

# Fermions, Scalars and Dark Energy

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with

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Presented at

Achievements and New Directions in Subatomic Physics:  
Workshop in Honour of Tony Thomas' 60th Birthday  
Adelaide, South Australia. February 15 - February 19, 2010.

# Outline

- 1 Background
  - The ultimate Copernican revolution
  - 14 Years ago in Adelaide .....
  - Is this related to Dark Energy?
- 2 Calculation of  $w$ 
  - $w$  as a function of fermion number density
  - Connecting data and parameters
- 3 Summary

## Type Ia Supernovae

- Discovery of the re-acceleration of the universe in 1998
- The acceleration of the time evolution of the scale parameter is given by the Friedmann Equation

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P) \quad (1)$$

- Introduce the parameter  $w$   
where  $p = w\rho_E$   
and  $\rho_E =$  total energy density
- Positive acceleration requires  $w < -1/3$
- $w = -1$  is Einstein's cosmological constant  $\Lambda$
- Most recent data and analysis (arXiv:0901.4804)  
 $1 + w = 0.013 \pm 0.067 \pm .11$   
assuming constant  $w$

# Dark Energy

- The *stuff* with  $w \approx -1$  is called **dark energy**, and it is about 70% of the present energy density of the universe.
- Only 5% is baryons, and the remaining 25% is **dark matter**
- Most of the universe is not the stuff of which we are made

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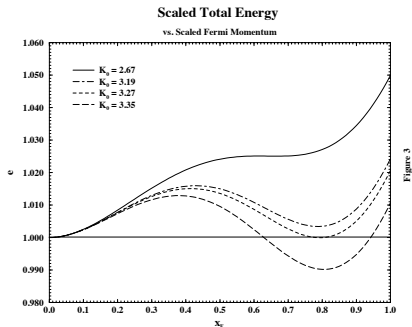
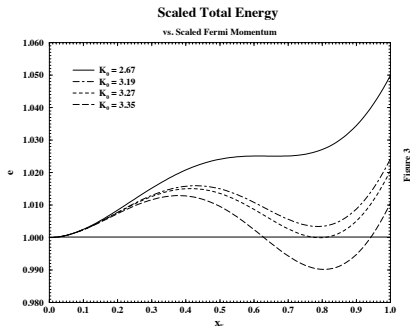


Figure 3

# The MSW Effect Without Matter

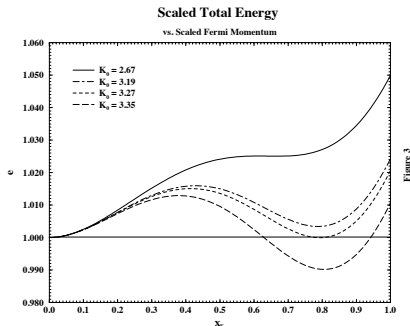
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- The total energy of a sea of interacting neutral fermions (neutrinos), mass  $m_0$ , Fermi energy  $k_F = m_0 x_F$ , as a function of  $x_F$  or density.

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- The total energy of a sea of interacting neutral fermions (neutrinos), mass  $m_0$ , Fermi energy  $k_F = m_0 x_F$ , as a function of  $x_F$  or density.
- Note the region of negative pressure.



# The basic fermion-scalar system

- Lagrangian:

Majorana Fermions (2 component)  $\Phi$  and scalars  $\zeta$

$$\begin{aligned}\mathcal{L} &= \frac{1}{2i}[\Phi^\dagger \sigma^\mu \partial_\mu \Phi - \partial_\mu \Phi^\dagger \sigma^\mu \Phi] + m_0 \frac{1}{2}[\Phi^T \sigma^2 \Phi + \Phi^\dagger \sigma^2 \Phi^*] \\ &+ \frac{1}{2}[\zeta(\partial^2 - m_\zeta^2)\zeta] \\ &+ \frac{1}{2}g[\Phi^T \sigma^2 \Phi + \Phi^\dagger \sigma^2 \Phi^*]\zeta\end{aligned}$$

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- the equations of motion

$$\begin{aligned} [\partial^2 + m_\zeta^2] \zeta &= -\frac{1}{2} g [\Phi^T \sigma^2 \Phi + \Phi^\dagger \sigma^2 \Phi^*] \\ [i\sigma^\mu \partial_\mu] \Phi &= m_0 \sigma^2 \Phi^* + g \zeta \sigma^2 \Phi^*. \end{aligned}$$

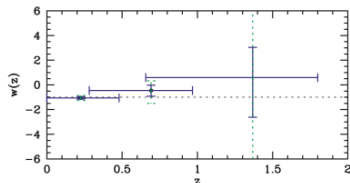
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Is  $w$  constant?



From Reiss et al, Ap J  
**659** (2007) 98

$w$  does not vary rapidly  
with  $z$

Is it flat or is it rising?

# Does the data constrain our model?

In our model we find

- $w$  is **not** constant

# Does the data constrain our model?

In our model we find

- $w$  is **not** constant
- **But** the variation of  $w$  can be made consistent with the present data

# The effective mass

- remember the equations of motion

$$\begin{aligned} \left[ \partial^2 + m_\zeta^2 \right] \zeta &= -\frac{1}{2} g [\Phi^T \sigma^2 \Phi + \Phi^\dagger \sigma^2 \Phi^*] \\ [i\sigma^\mu \partial_\mu] \Phi &= m_0 \sigma^2 \Phi^* + g \zeta \sigma^2 \Phi^*. \end{aligned}$$

- Clearly the scalar field induces an effective mass of the fermion  $m^* = m_0 + g\zeta$
- For a homogeneous filled Fermi sea, the expectation value of the Lorentz scalar density  $\mathcal{S} = \frac{1}{2} \{\Phi^T \sigma^2 \Phi + \Phi^\dagger \sigma^2 \Phi^*\}$  is simply related to the number density of the fermions and, with  $y = m^*/m_0$ ,

$$y = 1 - \{yK_0/2\} \left[ e_F x_F - y^2 \ln \left( \{e_F + x_F\}/y \right) \right],$$

with  $e_F = \sqrt{x_F^2 + y^2}$ , and  $K_0 = \{g^2(m_0)^2\}/\{\pi^2 m_\zeta^2\}$ .

# Calculating $w$

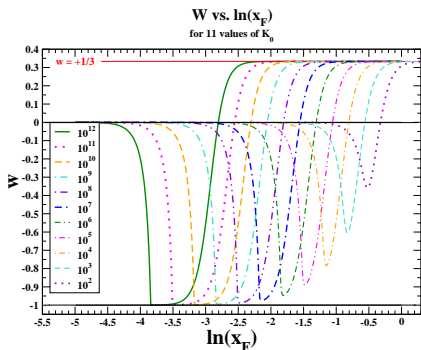
- Using the equations of the FLRW metric, with a scale parameter  $a$ ,  $1 + w = -\frac{1}{3} \frac{\partial \ln \rho_E}{\partial \ln a}$ .
- In our model, with the only contribution to  $\rho_E$  from the neutral fermions and the scalars, and  $\rho_E = em_n^{(0)} \rho_n$ , where  $\rho_n$  is the fermion number density,  $w = \frac{1}{3} \frac{\partial \ln(e)}{\partial \ln(x_F)}$ .
- 

$$w = \frac{1}{3} \frac{e_F K_0 x_F^3 - 3(2-y)(1-y)}{e_F K_0 x_F^3 + (2-y)(1-y)}$$

It follows immediately that  $w > -1$  in our model.

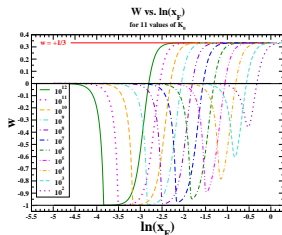


# The result



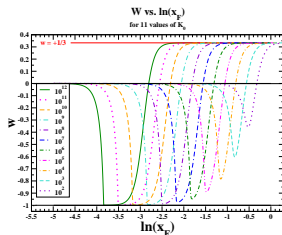
- For **very** small  $x_F$  (in fact  $x_F^3 \ll K_0^{-1}$ ),  $w \approx x_F^2/5 \rightarrow 0$
- For large  $x_F$ , the effective mass goes to zero, and  $w \rightarrow 1/3$ , the value for a relativistic gas of Fermions.
- There is a minimum of  $w$ ,  $w_{\min}$ , which gets closer to  $-1$  as  $K_0$  increases. It occurs at  $x_{\min} \approx (3.83/K_0)^{1/3}$  for  $K_0 > 10^6$

# How to fit the observed $w(z)$



- For  $x_F > x_{\min}$  and large  $K_0$ ,  $w(x_F)$  varies slowly with  $x_F$ . Push the “flat” region out to  $z \approx 1.0$
- For  $x_F < x_{\min}$  and large  $K_0$ ,  $w(x_F)$  rapidly approaches 0. Put this near or after now.
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- For  $K_0 > 10^6$ ,  $1 + w_{\min} \approx 2x_{\min}$ , and then the observed limit on  $1 + w$  requires  $x_{\min} < 0.06$
- Can we find a set of parameters for our model that do this?

# The parameters

The parameters of the model are

- the Fermion mass  $m_0$
- the Scalar mass  $m_\zeta$
- the coupling constant  $g$
- which enter in the dimensionless parameter
$$K_0 = \{g^2(m_0)^2\} / \{\pi^2 m_\zeta^2\}$$
- the density of the fermions, parameterized by the Fermi momentum in units of the Fermion mass,  $x_F$ , at some value of the LFRW scale parameter  $a = a_1$ , of the red shift  $z = z_1$

# Data

We know

- The limit on  $1 + w$ , if assumed constant:  $1 + w \leq 0.13$  at the  $1\sigma$  level
- The total dark energy density, if assumed constant:  
 $\rho_E^{(\text{obs})} = (2.4 \text{ meV})^4$
- **IF** we assume that the neutral fermions which generate dark energy are one of the species of active neutrinos, then their present number density is  $100 \text{ cm}^{-3}$

## Satisfying constraints

- From the calculated variation of  $w_{\min}$  with  $K_0$   
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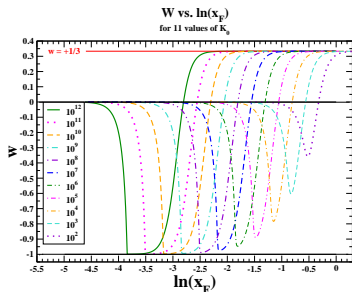
- whence  $m_0 \geq 160 \text{ meV}$
- with the above,  $x_F \approx (3.83 / K_0)^{1/3} \implies k_F < 2.49 \text{ meV}$

## With more analysis

- As  $x_F \propto (1 + z)$ , and we want  $w$  to stay near  $-1$  as  $z = 0.2 \rightarrow z = 1$  we need the minimum of  $w$  to be flat at its minimum for a range of a factor of 2 in  $x_F$ .

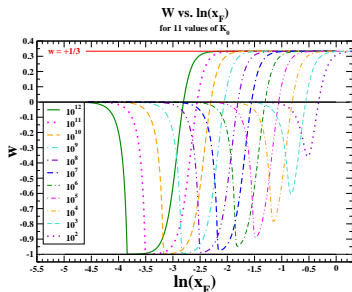
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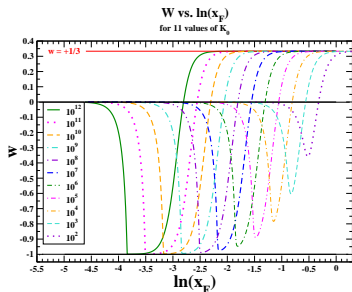
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- $K_0 > 10^9$  fits this
- $m_0 > 0.89$  eV and  $k_F < 1.40$  meV



# Accommodating LSP

- Take  $K_0 = 10^{54}$
- Then  $m_0 = 160$  GeV, and  $k_F = 25$   $\mu$ eV
- Remember  $K_0 = \{g^2(m_0)^2\}/\{\pi^2 m_\zeta^2\}$   
Take  $m_\zeta = (7 \times 10^9 \text{ lightyears})^{-1} \sim 3 \times 10^{-30}$  meV  
Then  $g^2/(4\pi) = 2.8 \times 10^{-34}$
- far too small to be constrained by terrestrial experiments

# Summary

- Our old neutrino cloud model finds a new life as a natural model for dark energy
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- Our old neutrino cloud model finds a new life as a natural model for dark energy
- Future experiments on SNIa and at LHC will tighten constraints on parameters
- Next questions
  - What is the extension of the standard model to include the new interactions?
  - Do “clouds” form in this version of the model?



# Thank you

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- To Tony, with best wishes for many happy returns from all the authors