Manuel E. Carrillo-Serrano

Collaborators: A. W. Thomas (ADL) I. C. Cloët (ANL)

April 8th 2014







 ρ^+ features

NJL Model

Constituent Quark Form Factors

 ρ^+ Form Factors

Results

Conclusions

◆□▶ ◆□▶ ◆三▶ ◆三▶ ◆□▶

Outline

ρ^+ features

NJL Model

Constituent Quark Form Factors

 ρ^+ Form Factors

Results

Conclusions

- イロト 4 団 ト 4 臣 ト 4 臣 - シッペー



 Up and down quark as constituents

< ロ > < 団 > < 三 > < 三 > < 三 > < 三 > < ○ < ○</p>



 Up and down quark as constituents



► Aligned quark spins → spin 1.



 Up and down quark as constituents



- ► Aligned quark spins → spin 1.
- ► Vector-isovector field.



 Up and down quark as constituents



► Three possible isospin states:

1.
$$\rho^+ = u\bar{d}$$

2. $\rho^- = \bar{u}d$
3. $\rho^0 = \frac{\left[u\bar{u} - d\bar{d}\right]}{\sqrt{2}}$

- ► Aligned quark spins → spin 1.
- ► Vector-isovector field.



 Up and down quark as constituents



- ► Aligned quark spins → spin 1.
- ► Vector-isovector field.

► Three possible isospin states:

1.
$$\rho^+ = u\bar{d}$$

2. $\rho^- = \bar{u}d$
3. $\rho^0 = \frac{\left[u\bar{u} - d\bar{d}\right]}{\sqrt{2}}$

- Charged ρ production:
 - ▶ e^+e^- anihilation
 - ▶ *τ*-lepton decays

200



 Up and down quark as constituents



- ► Aligned quark spins → spin 1.
- ► Vector-isovector field.

► Three possible isospin states:

1.
$$\rho^+ = u\bar{d}$$

2. $\rho^- = \bar{u}d$
3. $\rho^0 = \frac{\left[u\bar{u} - d\bar{d}\right]}{\sqrt{2}}$

- Charged ρ production:
 - e^+e^- anihilation
 - ▶ *τ*-lepton decays
- Resonances
 - 1. $\Gamma = 149.1 \pm 0.8$ MeV. $\sim \tau = 10^{-24}$ s.
 - 2. $M = 775.11 \pm 0.34$ MeV.

ペロ> ペロ> ペヨ> ペヨ> ペロ>

Outline

ρ^+ features

NJL Model

Constituent Quark Form Factors

 ρ^+ Form Factors

Results

Conclusions



- The Nambu-Jona-Lasinio (NJL) model simpler than full QCD.
- Keeps QCD symmetries $\rightarrow \chi$ sym. breaking $\rightarrow m_q$.

¹Rev. Mod. Phys., Vol. 64, No. 3, 1992.

ρ^+ features NJL Model Cons	stituent Quark Form Factors $ ho^+$	Form Factors	Results	Conclusions
NI II 1				

- The Nambu-Jona-Lasinio (NJL) model simpler than full QCD.
- Keeps QCD symmetries $\rightarrow \chi$ sym. breaking $\rightarrow m_q$.
- Initially N-N interaction $\rightarrow q\bar{q}$ interaction.

¹Rev. Mod. Phys., Vol. 64, No. 3, 1992.

- イロト (四) (三) (三) (三) (二)



NJL¹

- The Nambu-Jona-Lasinio (NJL) model simpler than full QCD.
- ► Keeps QCD symmetries $\rightarrow \chi$ sym. breaking $\rightarrow m_q$.
- Initially N-N interaction $\rightarrow q\bar{q}$ interaction.
- ▶ Main feature \rightarrow hadronic masses $\rightarrow (M_q) \rightarrow$ Gap Eq.



¹Rev. Mod. Phys., Vol. 64, No. 3, 1992.

6/23

- ▲ロ > ▲ 国 > ▲ 国 > ▲ 国 > 今 Q Q

The NJL Lagrangian $SU(2)^2$

Original SU(2) NJL Lagrangian

$$\mathcal{L}_{NJL} = G\left[\left(\bar{\psi} \psi \right)^2 + \left(\bar{\psi} i \gamma_5 \vec{\tau} \psi \right)^2 \right]$$
(1)

²Nucl. Phys. A. 587 (1995) 617-656

- ▲日 > ▲ 画 > ▲ 画 > ▲ 画 > ▲ の へ ()

The NJL Lagrangian $SU(2)^2$

Original SU(2) NJL Lagrangian

$$\mathcal{L}_{NJL} = G\left[\left(\bar{\psi}\psi \right)^2 + \left(\bar{\psi}i\gamma_5 \vec{\tau}\psi \right)^2 \right]$$
(1)

▶ In our model just SU(2) (Mesons) \rightarrow Fierz Transformations

$$\mathcal{L}_{NJL} = \frac{1}{2} G_{\pi} \left[\left(\bar{\psi} \psi \right)^2 - \left(\bar{\psi} \gamma_5 \vec{\tau} \psi \right)^2 \right] - \frac{1}{2} G_{\omega} \left(\bar{\psi} \gamma^{\mu} \psi \right)^2 - \frac{1}{2} G_{\rho} \left[\left(\bar{\psi} \gamma^{\mu} \vec{\tau} \psi \right)^2 + \left(\bar{\psi} \gamma^{\mu} \gamma_5 \vec{\tau} \psi \right)^2 \right]$$
(2)

► Feynman Rules.

²Nucl. Phys. A. 587 (1995) 617-656

- Point-like $\bar{q}q$ interaction \rightarrow must be regularized (Λ_{UV}) .
- ► Use proper time method.

- Point-like $\bar{q}q$ interaction \rightarrow must be regularized (Λ_{UV}) .
- ► Use proper time method.

$$\frac{1}{X^{n}} = \frac{1}{(n-1)} \int_{0}^{\infty} d\tau \tau^{n-1} e^{-\tau X}$$
$$\to \frac{1}{(n-1)} \int_{1/\Lambda_{UV}^{2}}^{1/\Lambda_{IR}^{2}} d\tau \tau^{n-1} e^{-\tau X}$$
(3)

 $\blacktriangleright \ X \to Feynman \ parametrization$

- Point-like $\bar{q}q$ interaction \rightarrow must be regularized (Λ_{UV}) .
- ► Use proper time method.

$$\frac{1}{X^{n}} = \frac{1}{(n-1)} \int_{0}^{\infty} d\tau \tau^{n-1} e^{-\tau X}$$
$$\to \frac{1}{(n-1)} \int_{1/\Lambda_{UV}^{2}}^{1/\Lambda_{IR}^{2}} d\tau \tau^{n-1} e^{-\tau X}$$
(3)

- $\blacktriangleright \ X \rightarrow Feynman \ parametrization$
- ► Infra-red cutoff, A_{IR}, simulates confinement (avoids hadrons decaying into free quarks).

- Point-like $\bar{q}q$ interaction \rightarrow must be regularized (Λ_{UV}) .
- ► Use proper time method.

$$\frac{1}{X^{n}} = \frac{1}{(n-1)} \int_{0}^{\infty} d\tau \tau^{n-1} e^{-\tau X}$$
$$\to \frac{1}{(n-1)} \int_{1/\Lambda_{UV}^{2}}^{1/\Lambda_{IR}^{2}} d\tau \tau^{n-1} e^{-\tau X}$$
(3)

- $\blacktriangleright \ X \to Feynman \ parametrization$
- ► Infra-red cutoff, A_{IR}, simulates confinement (avoids hadrons decaying into free quarks).
- $\Lambda_{IR} \rightarrow \text{QCD}$ scale Λ_{QCD} .

- Point-like $\bar{q}q$ interaction \rightarrow must be regularized (Λ_{UV}) .
- ► Use proper time method.

$$\frac{1}{X^{n}} = \frac{1}{(n-1)} \int_{0}^{\infty} d\tau \tau^{n-1} e^{-\tau X}$$
$$\to \frac{1}{(n-1)} \int_{1/\Lambda_{UV}^{2}}^{1/\Lambda_{IR}^{2}} d\tau \tau^{n-1} e^{-\tau X}$$
(3)

- $\blacktriangleright \ X \to Feynman \ parametrization$
- ► Infra-red cutoff, Λ_{IR} , simulates confinement (avoids hadrons decaying into free quarks).
- $\Lambda_{IR} \rightarrow \mathsf{QCD}$ scale Λ_{QCD} .
- $\Lambda_{UV} \rightarrow \text{pion decay constant } f_{\pi}$.

The Bethe Salpeter Equations for mesons in NJL

$$\mathcal{T}(p) = \int \frac{d^4k}{(2\pi)^4} \mathcal{K}S(p+k) S(k) \mathcal{T}(p)$$



 ρ^+ features

The Bethe Salpeter Equations for mesons in NJL

$$\mathcal{T}(p) = \int \frac{d^4k}{(2\pi)^4} \mathcal{K}S(p+k) S(k)\mathcal{T}(p)$$

► Solution depends on polarization bubble $(\Pi_{\rho(\pi)})$ on the diagram.



The Bethe Salpeter Equations for mesons in NJL

$$\mathcal{T}(p) = \int \frac{d^4k}{(2\pi)^4} \mathcal{K}S(p+k) S(k)\mathcal{T}(p)$$

- ► Solution depends on polarization bubble $(\Pi_{\rho(\pi)})$ on the diagram.
- ▶ Vertices depend on the nature of the meson → NJL Lagrangian (Feynman Rules)



► Solution reduced t-matrix pion.

$$\tau_{\pi} = \frac{4iG_{\pi}}{1 + 2G_{\pi}\Pi_{\pi} \left(p^2\right)}$$

(4)

► Solution reduced t-matrix pion.

$$\tau_{\pi} = \frac{4iG_{\pi}}{1 + 2G_{\pi}\Pi_{\pi} \left(p^2\right)}$$

• Mass defined at the pole position of τ_{π} :

(4)

► Solution reduced t-matrix pion.

$$\tau_{\pi} = \frac{4iG_{\pi}}{1 + 2G_{\pi}\Pi_{\pi} \left(p^2\right)}$$
(4)

- Mass defined at the pole position of τ_{π} :
- ▶ It can be Taylor expanded around M_{π} , such that

$$\tau_{\pi} = \frac{4iG_{\pi}}{1 + 2G_{\pi}\Pi_{\pi} (p^2)} \to -\frac{ig_{\pi}}{p^2 - M_{\pi}^2 + i\epsilon}$$
(5)

► Solution reduced t-matrix pion.

$$\tau_{\pi} = \frac{4iG_{\pi}}{1 + 2G_{\pi}\Pi_{\pi} \left(p^{2}\right)}$$
(4)

- Mass defined at the pole position of τ_{π} :
- ▶ It can be Taylor expanded around M_{π} , such that

$$\tau_{\pi} = \frac{4iG_{\pi}}{1 + 2G_{\pi}\Pi_{\pi} \left(p^{2}\right)} \to -\frac{ig_{\pi}}{p^{2} - M_{\pi}^{2} + i\epsilon}$$
(5)

with

$$g_{\pi} = 2 \left(\left. \frac{\partial \Pi_{\pi}}{\partial p^2} \right|_{p^2 = M_{\pi}^2} \right)^{-1} \tag{6}$$

DQ C

• Same for $\rho \rightarrow$ include Lorentz structure.

Outline

 ρ^+ features

NJL Model

Constituent Quark Form Factors

 ρ^+ Form Factors

Results

Conclusions

• Vector coupling γ^{μ} photon to quark \rightarrow Vector meson



• Vector coupling γ^{μ} photon to quark \rightarrow Vector meson



 Two options: isoscalar (ω) or isovector (ρ) and the contribution to Quark-photon vertex

• Vector coupling γ^{μ} photon to quark \rightarrow Vector meson

 Two options: isoscalar (ω) or isovector (ρ) and the contribution to Quark-photon vertex

$$\Lambda^{\mu}_{\gamma Q}(p,p') = \frac{1}{6} \Lambda^{\mu}_{\omega}(p,p') + \frac{\tau_3}{2} \Lambda^{\mu}_{\rho}(p,p')$$
(7)

where $\Lambda^{\mu}_{\rho(\omega)}=\gamma^{\mu}F_{1\rho(\omega)}\left(q^{2}\right)+\frac{i\sigma^{\mu\nu}q_{\nu}}{2M}F_{2\rho(\omega)}\left(q^{2}\right)$

• Vector coupling γ^{μ} photon to quark \rightarrow Vector meson

 Two options: isoscalar (ω) or isovector (ρ) and the contribution to Quark-photon vertex

$$\Lambda^{\mu}_{\gamma Q}(p,p') = \frac{1}{6} \Lambda^{\mu}_{\omega}(p,p') + \frac{\tau_3}{2} \Lambda^{\mu}_{\rho}(p,p')$$
(7)

where $\Lambda^{\mu}_{\rho(\omega)}=\gamma^{\mu}F_{1\rho(\omega)}\left(q^{2}\right)+\frac{i\sigma^{\mu\nu}q_{\nu}}{2M}F_{2\rho(\omega)}\left(q^{2}\right)$

• Vector coupling γ^{μ} photon to quark \rightarrow Vector meson

 Two options: isoscalar (ω) or isovector (ρ) and the contribution to Quark-photon vertex

$$\Lambda^{\mu}_{\gamma Q}(p,p') = \frac{1}{6} \Lambda^{\mu}_{\omega}(p,p') + \frac{\tau_3}{2} \Lambda^{\mu}_{\rho}(p,p')$$
(7)

where $\Lambda^{\mu}_{\rho(\omega)} = \gamma^{\mu} F_{1\rho(\omega)} \left(q^2\right) + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M} F_{2\rho(\omega)} \left(q^2\right)$

► Both depend on $\Pi_{\rho(\omega)}$. $F_{1\rho(\omega)} = \frac{1}{1+2G_{\rho(\omega)}\Pi_{\rho(\omega)}(q^2)}$ $F_{2\rho(\omega)} = 0$

- Vector coupling γ^{μ} photon to quark \rightarrow Vector meson

 Two options: isoscalar (ω) or isovector (ρ) and the contribution to Quark-photon vertex

$$\Lambda^{\mu}_{\gamma Q}(p,p') = \frac{1}{6} \Lambda^{\mu}_{\omega}(p,p') + \frac{\tau_3}{2} \Lambda^{\mu}_{\rho}(p,p')$$
(7)

where $\Lambda^{\mu}_{\rho(\omega)} = \gamma^{\mu} F_{1\rho(\omega)} \left(q^2\right) + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M} F_{2\rho(\omega)} \left(q^2\right)$

- Both depend on $\Pi_{\rho(\omega)}$. $F_{1\rho(\omega)} = \frac{1}{1+2G_{\rho(\omega)}\Pi_{\rho(\omega)}(q^2)}$ $F_{2\rho(\omega)} = 0$
- Use pole approximation: $F_{1\rho(\omega)} = \frac{1}{1+Q^2/m_{\rho(\omega)}^2}$.

Pion Cloud

► Solving following diagrams.



 Includes: Z (quark wave function renormalization), π FF. and reduced pion t-matrix.

Pion Cloud

Solving following diagrams.



- Includes: Z (quark wave function renormalization), π FF. and reduced pion t-matrix.
- Dressed Quark sector Form Factors given by

$$F_{1U(D)} = Z \left[\frac{1}{6} F_{1\omega} + (-) \frac{1}{2} F_{1\rho} \right] + \left[F_{1\omega} - (+) F_{1\rho} \right] f_1^q + (-) F_{1\rho} f_1^\pi$$
(8)

$$F_{2U(\mathbf{D})} = [F_{1\omega} - (+)F_{1\rho}] f_2^{(q)} + (-)F_{1\rho} f_2^{\pi}$$
(9)

Outline

 ρ^+ features

NJL Model

Constituent Quark Form Factors

ρ^+ Form Factors

Results

Conclusions

< ロ > < 団 > < 三 > < 三 > < 三 > < 三 > < ○ < ○</p>





► Three Form Factors:

$$j_{\rho}^{\mu,\alpha\beta}\left(p,p'\right) = \left[g^{\alpha\beta}F_{1\rho}\left(Q^{2}\right) - \frac{q^{\alpha}q^{\beta}}{2M_{\rho}^{2}}F_{2\rho}\left(Q^{2}\right)\right]\left(p'+p\right)^{\mu} - \left(q^{\alpha}g^{\mu\beta} - q^{\beta}g^{\mu\alpha}\right)F_{3\rho}\left(Q^{2}\right)$$
(10)

with



► Three Form Factors:

$$j_{\rho}^{\mu,\alpha\beta}\left(p,p'\right) = \left[g^{\alpha\beta}F_{1\rho}\left(Q^{2}\right) - \frac{q^{\alpha}q^{\beta}}{2M_{\rho}^{2}}F_{2\rho}\left(Q^{2}\right)\right]\left(p'+p\right)^{\mu} - \left(q^{\alpha}g^{\mu\beta} - q^{\beta}g^{\mu\alpha}\right)F_{3\rho}\left(Q^{2}\right)$$
(10)

with

$$F_{i\rho}(Q^{2}) = [F_{1U}(Q^{2}) - F_{1D}(Q^{2})] f_{i}^{V}(Q^{2}) + [F_{2U}(Q^{2}) - F_{2D}(Q^{2})] f_{i}^{T}(Q^{2}).$$
(11)

i = 1, 2, 3.

► Sachs-like charge, magnetic and quadrupole Form Factors

► Sachs-like charge, magnetic and quadrupole Form Factors

$$G_C(Q^2) = F_1(Q^2) + \frac{2}{3}\eta G_Q(Q^2), \qquad (12)$$

$$G_M\left(Q^2\right) = F_3\left(Q^2\right),\tag{13}$$

$$G_Q(Q^2) = F_1(Q^2) + (1+\eta)F_2(Q^2) - F_3(Q^2).$$
(14)

$$\eta = rac{Q^2}{4M_
ho^2}$$

► Sachs-like charge, magnetic and quadrupole Form Factors

$$G_C(Q^2) = F_1(Q^2) + \frac{2}{3}\eta G_Q(Q^2), \qquad (12)$$

$$G_M\left(Q^2\right) = F_3\left(Q^2\right),\tag{13}$$

$$G_Q(Q^2) = F_1(Q^2) + (1+\eta)F_2(Q^2) - F_3(Q^2).$$
(14)

$$\eta = \frac{Q^2}{4M_{\rho}^2}$$

 \blacktriangleright Charged squared radius $\left< r^2 \right>$ and magnetic moment μ_ρ

$$\left\langle r^{2}\right\rangle = -6 \left. \frac{\partial G_{C}\left(Q^{2}\right)}{\partial Q^{2}} \right|_{Q^{2}=0}$$
(15)

$$\mu_{\rho}\left(\mu_{N}\right) = \left[\frac{m_{N}}{m_{\rho}\left(m_{\pi}\right)}G_{M}\left(0,m_{\pi}\right)\right]\mu_{N}$$
(16)

◆□▶ ◆□▶ ◆三▶ ◆三▶ ◆□▶

Outline

 ρ^+ features

NJL Model

Constituent Quark Form Factors

 ρ^+ Form Factors

Results

Conclusions

- イロト 4 団 ト 4 臣 ト 4 臣 - シッペー

Form Factor Results

 ρ^+ Form Factors Vs Q^2



QCDSF PoS LAT2008 (2008) 051 [arXiv:hep-lat/0611029v1]



nan

 $\langle r^2 \rangle$ Lattice at $Q^2 = 0.22$ Mev (CSSM): Phys. Rev. **D** 75, 094504, 2007.







Outline

 ρ^+ features

NJL Model

Constituent Quark Form Factors

 ρ^+ Form Factors

Results

Conclusions



Conclusions

- Results for $\langle r^2 \rangle$ and $\mu_{
 ho}$ promising.
- ► Look forward to more detailed Lattice results.
- Ability to calculate FF for many $m_\pi^2 \rightarrow$ Lattice.
- Look forward to experiments?