

The role of strange quarks in nucleon structure

Ross Young University of Adelaide

Achievements and New Directions in Subatomic Physics: Workshop in Honour of Tony Thomas's 60th Birthday

University of Adelaide



Chiral Extrapolation



Chiral Extrapolation







Differences described by chiral loops







Physics!!

- Let's use these new tools to learn something new
- How about strangeness?

Electromagnetic form factors characterise the charge and magnetisation distribution in the nucleon

 $G_M(Q^2)$ $G_E(Q^2)$

Electromagnetic form factors characterise the charge and magnetisation distribution in the nucleon

$$G_E(Q^2)$$
 $G_M(Q^2)$

Measure total response from all quarks

Proton
$$G_{E,M}^{p} = +\frac{2}{3}G_{E,M}^{u} - \frac{1}{3}G_{E,M}^{d} - \frac{1}{3}G_{E,M}^{s}$$

Electromagnetic form factors characterise the charge and magnetisation distribution in the nucleon



Electromagnetic form factors characterise the charge and magnetisation distribution in the nucleon



Charge symmetry: *up* quark in proton = *down* quark in neutron, ...

Electromagnetic form factors characterise the charge and magnetisation distribution in the nucleon

$$G_{E}(Q^{2}) \qquad G_{M}(Q^{2})$$
Measure total response from all quarks
Proton
$$G_{E,M}^{p} = +\frac{2}{3}G_{E,M}^{u} - \frac{1}{3}G_{E,M}^{d} - \frac{1}{3}G_{E,M}^{s}$$
Charge symmetry: *up* quark in proton = *down* quark in neutron, ...
Neutron
$$G_{E,M}^{n} = -\frac{1}{3}G_{E,M}^{u} + \frac{2}{3}G_{E,M}^{d} - \frac{1}{3}G_{E,M}^{s}$$

Electromagnetic form factors characterise the charge and magnetisation distribution in the nucleon

$$G_{E}(Q^{2}) \qquad G_{M}(Q^{2})$$
Measure total response from all quarks
Proton
$$G_{E,M}^{p} = +\frac{2}{3}G_{E,M}^{u} - \frac{1}{3}G_{E,M}^{d} - \frac{1}{3}G_{E,M}^{s}$$
Charge symmetry: *up* quark in proton = *down* quark in neutron, ...
Neutron
$$G_{E,M}^{n} = -\frac{1}{3}G_{E,M}^{u} + \frac{2}{3}G_{E,M}^{d} - \frac{1}{3}G_{E,M}^{s}$$
2 Equations - 3 Unkowns!

$$G_{E,M}^{p,Z} = \left(1 - \frac{8}{3}\sin^2\theta_W\right)G_{E,M}^u + \left(-1 + \frac{4}{3}\sin^2\theta_W\right)G_{E,M}^d + \left(-1 + \frac{4}{3}\sin^2\theta_W\right)G_{E,M}^s$$

Electroweak couplings differ from usual charges

Weak mixing angle: $\sin^2 \theta_W$

$$G_{E,M}^{p,Z} = \left(1 - \frac{8}{3}\sin^2\theta_W\right)G_{E,M}^u + \left(-1 + \frac{4}{3}\sin^2\theta_W\right)G_{E,M}^d + \left(-1 + \frac{4}{3}\sin^2\theta_W\right)G_{E,M}^s$$

Electroweak couplings differ from usual charges

Weak mixing angle: $\sin^2 \theta_W$

$$G_{E,M}^{p} = +\frac{2}{3}G_{E,M}^{u} - \frac{1}{3}G_{E,M}^{d} - \frac{1}{3}G_{E,M}^{s}$$
$$G_{E,M}^{n} = -\frac{1}{3}G_{E,M}^{u} + \frac{2}{3}G_{E,M}^{d} - \frac{1}{3}G_{E,M}^{s}$$

3 Equations — 3 Unkowns Can isolate strangeness!





 $3O_N = 2p + n - u^p$



$$3O_N = 2p + n - \frac{u}{u^{\Sigma}} (\Sigma^+ - \Sigma^-)$$



$$3O_N = 2p + n - \frac{u}{u^{\Sigma}} (\Sigma^+ - \Sigma^-)$$

$$3O_N = p + 2n - u^n$$

$$\Xi^0 - \Xi^- = u^{\Xi}$$

$$3O_N = p + 2n - \frac{u^n}{u^{\Xi}} (\Xi^0 - \Xi^-)$$



 $3O_N = p + 2n - u^n$ $\Xi^0 - \Xi^- = u^\Xi$ $3O_N = p + 2n - \frac{u^n}{u^\Xi} \Xi^0 - \Xi^-)$

Constraint on GMs









u-quark in the Sigma



Magnetic strangeness result



 $G_M^s = -0.046 \pm 0.022 \mu_N$

Limited hyperon info, take absolute values from lattice





















Experimental Status (2005)

- Conglomerate of world's strangeness measurements
- Uses constraint from best theoretical estimate of anapole form factor *Zhu et al*.

 $Q^2 = 0.1 \,\mathrm{GeV}^2$



Strange Magnetic

Global Analysis

• Extract the anapole contribution from experiment

 $\tilde{G}_A^N = \tilde{g}_A^N (1 + Q^2 / \Lambda^2)^{-2}$

- Fit strangeness to measured asymmetries
 - Consistent treatment of electromagnetic form factors and radiative corrections
- Use all available data for $Q^2 < 0.3 \,{
 m GeV}^2$
- Taylor expansion of strangeness

$$G_E^s = \rho_s Q^2 + \rho'_s Q^4 + \dots$$
$$G_M^s = \mu_s + \mu'_s Q^2 + \dots$$





Global Analysis

• Extract the anapole contribution from experiment

 $\tilde{G}_A^N = \tilde{g}_A^N (1 + Q^2 / \Lambda^2)^{-2}$

- Fit strangeness to measured asymmetries
 - Consistent treatment of electromagnetic form factors and radiative corrections
- Use all available data for $Q^2 < 0.3 \,\mathrm{GeV}^2$
- Taylor expansion of strangeness

$$G_E^s = \rho_s Q^2 + \rho'_s Q^4 + \dots$$
$$G_M^s = \mu_s + \mu'_s Q^2 + \dots$$





Strangeness form factors



$Q^2 = 0.1 \,\mathrm{GeV}^2$

New HAPPEX Measurement

• Excellent agreement with global analysis



Strange Magnetic

Combined Analysis

- Combined constraints on current knowledge of strangeness content.
- Strangeness is small!
- 95% confidence:
 < 5% charge radius
 < 6% magnetic moment
- In support of theory estimate



Strange Magnetic

Searching for new physics



Constrained by low-energy data!

 $C_{1u} \sim Q_e^W Q_u^W$ $C_{1d} \sim Q_e^W Q_d^W$

 $\mathcal{L}_{\rm SM}^{\rm PV} = -\frac{G_F}{\sqrt{2}} \bar{e} \gamma_{\mu} \gamma_5 e \sum_{q} C_{1q}^{\rm SM} \bar{q} \gamma^{\mu} q$

Proton weak charge extrapolation



Current knowledge and future: Q-weak



Q-weak: Assuming SM



Q-weak: Assuming SM



Q-weak constrains new physics to beyond 2 TeV

Strange condensate in the nucleon

• Gell-Mann–Okubo Relation and Pion-Nucleon sigma term $M_N, M_\Lambda, M_\Sigma, M_\Xi$

$m_s \langle N | \bar{s}s | N \rangle \simeq 335 \pm 132 \,\mathrm{MeV}$

Nelson & Kaplan PLB(1987)

 $\sim M_N^{phys} - M_N^{SU(3)chiral limit}$

Strange condensate in the nucleon

• Gell-Mann–Okubo Relation and Pion-Nucleon sigma term $M_N, M_\Lambda, M_\Sigma, M_\Xi$

$m_s \langle N | \bar{s}s | N \rangle \simeq 335 \pm 132 \,\mathrm{MeV}$ Nelson & Kaplan PLB(1987)

 $\sim M_N^{phys} - M_N^{SU(3)chiral \ limit}$

QCD Lagrangian $\sim \ldots \bar{s}(D + m_s)s$

$$m_s \langle N | \bar{s}s | N \rangle = m_s \frac{\partial M_N}{\partial m_s}$$

evaluated at physical point!

Strange condensate in the nucleon

• Gell-Mann–Okubo Relation and Pion-Nucleon sigma term $M_N, M_\Lambda, M_\Sigma, M_\Xi$

 $m_s \langle N | \bar{s}s | N \rangle \simeq 335 \pm 132 \,\mathrm{MeV}$

Nelson & Kaplan PLB(1987)

 $\sim M_N^{phys} - M_N^{SU(3)chiral limit}$

QCD Lagrangian $\sim \dots \bar{s}(D + m_s)s$ $m_s \langle N | \bar{s}s | N \rangle = m_s \frac{\partial M_N}{\partial m_s}$ evaluation

evaluated at *physical* point!

Improved Effective Field Theory estimate

$$m_s \frac{\partial M_N}{\partial m_s} = 113 \pm 108 \,\mathrm{MeV}$$

Borasoy & Meissner (1997)

Fit to 8 LHPC points





Fit to 8 PACS-CS points

PACS-CS: 2+1-flavour simulation; different action discretization to LHPC



PACS-CS have an additional run with a different strange quark mass



Quantifying Sources of Uncertainty

Nucleon Mass (GeV)

Discretisation

 LHPC
 0.945 ± 0.029

 PACS-CS
 0.954 ± 0.042

Extrapolated baryon masses and fit parameters (LECs) in agreement — good news for us and lattice practitioners

RegulatorDipole0.9410Sharp0.9452

Small dependence on choice of regulator — similarly for other functional forms

Source	MeV	
Statistical	23.6	
Discretisation	4.2	
Model	3.1	
Regulator	2.1	
f_{π} (5%)	0.7	
F (15%)	1.3	
D (15%)	1.3	
C (15%)	0.9	
Δ_{10-8} (15%)	0.4	



Baryon Sigma Terms

 $\bar{\sigma}_{Bq} = \frac{m_q}{M_B} \frac{\partial M_B}{\partial m_a}$

	N	Λ	\sum	[I]
$\bar{\sigma}_{Bl}$	0.050(9)(1)(3)	0.028(4)(1)(2)	0.0212(27)(1)(17)	0.0100(10)(0)(4)
$\bar{\sigma}_{Bs}$	0.033(16)(4)(2)	0.144(15)(10)(2)	0.187(15)(3)(4)	0.244(15)(12)(2)



πN Sigma Term (Expt):GL: Gasser & Leutwyler (1991)GW: Pavan et al. (2001)

Octet Masses & Breaking: Gasser (1981) NK: Nelson & Kaplan (1987) BM: Borasoy & Meissner (1997)

3-flavour Lattice QCD: YT: Young & Thomas (2009) TF: Toussaint & Freeman (2009)

We determine precisely *both* the light and strange quark sigma terms



Spin-independent neutralino cross sections

- Ellis, Olive & Savage, PRD(2008)
 - Constrained Minimal Supersymmetric Standard Model (CMSSM)
 - Neutralino as dark matter candidate
 - Scalar contact interaction

$$\mathcal{L}_{SI} = \sum_{i} \alpha_{3i} \bar{\chi} \chi \bar{q}_i q_i$$

$$\sigma_{SI}^p \propto |f_p|^2$$
$$\frac{f_p}{M_p} = \sum_{q=u,d,s} \bar{\sigma}_{pq} \frac{\alpha_{3q}}{m_q} + \frac{2}{27} f_{TG}^p \sum_{q=c,b,t} \frac{\alpha_{3q}}{m_q}$$

 $f^p_{TG} = 1 - \sum_{q=u,d,s} \bar{\sigma}_{pq} \qquad \begin{array}{l} \mbox{Trace anomaly:} \\ \mbox{Shifman, Vainstein} \end{array}$

Trace anomaly: Shifman, Vainstein & Zakharov, PLB(1978)

Uncertainty dominated by knowledge of light-quark sigma terms



Updated cross sections for benchmark models



Ellis, Olive & Savage

Strong dependence on sigma term from poorly known strangeness

Giedt, Thomas & Young, PRL (2009)

Tremendous advance in precision from new lattice QCD results

Nuclear physics & lattice QCD can help discriminate supersymmetry scenarios



Summary

- Strange quarks play a relatively small role in the nucleon structure
- "Spin-offs": precision electroweak physics, dark matter scattering

Summary

- Strange quarks play a relatively small role in the nucleon structure
- "Spin-offs": precision electroweak physics, dark matter scattering





Anapole Form Factor

- Backward angle measurements have increased sensitivity to axial form factor.
- Anapole corresponds to an electroweak correction to the proton structure
- Anapole form factor is not measured in any other process



Axial+Anapole

Anapole Form Factor

- Backward angle measurements have increased sensitivity to axial form factor.
- Anapole corresponds to an electroweak correction to the proton structure
- Anapole form factor is not measured in any other process



The axial/anapole term

