

Five-quark components and breathing mode for baryons

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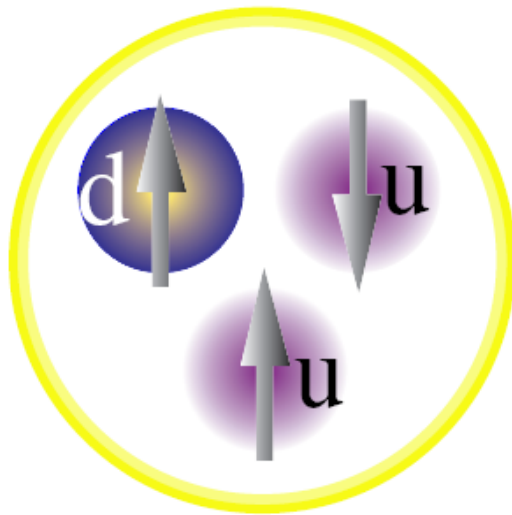
Outline:

- **Introduction -- 5-quark components in the proton**
- **New scheme for $N^*(1535)$ and its $1/2^-$ nonet partners with large 5-quark components and breathing mode**
- **Evidence for the predicted $\Sigma^*(1/2^-)$**
- **Extension to hidden charm and beauty**
- **Conclusion**

1. Introduction: 5-quark components in the proton

Classical picture of the proton

Constituent Quarks



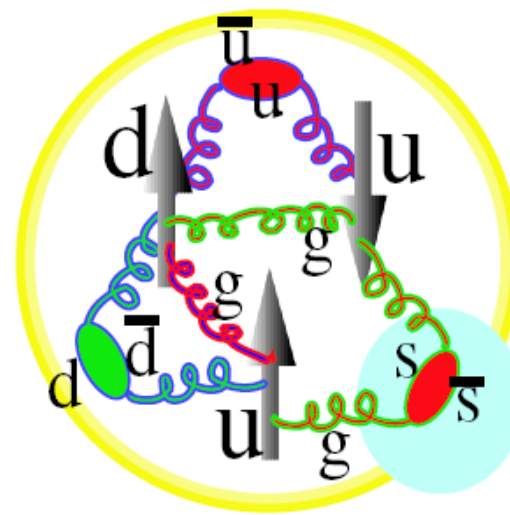
($Q^2 = 0 \text{ GeV}^2$)

baryon octet

masses, magn. momenta

1964-1974

Parton Distributions



($Q^2 > 1 \text{ GeV}^2$)

structure functions

momentum, spin

$$\bar{u}(x) = \bar{d}(x), \quad \bar{s}(x) = s(x)$$

1974-1992

Flavor asymmetry of light quarks in the nucleon sea

Deep Inelastic Scattering (DIS) + Drell-Yan (DY) process

$$\rightarrow \quad \bar{d} - \bar{u} \sim 0.12 \quad \text{for a proton}$$

Garvey&Peng, *Prog. Part. Nucl. Phys.*47, 203 (2001)

Table 1. Values of the integral $\int_0^1 [\bar{d}(x) - \bar{u}(x)]dx$ determined from the DIS, semi-inclusive DIS, and Drell-Yan experiments.

Experiment	$\langle Q^2 \rangle$ (GeV ² /c ²)	$\int_0^1 [\bar{d}(x) - \bar{u}(x)]dx$
NMC/DIS	4.0	0.147 ± 0.039
HERMES/SIDIS	2.3	0.16 ± 0.03
FNAL E866/DY	54.0	0.118 ± 0.012

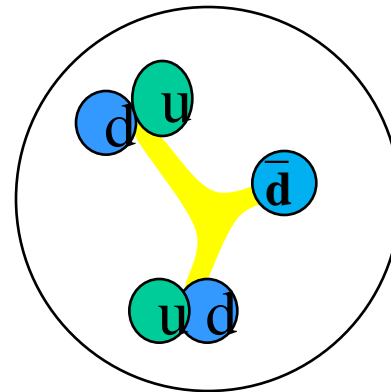
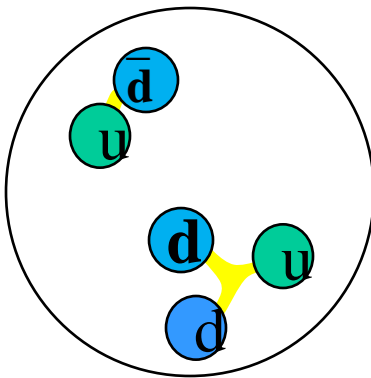
Two major theoretical schemes for $\bar{d} - \bar{u} \sim 0.12$

Meson cloud picture: Thomas, Speth, Henley, Meissner, Miller, Weise, Oset, Brodsky, Ma, ...

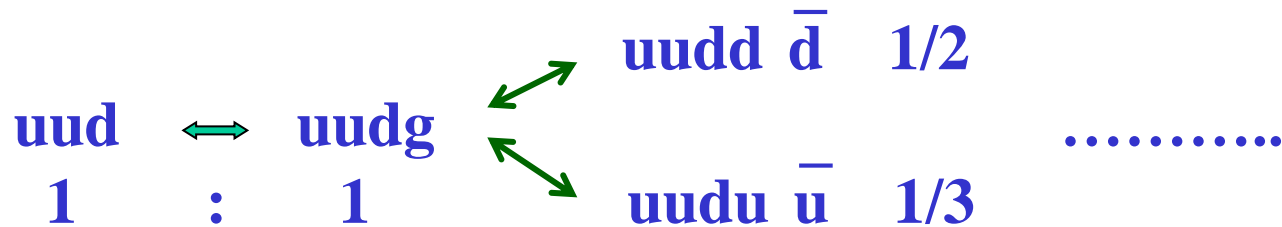
$$|p\rangle \sim |uud\rangle + \varepsilon_1 |n(udd)\pi^+(\bar{d}u)\rangle + \varepsilon_2 |\Delta^{++}(uuu)\pi^-(\bar{u}d)\rangle + \varepsilon' |\Lambda(uds)K^+(\bar{s}u)\rangle + \dots$$

Penta-quark picture: Riska, Zou, Zhu, ...

$$|p\rangle \sim |uud\rangle + \varepsilon_1 |[ud][ud]\bar{d}\rangle + \varepsilon' |[ud][us]\bar{s}\rangle + \dots$$



Detailed balance model : Zhang, Ma, Zou, Yang, Alberg, Henley



$p = 0.168 (uud) + 0.168 (uudg) + 0.084 (uudd \bar{d}) + 0.056 (uudu \bar{u})$
 $+ 0.084 (uudgg) + \dots$ $\bar{d} - \bar{u} \sim 0.124$

$(uud+ng)$ 50% $(uudd \bar{d}+ng)$ 22.4% $(uudu \bar{u} +ng)$ 15.0%

With ~25% $\bar{q}qqqq$ components in the proton, the “spin crisis” and single spin asymmetry may also be naturally explained.

An-Riska-Zou, PRC73 (2006) 035207; F.X.Wei, B.S.Zou, hep-ph/0807.2324

$$\Delta_u = 0.85 \pm 0.17$$

$$\Delta_d = -(0.33 \sim 0.56)$$

$$\Delta_u = \frac{4}{3} |A_{3q}|^2$$

$$\Delta_d = -\frac{1}{3}(1 - P_{s\bar{s}})$$

$$\Delta L_q = \frac{4}{3}(P_{d\bar{d}} + \bar{P}_{s\bar{s}})$$

**We must go beyond the simple 3q models,
meson cloud vs penta-quark not settled yet.**

2. New scheme for $N^*(1535)$ and its $1/2^-$ nonet partners

- Mass order reverse problem for the lowest excited baryons

$uud (L=1) 1/2^- \sim N^*(1535)$ **should be the lowest**

$uud (n=1) 1/2^+ \sim N^*(1440)$

$uds (L=1) 1/2^- \sim \Lambda^*(1405)$

harmonic oscillator $(2n + L + 3/2) \hbar\omega$

- Strange decays of $N^*(1535)$: **PDG \rightarrow large $g_{N^*N\eta}$**

$J/\psi \rightarrow \bar{p}N^* \rightarrow \bar{p} (K\Lambda) / \bar{p} (p\eta) \rightarrow$ **large $g_{N^*K\Lambda}$**

Liu&Zou, PRL96 (2006) 042002; Geng,Oset,Zou&Doring, PRC79 (2009) 025203

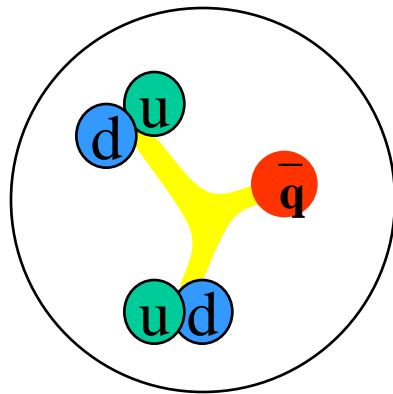
$\gamma p \rightarrow p\eta' \text{ \& } pp \rightarrow pp\eta' \rightarrow$ **large $g_{N^*N\eta'}$**

M.Dugger et al., PRL96 (2006) 062001; Cao&Lee, PRC78(2008) 035207

$\pi^- p \rightarrow n\phi \text{ \& } pp \rightarrow pp\phi \text{ \& } pn \rightarrow d\phi \rightarrow$ **large $g_{N^*N\phi}$**

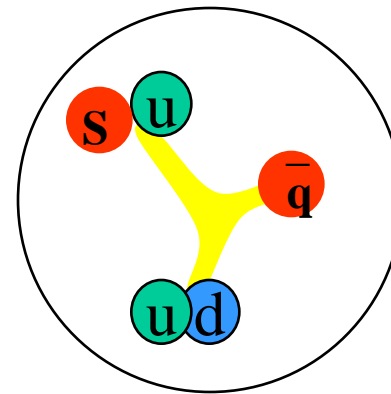
Xie, Zou & Chiang, PRC77(2008)015206; Cao, Xie, Zou & Xu, PRC80(2009)025203

New Scheme for $N^*(1535)$ and its $1/2^-$ nonet partners



$$\bar{q} \quad 1/2^+$$

$$\left. \begin{array}{l} [ud] \\ [ud] \end{array} \right\} L=1$$



$$\bar{q} \quad 1/2^-$$

$$\left. \begin{array}{l} [ud] \\ [us] \end{array} \right\} L=0$$

Zhang et al, hep-ph/0403210

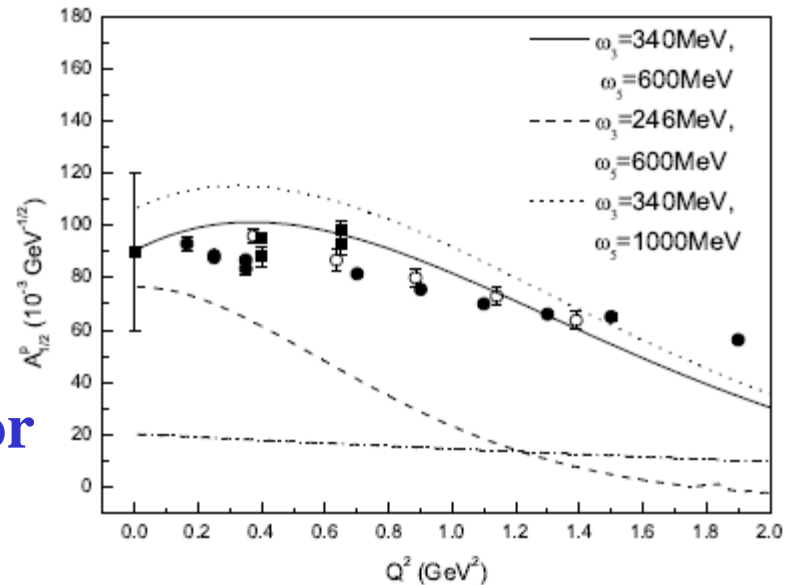
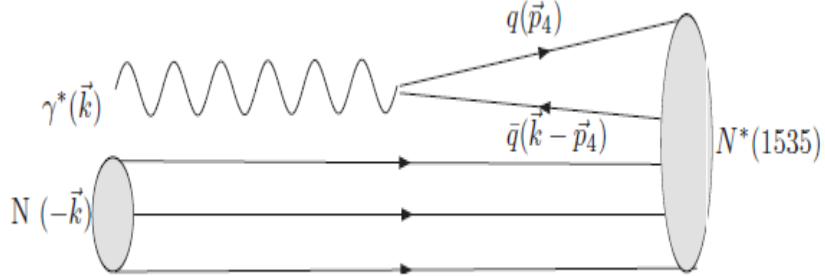
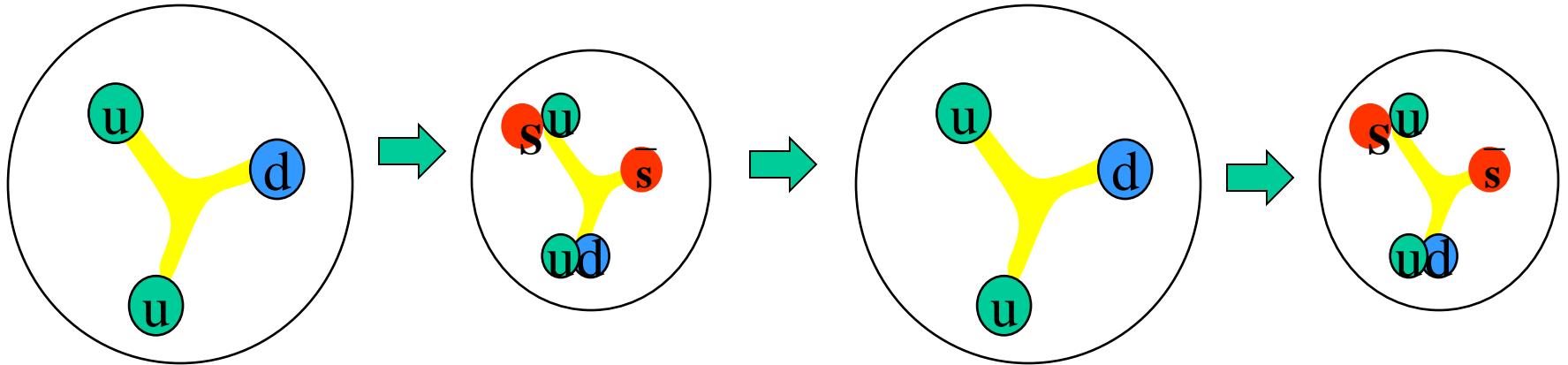
$$N^*(1535) \sim uud (L=1) + \varepsilon [ud][us] \bar{s} + \dots$$

$$N^*(1440) \sim uud (n=1) + \xi [ud][ud] \bar{d} + \dots$$

$$\Lambda^*(1405) \sim uds (L=1) + \varepsilon [ud][su] \bar{u} + \dots$$

$N^*(1535)$: $[ud][us] \bar{s} \rightarrow$ larger coupling to $N\eta$, $N\eta'$, $N\phi$ & $K\Lambda$, weaker to $N\pi$ & $K\Sigma$, and heavier !

The breathing mode for the $N^*(1535)$

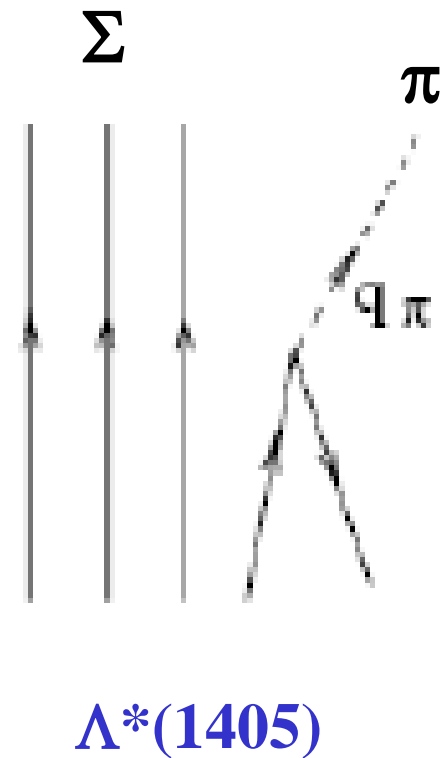
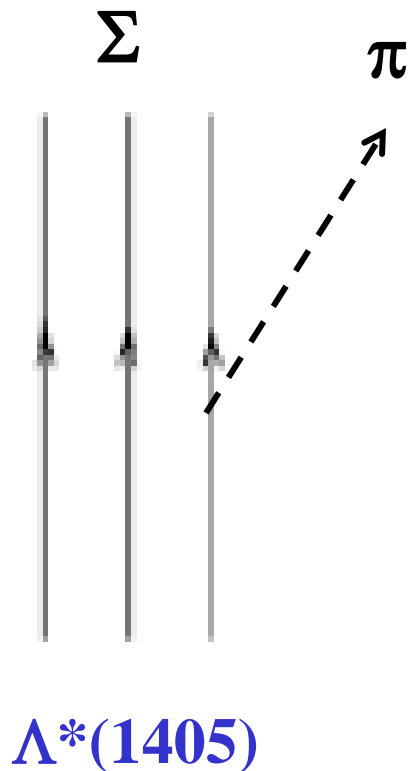


Important role for N^* EM form factor

An & Zou, EPJA39(2009)195

**50% 5q components in $\Lambda^*(1405)$
to reproduce $\Gamma(\Lambda^* \rightarrow \Sigma\pi) = 50 \text{ MeV}$**

An, Saghai, Yuan, He, PRC81(2010)045203



The new scheme for the $1/2^-$ nonet predicts:

$$\Lambda^* \quad [us][ds] \bar{s} \quad \sim \quad 1575 \text{ MeV}$$

$$\Sigma^* \quad [us][du] \bar{d} \quad \sim \quad 1360 \text{ MeV}$$

$$\Xi^* \quad [us][ds] \bar{u} \quad \sim \quad 1520 \text{ MeV}$$

Prediction of other unquenched models:

(1) 5-quark model Helminen & Riska, NPA699(2002)624

$$\Sigma^*(1/2^-) \sim \Lambda^*(1/2^-)$$

(2) K Λ -K Σ dynamics Weise, Oset et al.

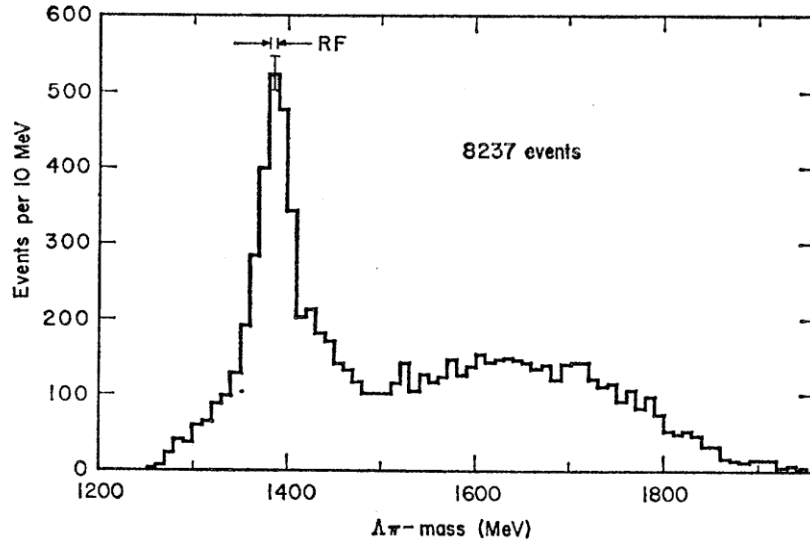
broad non-resonant $\Sigma^*(1/2^-)$ structure

Jido-Oset et al , NPA725(2003)181

Important to look for the $\Sigma^*(1/2^-)$ around 1380 MeV !

3. Evidence for the predicted $\Sigma^*(1/2^-)$

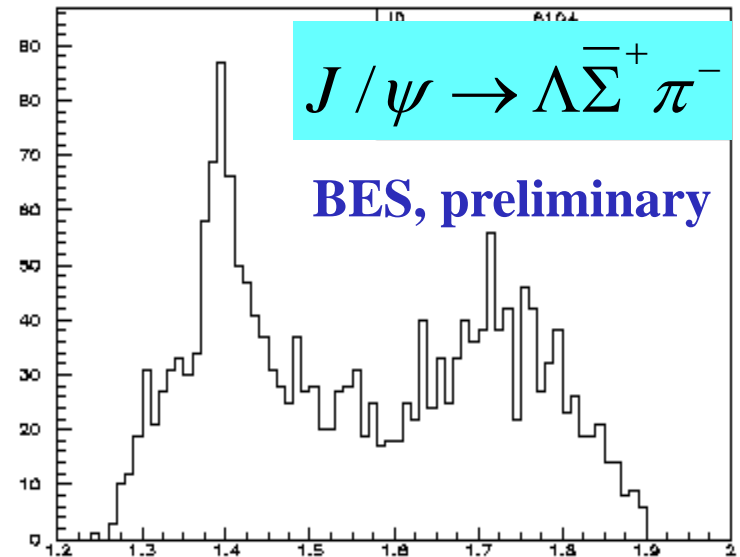
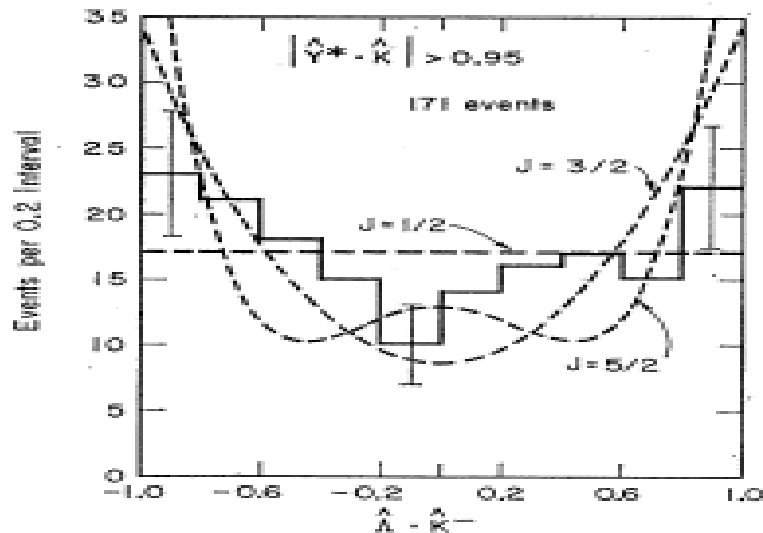
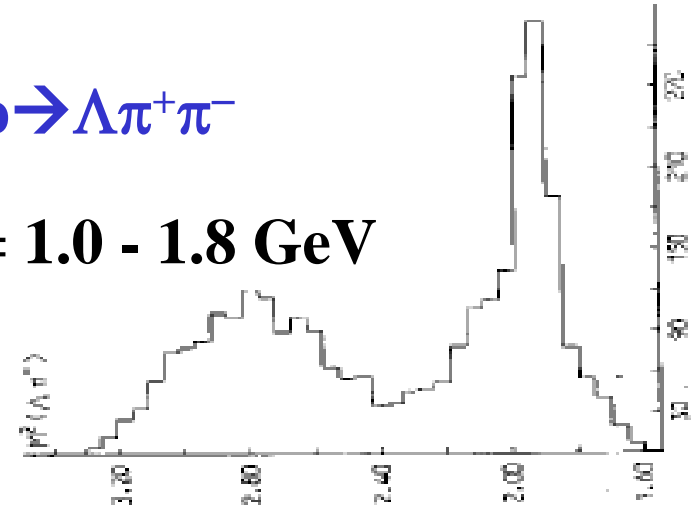
Huwe, PR181(1969)1824



Cameron et al., NPB143(1978)189

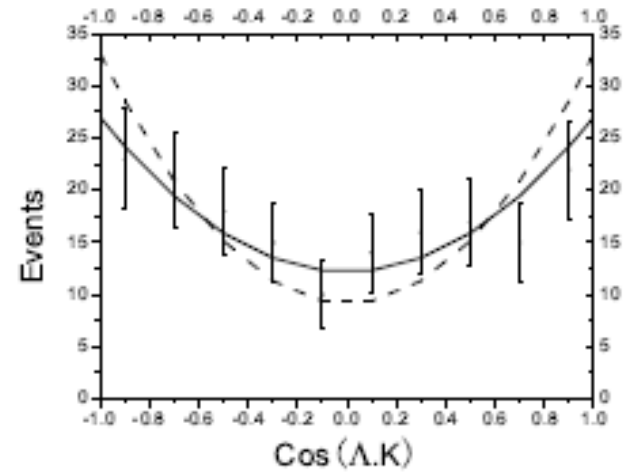
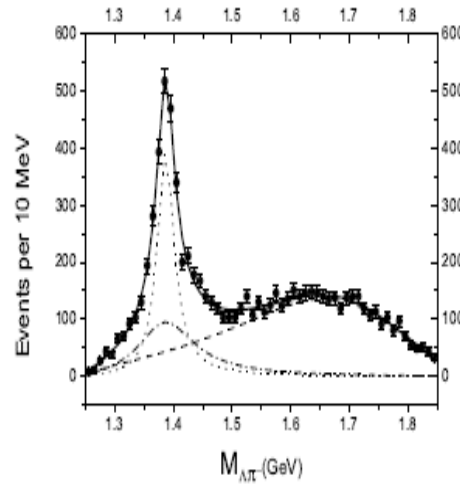
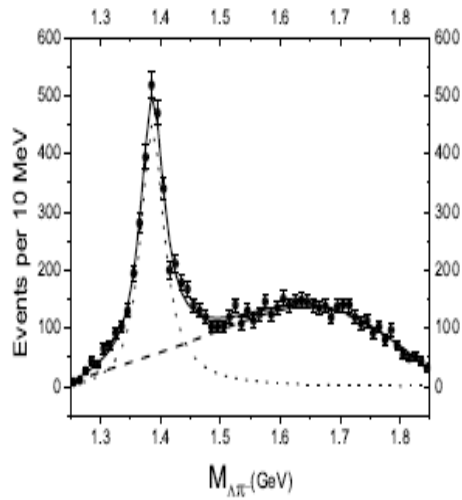
$K^- p \rightarrow \Lambda \pi^+ \pi^-$

$P_K = 1.0 - 1.8 \text{ GeV}$



BES, NSTAR04

$M_{\pi\Lambda}$



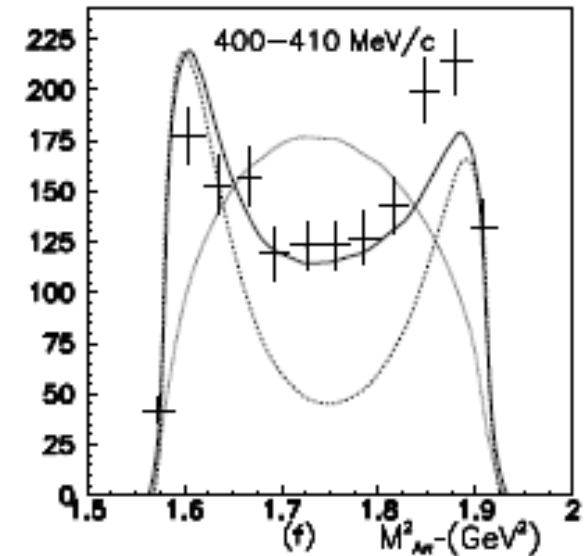
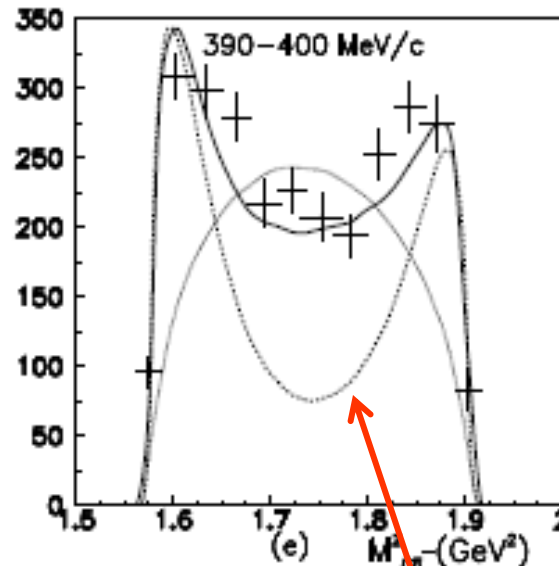
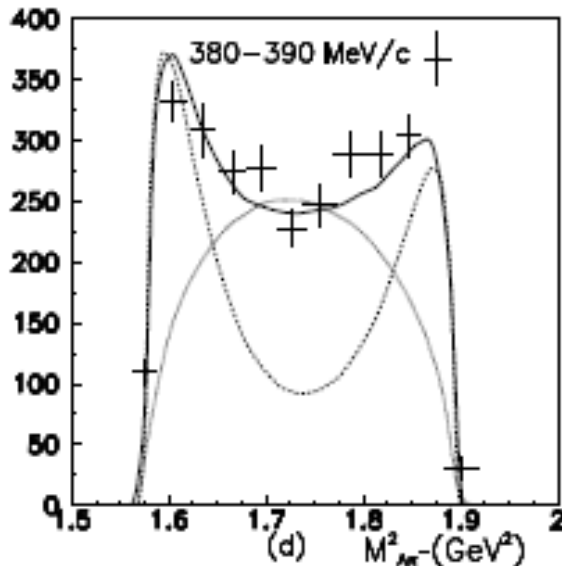
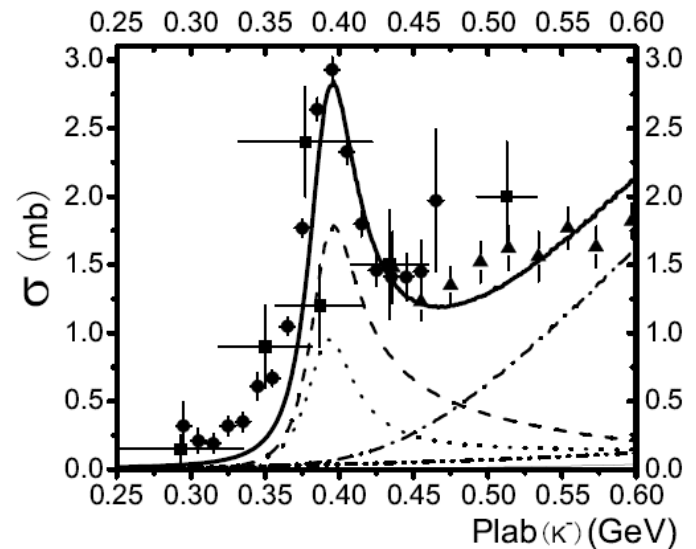
	$M_{\Sigma^*(3/2)}$	$\Gamma_{\Sigma^*(3/2)}$	$M_{\Sigma^*(1/2)}$	$\Gamma_{\Sigma^*(1/2)}$	χ^2/ndf (Fig.1)	χ^2/ndf (Fig.2)
Fit1	1385.3 ± 0.7	46.9 ± 2.5			68.5/54	10.1/9
Fit2	$1386.1^{+1.1}_{-0.9}$	$34.9^{+5.1}_{-4.9}$	$1381.3^{+4.9}_{-8.3}$	$118.6^{+55.2}_{-35.1}$	58.0/51	3.2/9

J.J.Wu, S.Dulat, B.S.Zou, PRD80 (2009) 017503

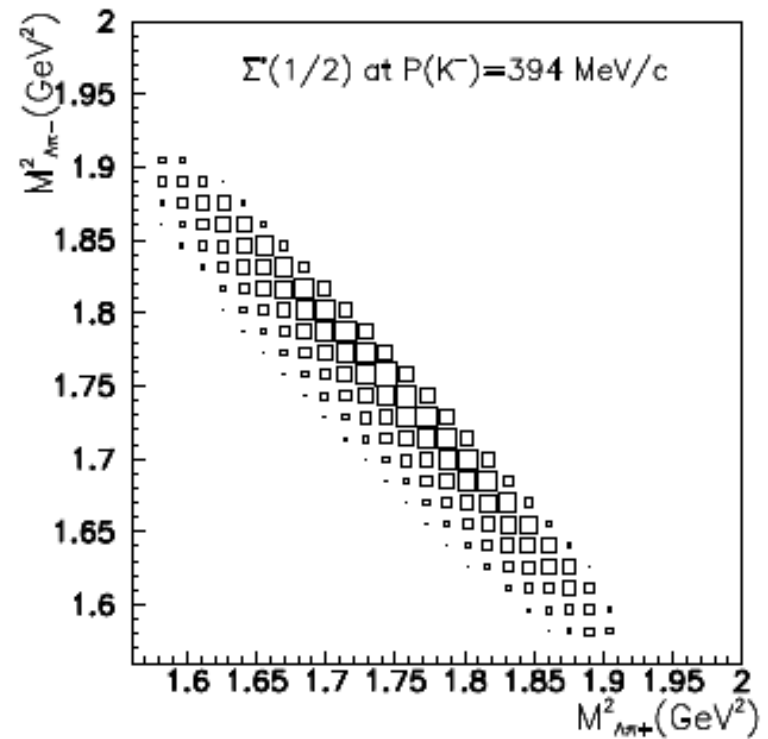
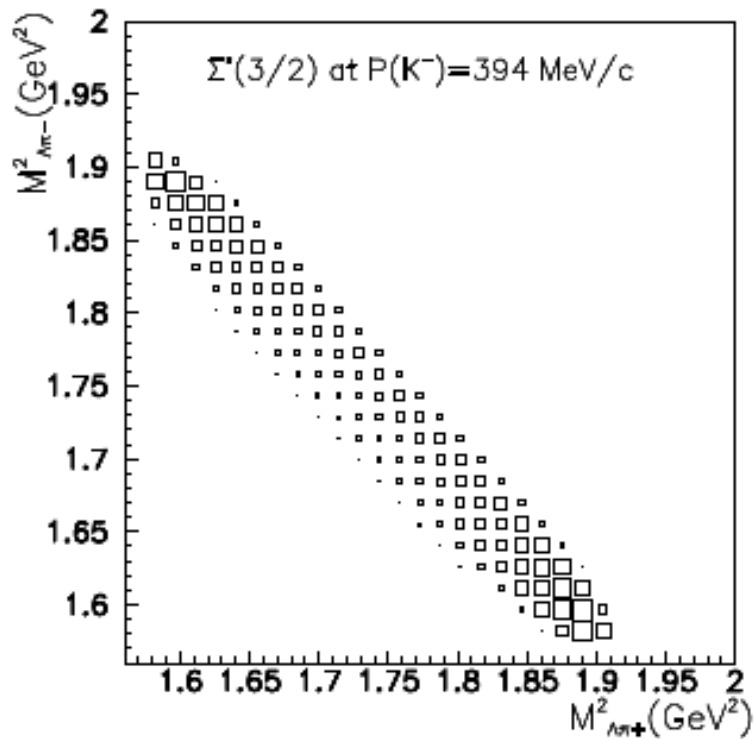
$$K^- p \rightarrow \Lambda^* \rightarrow \Sigma_{3/2}^{*-} \pi^+ \rightarrow \Lambda \pi^+ \pi^-$$

$$K^- p \rightarrow \Lambda^* \rightarrow \Sigma_{1/2}^{*-} \pi^+ \rightarrow \Lambda \pi^+ \pi^-$$

$$P_K \approx 0.4 \text{ GeV}$$



$\Sigma^*(3/2^+)$ only



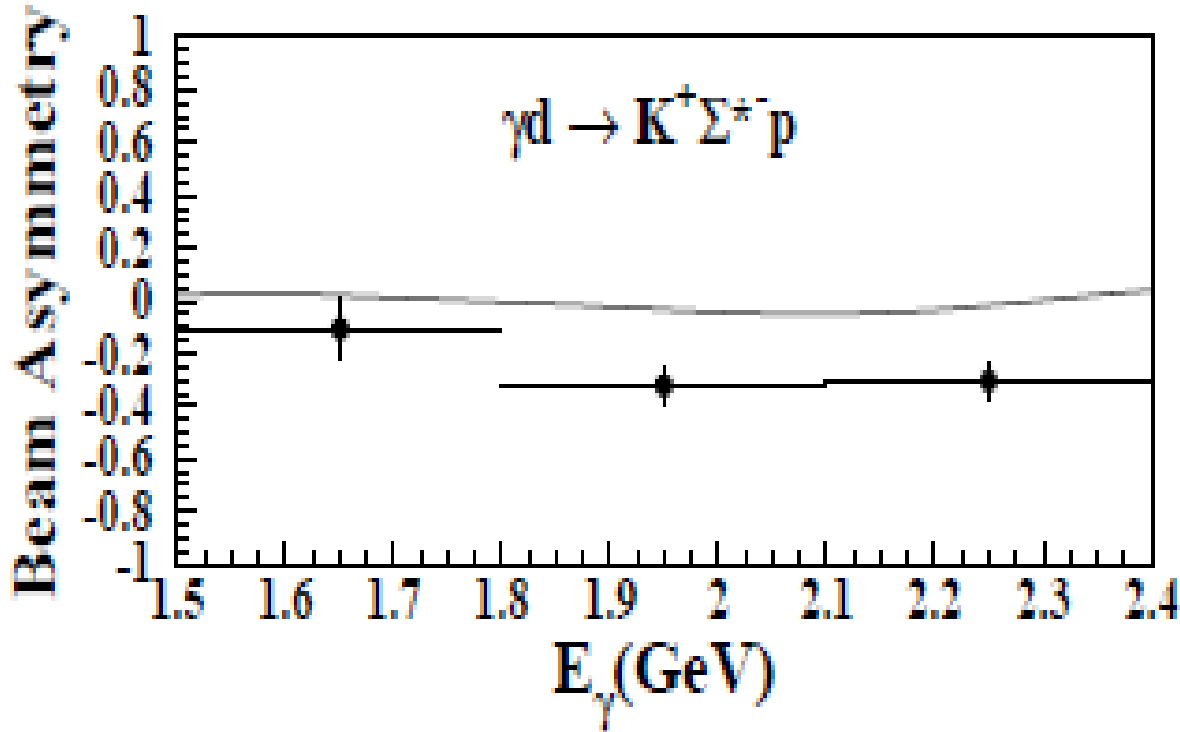
$\Sigma^*(3/2^+)$ & $\Sigma^*(1/2^-)$ \rightarrow different Dalitz plots & mass spectra

Both are needed to reproduce the data !

Other evidence: failed to reproduce data with $\Sigma^*(1385)$

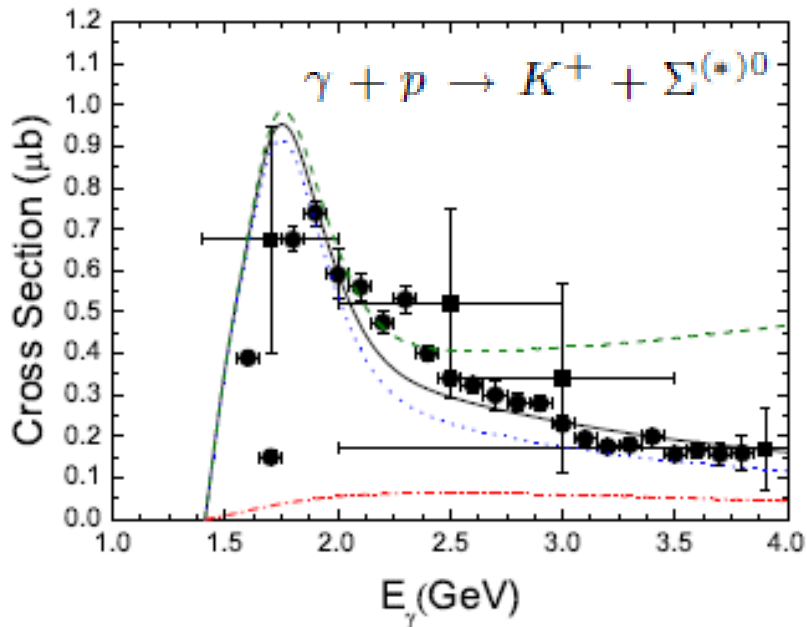
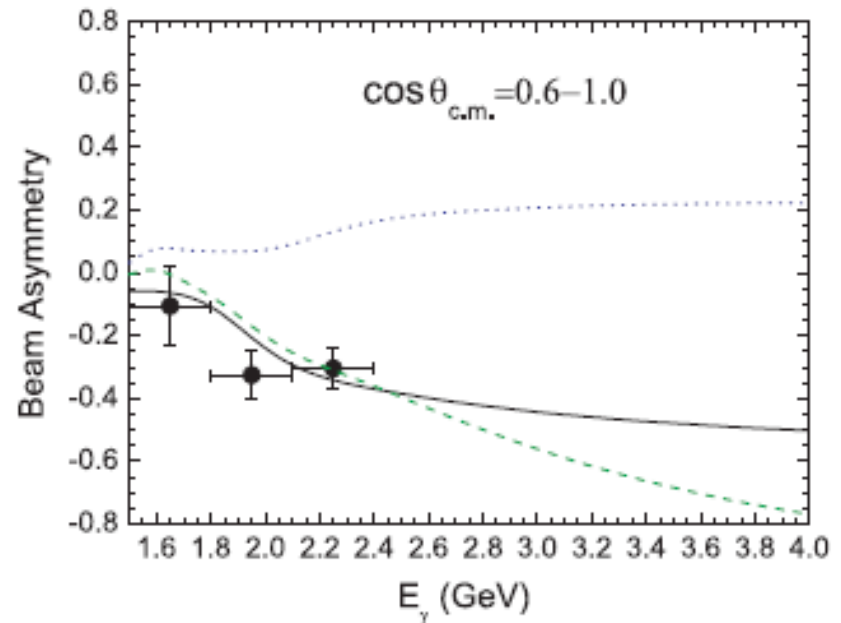
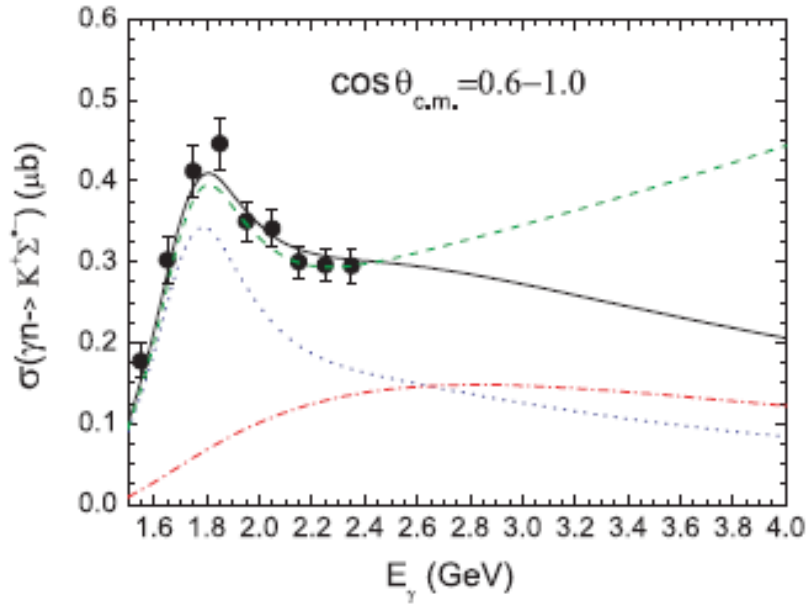
LEPS, PRL102(2009)012501

Y. Oh, C. M. Ko, and K. Nakayama, PRC77(2008) 045204



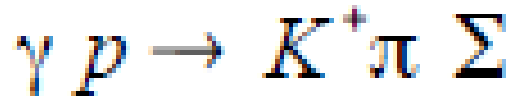
Something new ? $\Sigma^*(1/2^-)$?

P.Gao, J.J.Wu, B.S.Zou, Phys. Rev. C 81 (2010) 055203

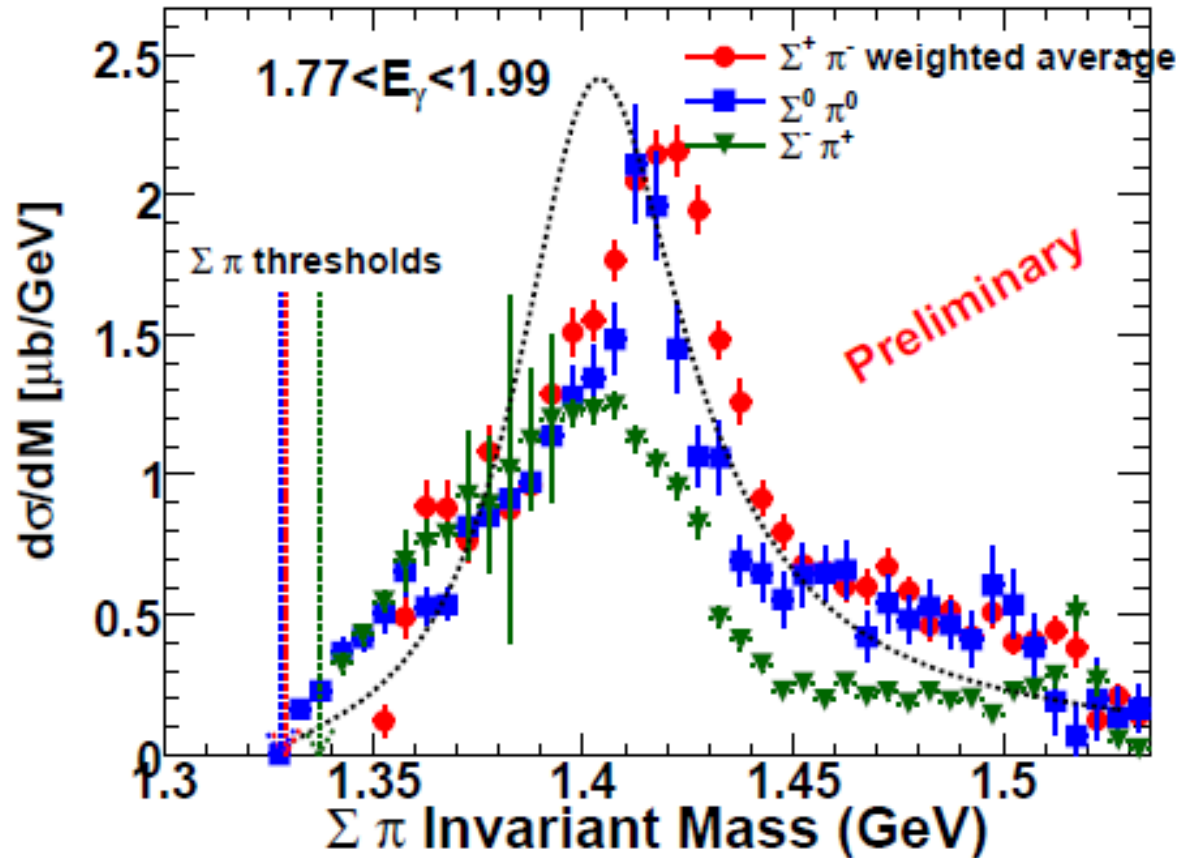


dot lines: $\Sigma^{*}(3/2^+)$ with $h=1.00$
dashed : $\Sigma^{*}(3/2^+)$ with $h=1.11$
solid: including $\Sigma^{*}(1/2^-)$

P.Gao, J.J.Wu, B.S.Zou,
Phys. Rev. C 81 (2010) 055203



R.Schumacher, K.Moriya



$J^P=1/2^-$ $I=1$ is needed besides $\Lambda^*(1405)$!

$$\frac{d\sigma(\pi^+\Sigma^-)}{dM_\gamma} \propto \frac{1}{2}|T^{(0)}|^2 + \frac{1}{3}|T^{(0)}|^2 + \frac{2}{\sqrt{6}}\text{Re}(T^{(0)}T^{(1)*}) + O(T^{(2)})$$

$$\frac{d\sigma(\pi^-\Sigma^+)}{dM_\gamma} \propto \frac{1}{2}|T^{(0)}|^2 + \frac{1}{3}|T^{(0)}|^2 - \frac{2}{\sqrt{6}}\text{Re}(T^{(0)}T^{(1)*}) + O(T^{(2)})$$

$$\frac{d\sigma(\pi^0\Sigma^0)}{dM_\gamma} \propto \frac{1}{3}|T^{(0)}|^2 + O(T^{(2)})$$

J/ψ decay

branching ratio * 10⁴

$\bar{p} \Delta(1232)^+$	3/2+	< 1	} SU(3) breaking
$\bar{\Sigma}^- \Sigma(1385)^+$		3.1 ± 0.5	
$\bar{\Xi}^+ \Xi(1530)^-$		5.9 ± 1.5	
$\bar{p} N^*(1535)^+$	1/2-	10 ± 3	} SU(3) allowed
$\bar{\Sigma}^- \Sigma(1360)^+$?	
$\bar{\Xi}^+ \Xi(1520)^-$?	

It is very important to check whether under the $\Sigma(1385)$ and $\Xi(1520)$ peaks there are $1/2^-$ components ?

4. Extension to hidden charm and beauty

4-quark components in mesons

$\bar{q}q \ ^3S_1$ nonet

$\phi(1020) \quad \bar{s}s$

$K(892) \quad \bar{s}d$

$\omega(782) \quad \bar{u}u + \bar{d}d$

$\rho(770) \quad \bar{u}u - \bar{d}d$

$\bar{q}q \ ^3P_0$ or \bar{q}^2q^2 nonet ?

$a_0(980) \quad \bar{u}u - \bar{d}d, \quad [\bar{u}s][us] - [\bar{d}s][ds]$

$f_0(980) \quad \bar{s}s, \quad [\bar{u}s][us] + [\bar{d}s][ds]$

$\kappa(800) \quad \bar{s}d, \quad [\bar{s}u][ud]$

$f_0(600) \quad \bar{u}u + \bar{d}d, \quad [\bar{u}d][ud]$

$D_{s0}^*(2317) \sim \bar{s}c (L=1) + [\bar{q}s][qc] + DK + \dots$

$D_{s1}^*(2460) \sim \bar{s}c (L=1) + D^*K + \dots$

$X(3872) \sim \bar{c}c (L=1) + [\bar{q}c][qc] + D^*D + \dots$

Many other proposed dynamically generated states

Problem:

None of them can be clearly distinguished from qqq or $\bar{q}q$ due to tunable ingredients and possible large mixing of various configurations

Solution: Extension to hidden charm and beauty for baryons

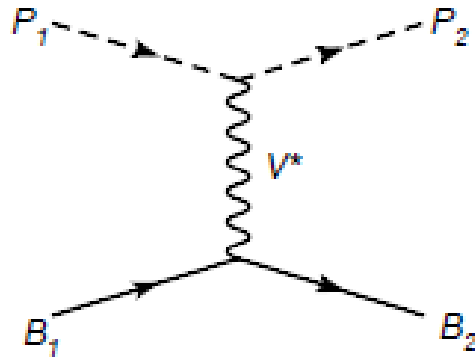
$N^*(1535)$ $\bar{s}suud$

$N^*(4260)$ $\bar{c}cuud$ J.J.Wu, R.Molina, E.Oset, B.S.Zou.
arXiv:1007.0573[nucl-th]

$N^*(11050)$ $\bar{b}buud$ J.J.Wu, B.S.Zou. to be submitted.

$K\Sigma, K\rho \rightarrow \bar{D}\Sigma_c, \bar{D}_s\Lambda_c \rightarrow B\Sigma_b, B_s\Lambda_b$ bound states

J.J.Wu, R.Molina, E.Oset, B.S.Zou. arXiv:1007.0573[nucl-th]



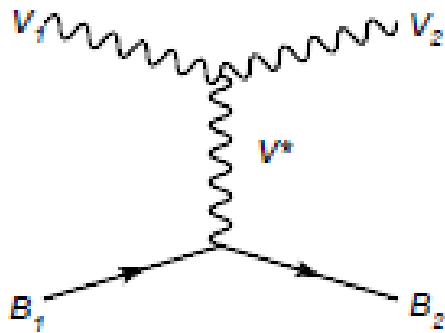
$$\mathcal{L}_{VVV} = ig\langle V^\mu [V^\nu, \partial_\mu V_\nu] \rangle$$

$$\mathcal{L}_{PPV} = -ig\langle V^\mu [P, \partial_\mu P] \rangle$$

$$\mathcal{L}_{BBV} = g(\langle \bar{B}\gamma_\mu [V^\mu, B] \rangle + \langle \bar{B}\gamma_\mu B \rangle \langle V^\mu \rangle)$$

$$V_{ab}(P_1 B_1 \rightarrow P_2 B_2) = \frac{C_{ab}}{4f^2} (E_{P_1} + E_{P_2}),$$

$$V_{ab}(V_1 B_1 \rightarrow V_2 B_2) = \frac{C_{ab}}{4f^2} (E_{V_1} + E_{V_2}) \vec{\epsilon}_1 \cdot \vec{\epsilon}_2,$$



$$T = [1 - VG]^{-1}V$$

$$T_{ab} = \frac{g_a g_b}{\sqrt{s} - z_R}$$

	(I, S)	z_R (MeV)	g_a		
N^*	$(1/2, 0)$		$\bar{D}\Sigma_c$	$\bar{D}\Lambda_c^+$	
		4269	2.85	0	
Λ^*	$(0, -1)$		$\bar{D}_s\Lambda_c^+$	$\bar{D}\Xi_c$	$\bar{D}\Xi'_c$
		4213	1.37	3.25	0
		4403	0	0	2.64

TABLE III: Pole positions z_R and coupling constants g_a for the states from $PB \rightarrow PB$.

	(I, S)	z_R (MeV)	g_a		
N^*	$(1/2, 0)$		$\bar{D}^*\Sigma_c$	$\bar{D}^*\Lambda_c^+$	
		4418	2.75	0	
Λ^*	$(0, -1)$		$\bar{D}_s^*\Lambda_c^+$	$\bar{D}^*\Xi_c$	$\bar{D}^*\Xi'_c$
		4370	1.23	3.14	0
		4550	0	0	2.53

TABLE IV: Pole position and coupling constants for the bound states from $VB \rightarrow VB$.

	(I, S)	M	Γ	Γ_i					
N^*	$(1/2, 0)$			πN	ηN	$\eta' N$	$K\Sigma$	$\eta_c N$	
		4261	56.9	3.8	8.1	3.9	17.0	23.4	
Λ^*	$(0, -1)$			$K N$	$\pi\Sigma$	$\eta\Lambda$	$\eta'\Lambda$	$K\Xi$	$\eta_c\Lambda$
		4209	32.4	15.8	2.9	3.2	1.7	2.4	5.8
		4394	43.3	0	10.6	7.1	3.3	5.8	16.3

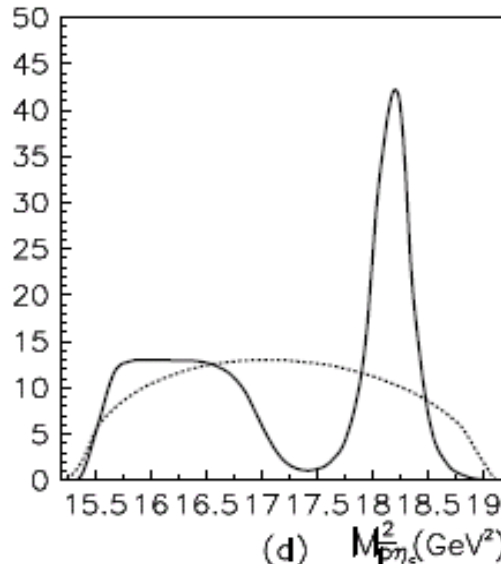
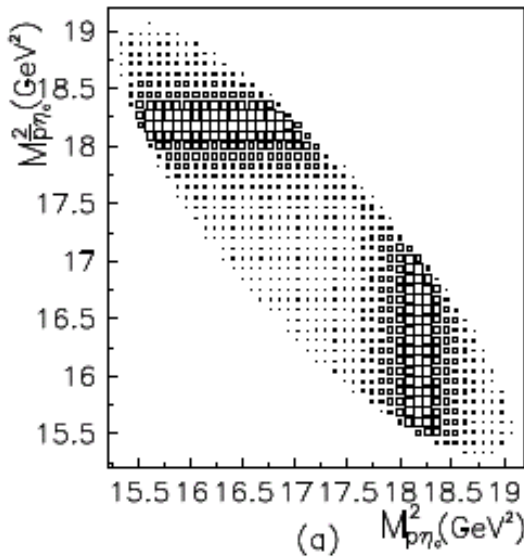
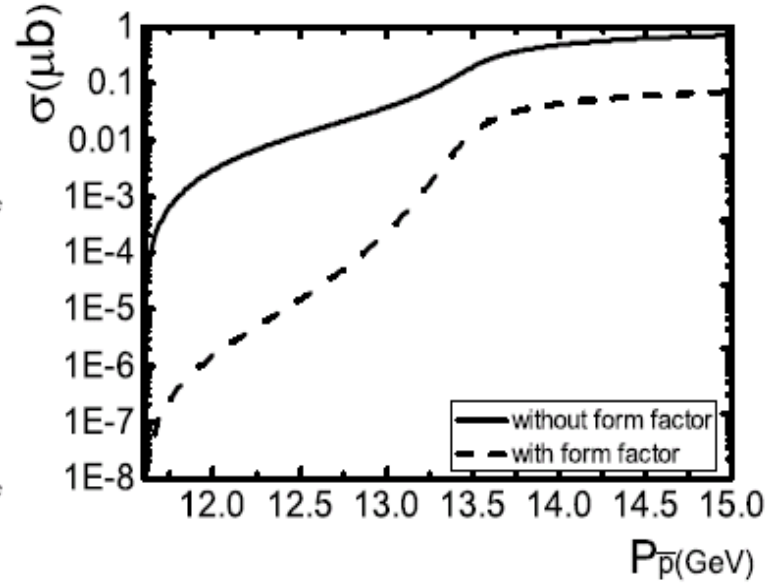
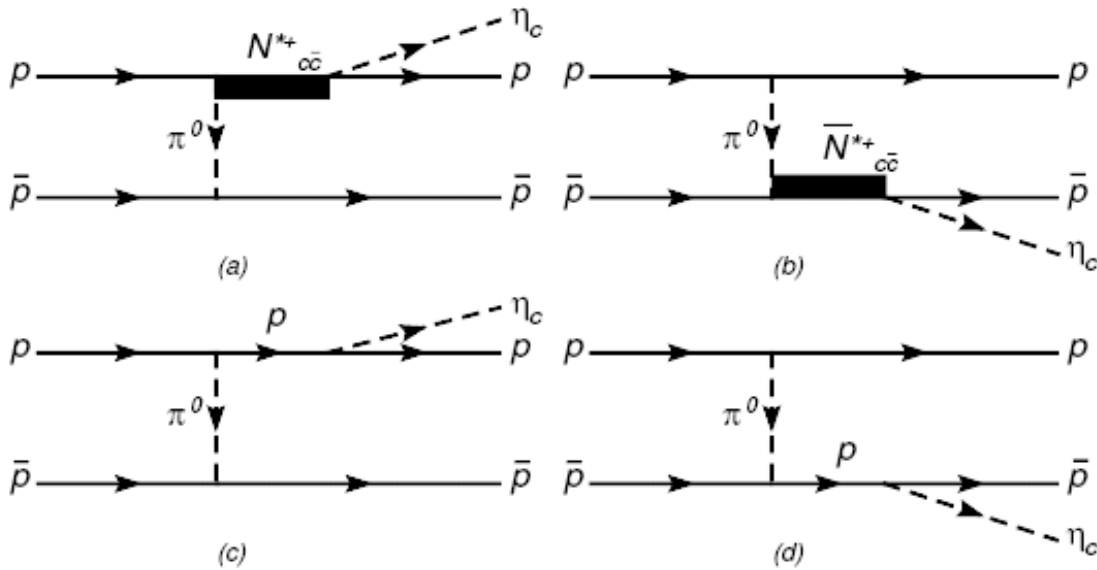
TABLE V: Mass (M), total width (Γ), and the partial decay width (Γ_i) for the states from $PB \rightarrow PB$, with units in MeV.

	(I, S)	M	Γ	Γ_i					
N^*	$(1/2, 0)$			ρN	ωN	$K^*\Sigma$		$J/\psi N$	
		4412	47.3	3.2	10.4	13.7		19.2	
Λ^*	$(0, -1)$			$K^* N$	$\rho\Sigma$	$\omega\Lambda$	$\phi\Lambda$	$K^*\Xi$	$J/\psi\Lambda$
		4368	28.0	13.9	3.1	0.3	4.0	1.8	5.4
		4544	36.6	0	8.8	9.1	0	5.0	13.8

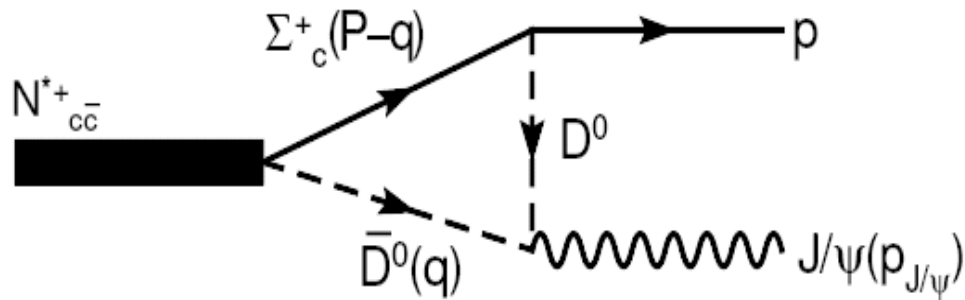
TABLE VI: Mass (M), total width (Γ), and the partial decay width (Γ_i) for the states from $VB \rightarrow VB$ with units in MeV.

Super-heavy narrow N^* and Λ^* with hidden charm !
Definitely not qqq states !

Prediction for PANDA



$\bar{p} p \rightarrow \bar{p} p \eta_c$
0.07 -- 0.7 μb



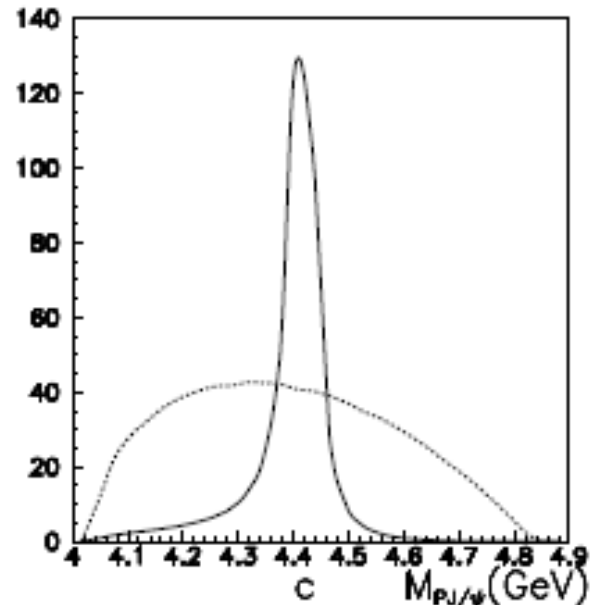
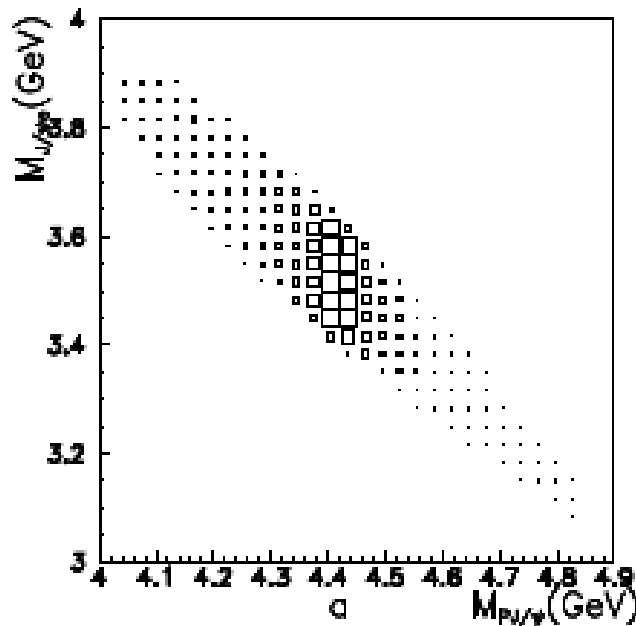
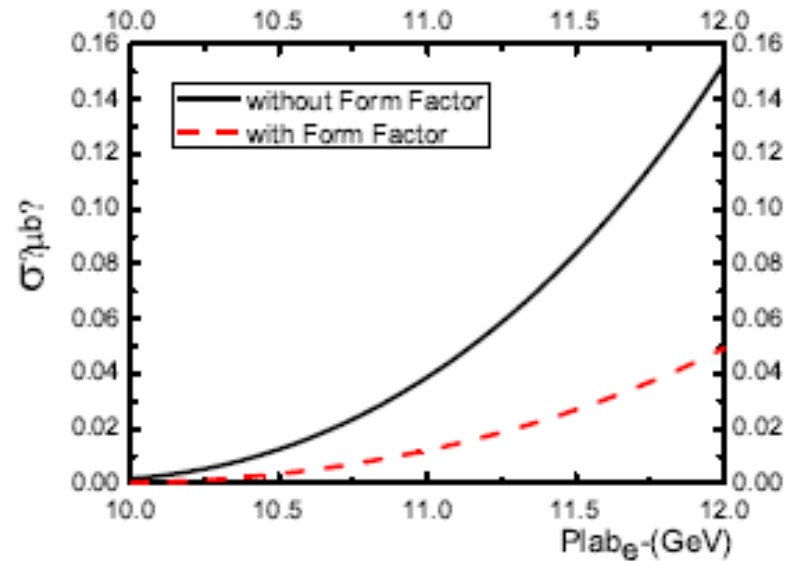
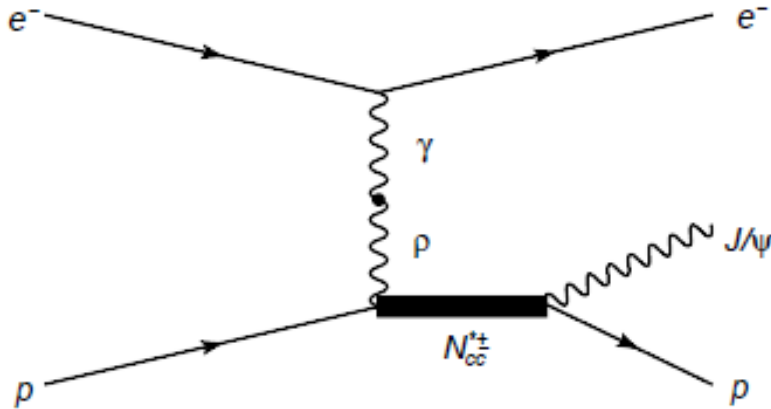
3 orders of magnitude smaller than $N^* \rightarrow p\eta_c$

$$\bar{p}p \rightarrow \bar{p}pJ/\psi \sim 0.03 \text{ nb}$$

~ 250 events per day at PANDA/FAIR by $L=10^{31} \text{ cm}^{-2}\text{s}^{-1}$

These Super-heavy narrow N^* and Λ^* can be found at PANDA !

Prediction for 12GeV@JLab



Conclusion I

- **Meson-cloud vs diquark cluster for $\bar{d} - \bar{u} \sim 0.12$**
 - **Predictions for the strangeness in the proton:**
 - meson cloud : $\Delta s < 0$, $\mu_s < 0$, $r_s < 0$
 - diquark cluster : $\Delta s < 0$, $\mu_s > 0$, $r_s > 0$
 - **$\bar{q}qqqq$ in S-state more favorable than qqq with $L=1$!
& $qqqq$ in S-state more favorable than qq with $L=1$!**
 - $1/2^-$ baryon nonet $\sim \bar{q}q^2q^2$ state + ...**
 - 0^+ meson octet $\sim \bar{q}^2q^2$ state + ...**
- multiquark components are important for hadrons!**

Conclusion II

- Quenched quark models and unquenched models give very distinctive predictions for $\Sigma^*(1/2^-)$;
- Possible existence of a $\Sigma^*(1/2^-)$ around 1380 MeV: evidence needs confirmation ;
relevant to **Kp, Kpp interactions or bound states**
- It should be checked by forthcoming experiments :

$K^- p \rightarrow \pi \Sigma^*, \Sigma^* \rightarrow \Lambda \pi, \Sigma \pi$ @ JPARC

$\gamma N \rightarrow K^+ \Sigma^*, \Sigma^* \rightarrow \Lambda \pi, \Sigma \pi$ @ JLab, Spring-8, ELSA

$\psi \rightarrow \bar{\Sigma} \Sigma^*, \Sigma^* \rightarrow \Lambda \pi, \Sigma \pi$ @ BESIII

Conclusion III

- Super heavy narrow N^* and Λ^* are predicted to exist
 $\bar{D}\Sigma_c, \bar{D}_s\Lambda_c \rightarrow B\Sigma_b, B_s\Lambda_b$ bound states
- They are definitely not qqq baryons
- They can be looked for at 12GeV@Jlab and PANDA
maybe also at RHIC , EIC?