

The $\Lambda(1405)$ Resonance in Full-QCD

Benjamin J. Menadue Waseem Kamleh
Derek B. Leinweber M. S. Mahbub

Special Research Centre for the Subatomic Structure of Matter
School of Chemistry & Physics, University of Adelaide, Australia, 5005

8th Circum-Pan-Pacific Symposium
on High Energy Spin Physics

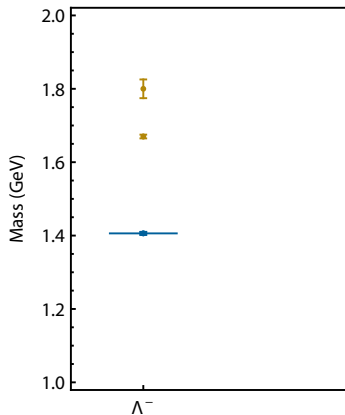


THE UNIVERSITY
of ADELAIDE



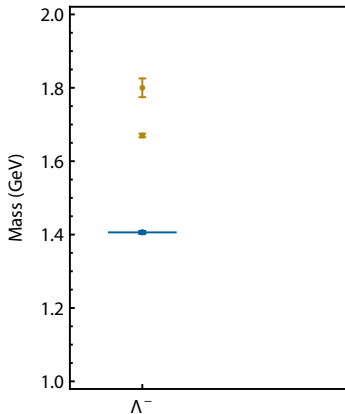
The $\Lambda(1405)$

- The negative-parity ground state of the Lambda has a mass of 1406 ± 4 MeV.



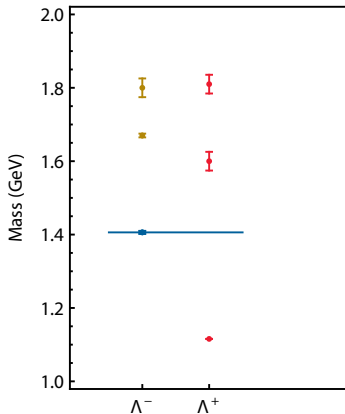
The $\Lambda(1405)$

- The negative-parity ground state of the Lambda has a mass of 1406 ± 4 MeV.
- Such a low mass is puzzling:



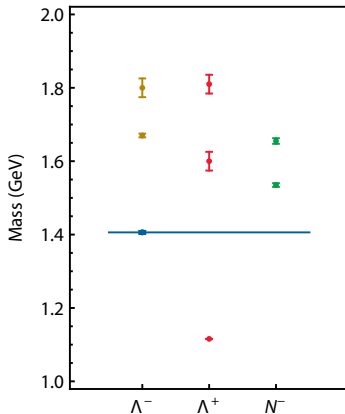
The $\Lambda(1405)$

- The negative-parity ground state of the Lambda has a mass of 1406 ± 4 MeV.
- Such a low mass is puzzling:
 - lies lower than the $\Lambda(1600)$, but has negative parity.



The $\Lambda(1405)$

- The negative-parity ground state of the Lambda has a mass of 1406 ± 4 MeV.
- Such a low mass is puzzling:
 - lies lower than the $\Lambda(1600)$, but has negative parity.
 - lies lower than the $N(1535)$, but has a valence strange quark.



The $\Lambda(1405)$

- Lattice QCD has so far been unable to isolate such a low-lying state.

The $\Lambda(1405)$

- Lattice QCD has so far been unable to isolate such a low-lying state.
- The Roper resonance of the nucleon is also abnormally low-lying, and Lattice QCD has had similar trouble in isolating it.

The $\Lambda(1405)$

- Lattice QCD has so far been unable to isolate such a low-lying state.
- The Roper resonance of the nucleon is also abnormally low-lying, and Lattice QCD has had similar trouble in isolating it.
- The CSSM Lattice Collaboration has developed a technique that has successfully isolated the Roper.

M. Selim Mahbub, et al., arXiv:1011.5724

M. S. Mahbub, et al., PoS Lattice 2010, 112, arXiv:1011.0480

The $\Lambda(1405)$

- We apply the same techniques to the Lambda in an attempt to isolate the $\Lambda(1405)$.

The $\Lambda(1405)$

- We apply the same techniques to the Lambda in an attempt to isolate the $\Lambda(1405)$.
- Last year we showed that such an analysis is necessary to isolate this state.

BM, et al., AIP Conf. Proc., **1354**, 213 (2011), arXiv:1102.3492

Outline

1 Technique

- Variational Analysis
- Lattice Details

2 Results

- Common Interpolator Results
- Octet & Singlet Flavour-Symmetry Results

3 Conclusions

Variational Analysis

- 1 Construct a correlation matrix of cross-correlation functions from various interpolating operators.

Variational Analysis

- 1 Construct a correlation matrix of cross-correlation functions from various interpolating operators.
- 2 Use eigenanalysis to project out correlation functions for individual baryon states.

Variational Analysis

- 1 Construct a correlation matrix of cross-correlation functions from various interpolating operators.
- 2 Use eigenanalysis to project out correlation functions for individual baryon states.
- 3 Analyse these projected correlation functions using normal techniques.

Correlation Matrices

- Take a set of N operators $\chi_i(\mathbf{x}, t)$ that couple to the baryon we interested in.

Correlation Matrices

- Take a set of N operators $\chi_i(\mathbf{x}, t)$ that couple to the baryon we interested in.
- Calculate the $N \times N$ matrix of zero-momentum, parity-projected cross correlation functions from these operators:

$$\begin{aligned} G_{ij}^{\pm}(t) &= \sum_{\mathbf{x}} \text{tr} \left(\Gamma_{\pm} \langle \Omega | \chi_i(\mathbf{x}, t) \bar{\chi}_j(\mathbf{x}, 0) | \Omega \rangle \right) \\ &= \sum_{\alpha=1}^N \lambda_i^{\alpha} \bar{\lambda}_j^{\alpha} e^{-m_{\alpha} t}. \end{aligned}$$

Correlation Matrices

- Take a set of N operators $\chi_i(\mathbf{x}, t)$ that couple to the baryon we interested in.
- Calculate the $N \times N$ matrix of zero-momentum, parity-projected cross correlation functions from these operators:

$$\begin{aligned}
 G_{ij}^{\pm}(t) &= \sum_{\mathbf{x}} \text{tr} \left(\Gamma_{\pm} \langle \Omega | \chi_i(\mathbf{x}, t) \bar{\chi}_j(\mathbf{x}, 0) | \Omega \rangle \right) \\
 &= \sum_{\alpha=1}^N \lambda_i^{\alpha} \bar{\lambda}_j^{\alpha} e^{-m_{\alpha} t}.
 \end{aligned}$$

- Construct a set of N “perfect” operators $\varphi_{\alpha}(\mathbf{x}, t)$ that completely isolate the N lowest states from linear combinations of our original operators:

$$\chi_i = \sum_{\alpha} v_i^{\alpha} \varphi_{\alpha} \quad \text{and} \quad \bar{\chi}_j = \sum_{\alpha} u_j^{\alpha} \varphi_{\alpha}.$$

Correlation Matrices

- With a bit of algebraic manipulation, can show that u^α and v^α are the right- and left-eigenvectors of $G^\pm(t)^{-1}G^\pm(t + \Delta t)$, with eigenvalue $e^{-m_\alpha \Delta t}$.

Correlation Matrices

- With a bit of algebraic manipulation, can show that u^α and v^α are the right- and left-eigenvectors of $G^\pm(t)^{-1}G^\pm(t + \Delta t)$, with eigenvalue $e^{-m_\alpha \Delta t}$.
- Moreover, the quadratic form $v^\alpha G^\pm(t)u^\beta \propto \delta^{\alpha\beta} e^{-m_\alpha t}$ has t -dependence only in a single exponential term.

Correlation Matrices

- With a bit of algebraic manipulation, can show that u^α and v^α are the right- and left-eigenvectors of $G^\pm(t)^{-1}G^\pm(t + \Delta t)$, with eigenvalue $e^{-m_\alpha \Delta t}$.
- Moreover, the quadratic form $v^\alpha G^\pm(t)u^\beta \propto \delta^{\alpha\beta} e^{-m_\alpha t}$ has t -dependence only in a single exponential term.
- Hence, defining $G_\alpha^\pm(t) := v^\alpha G^\pm(t)u^\alpha$, can extract the masses m_α using the usual

$$m_\alpha = \ln \left(\frac{G_\alpha^\pm(t)}{G_\alpha^\pm(t+1)} \right).$$

Available Operators

- There are quite a few operators for the Lambda baryon.

Available Operators

- There are quite a few operators for the Lambda baryon.
- Flavour symmetry gives three operators:

- Octet:

$$\chi_i^8 = \frac{1}{\sqrt{6}} \varepsilon_{abc} (2(u_a^T A_i d_b) B_i s_c + (u_a^T A_i s_b) B_i d_c - (d_a^T A_i s_b) B_i u_c),$$

- Singlet:

$$\chi_i^1 = -2\varepsilon_{abc} (-(u_a^T C \gamma_5 d_b) s_c + (u_a^T C \gamma_5 s_b) d_c - (d_a^T C \gamma_5 s_b) u_c),$$

- Common:

$$\chi_i^c = \frac{1}{\sqrt{2}} \varepsilon_{abc} ((u_a^T A_i s_b) B_i d_c - (d_a^T A_i s_b) B_i u_c),$$

- The i -index indicates the spin-structure of the operator:

$$(A_1, B_1) = (C \gamma_5, I), \quad (A_2, B_2) = (C, \gamma_5), \text{ and} \\ (A_4, B_4) = (C \gamma_5 \gamma_4, I).$$

Available Operators

- Similarly to the CSSM Collaboration's investigation of the Roper, we use gauge-invariant Gaussian smearing of the source and sink to further increase our operator basis.
 - We use 16, 35, 100, and 200 sweeps of smearing.

Available Operators

- Similarly to the CSSM Collaboration's investigation of the Roper, we use gauge-invariant Gaussian smearing of the source and sink to further increase our operator basis.
 - We use 16, 35, 100, and 200 sweeps of smearing.
- This gives us a total of 28 available operators.
 - There will not be enough signal to extract 28 states, but smaller subsets should give useful isolation of the lowest states.

Lattice Details

- We use the PACS-CS (2 + 1)-flavour full-QCD lattices, available from the ILDG.

PACS-CS Collaboration, Phys. Rev. D, **79**, 034503 (2009), arXiv:0807.1661

- Lattice size is $32^3 \times 64$, with a lattice spacing of 0.0907(13) fm.
- There are 5 light quark masses, with the strange quark mass held fixed.
- These correspond to pion masses of 623.40 ± 0.75 MeV down to 170.7 ± 2.1 MeV.

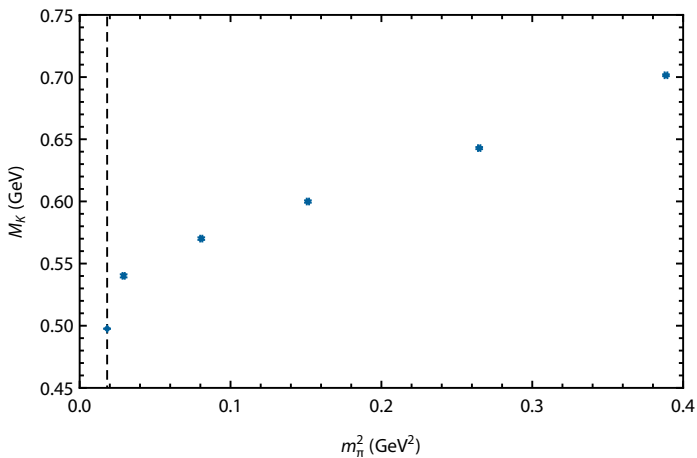
Lattice Details

- We use the PACS-CS (2 + 1)-flavour full-QCD lattices, available from the ILDG.

PACS-CS Collaboration, Phys. Rev. D, **79**, 034503 (2009), arXiv:0807.1661

- Lattice size is $32^3 \times 64$, with a lattice spacing of 0.0907(13) fm.
 - There are 5 light quark masses, with the strange quark mass held fixed.
 - These correspond to pion masses of 623.40 ± 0.75 MeV down to 170.7 ± 2.1 MeV.
- **PROBLEM:** The strange quark is too heavy!

Kaon Mass



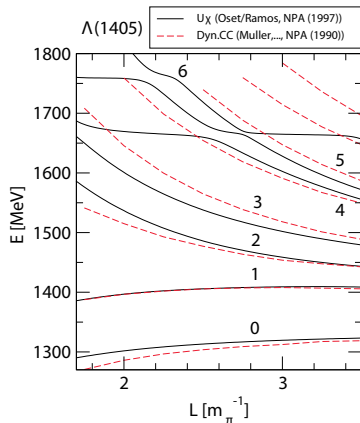
Implications of a Heavy Strange Quark

- Won't see energies of states moving directly into their physical values.

Implications of a Heavy Strange Quark

- Won't see energies of states moving directly into their physical values.
- However, don't have the avoided level crossings that are present in, e.g., the Roper resonance.

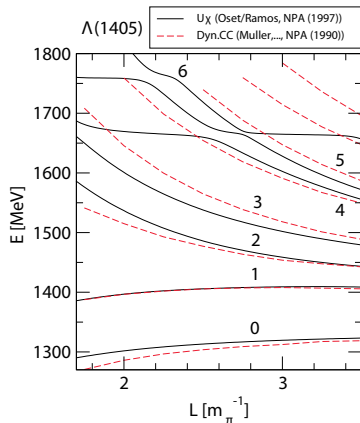
M. Döring, et al., NSTAR 2011



Implications of a Heavy Strange Quark

- Won't see energies of states moving directly into their physical values.
- However, don't have the avoided level crossings that are present in, e.g., the Roper resonance.
- We instead look for the correct arrangement of the states.

M. Döring, et al., NSTAR 2011



Results from Common Interpolating Fields

BM, Waseem Kamleh, Derek B. Leinweber, M. Selim Mahbub, in preparation

- We use the common interpolating fields χ_i^c to initially investigate the spectrum.
 - Gives an idea about how well our analysis can extract the low-lying states, without making assumptions about the flavour-symmetry properties.

Results from Common Interpolating Fields

BM, Waseem Kamleh, Derek B. Leinweber, M. Selim Mahbub, in preparation

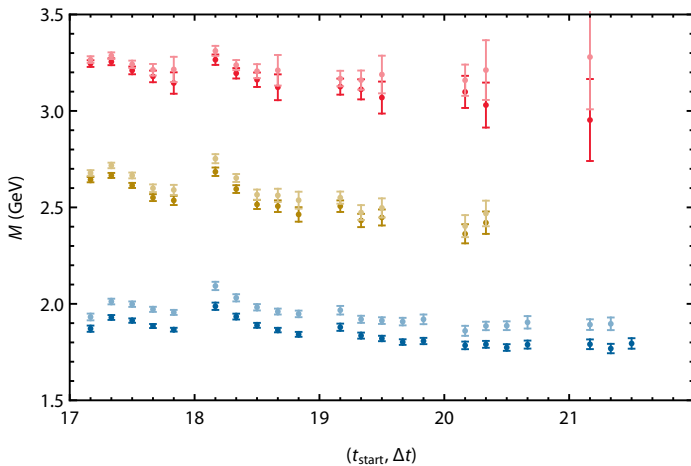
- We use the common interpolating fields χ_i^c to initially investigate the spectrum.
 - Gives an idea about how well our analysis can extract the low-lying states, without making assumptions about the flavour-symmetry properties.
- Need to determine which t_{start} and Δt , and which matrix basis to use.

Results from Common Interpolating Fields

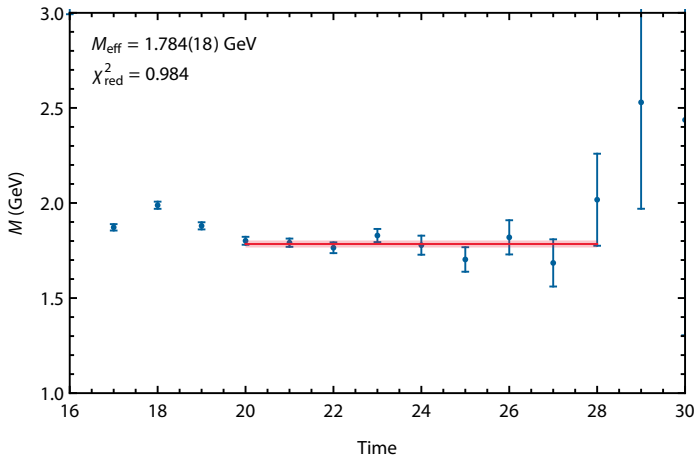
BM, Waseem Kamleh, Derek B. Leinweber, M. Selim Mahbub, in preparation

- We use the common interpolating fields χ_i^c to initially investigate the spectrum.
 - Gives an idea about how well our analysis can extract the low-lying states, without making assumptions about the flavour-symmetry properties.
- Need to determine which t_{start} and Δt , and which matrix basis to use.
 - Pick $\kappa_{u,d} = 0.13727$ ($m_\pi = 514.6 \pm 0.7$ MeV).
 - Use 6×6 basis from χ_1^c and χ_2^c with 16, 100, and 200 sweeps of smearing to investigate t_{start} and Δt .

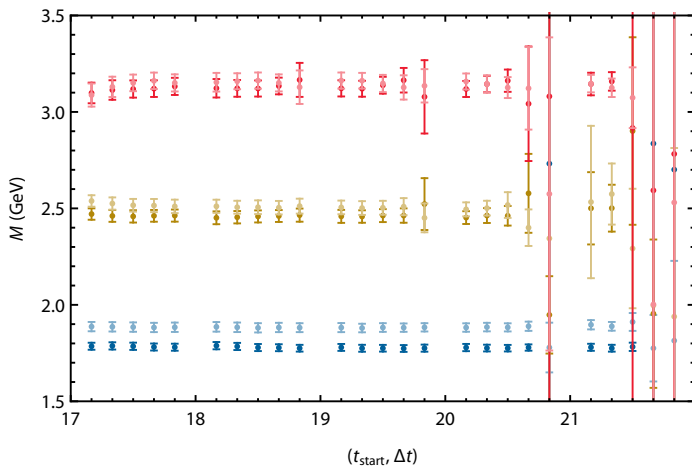
Eigenvalues



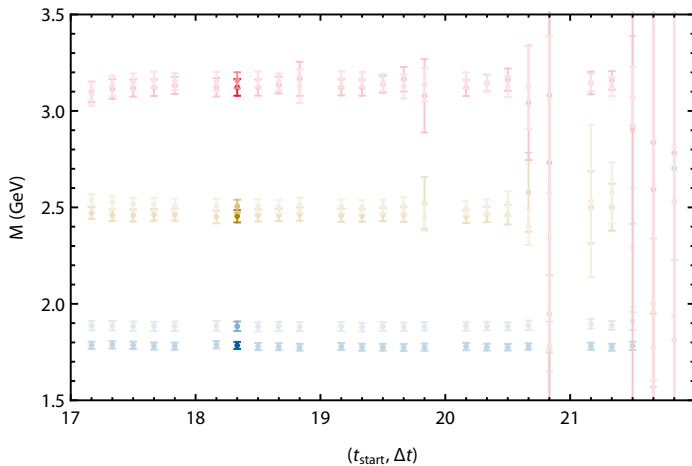
Ground State Effective Mass



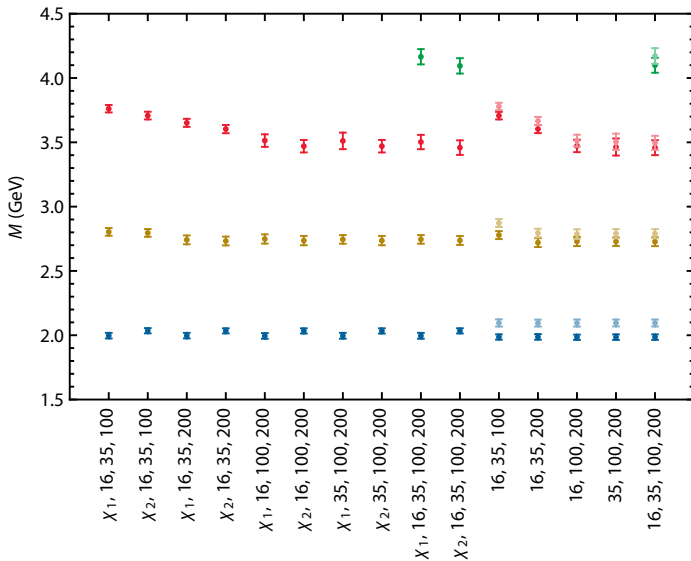
Fits



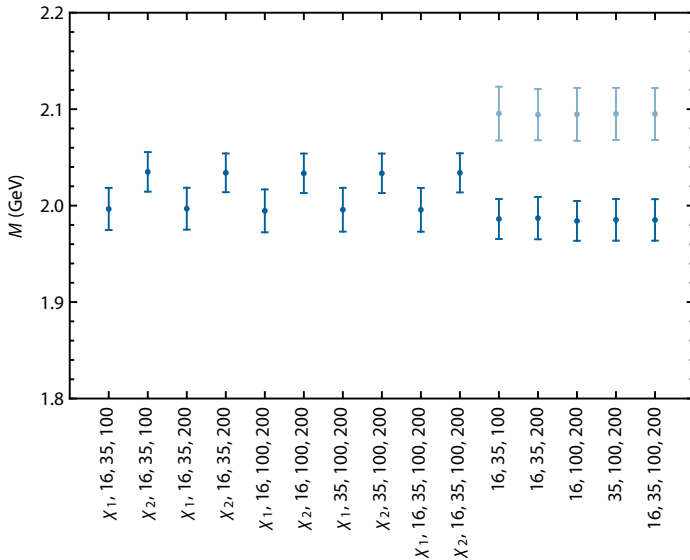
Fits



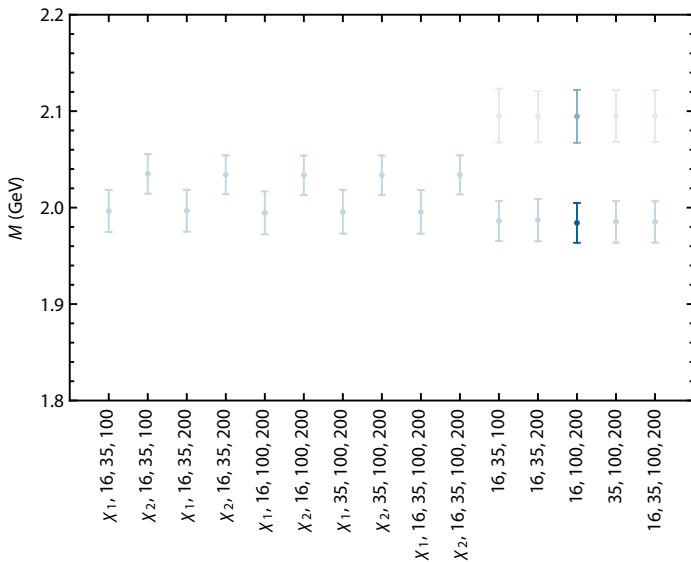
Comparison of Bases



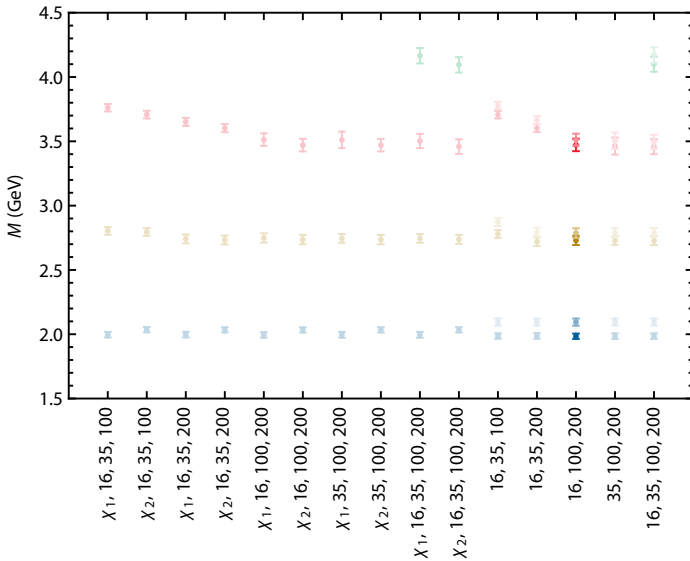
Comparison of Bases



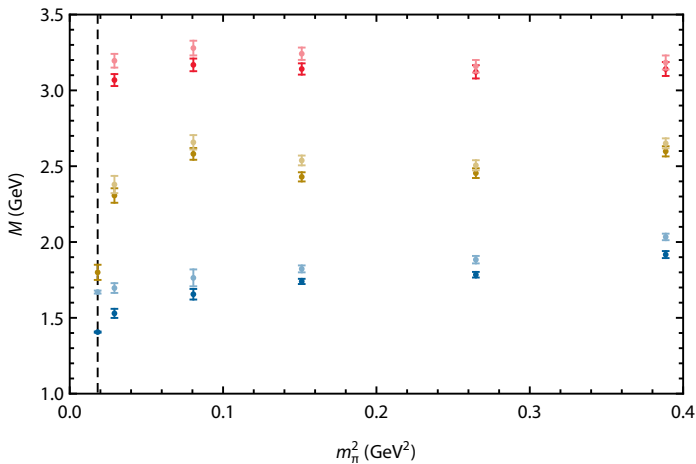
Comparison of Bases



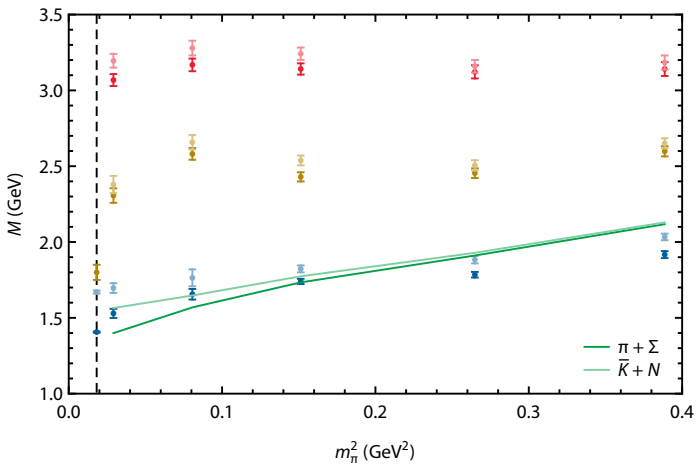
Comparison of Bases



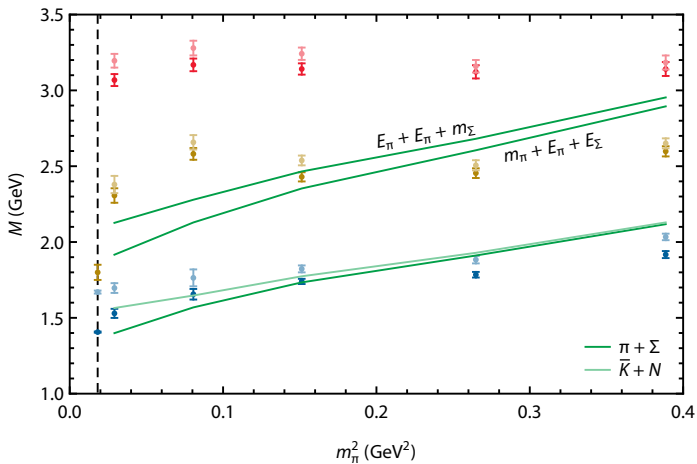
All States



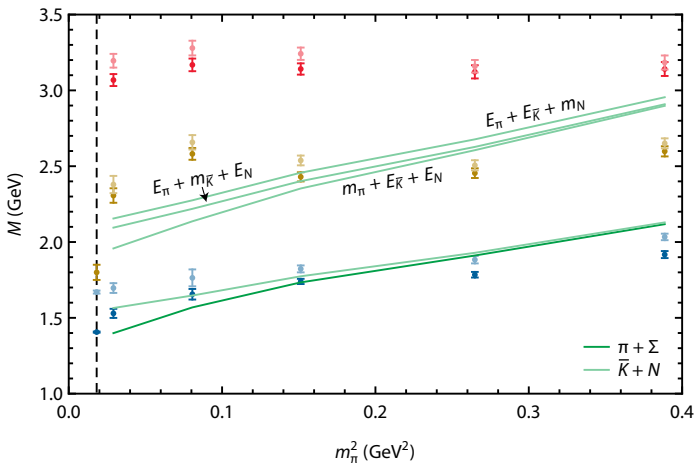
All States



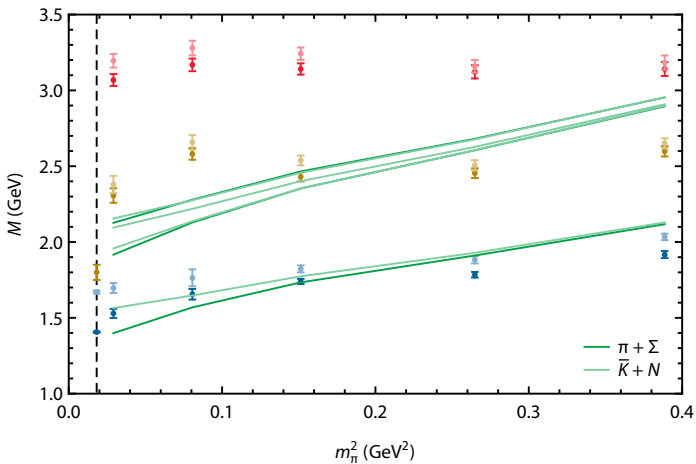
All States



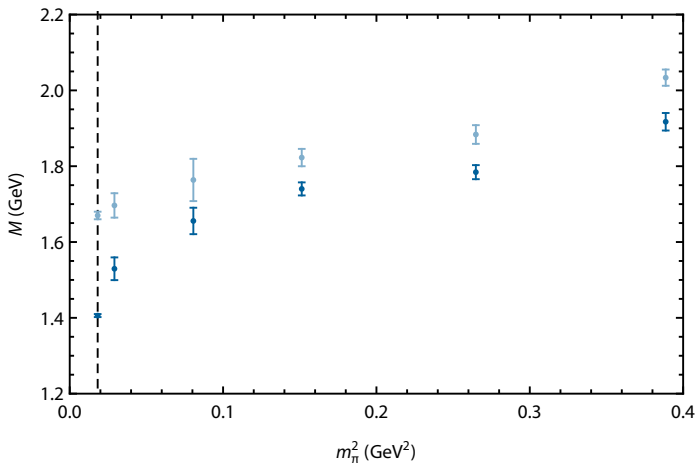
All States



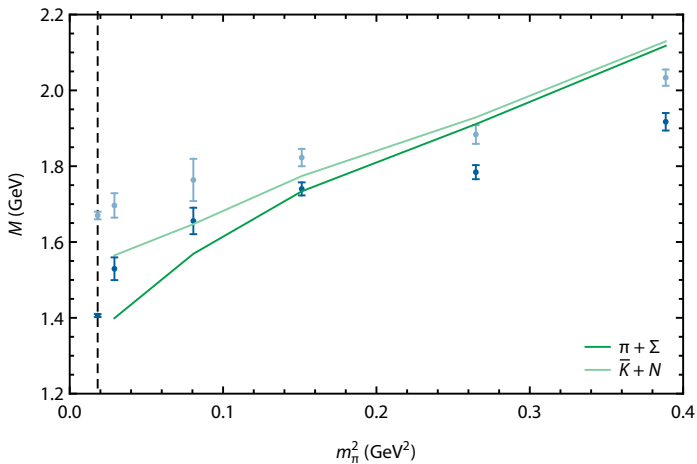
All States



Lowest Two States



Lowest Two States



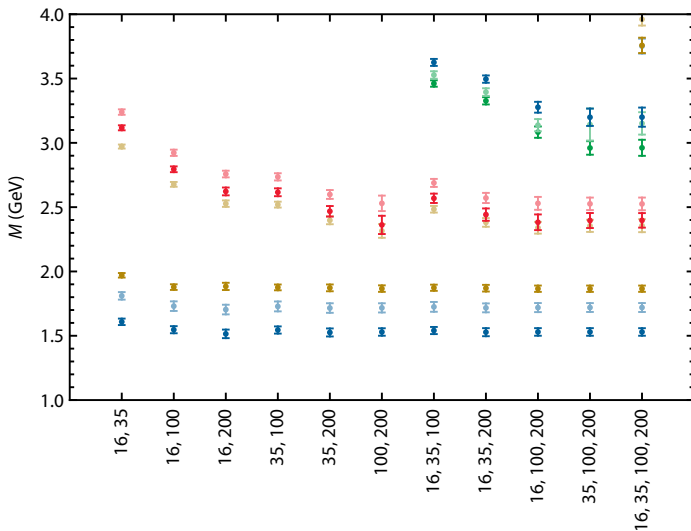
Octet and Singlet Flavour Symmetry

- Now that we have isolated the $\Lambda(1405)$, we investigate its properties.
- To begin with, we can extend the analysis to include the octet and singlet flavour-symmetric operators.

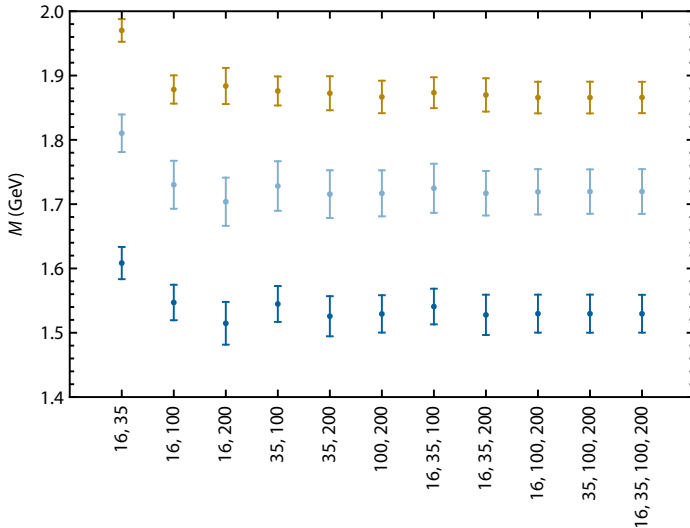
Octet and Singlet Flavour Symmetry

- Now that we have isolated the $\Lambda(1405)$, we investigate its properties.
- To begin with, we can extend the analysis to include the octet and singlet flavour-symmetric operators.
- So far, looked at lightest quark mass ($m_\pi = 170.7 \pm 2.1$ MeV).
 - Use $\chi_{1,2}^8$ and χ^1 with various sets of smearing.

Comparison of Bases – All States



Comparison of Bases – Lowest Three States



Conclusions

- Now have a method for isolating the $\Lambda(1405)$ in Lattice QCD.

Conclusions

- Now have a method for isolating the $\Lambda(1405)$ in Lattice QCD.
- Using this, we can investigate the properties of this unusual resonance.

Conclusions

- Now have a method for isolating the $\Lambda(1405)$ in Lattice QCD.
- Using this, we can investigate the properties of this unusual resonance.
 - Can attempt to identify the flavour-symmetry associated with each state.

Conclusions

- Now have a method for isolating the $\Lambda(1405)$ in Lattice QCD.
- Using this, we can investigate the properties of this unusual resonance.
 - Can attempt to identify the flavour-symmetry associated with each state.
 - Have also begun a form factor study of the $\Lambda(1405)$.

Acknowledgements

This research was undertaken on the NCI National Facility in Canberra, Australia, which is supported by the Australian Commonwealth Government.