rho}{\rho_0} &= f_3 \left(\frac{\rho}{\rho_0} \right) \equiv \frac{(\alpha_p^0 \gamma_s)^{3/2}}{\alpha_s^{02}}
 \end{aligned}$$

MEDIUM MODIFICATIONS

Nucleon in nuclear matter

- Isoscalar mass

$$m_{N,s}^* = M_S^* + \frac{3}{8\Lambda^*} + \frac{\Lambda^*}{2} \left(a^{*2} + \frac{\Lambda_{\text{env}}^{*2}}{\Lambda^{*2}} \right)$$

- Isovector mass

$$\Delta m_{np}^* = a^* + \frac{\Lambda_{\text{env}}^*}{\Lambda^*}$$

- Mass of the nucleon

$$m_{n,p}^* = m_{N,s}^* - \Delta m_{np}^* T_3$$

NUCLEAR MATTER

The binding-energy-formula terms in the present model

$$\varepsilon(A, Z) = -a_V + a_S \frac{(N - Z)^2}{A^2} + \dots$$

We are ready
to reproduce

- Volume term
 - Infinite and asymmetric nuclear matter
- Asymmetry term
 - Isospin asymmetric environment
- Surface and Coulomb terms
 - Nucleons in a finite volume
- Finite nuclei properties
 - Local density approximation

NUCLEAR MATTER

Volume term and Symmetry energy

- At infinite nuclear matter approximation the binding energy per nucleon takes the form

$$\varepsilon(\lambda, \delta) = \varepsilon_V(\lambda) + \varepsilon_S \delta^2 + O(\delta^4) \equiv \varepsilon_V(\lambda) + \varepsilon_A(\lambda, \delta)$$

- λ is normalized nuclear matter density
 - δ is asymmetry parameter
 - ε_S is symmetry energy
- In our model
 - Symmetric matter $\varepsilon_V(\lambda) = m_{N,s}^*(\lambda, 0) - m_N^{\text{free}}$
 - Asymmetric matter

$$\begin{aligned} \varepsilon_A(\lambda, \delta) &= \varepsilon(\lambda, \delta) - \varepsilon_V(\lambda) \\ &= m_{N,s}^*(\lambda, \delta) - m_{N,s}^*(\lambda, 0) + m_{N,V}^*(\lambda, \delta)\delta \end{aligned}$$

NUCLEAR MATTER

Nuclear matter properties

- Symmetric matter properties (pressure, compressibility and third derivative)

$$p = \rho_0 \lambda^2 \left. \frac{\partial \varepsilon_V(\lambda)}{\partial \lambda} \right|_{\lambda=1}, \quad K_0 = 9 \rho^2 \left. \frac{\partial^2 \varepsilon_V(\lambda)}{\partial \rho^2} \right|_{\rho=\rho_0}, \quad Q = 27 \lambda^3 \left. \frac{\partial^3 \varepsilon_V(\lambda)}{\partial \lambda^3} \right|_{\lambda=1}$$

- Symmetry energy properties (coefficient, slop and curvature)

$$\varepsilon_s(\lambda) = \varepsilon_s(1) + \frac{L_s}{3}(\lambda - 1) + \frac{K_s}{18}(\lambda - 1)^2 + \dots$$

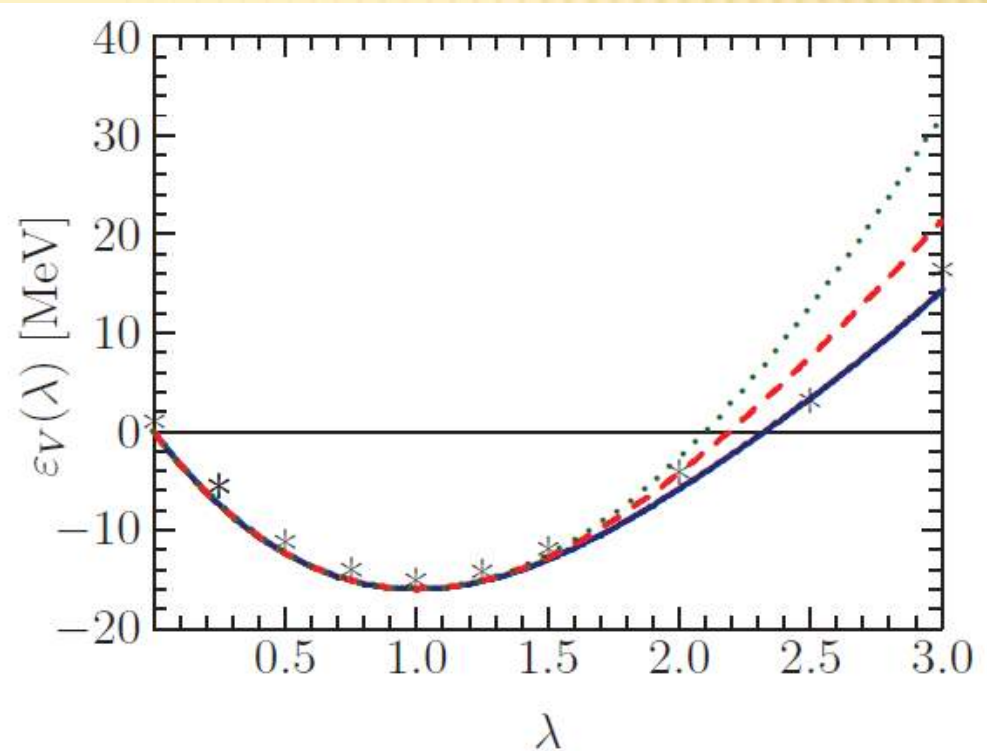
SYMMETRIC MATTER

Volume energy [UY, PRC88 (2013)]

- Set I – solid
- Set II – dashed
- Set III – dotted

For comparison: Akmal-Pandharipande-Ravenhall (APR) predictions [PRC 58, 1804 (1998)] are given by stars.

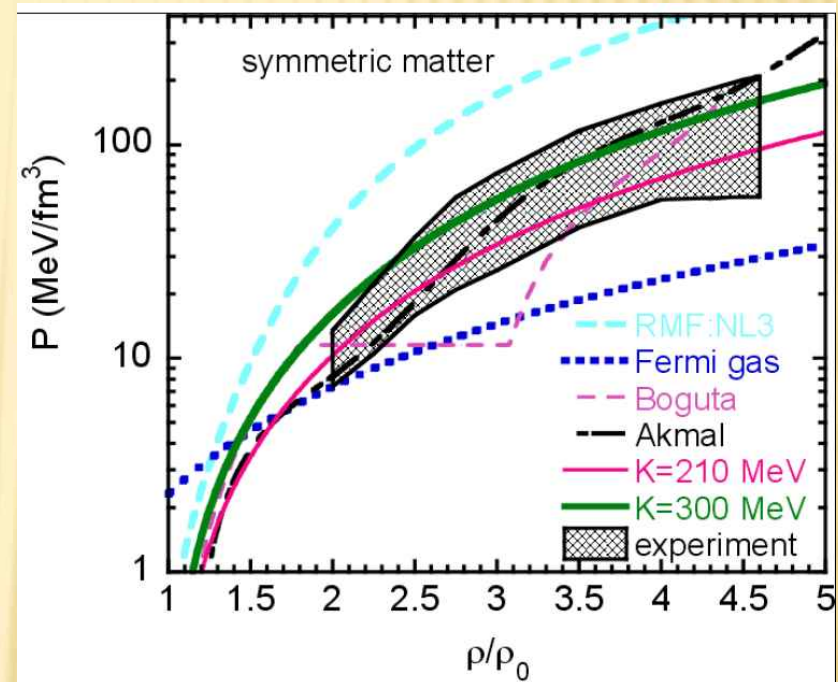
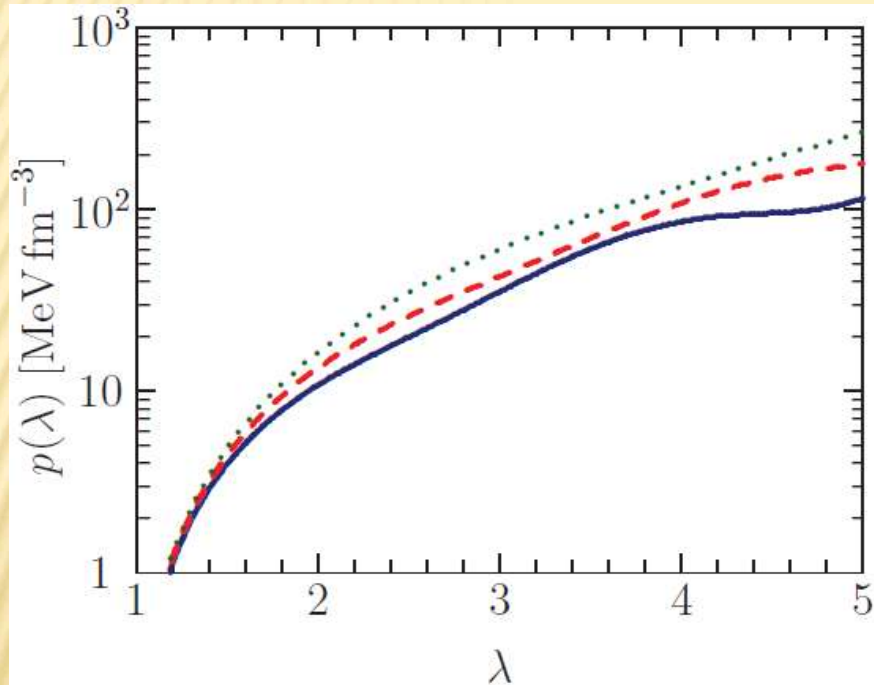
(From arigonna 2 body interactions + 3 body interactions)



Set	C_1	C_2	C_3	$\epsilon_V(\rho_0)$ (MeV)	K_0 (MeV)	Q (MeV)
I	-0.279	0.737	1.782	-16	240	-410
II	-0.273	0.643	1.858	-16	250	-279
III	-0.277	0.486	2.124	-16	260	-178

SYMMETRIC MATTER

Pressure [UY, PRC88 (2013)]



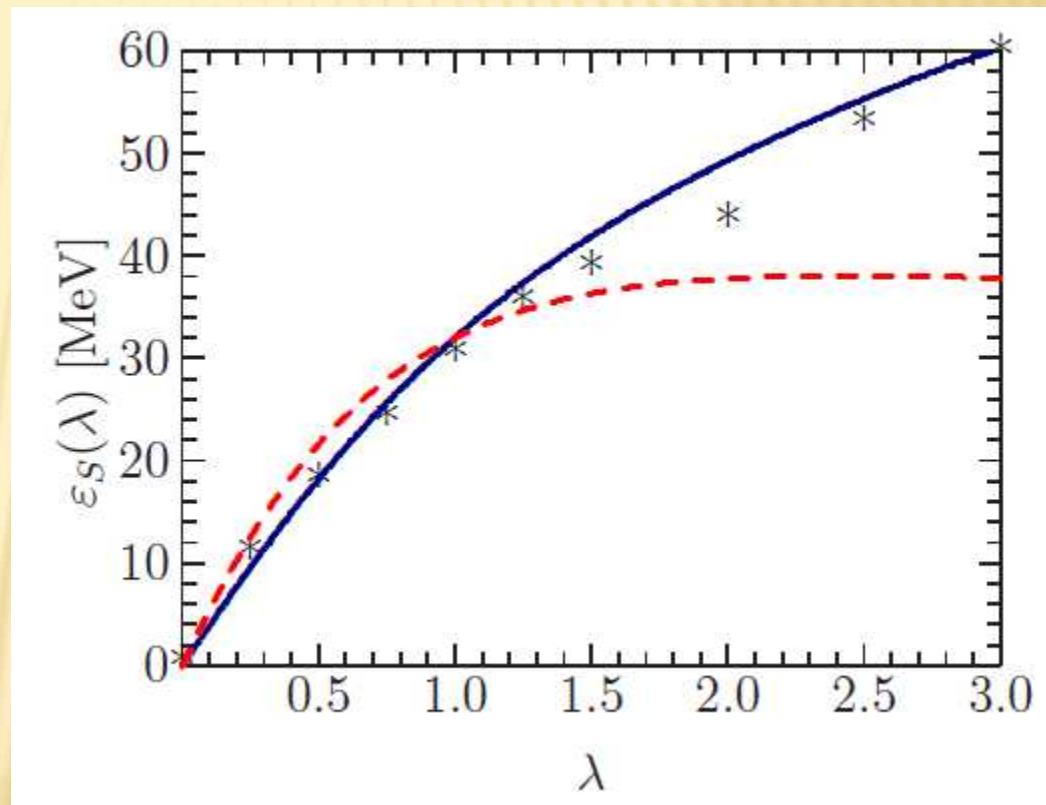
For comparison: Right figure from
Danielewicz- Lacey-Lynch, Science 298, 1592 (2002).
(Deduced from experimental flow data and simulations studies)

ASYMMETRIC MATTER

Symmetry energy

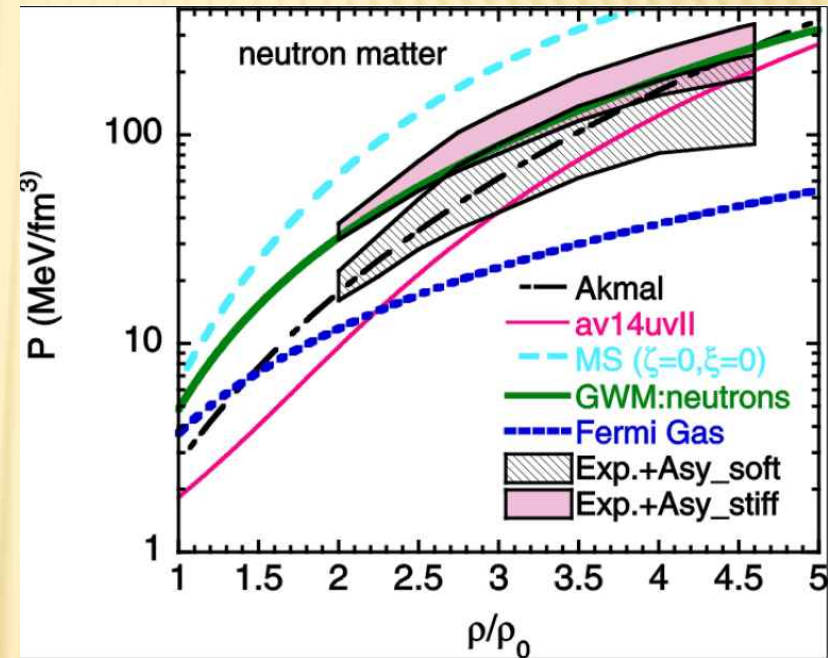
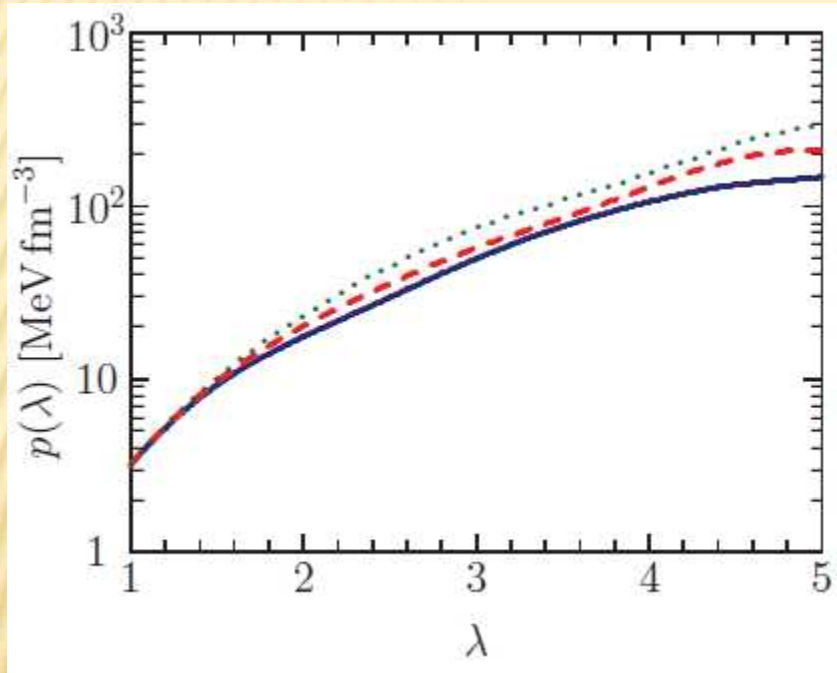
- Solid $L_s = 70 \text{ MeV}$
- Dashed $L_s = 40 \text{ MeV}$

For comparison: Akmal-Pandharipande-Ravenhall (APR) predictions [PRC 58, 1804 (1998)] are given by stars.
(From arigonna 2 body interactions + 3 body interactions)



ASYMMETRIC MATTER

Pressure in neutron matter [UY, PRC88 (2013)]



For comparison: Right figure from Danielewicz- Lacey-Lynch, Science 298, 1592 (2002).
(Deduced from experimental flow data and simulations studies)

ASYMMETRIC MATTER

Low density behavior of symmetry energy

For comparison:

Trippa-Colo-Vigezzi

[PRC 77, 061304 (2008)];

From analysis of GDR
(208Pb).

$$23.3 < \varepsilon_s(\rho = 0.1\text{fm}^{-3}) < 24.9 \text{ MeV}$$

$\varepsilon_s(\rho_0)$ [MeV]	L_S [MeV]	K_S [MeV]	K_τ [MeV]	$K_{0,2}$ [MeV]	$\varepsilon_s(0.1\text{fm}^{-3})$ [MeV]
32	40	-181	-301	-257	25.15
32	50	-160	-310	-254	24.15
32	60	-126	-306	-239	23.22
32	70	-80	-290	-211	22.37
32	80	-21	-261	-172	21.57
32	90	50	-220	-119	20.82
32	100	134	-166	-55	20.13

Consequently one can
predict in this model:

$$K_\tau = K_S - 6L_S$$

$$K_{0,2} = K_\tau - \frac{Q}{K_0} L_S$$

NUCLEAR MATTER

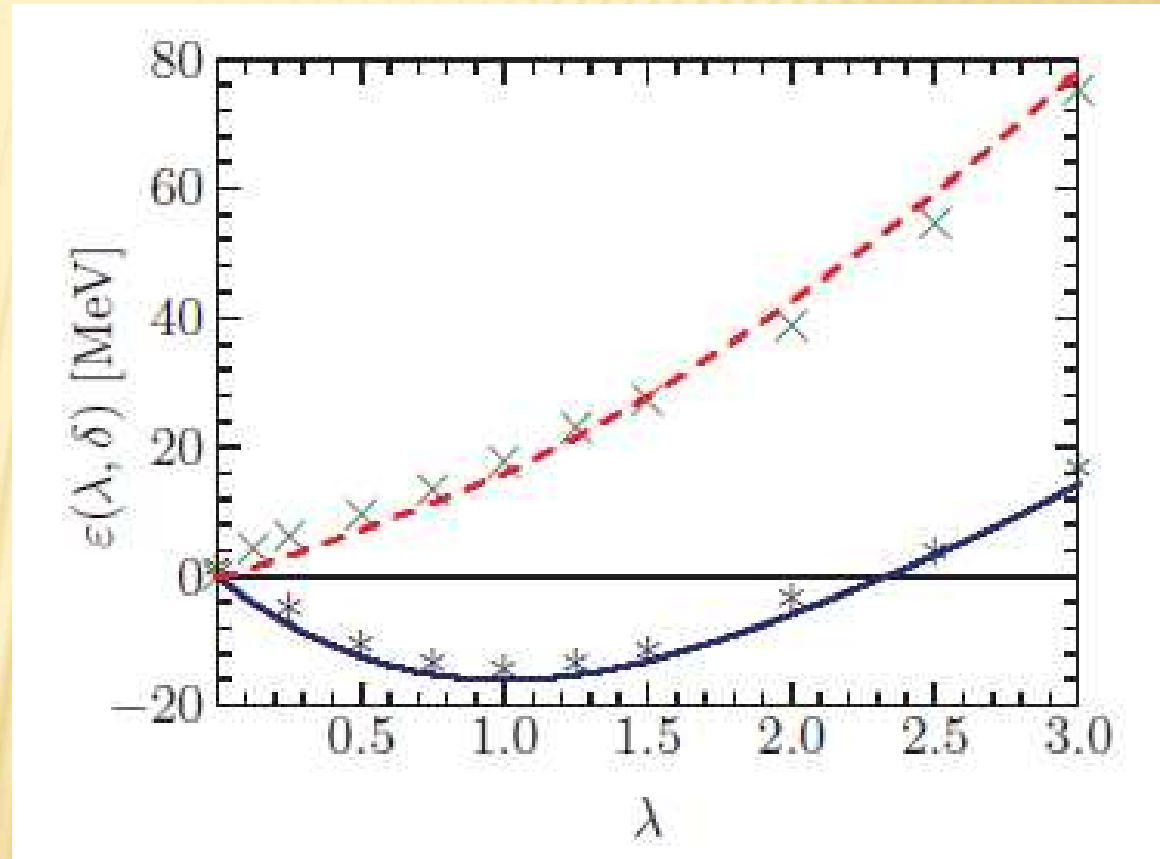
For completeness [UY, PRC88 (2013)]

- Symmetric matter – solid
- Neutron matter –dashed

For comparison:

APR predictions

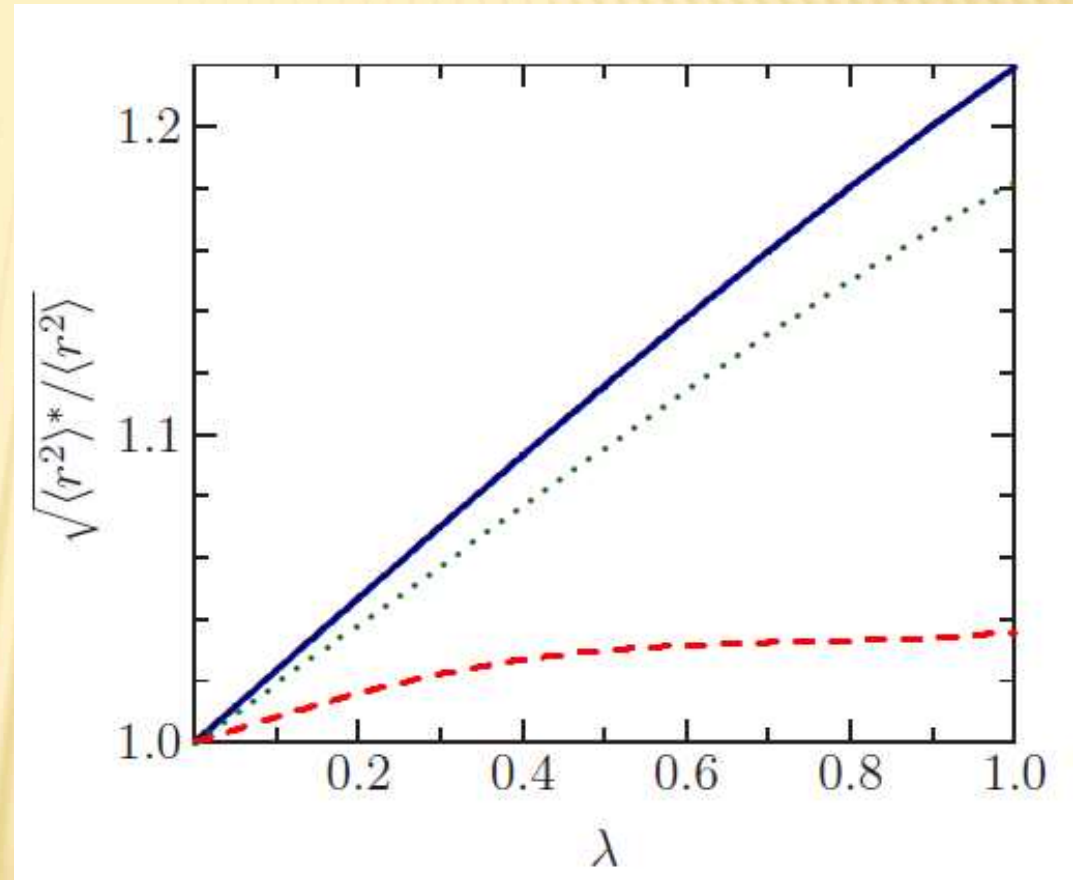
[PRC 58, 1804 (1998)] are given by stars.



HADRON PROPERTIES IN NUCLEAR MATTER

Mean square radii of nucleons

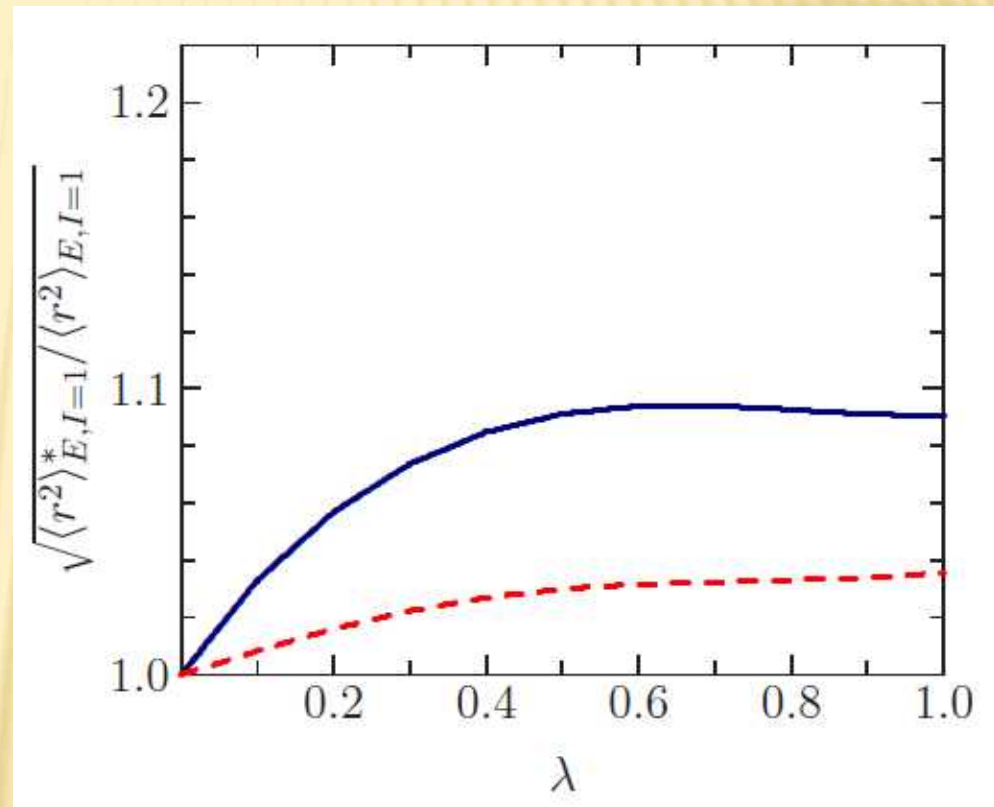
- ❑ Electric-isoscalar – solid curve
- ❑ Electric-isovector – dashed curve
- ❑ Magnetic-isoscalar – dashed curve
- ❑ Magnetic-isovector – dotted curve



HADRON PROPERTIES IN NUCLEAR MATTER

Mean square radii of nucleons

- Neutron matter – solid curve
- Symmetric matter – dashed curve



HADRON PROPERTIES IN NUCLEAR MATTER

Low energy constants in nuclear matter at ρ_0

$$F_{\pi,\tau} \rightarrow F_{\pi,\tau}^*, \quad F_{\pi,S} \rightarrow F_{\pi,S}^*$$

	Present model	ChPT [1]	QCD sum rules [2]
$F_{\pi,t}^* / F_{\pi}$	0.37	0.74	0.79
$F_{\pi,S}^* / F_{\pi}$	0.72	< 0	0.78

[1] U. Meissner, J. Oller, A. Wirzba, Annals Phys. 297 (2002) 27.

[2] H. Kim, M. Oka, NPA720 (2003) 368.

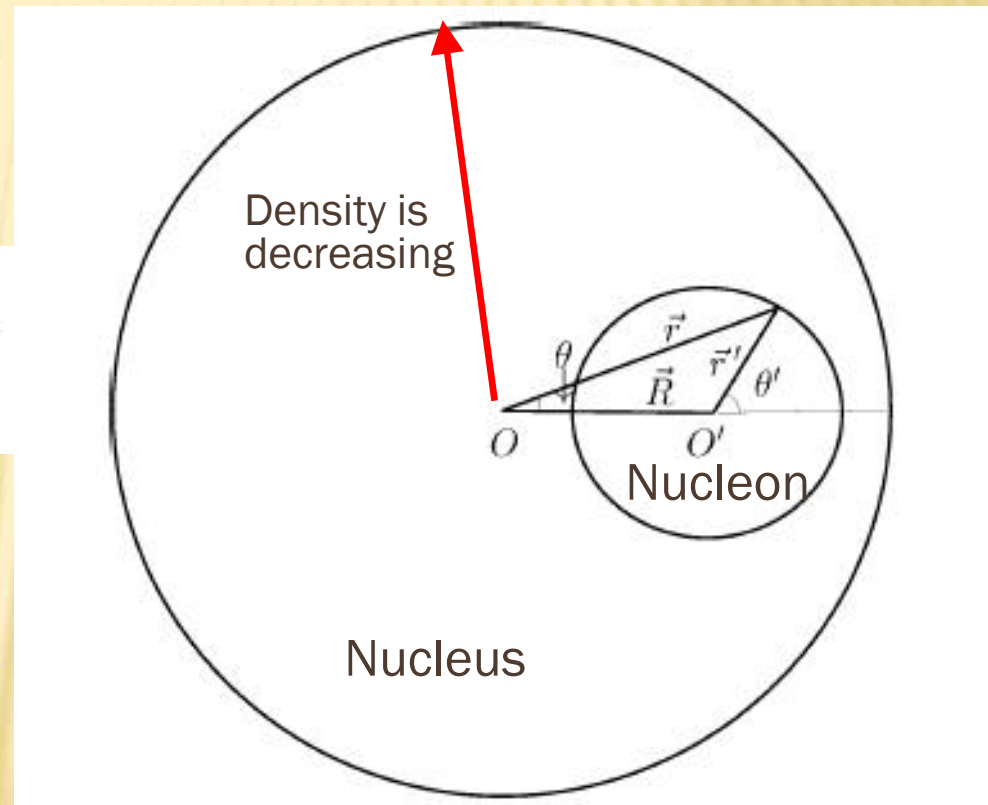
FINITE NUCLEI

Nucleon in a nucleus (local density approximation)

- The problem is coupled partial differential equations

$$f(F_{\bar{r}\bar{r}}, F_{\theta\theta}, F_{\bar{r}}, F_{\theta}, \Theta_{\theta}, F, \Theta) = 0,$$

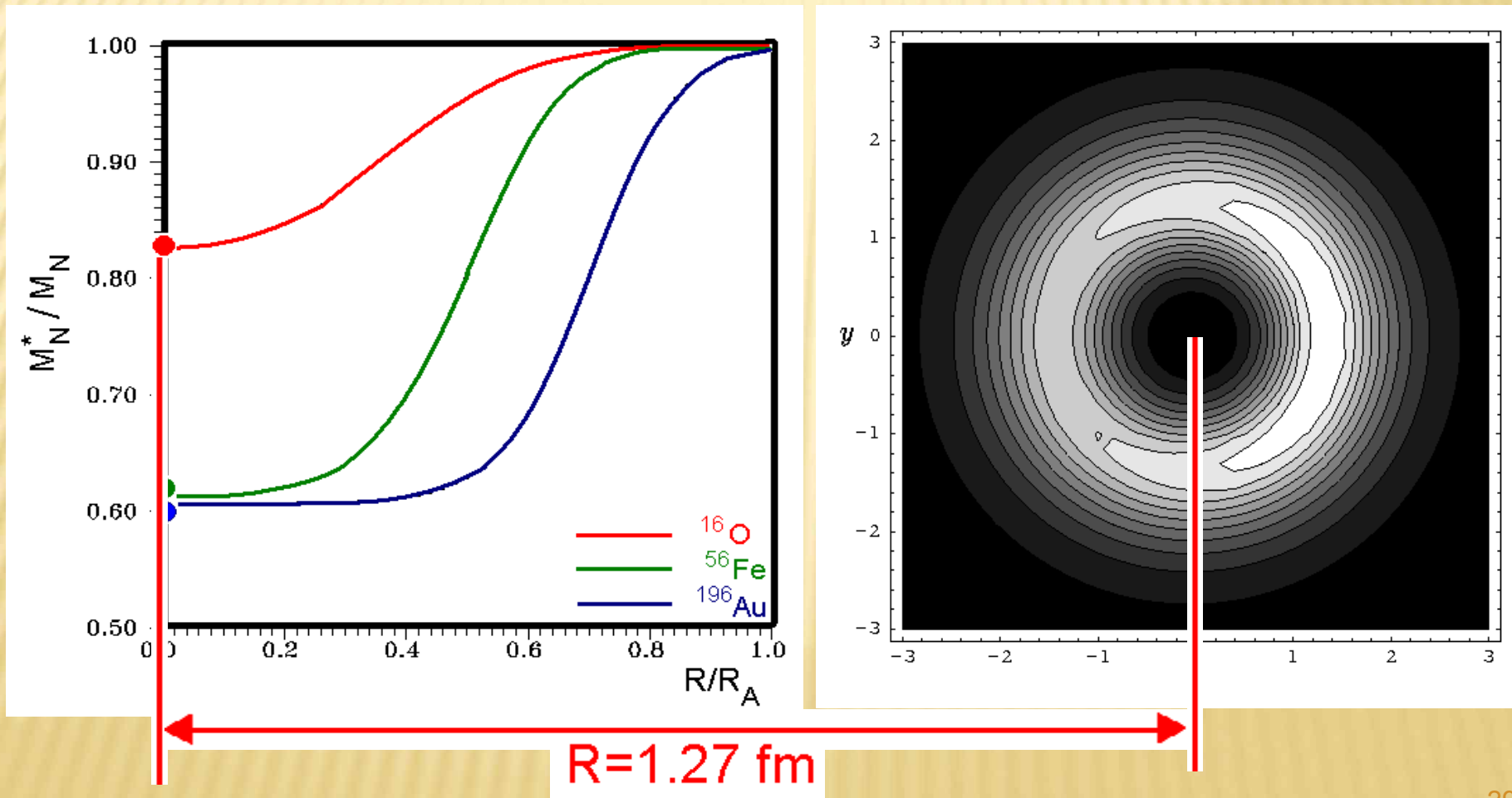
$$g(\Theta_{\theta\theta}, \Theta_{\theta}, F_{\bar{r}}, F_{\theta}, \Theta, F) = 0,$$



FINITE NUCLEI

Nucleon may deform [UY, *et al.*, NPA200 (2002) 403]

- The core modifications are not taken into account



SUMMARY

- ❑ Within the applicability range, the model describes
 - ❑ the single hadrons properties
 - in separate state
 - in the community of their partners
 - ❑ as well as the properties of that whole community at same footing

OUTLOOK

Extensions and applicability of the approach

- ❑ Nucleon tomography in nuclear matter
[H.Ch. Kim, UY, PLB726 (2013), arXiv:1304.5926]
- ❑ NN interactions in nuclear matter
- ❑ Neutron stars
- ❑ Finite nuclei properties
 - ❑ Mirror nuclei
 - ❑ Exotic nuclei
 - ❑ Halo nuclei
- ❑ Nucleon-knock out reactions
- ❑ Vector mesons in nuclear matter
[J.H.Jung, UY, H.Ch.Kim, PLB 723 (2013), arXiv:1212.4616]

Thank you for your attention!