

# T(R)OPICAL QCD 2010

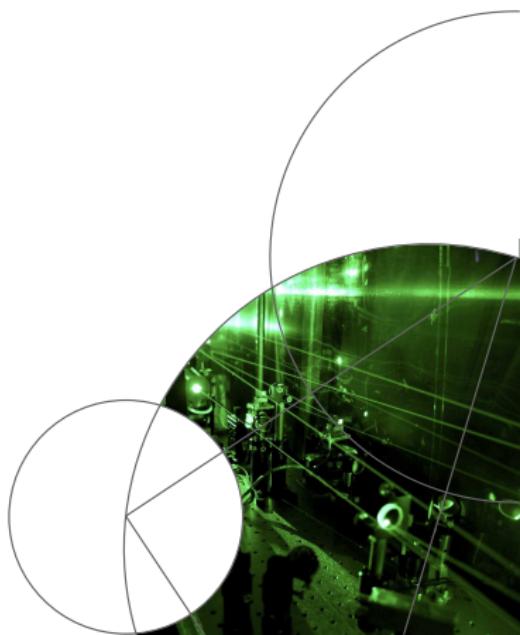


## The Radius of the Proton: Size Does Matter.

Jonathan D. Carroll

Centre for the Subatomic Structure of Matter

September 27, 2010



# Outline

## ① The Puzzle

The Publication

The Experiment

The Original Calculations

## ② The Investigation

The New Calculations

## ③ The Results



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## ① The Puzzle

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# The Publication

Pohl *et al.* Nature 466, 213-216 (8 July 2010)



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$$\text{CODATA: } r_p = 0.8768(69) \text{ fm}$$

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## Conclusion

*“... obtain a new value of the Rydberg constant,  
 $R_\infty = 10,973,731.568160(16) \text{ m}^{-1}$ .*

*This is 2110 kHz/c or 4.9σ away from the CODATA value, but 4.6 times more precise.”*

*“Finally, the origin of the discrepancy with the HFR data could originate from wrong or missing QED terms or from unexpectedly large contributions of yet uncalculated higher order terms.”*

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# The Experiment

## Muonic Hydrogen

The experiment focussed on muonic hydrogen ( $\mu p$ )

$$\frac{m_\mu}{m_e} \sim 200$$

so the wavefunction overlaps with the proton  
 $(m_\mu/m_e)^3 \sim 10^7$  times stronger than that of the electron  
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$2P_{1/2}$  —————

$\Delta E_{\text{Lamb}}^{2S-2P}$

$2S_{1/2}$  —————

The Lamb shift is the splitting of the degenerate  $2S_{1/2}$  and  $2P_{1/2}$  eigenstates, due to vacuum polarization

$$V_{VP}(r) = -\frac{Z\alpha}{r} \frac{\alpha}{3\pi} \int_4^\infty \frac{d(q^2)}{q^2} \frac{-m_e \alpha r}{1 - \frac{4}{q^2}} \left( 1 + \frac{2}{T^2} \right)$$

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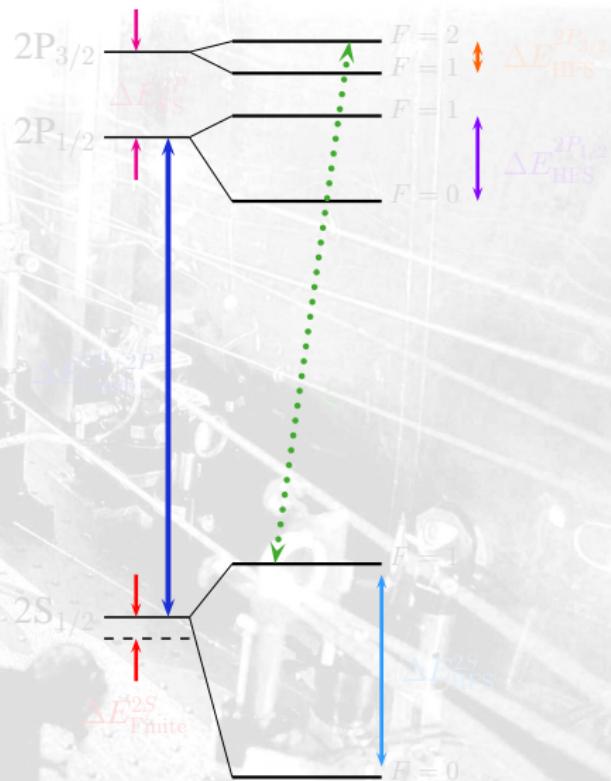
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## Muonic Hydrogen

The measured transition corresponds to the following sum of effects:

$$\begin{aligned}\xi &= E_{2P, F=2} - E_{2S, F=1} \\ &= \Delta E_{\text{Lamb}}^{2S} + \Delta E_{\text{Finite}}^{2S} \\ &\quad + \Delta E_{\text{FS}}^{2P} \\ &\quad + \frac{3}{8} \Delta E_{\text{HFS}}^{2P_{3/2}} - \frac{1}{4} \Delta E_{\text{HFS}}^{2S}\end{aligned}$$

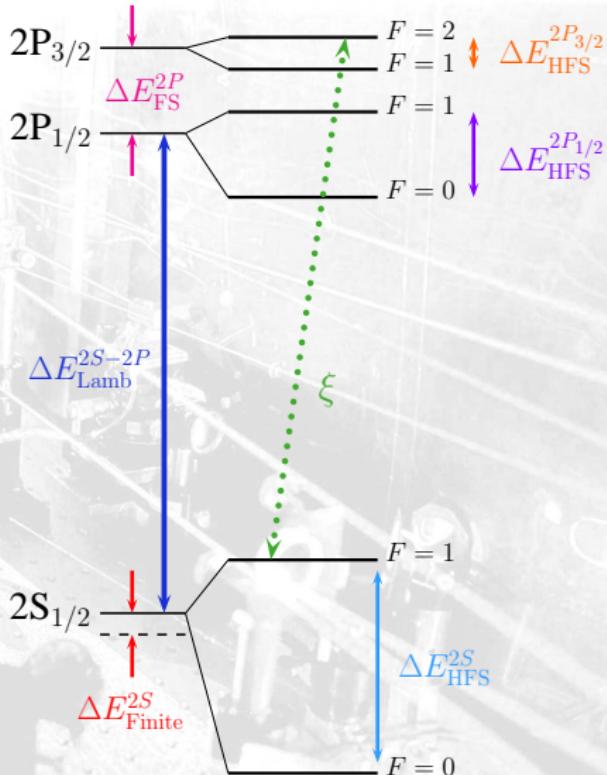


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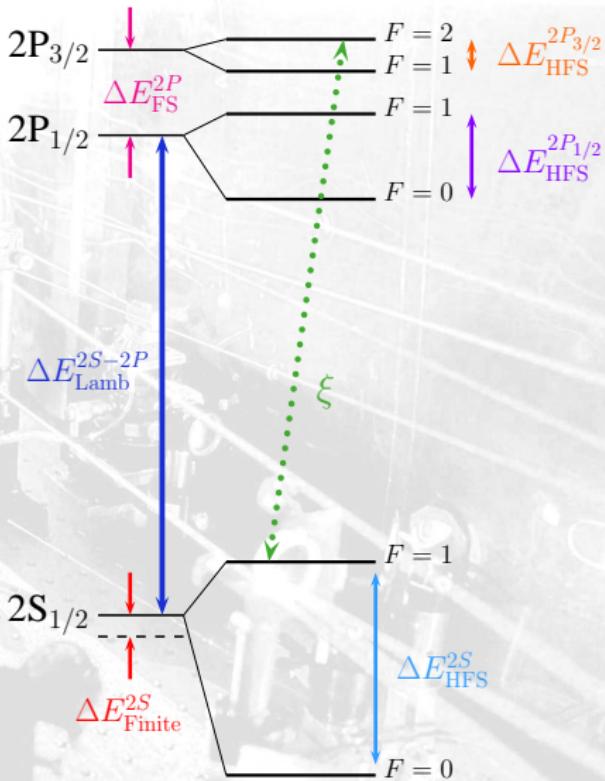


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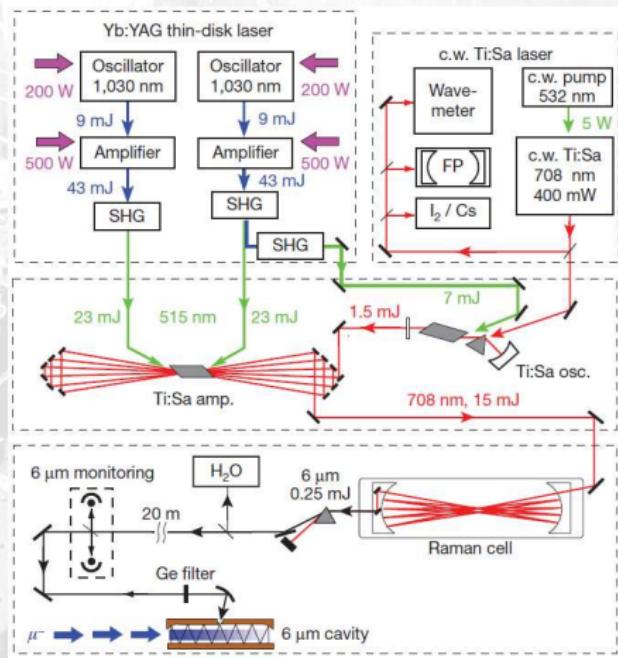
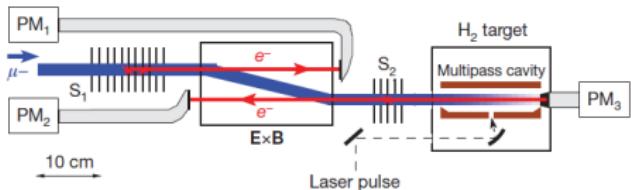
$$\begin{aligned}\xi &= E_{2P_{3/2}^F=2} - E_{2S_{1/2}^F=1} \\ &= \Delta E_{\text{Lamb}}^{2S-2P} + \Delta E_{\text{Finite}}^{2S} \\ &\quad + \Delta E_{FS}^{2P} \\ &\quad + \frac{3}{8} \Delta E_{HFS}^{2P_{3/2}} - \frac{1}{4} \Delta E_{HFS}^{2S}.\end{aligned}$$



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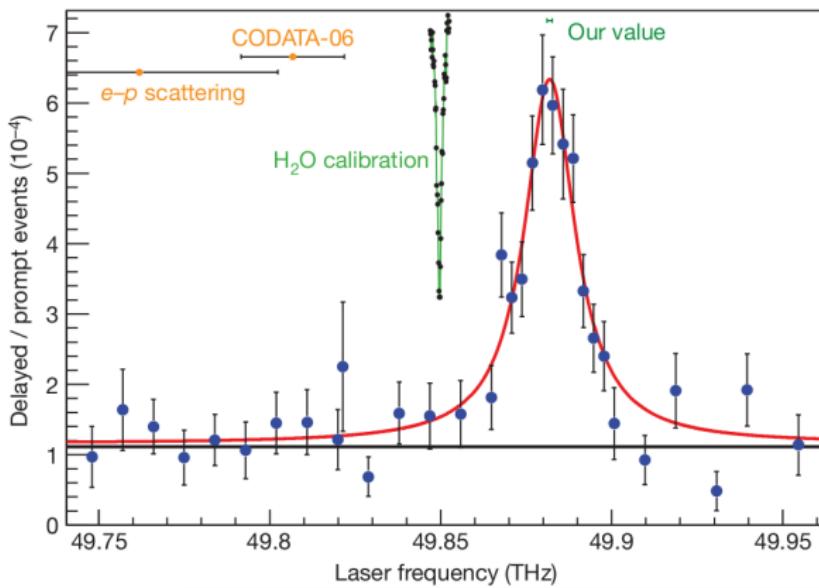
The experiment was performed at the  $\pi E5$  beam-line of the proton accelerator at PSI in Switzerland.



# The Experiment

## Muonic Hydrogen

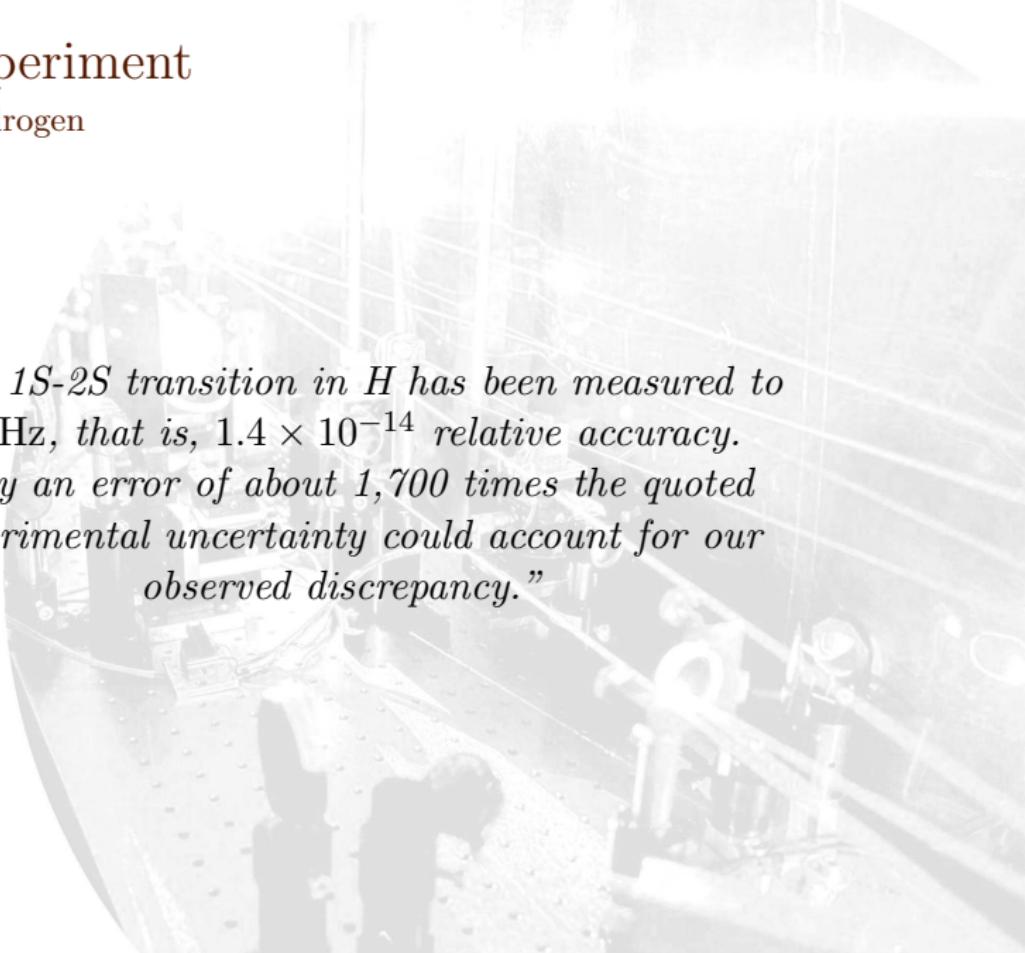
The data is clearly in disagreement with previous measurements and the world average



# The Experiment

## Muonic Hydrogen

*“The  $1S$ - $2S$  transition in  $H$  has been measured to 34 Hz, that is,  $1.4 \times 10^{-14}$  relative accuracy. Only an error of about 1,700 times the quoted experimental uncertainty could account for our observed discrepancy.”*



# The Experiment

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*“This new value of the proton radius  
 $r_p = 0.84184(67) \text{ fm}$   
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# The Original Calculations

## The Shift

The extraction of the proton radius from the shift involves comparing the measured shift to the following sum of contributions

$$\begin{aligned}\xi &= 206.2949(32) \text{ meV} \\ &= 206.0573(45) - 5.2262 \langle r_p \rangle^2 + 0.0347 \langle r_p \rangle^3 \text{ meV}\end{aligned}$$

The radius-dependence is entirely contained in the finite-size contributions.

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## The Approximations

To calculate the theoretical shift corresponding to the measured transition, **Borie, Pachucki, Martynenko, et al.** have used perturbation theory with non-relativistic wavefunctions to predict the size of the contributing effects, including relativistic effects

$$\Delta E_{V'}^{nlm} \sim \int_0^{\infty} V'(r) |\Psi_{\text{Schr}\ddot{\text{o}}\text{d.}}^{nlm}(r)|^2 dr$$

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To calculate the theoretical shift corresponding to the measured transition, **Borie, Pachucki, Martynenko, et al.** have used perturbation theory with non-relativistic wavefunctions to predict the size of the contributing effects, including relativistic effects.

$$V(r) = -\frac{Z\alpha}{r} + V'(r)$$

$$\Delta E_{V'}^{nlm} \sim \int_0^\infty V'(r) |\Psi_{\text{Schröd.}}^{nlm}(r)|^2 d^3r$$

# The Original Calculations

## Lamb Shift

The Lamb shift value(s) of Borie is(are) cited in the Nature article.

$$\Delta E_{\text{Lamb}}^{nlm} \sim \int_0^{\infty} V_{VP}(r) |\Psi_{\text{Schröd.}}^{nlm}(r)|^2 d^3r$$

$$\Delta E_{\text{Lamb}}^{2S-2P} \sim \int_0^{\infty} V_{VP}(r) |\Psi_{\text{Schröd.}}^{2S}(r)|^2 d^3r$$

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Still need to add recoil, radiative, and higher-order terms.

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26 other

Still need to add ~~recoil, radiative, and higher order~~ terms.

#	Contribution	Ref.	Our selection Value	Unc.	Pachucki <sup>1-3</sup> Value	Unc.	Borie <sup>5</sup> Value	Unc.
1	NR One loop electron VP	1,2			205.0074			
2	Relativistic correction (corrected)	1-3,5			0.0169			
3	Relativistic one loop VP	5	205.0282				205.0282	
4	NR two-loop electron VP	5,14	1.5081		1.5079		1.5081	
5	Polarization insertion in two Coulomb lines	1,2,5	0.1509		0.1509		0.1510	
6	NR three-loop electron VP	11	0.00529					
7	Polarisation insertion in two and three Coulomb lines (corrected)	11,12	0.00223					
8	Three-loop VP (total, uncorrected)				0.0076		0.00761	
9	Wichmann-Kroll	5,15,16	-0.00103				-0.00103	
10	Light by light electron loop contribution (Virtual Delbrück scattering)	6	0.00135	0.00135			0.00135	0.00015
11	Radiative photon and electron polarization in the Coulomb line $\alpha^2(Z\alpha)^4$	1,2	-0.00500	0.0010	-0.006	0.001	-0.005	
12	Electron loop in the radiative photon of order $\alpha^2(Z\alpha)^4$	17-19	-0.00150					
13	Mixed electron and muon loops	20	0.00007				0.00007	
14	Hadronic polarization $\alpha(Z\alpha)^4 m_r$	21-23	0.01077	0.00038	0.0113	0.0003	0.011	0.002
15	Hadronic polarization $\alpha(Z\alpha)^5 m_r$	22,23	0.000047					
16	Hadronic polarization in the radiative photon $\alpha^2(Z\alpha)^4 m_r$	22,23	-0.000015					
17	Recoil contribution	24	0.05750		0.0575		0.0575	
18	Recoil finite size	5	0.01300	0.001			0.013	0.001
19	Recoil correction to VP	5	-0.00410				-0.0041	
20	Radiative corrections of order $\alpha^n(Z\alpha)^k m_r$	2,7	-0.66770		-0.6677		-0.66788	
21	Muon Lamb shift 4th order	5	-0.00169				-0.00169	
22	Recoil corrections of order $\alpha(Z\alpha)^5 \frac{m}{M} m_r$	2,5-7	-0.04497		-0.045		-0.04497	
23	Recoil of order $\alpha^6$	2	0.00030		0.0003			
24	Radiative recoil corrections of order $\alpha(Z\alpha)^n \frac{m}{M} m_r$	1,2,7	-0.00960		-0.0099		-0.0096	
25	Nuclear structure correction of order $(Z\alpha)^5$ (Proton polarizability contribution)	2,5,22,25	0.015	0.004	0.012	0.002	0.015	0.004
26	Polarization operator induced correction to nuclear polarizability $\alpha(Z\alpha)^5 m_r$	23	0.00019					
27	Radiative photon induced correction to nuclear polarizability $\alpha(Z\alpha)^5 m_r$	23	-0.00001					
	Sum		206.0573	0.0045	206.0432	0.0023	206.05856	0.0046

# The Original Calculations

## Finite-Size Correction

The finite-size correction value(s) of Borie is(are) cited in the Nature article.

$$\Delta E_{\text{finite}} = -\frac{2Z\alpha}{3} \left( \frac{Z\alpha\mu}{2} \right)^3 \left[ \langle r_p^2 \rangle - \frac{Z\alpha\mu}{2} \langle r_p^3 \rangle \right]$$

# The Original Calculations

## 2P Fine Structure

The 2P fine structure shift value(s) of Martynenko is(are) cited in the Nature article.

$$\Delta E_{FS}^{2P} = \frac{\mu^3 (Z\alpha)^4}{32m_\mu^2} \left( 1 + \frac{m_\mu}{2m_p} \right)$$

# The Original Calculations

## 2S Hyperfine Structure

The 2S hyperfine structure shift value(s) of Martynenko  
is(are) cited in the Nature article.

$$\Delta E_{HFS}^{2S_{1/2}} = \frac{1}{3}(Z\alpha)^4 \frac{\mu^3}{m_\mu m_p} (1 + \kappa)$$

# The Original Calculations

## 2P Hyperfine Structure

The 2P hyperfine structure shift value(s) of Martynenko is(are) cited in the Nature article.

$$\Delta E_{HFS}^{2P_{3/2}} = (Z\alpha)^4 \frac{\mu^3(1 + \kappa)}{3m_\mu m_p} \left[ \frac{2}{15} - \frac{a_\mu}{30} + \frac{m_\mu(1 + 2\kappa)}{2m_p(1 + \kappa)} \right]$$

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# The New Calculations

## Dirac Equation

To calculate the shifts exactly, we can use the Dirac Equation with the appropriate potential to calculate the perturbed wavefunctions.

$$U_\alpha(x) = E_{\alpha m n t}(x) = \begin{bmatrix} G_{\alpha m}(r) & F_{\alpha m}(r) \\ -F_{\alpha m}(r) & G_{\alpha m}(r) \end{bmatrix}$$
$$\int_0^\infty dr \left( |G_{\alpha m}(r)|^2 + |F_{\alpha m}(r)|^2 \right) = 1$$

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To calculate the shifts exactly, we can use the Dirac Equation with the appropriate potential to calculate the perturbed wavefunctions.

$$U_\alpha(x) = U_{n\kappa m t}(x) = \begin{bmatrix} iG_{n\kappa t}(r)/r & \Phi_{\kappa m}\eta_t \\ -F_{n\kappa t}(r)/r & \Phi_{-\kappa m}\eta_t \end{bmatrix}$$

$$\int_0^\infty dr (|G_\alpha(r)|^2 + |F_\alpha(r)|^2) = 1$$

# The New Calculations

## Dirac Equation

$$\frac{d}{dr} G_\alpha(r) = -\frac{\kappa}{r} G_\alpha(r) + [\epsilon_\alpha + \mu - V(r)] F_\alpha(r)$$

$$\frac{d}{dr} F_\alpha(r) = \frac{\kappa}{r} F_\alpha(r) - [\epsilon_\alpha - \mu - V(r)] G_\alpha(r)$$

$\Rightarrow$

$$\frac{d}{dr} G_\alpha(r) = -C_\alpha(r) + [\lambda_\alpha + 2\mu - V(r)] F_\alpha(r)$$

$$\frac{d}{dr} F_\alpha(r) = \frac{\kappa}{r} F_\alpha(r) - [\lambda_\alpha - V(r)] G_\alpha(r)$$

$$\kappa = \frac{M_p m_\mu}{M_p + m_\mu}, \quad \lambda_\alpha = \epsilon_\alpha - \mu$$

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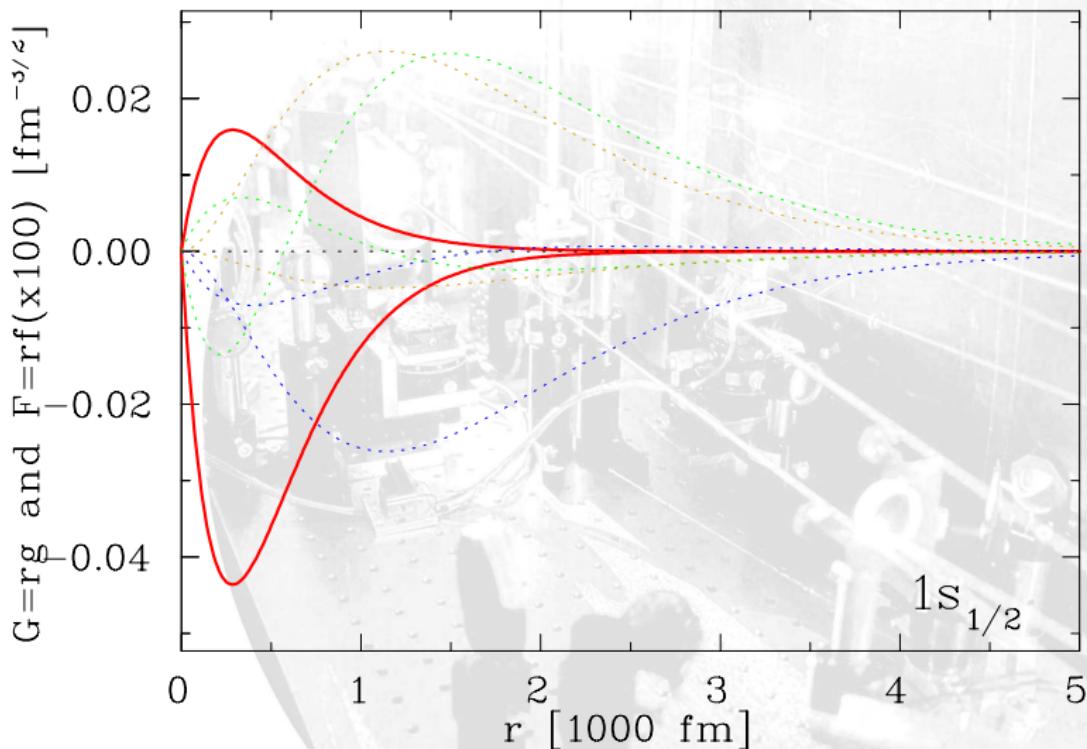
$$\frac{d}{dr}G_\alpha(r) = -\frac{\kappa}{r}G_\alpha(r) + [\lambda_\alpha + 2\mu - V(r)] F_\alpha(r)$$

$$\frac{d}{dr}F_\alpha(r) = \frac{\kappa}{r}F_\alpha(r) - [\lambda_\alpha - V(r)] G_\alpha(r)$$

$$\mu = \frac{M_p m_\mu}{M_p + m_\mu}; \quad \lambda_\alpha = \epsilon_\alpha - \mu$$

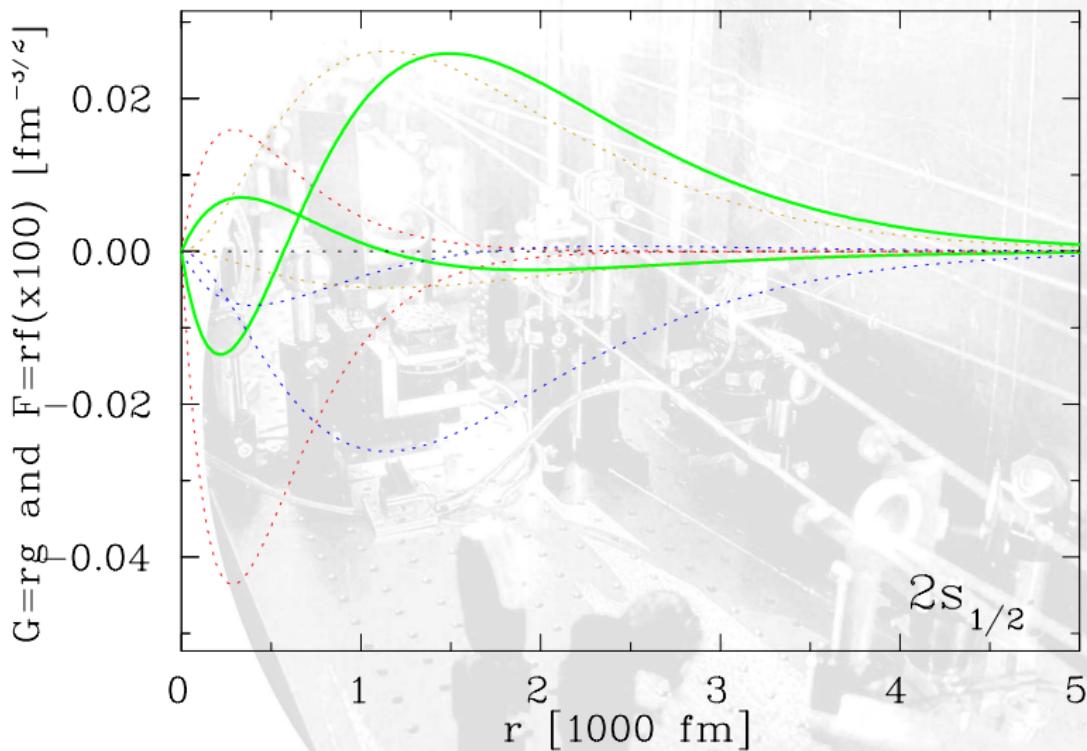
# The New Calculations

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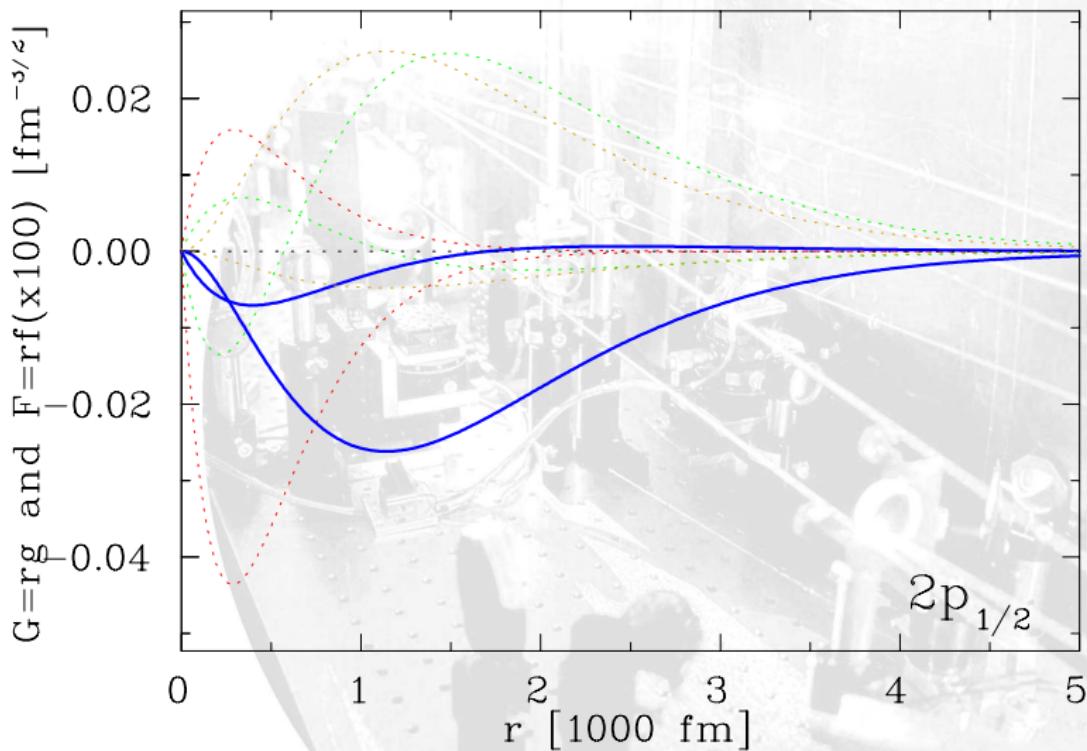
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## Dirac Equation

 $2s_{1/2}$

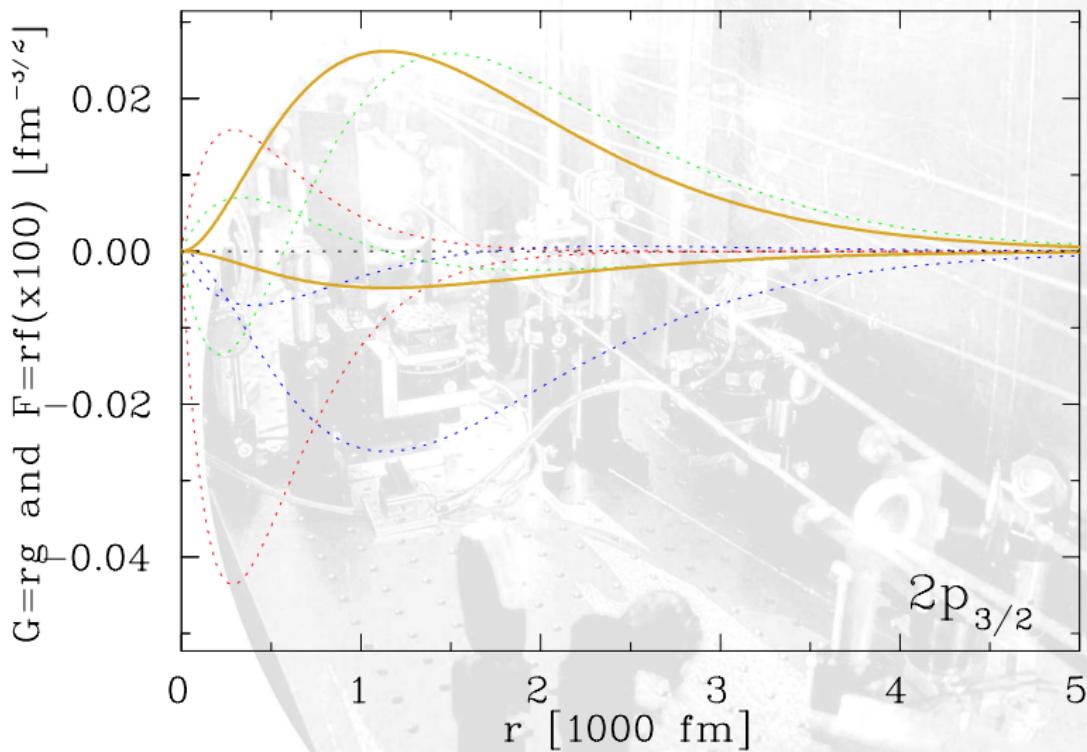
# The New Calculations

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## Dirac Equation



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## Dirac Equation

The (shifted) eigenvalues can be reliably reproduced by using the point-Coulomb potential

$$V(r) = -\frac{Z\alpha}{r}$$

and the unperturbed eigenvalues are known exactly

$$\lambda_\alpha = \epsilon_\alpha - \mu = \mu \left( \left[ 1 + \frac{Z^2 \alpha^2}{(n_\alpha - |\kappa_\alpha| + \sqrt{\kappa_\alpha^2 - Z^2 \alpha^2})^2} \right]^{-\frac{1}{2}} - 1 \right)$$

so we can check the accuracy and precision of our code...

# The New Calculations

## Dirac Equation

$1S_{1/2} : \lambda = -2.5285267981975594 \text{ keV}$

$2S_{1/2} : \lambda = -0.6321338034582128 \text{ keV}$

$2P_{1/2} : \lambda = -0.6321338034582128 \text{ keV}$

$2P_{3/2} : \lambda = -0.6321253878229215 \text{ keV}$

# The New Calculations

## Dirac Equation

$1S_{1/2} : \lambda = -2.5285267981975594$  keV

$1S_{1/2} : \lambda = -2.5285267981975\textcolor{red}{605}$  keV

$2S_{1/2} : \lambda = -0.6321338034582128$  keV

$2S_{1/2} : \lambda = -0.6321338034\textcolor{red}{390614}$  keV

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$$2P_{3/2} : \lambda = -0.632125387822921\textcolor{red}{3} \text{ keV} \leftarrow \delta = 0.2 \text{ peV !}$$

# The New Calculations

## Lamb Shift

We can calculate the Lamb shift by including the vacuum polarization potential such that

$$V(r) = -\frac{Z\alpha}{r} - \frac{Z\alpha}{r} \frac{\alpha}{3\pi} \int_4^\infty \frac{d(q^2)}{q^2} e^{-m_{eq}r} \sqrt{1 - \frac{4}{q^2}} \left(1 + \frac{2}{q^2}\right)$$

and calculating a converged wavefunction, then taking the difference in eigenvalues.

# The New Calculations

## Lamb Shift

We can calculate the Lamb shift by including the vacuum polarization potential such that

$$V(r) = -\frac{Z\alpha}{r} - \frac{Z\alpha}{r} \frac{\alpha}{3\pi} \int_4^\infty \frac{d(q^2)}{q^2} e^{-m_e qr} \sqrt{1 - \frac{4}{q^2}} \left(1 + \frac{2}{q^2}\right)$$

and calculating a converged wavefunction, then taking the difference in eigenvalues.

# The New Calculations

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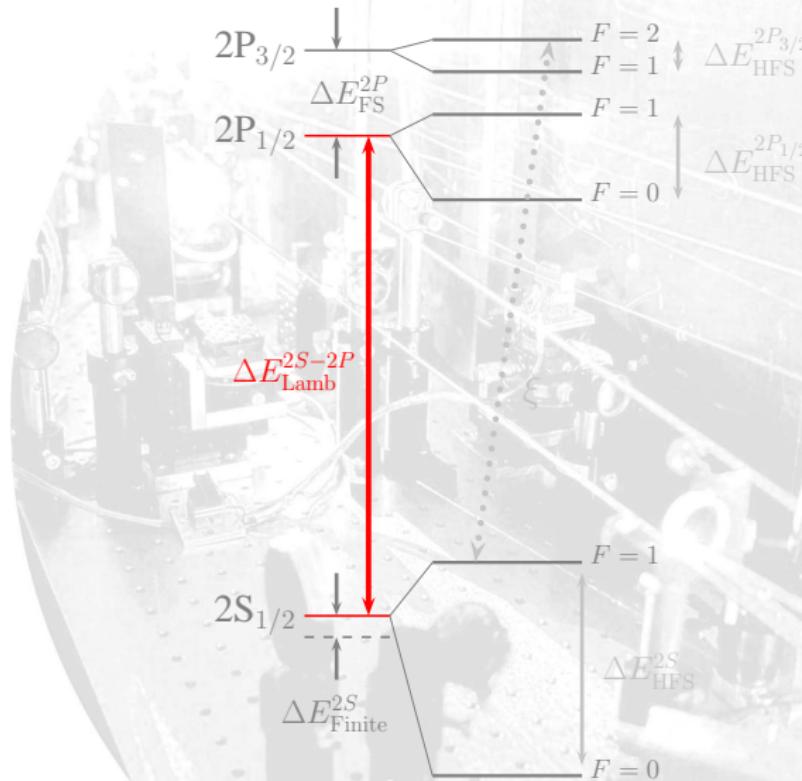
## Lamb Shift

This is **almost** the Lamb shift... still need to add the **24** additional terms.

$$\text{Full Calculation} = \underbrace{\begin{array}{l} \text{Relativistic one-loop VP} \\ + \text{Polarization insertion} \\ \text{in 2 and 3 Coulomb lines} \end{array}}_{\text{Perturbative Calculation}}$$

# The New Calculations

## Lamb Shift



# The New Calculations

## More Approximations

$$\Psi_{\text{Schröd.}} \neq \Psi_{\text{Dirac}}$$

But even this is not perfect.

There are still recoil corrections to include to approximate the Bethe-Salpeter equation.

For now we will assume these are unaffected by the vacuum polarization.

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## Finite-Size Correction

The point-Coulomb potential assumes that the proton is a point-particle.

This is not an accurate assumption to make, so we use the finite-size Coulomb potential instead

$$V_C(r) = -\frac{Z\alpha}{r} \rightarrow -Z\alpha \int \frac{\rho(r')}{|\vec{r} - \vec{r}'|} d\tau'$$

Since this involves the charge density, this leads to a radius-dependent quantity

$$\rho(r) = \frac{n}{8\pi} e^{-nr}; \quad n = \sqrt{12/r_p^2}$$

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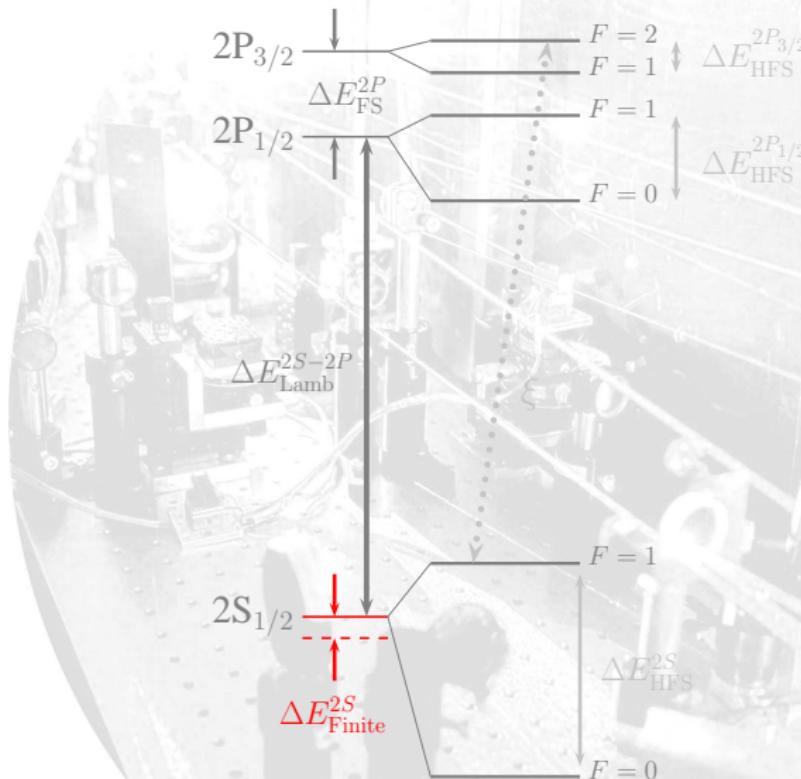
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# The New Calculations

## Finite-Size Correction



# The New Calculations

## 2p Fine Structure

Subtracting the exact eigenvalues of the  $2P_{1/2}$  and  $2P_{3/2}$  gives the fine structure splitting precisely.

$$\Delta E_{FS}^{2P} = \epsilon_{2P_{3/2}} - \epsilon_{2P_{1/2}}$$

which we can also calculate in the presence of the vacuum polarizations

$$\frac{Ze}{r} + VVP(r)$$

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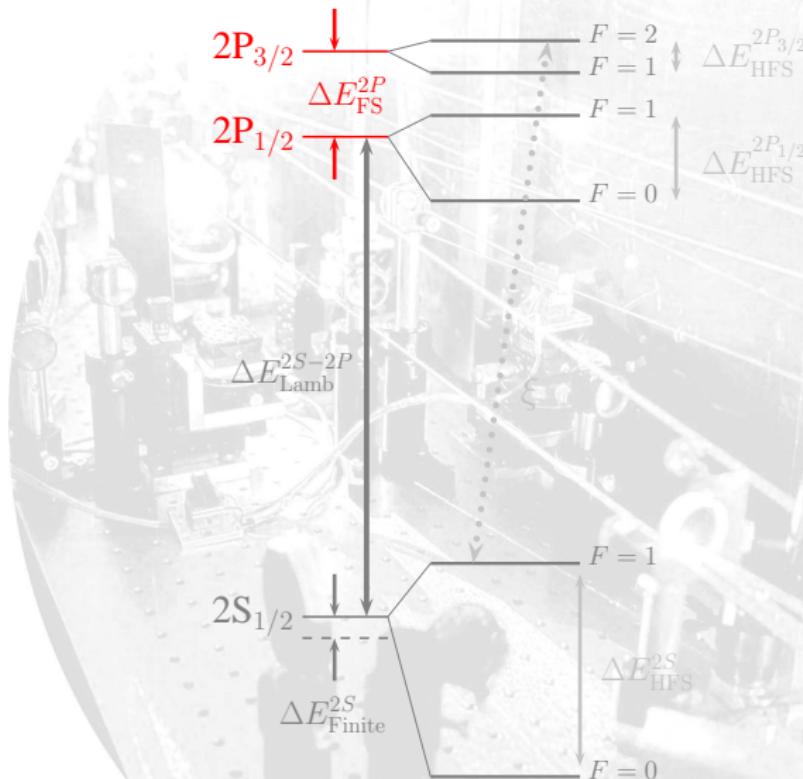
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# The New Calculations

## 2p Fine Structure



# The New Calculations

## Hyperfine Structure

The hyperfine structure is a measure of the  $\ell \cdot \sigma$  coupling.  
The Hamiltonian is

$$\mathcal{H} = 2\beta\gamma\hbar \frac{\ell(\ell+1)}{j(j+1)} \left\langle \frac{1}{r^3} \right\rangle \mathbf{I} \cdot \mathbf{J} + A_F \mathbf{I} \cdot \mathbf{S}$$

$$A_F = \frac{16\pi}{3} \beta\gamma\hbar |\psi(0)|^2$$

$$\mathbf{I} \cdot \mathbf{J} = \frac{1}{2} [F(F+1) - I(I+1) - j(j+1)]$$

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# The New Calculations

## 2s Hyperfine Structure

For  $\ell = 0$  the contact term in the Hamiltonian is non-zero, while the dipole term vanishes

$$E_{HFS}^{2S} = \frac{16\pi}{3} \beta \gamma \hbar |\psi(0)|^2 \mathbf{I} \cdot \mathbf{S}$$

$\Rightarrow$

$$\Delta E_{HFS}^{2S(F=1-F=0)} = \frac{16\pi}{3} \beta \gamma \hbar |\psi(0)|^2$$

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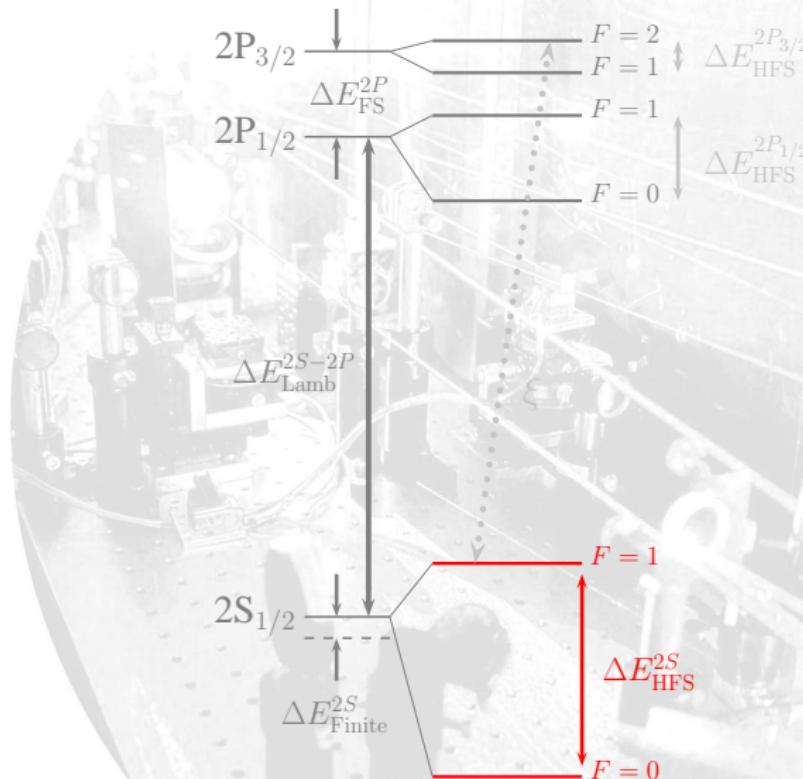
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# The New Calculations

## 2s Hyperfine Structure



# The New Calculations

## 2p Hyperfine Structure

For  $\ell \neq 0$  the dipole term in the Hamiltonian is non-zero, while the contact term vanishes

$$E_{HFS}^{2P} = 2\beta\gamma\hbar \frac{\ell(\ell+1)}{j(j+1)} \left\langle \frac{1}{r^3} \right\rangle \mathbf{I} \cdot \mathbf{J}$$

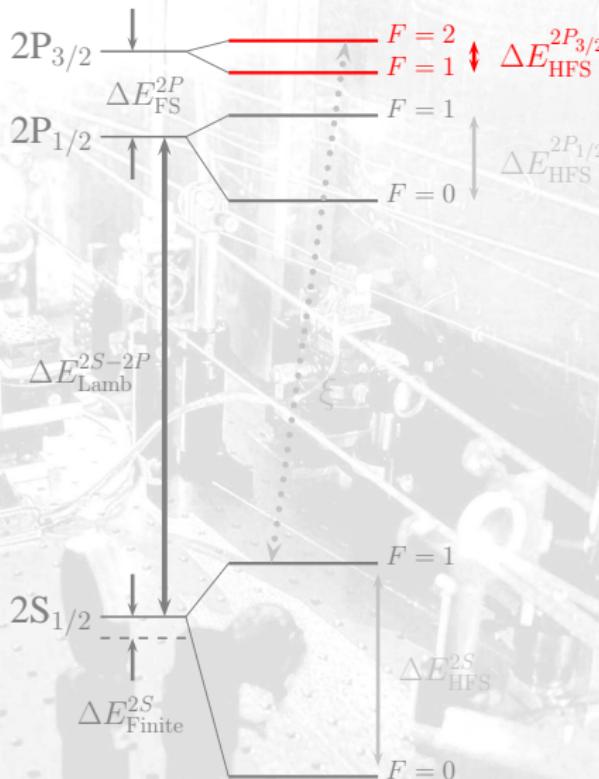
$\Rightarrow$

$$\Delta E_{HFS}^{2P_{3/2}(F=2-F=1)} = 2\beta\gamma\hbar \frac{\ell(\ell+1)}{j(j+1)} \left\langle \frac{1}{r^3} \right\rangle \times [F(F+1) - I(I+1) - j(j+1)]$$

TO BE COMPLETED!

# The New Calculations

## 2p Hyperfine Structure



# The New Calculations

## Vacuum Polarization Finite-Size Correction

Since we are solving for exact wavefunctions, we can also add additional terms to the potential. For example, we can investigate the finite-size effect of the vacuum polarization

$$V_{VP}(r) \rightarrow -\frac{2Z\alpha^2}{3\pi} \int \frac{\rho(r')}{|\vec{r} - \vec{r}'|} Z_0(|\vec{r} - \vec{r}'|) d\tau'$$

$$Z_n(|\vec{r}|) = \int_1^\infty e^{-\frac{2}{\lambda}|\vec{r}|\xi} \left(1 + \frac{1}{2\xi^2}\right) \frac{(1+\xi)^{\frac{1}{2}}}{\xi^n \xi^2} d\xi$$

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TO BE COMPLETED!

# Outline

## ① The Puzzle

The Publication

The Experiment

The Original Calculations

## ② The Investigation

The New Calculations

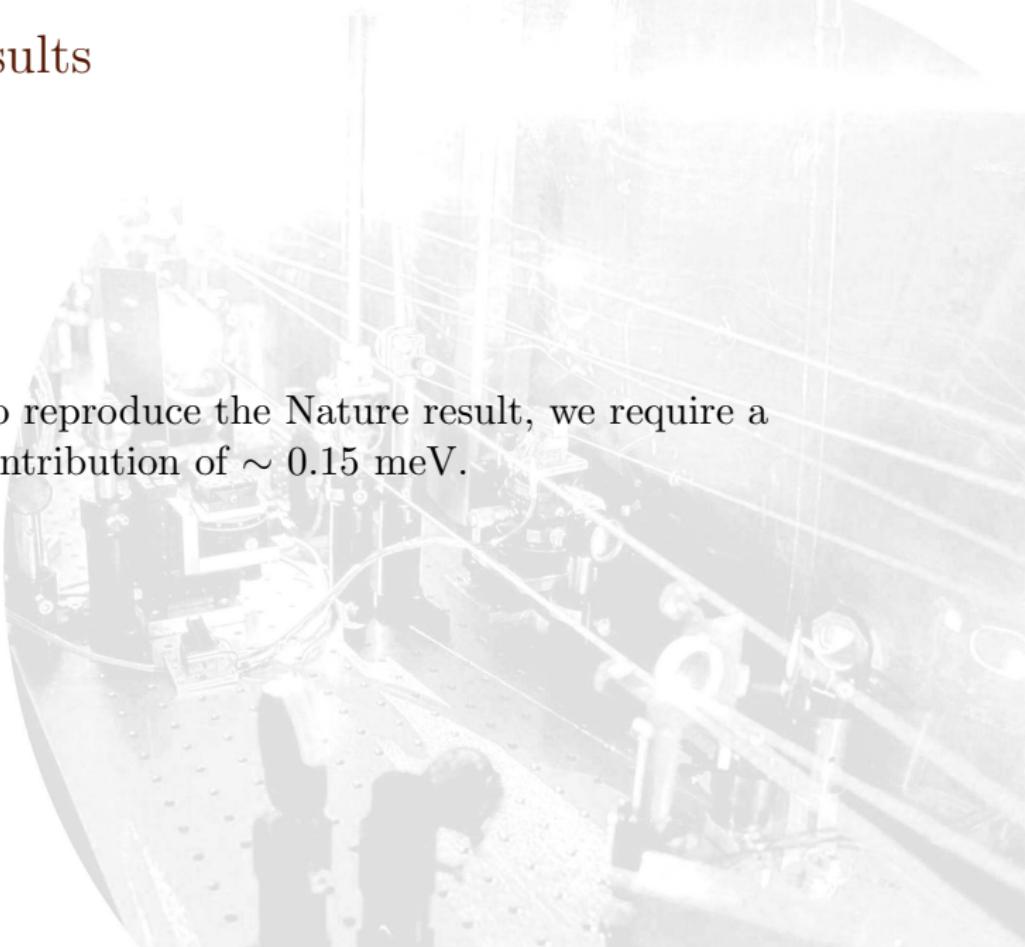
## ③ The Results



# The Results

## The Data

In order to reproduce the Nature result, we require a missing contribution of  $\sim 0.15$  meV.



I	Contribution	Carroll	Nature	Other	$ \delta $	Precision
	$E(1S_{1/2}) : V = V_C$	-2.5285267981975605 keV		-2.5285267981975594 keV [exact]	1.1 peV	QP
	$E(1S_{1/2}) : V = V_C + V_{VP}$	-2.530427808554 keV				DP
	$E(2S_{1/2}) : V = V_C$	-0.6321338034390614 keV		-0.6321338034582128 keV [exact]	19.15 neV	QP
	$E(2S_{1/2}) : V = V_C + V_{VP}$	-0.632353566988 keV				DP
	$E(2P_{1/2}) : V = V_C$	-0.632133803154 keV		-0.632133803457 keV [exact]	0.30 $\mu$ eV	DP
	$E(2P_{1/2}) : V = V_C + V_{VP}$	-0.632148387529 keV				DP
	$E(2P_{3/2}) : V = V_C$	-0.6321253878229213 keV		-0.6321253878229215 keV [exact]	0.2 peV	QP
	$E(2P_{3/2}) : V = V_C + V_{VP}$	-0.632139967131 keV				DP
	NR One loop electron VP	205.008 meV		205.0074 meV [Pachucki]	0.6 $\mu$ eV	DP
	Relativistic one loop VP	205.029 meV	205.0282 meV	205.0282 meV [Borie]	0.8 $\mu$ eV	DP
	+Polarization insertion in two Coulomb lines		+0.1509 meV			DP
	+Polarisation insertion in two and three Coulomb lines		+0.00223 meV			DP
	$= \Delta E(2P_{1/2}) - \Delta E(2S_{1/2}) : V = V_C + V_{VP}$	205.17953 meV	205.18133 meV		1.8 $\mu$ eV	DP
!	$\langle r_p^2 \rangle$ nuclear size correction	5.22047 $\langle r_p^2 \rangle$	5.22619 $\langle r_p^2 \rangle$		5.7e-3 $\langle r_p^2 \rangle$	DP
	$\langle r_p^3 \rangle$ nuclear size correction	0.0484745 $\langle r_p^3 \rangle$	0.0347 $\langle r_p^3 \rangle$		1.4e-2 $\langle r_p^3 \rangle$	DP
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	$\Delta E_{FS}^{2P} : V = V_C + V_{VP}$	8.420398 meV				DP
	(NR) $\Delta E_{HFS}^{2S}$	22.8053 meV		22.8054 meV [Martynenko]	0.1 $\mu$ eV	DP
	+Relativistic correction $\mathcal{O}(\alpha^6)$		+0.0026 meV			DP
	+One-loop $e^-$ VP $\mathcal{O}(\alpha^6), \mathcal{O}(\alpha^6)$		+0.0482 meV			DP
	+VP corrections in second-order PT $\mathcal{O}(\alpha^6), \mathcal{O}(\alpha^6)$		+0.00746 meV			DP
	$= \Delta E_{HFS}^{2S} : V = V_C + V_{VP}$	22.8988 meV		22.9308 meV	32.5 $\mu$ eV	DP
	$\Delta E_{HFS}^{2S} : V = V_C$	22.8182 meV				DP
	$\Delta E_{HFS}^{2P_{3/2}} : V = V_C$	? meV	3.392588 meV	3.392588 meV [Martynenko]		DP
	$\Delta E_{HFS}^{2P_{3/2}}$	: $V = V_C + V_{VP}$	? meV			DP

16 quantities reconstructed.  
No significant deviations.

TOXICITY INDEX

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	$\Delta E_{HFS}^{2P_{3/2}} : V = V_C + V_{VP}$	? meV				DP

~ 16 quantities re-calculated.

No significant deviations from approximations.

# The Results

## The Data

Lamb Shift:

$$\begin{aligned} & \text{Relativistic one-loop VP} \\ & + \text{Polarization insertion in} \\ & \quad 2 \text{ and } 3 \text{ Coulomb lines} \\ & = 205.18133 \text{ meV} \end{aligned}$$

---

$$\begin{aligned} \text{Full calculation} &= 205.1\textcolor{red}{7953} \text{ meV} \end{aligned}$$

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## The Data

Finite Size Correction:

$$\begin{aligned} \text{coefficient } & \langle r_p \rangle^2 \\ = & 5.22619 \end{aligned}$$

---

$$\begin{aligned} \text{coefficient } & \langle r_p \rangle^3 \\ = & 0.0347 \end{aligned}$$

---

$$\begin{aligned} \text{Full calculation} & = 5.22\textcolor{red}{047} \end{aligned}$$

$$\begin{aligned} & = 0.0\textcolor{red}{484745} \end{aligned}$$

# The Results

## The Data

2P Fine Structure:

$$\begin{aligned} &\text{Dirac + Vacuum Polarization} \\ &= 8.42064 \text{ meV} \end{aligned}$$

---

$$\begin{aligned} \text{Full calculation} &= 8.420\textcolor{red}{398} \end{aligned}$$

# The Results

## The Data

$2S_{1/2}$  Hyperfine Structure:

$$\begin{aligned} & \text{Non-relativistic shift} \\ & + \text{Relativistic correction} \\ & + \text{One-loop } e^- \text{ VP } \mathcal{O}(\alpha^5), \mathcal{O}(\alpha^6) \\ & + \text{VP corrections in second-order PT } \mathcal{O}(\alpha^6) \\ & = 22.9308 \text{ meV} \end{aligned}$$

---

Full calculation  $= 22.8988 \text{ meV}$

# The Results

## The Data

$2P_{3/2}$  Hyperfine Structure:

Non-relativistic shift

?

?

$$= 3.392588 \text{ meV}$$

Full calculation

$$= ? \text{ meV}$$

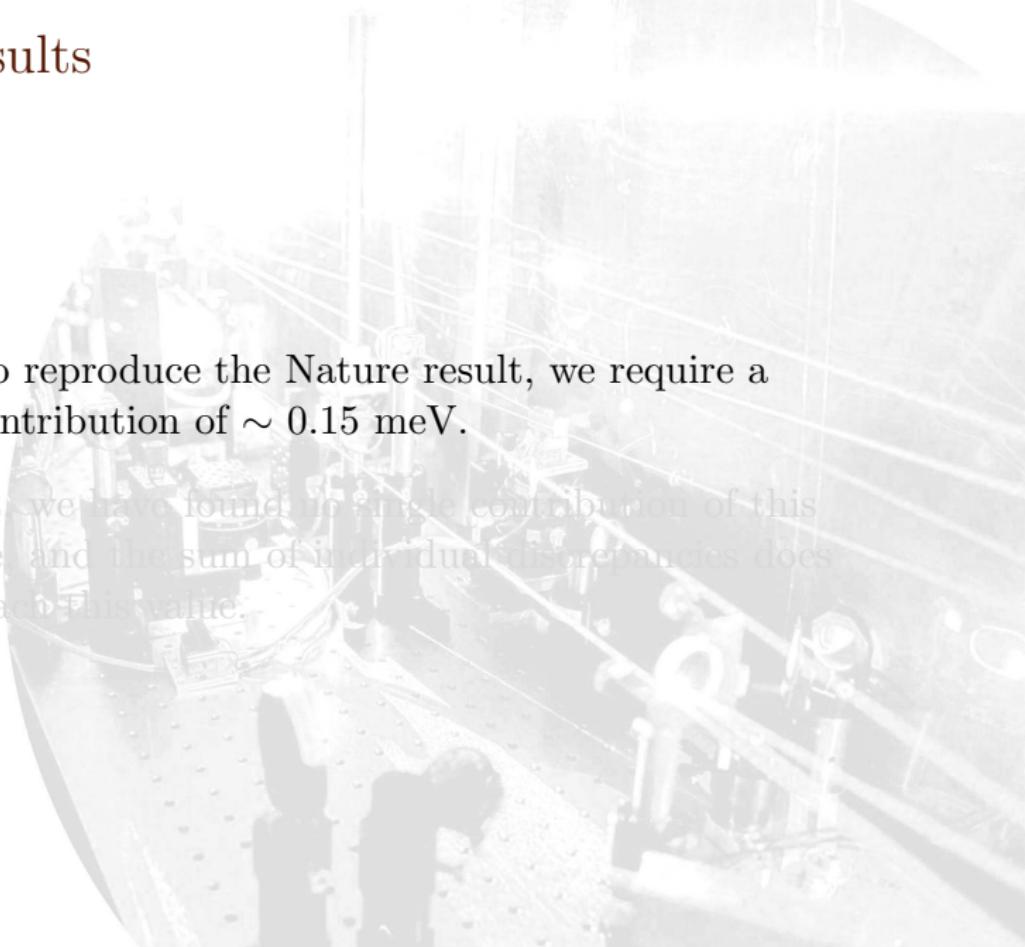
TO BE COMPLETED!

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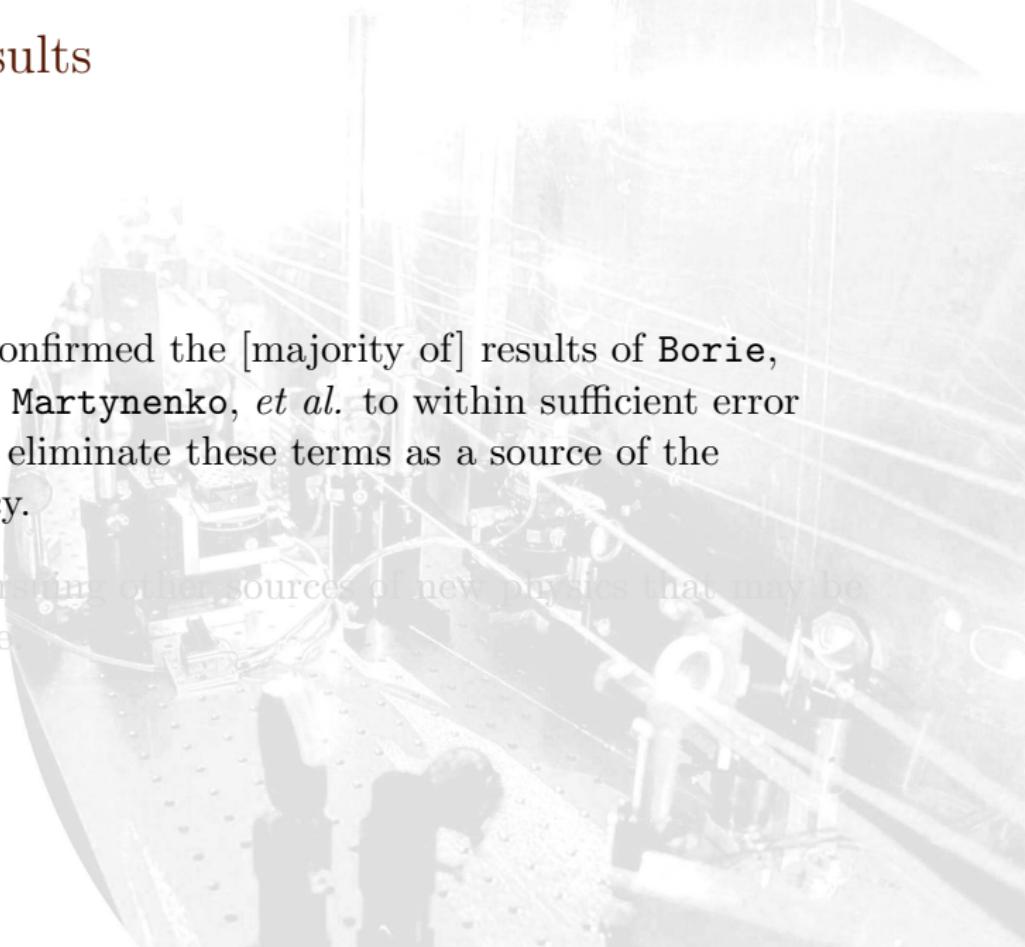
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We have confirmed the [majority of] results of **Borie, Pachucki, Martynenko, et al.** to within sufficient error such as to eliminate these terms as a source of the discrepancy.

We are pursuing other sources of new physics that may be responsible.



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# Further Reading I



R. Pohl et al.

The size of the proton

Nature Vol 466 (8 July 2010)

[doi:10.1038/nature09250]



E. Borie

Lamb shift in muonic hydrogen

Phys.Rev.A71:032508 (2005)



K. Pachucki

Theory of the Lamb shift in muonic hydrogen

Phys.Rev.A53:2092-2100 (1996)



A. P. Martynenko

2S Hyperfine Splitting of Muonic Hydrogen

Phys.Rev.A71:022506 (2005)

## Further Reading II



A. P. Martynenko

Fine and hyperfine structure of P-wave levels in  
muonic hydrogen

Phys.At.Nucl.71:125-135 (2008)



J. Rafelski, L. P. Fulcher

Fermions And Bosons Interacting With Arbitrarily  
Strong External Fields

Phys.Rep.38:227-361 (1978)

# Cheers, Everyone!

*The spectrum of the hydrogen atom has proved  
to be the Rosetta stone of modern physics.*

— T.W. Hänsch

