

**THE STATUS  $\alpha_s$  FROM THE LATTICE AND  
HADRONIC  $\tau$  DECAYS**

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*Achievements and New Directions in Subatomic Physics*

*(Tony 60-Fest)*

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## OUTLINE

- *Context (including lattice vs.  $\tau$  decay tension)*
- *Recent updates of UKQCD/HPQCD lattice approach*
- *New results on the hadronic  $\tau$  decay determination*
- *Prospects/issues*

## THANKS/COLLABORATORS

- CSSM version lattice  $\alpha_s$ : Derek Leinweber, Peter Moran and Andre Sternbeck
- $\tau$  decay  $\alpha_s$ : Tzahi Yavin
- work in progress on duality violation in  $\tau$   $\alpha_s$  (etc.): Maarten Golterman, Santi Peris, Oscar Cata
- CSSM in general for its terrific physics atmosphere!!

## CONTEXT ETC.

- HPQCD/UKQCD, PRL95 (2005) 052002: perturbative analysis of UV-sensitive lattice observables [dominant input to PDG08 assessment  $\alpha_s(M_Z) = 0.1176(20)$ ]

$$[\alpha_s(M_Z)]_{latt} = 0.1170(12)$$

- Conventional ALEPH, OPAL [e.g., EPJC56 (2008) 305]: “(k,m) spectral weight” hadronic  $\tau$  determination as of mid-2008:

$$[\alpha_s(M_Z)]_{\tau} = 0.1212(11)$$

- c.f. recent experimental determinations

Source	$\alpha_s(M_Z)$
Global EW fit	0.1193(28)
H1+ZEUS NLO inclusive jets	0.1198(32)
H1 high- $Q^2$ NLO jets	0.1182(45)
Non-singlet structure functions	0.1142(23)
NNLO+NLLA LEP event shapes	0.1224(39)
NNLO+NLLA JADE event shapes	0.1172(51)
NLO inclusive jets, $p\bar{p}$	0.1161(48)
ZEUS NLO inclusive jets, $\gamma p$	0.1223(38)
NNLL ALEPH+OPAL thrust distributions	0.1172(21)
$\Gamma[\Upsilon(1s) \rightarrow \gamma X]/\Gamma[\Upsilon(1s) \rightarrow X]$	0.1190(60)

NOTE: expt'l det'n errors large c.f. nominal lattice,  $\tau$

Excluding  $\tau$ , lattice input, Bethke [0908.1135] average

$$\alpha_s(M_Z) = 0.1184(7) \rightarrow 0.1179(13)$$

## UPDATES OF HPQCD LATTICE APPROACH

- Based on perturbative analyses of observables,  $O_k$ , measured on MILC (asqtad)  $n_f = 2 + 1$  ensembles

- $O(\alpha_s^3)$   $D = 0$  ( $m_q = 0$ ) expansion

$$[O_k]_{D=0} = D_k \alpha_T(Q_k) \left[ 1 + c_1^{(k)} \alpha_T(Q_k) + c_2^{(k)} \alpha_T^2(Q_k) + \dots \right]$$

with  $Q_k = d_k/a$  the BLM scale for  $O_k$

- $D_k, c_1^{(k)}, c_2^{(k)}, d_k$ : Q. Mason et al. 3-loop lattice PT

- Original HPQCD/UKQCD analysis [PRL 95 (2005) 052002]:  $a \sim 0.18, 0.12, 0.09$  fm ensembles
- HPQCD [PRD78 (2008) 114507], CSSM [PRD78 (2008) 114504] updates add new  $a \sim 0.15, 0.06$  fm ensembles, one  $(am_\ell, am_s)$   $a \sim 0.045$  fm ensemble (HPQCD only) (results dominated by finer ensembles)
- $m_q$ -dependent NP contributions: linear  $m_q$  extrapolation/subtraction
- $m_q$ -independent NP: estimate/subtract via LO  $\langle aG^2 \rangle$  (+ fitted  $D > 4$  for more long-distance-sensitive observables in 2008 HPQCD)

## Some relevant details

- $D = 0$  to  $O(\alpha_s^3)$  insufficient to account for observed scale dependence  $\Rightarrow$  **MUST fit additional HO term(s)**
- 2008 HPQCD, CSSM: different  $D = 0$  expansion parameter choices  $\Rightarrow$  different (complementary) handling of residual HO perturbative uncertainties
- $m_q \rightarrow 0$  extrapolation very reliable:
  - many  $(am_\ell, am_s)$  for  $a \sim 0.12$  fm, very good linearity (plus good linearity for other  $a$  as well)
  - extrapolation very stable to added non-linear terms



- Re  $m_q$ -independent NP subtraction:
  - $\langle aG^2 \rangle = 0 \pm 0.012 \text{ GeV}^4$  (HPQCD), with independent fit for each  $O_k$
  - $\langle aG^2 \rangle = 0.009 \pm 0.007 \text{ GeV}^4$  (CSSM), common input for all  $O_k$  to identify small NP cases
  - estimated  $D = 4$  correction tiny for shortest-distance-sensitive observables (e.g.,  $\log(W_{11})$ ,  $\log(W_{12})$ )
  - After fitted  $m_q$ -independent NP subtractions, HPQCD observables with LARGE estimated  $D = 4$  corrections yield  $\alpha_s$  in good agreement with  $\log(W_{11})$  etc.

- **COMPARISON OF HPQCD, CSSM RESULTS**
  - Results for a selection of three least-NP and four most-NP observables
  - $\delta_{D=4} \equiv$  fractional change from scale dependence of “raw” observable to that of  $m_q$ -independent NP-subtracted version between  $a \sim 0.12$  and  $\sim 0.06$  fm ( $\langle aG^2 \rangle = 0.009 \text{ GeV}^4$  as input)
  - **NOTE: re estimated NP  $D = 4$  corrections**
    - \* corrections far and away the largest for the 3 HPQCD “outliers”
    - \* despite *large* corrections,  $\alpha_s$  agree with results from observables where NP corrections negligible

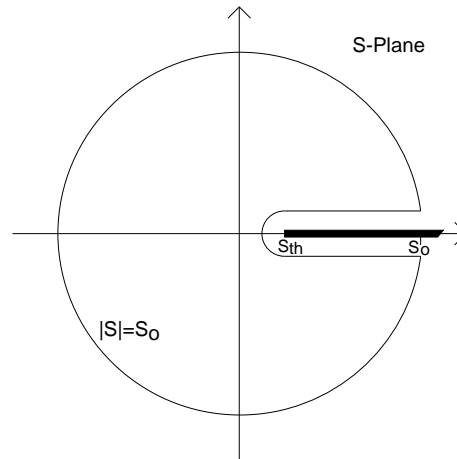
–  $\delta_{D=4}$  and resulting  $\alpha_s(M_Z)$  values

$O_k$	$\alpha_s(M_Z)$ (HPQCD)	$\alpha_s(M_Z)$ (CSSM)	$\delta_{D=4}$
$\log(W_{11})$	0.1185(8)	0.1190(11)	0.7%
$\log(W_{12})$	0.1185(8)	0.1191(11)	2.0%
$\log\left(\frac{W_{12}}{u_0^6}\right)$	0.1183(7)	0.1191(11)	5.2%
$\log\left(\frac{W_{11}W_{22}}{W_{12}^2}\right)$	0.1185(9)	N/A	32%
$\log\left(\frac{W_{23}}{u_0^{10}}\right)$	0.1176(9)	N/A	53%
$\log\left(\frac{W_{14}}{W_{23}}\right)$	0.1171(11)	N/A	79%
$\log\left(\frac{W_{11}W_{23}}{W_{12}W_{13}}\right)$	0.1174(9)	N/A	92%

## THE HADRONIC $\tau$ DETERMINATION

- Based on FESRs for  $\Pi_{T;ud}^{(0+1)}$ ,  $T = V, A, V + A$

$$\int_0^{s_0} w(s) \rho_{T;ud}^{(0+1)}(s) ds = -\frac{1}{2\pi i} \oint_{|s|=s_0} w(s) \Pi_{T;ud}^{(0+1)}(s) ds$$



- valid for any  $s_0$ , analytic  $w(s)$
- LHS: data; RHS: OPE (hence  $\alpha_s$ ) for  $s_0 \gg \Lambda_{QCD}^2$

- The spectral integrals

- V, A, I = 1 spectral function  $\rho_{V/A;ud}^{(J)=(0+1)}(s)$  from experimental differential decay distributions  $\frac{dR_{V/A;ud}}{ds}$ , with  $R_{V/A;ud} \equiv \frac{\Gamma[\tau \rightarrow \nu_\tau \text{ hadrons}_{V/A;ud}(\gamma)]}{\Gamma[\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e(\gamma)]}$
- $\Rightarrow$  experimental access to generic  $(J) = (0 + 1)$ ;  $w(s)$ -weighted,  $0 < s \leq s_0 \leq m_\tau^2$  spectral integrals

$$I_{spec;T}^w(s_0) = \int_0^{s_0} ds w(s) \rho_{T;ud}^{(0+1)}(s)$$

- The OPE side:

- $D = 0$ : fixed by  $\alpha_s$  (known to 5 loops); strongly dominant for  $s_0 \gtrsim 2 \text{ GeV}^2$

- $D = 2$ :  $\propto (m_d \pm m_u)^2$ , hence negligible

- $D = 4$ : fixed by  $\langle aG^2 \rangle$ ,  $\langle m_\ell \bar{\ell} \ell \rangle$ ,  $\langle m_s \bar{s} s \rangle$

- $D = 6, 8, \dots$ :

- \* not known phenomenologically, hence fitted to data (or guesstimated)

- \* for  $\sim 1\%$   $\alpha_s(M_Z)$  determination need integrated  $D > 4$  to  $\lesssim 0.5\%$  of  $D = 0$

– More on fitting the  $D > 4$  contributions

\*  $w(y) = \sum_{m=0} b_m y^m$ ,  $y = s/s_0$  to distinguish contributions with different  $D$  (differing  $s_0$  dependence)

\* integrated  $D = 2k + 2 \geq 2$  contribution  $\Leftrightarrow b_k \neq 0$  (up to  $O[\alpha_s^2(m_T^2)]$ )  $\Rightarrow$  contributions up to  $D_{max} = 2N + 2$  for degree  $N$   $w(y)$

\* integrated  $D = 2k + 2$  contributions  $\propto 1/s_0^k$

$$\frac{-1}{2\pi i} \oint_{|s|=s_0} ds w(y) \sum_{D>4} \frac{C_D}{Q^D} = \sum_{k \geq 2} (-1)^k \frac{b_k C_{2k+2}}{s_0^k}$$

## Summary of recent $\tau$ -based determinations

- Differences in 6-loop  $D = 0$  Adler function coeff,  $d_5$ ;  $D = 0$  series integral prescription;  $D > 4$  treatment
- Duality violation typically assumed negligible

Source	$d_5$	$D > 4$ self-consistency	PT scheme	$\alpha_s(M_Z^2)$
BCK08	275	No	$\frac{1}{2}(\text{FO} + \text{CI})$	0.1202(19)
ALEPH08	383	No	CI	0.1211(11)
BJ08	283	No	FO	0.1185(14)
	283	No	model	0.1179(8)
MY08	275	Yes	CI	0.1187(16)
N09	0	partly	$\frac{1}{2}(\text{FO} + \text{CI})$	0.1192(10)
M09	400	No	$\frac{1}{2}(\text{RC} + \text{CI})$	0.1213(11)
CF09	283	No	modified CI	0.1186(13)



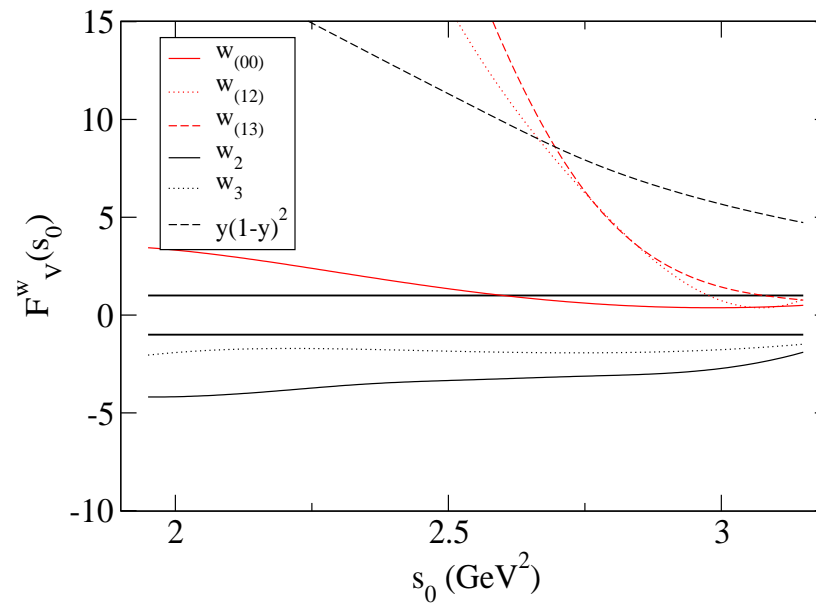
## THE ALEPH, OPAL (AND RELATED) ANALYSES

- $w_{(00)}(y) = 1 - 3y^2 + 2y^3 \Rightarrow$  OPE up to  $D = 6, 8$
- $\Gamma[\tau \rightarrow hadrons_{ud} \nu_\tau]$  alone ( $\leftrightarrow I_{spec;V+A}^{w(00)}(m_\tau^2)$ ) insufficient to fix  $\alpha_s, C_6, C_8$
- ALEPH, OPAL approach
  - add  $s_0 = m_\tau^2, (km) = (10), (11), (12), (13)$  “spectral weight” FESRs [ $w(y) \rightarrow y^m (1-y)^k w_{(00)}(y)$ ]
  - neglect (in ppl present)  $D = 10, \dots, 16$  contribs
  - $\alpha_s, \langle aG^2 \rangle, C_6, C_8$  fitted to 5 integral set

- NOTE: ALEPH  $C_6, C_8$  input to most other analyses
- Potential problem: single  $s_0 (= m_\tau^2) \Rightarrow D > 8$  (if non-negligible) distort  $D = 0, 4, 6, 8$  fit parameters
- Test for possible symptoms (systematic  $s_0$ -dependence problems) using “fit qualities”

$$F_T^w(s_0) \equiv \left[ I_{spec;T}^w(s_0) - I_{OPE;T}^w(s_0) \right] / \delta I_{spec;T}^w(s_0)$$

- FIGURE:  $F_V^w(s_0)$  for ALEPH data, OPE fit, and 3  $w_{(k,m)}$  used in ALEPH/OPAL fit, PLUS 3 other degree 3  $w(y)$  (to provide independent  $C_{6,8}$  tests)



- OPE-spectral mismatch  $\Rightarrow$  either a problem with assumption that  $D > 8$  negligible, or OPE breakdown (either way a problem for extracted  $\alpha_s$ )

## A MODIFIED ANALYSIS

- V, A and V+A,  $w_N(y) \equiv 1 - \frac{N}{N-1}y + \frac{1}{N-1}y^N$  FESRs  
[KM,T. Yavin, PRD78 (2008) 094020 (arXiv:0807.0650)]
- single unsuppressed  $D = 2N + 2 > 4$  contrib ( $N \geq 2$ ),  
 $(-1)^N C_{2N+2} / [(N-1)s_0^N]$
- $1/s_0^{N+1}$  scaling c.f.  $D = 0 \Rightarrow$  joint  $\alpha_s, C_{2N+2}$  fit
- $1/(N-1)$   $D = 2N + 2$  suppression, no  $D = 0$  suppression  $\Rightarrow$  MUCH better  $\alpha_s$  emphasis than  $w_{(k,m)}$  set

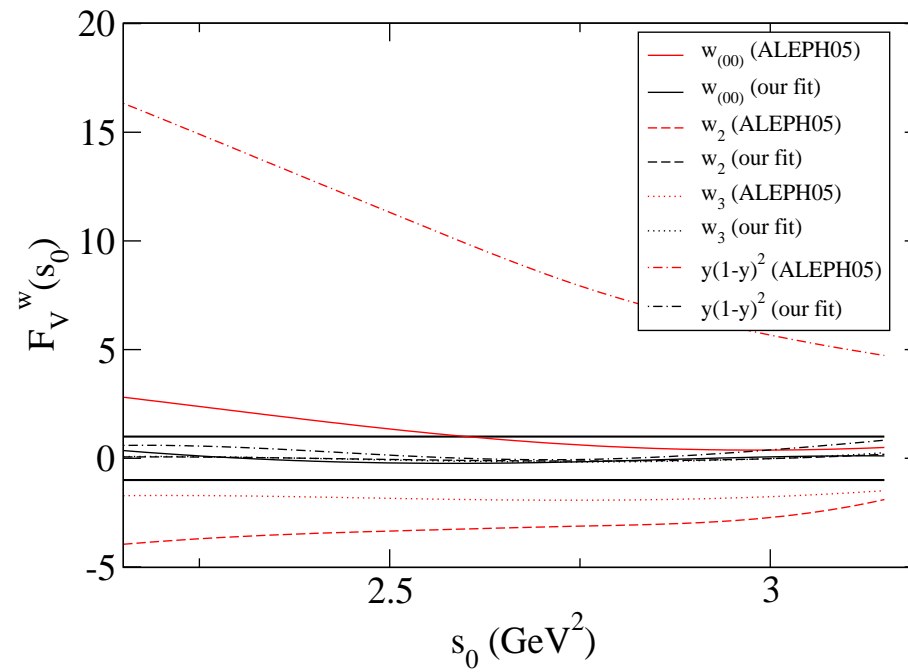
## RESULTS

- Results for  $\alpha_s(m_\tau^2)$  using the CIPT  $D = 0$  prescription

$w(y)$	ALEPH V+A	OPAL V+A
$w_2$	0.320(5)(12)	0.322(7)(12)
$w_3$	0.320(5)(12)	0.322(7)(12)
$w_4$	0.320(5)(12)	0.322(7)(12)
$w_5$	0.320(5)(12)	0.322(7)(12)
$w_6$	0.320(5)(12)	0.322(8)(12)

$w(y)$	ALEPH V	ALEPH A	ALEPH V+A
$w_2$	0.321(7)(12)	0.319(6)(12)	0.320(5)(12)
$w_3$	0.321(7)(12)	0.319(6)(12)	0.320(5)(12)
$w_4$	0.321(7)(12)	0.319(6)(12)	0.320(5)(12)
$w_5$	0.321(7)(12)	0.319(6)(12)	0.320(5)(12)
$w_6$	0.321(7)(12)	0.319(6)(12)	0.320(5)(12)

- Much improved  $F_V^w(s_0)$  for  $w = w_N$  c.f.  $w = w_{(k,m)}$



- CIPT  $w_2, \dots, w_6$  fit values consistent to  $\pm 0.0001$
- Averaging ALEPH and OPAL based results with non-normalization component of error  $\Rightarrow$

$$\alpha_s^{(n_f=3)}(m_\tau) = 0.3209(46)_{exp}(118)_{th}$$

- standard self-consistent combination of 4-loop running, 3-loop matching at flavor thresholds  $\Rightarrow$

$$\alpha_s^{(n_f=5)}(M_Z) = 0.1187(3)_{evol}(6)_{exp}(15)_{th}$$

## CONCLUSIONS/SUMMARY/PROSPECTS

- Lattice ( $\log(W_{11})$  to be specific) and  $\tau$  determinations now in excellent agreement

$$[\alpha_s(M_Z)]_{latt} = 0.1185(8), 0.1190(11)$$

$$[\alpha_s(M_Z)]_{\tau} = 0.1187(16)$$

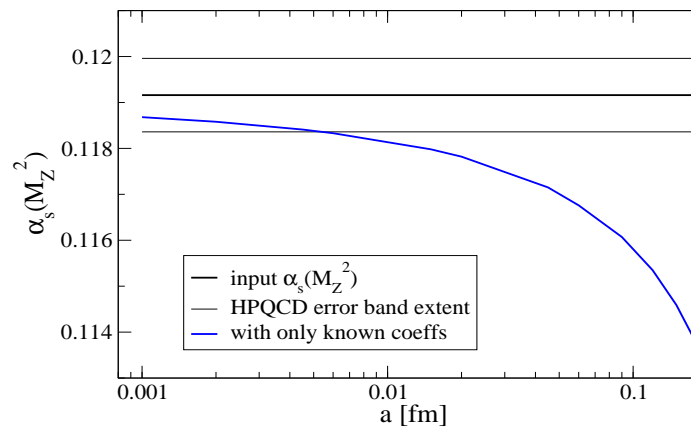
- Future prospects:
  - Significant improvement to lattice errors difficult
  - Some improvement in  $\tau$  decay analysis probable



## MORE ON FUTURE PROSPECTS

- The lattice analysis case:
  - further self-consistency checks from additional  $a \sim 0.045$  fm MILC ensembles, BUT  $a$  small enough to avoid fitting HO  $D = 0$  coefficients impractical

$\alpha_s(M_Z^2)$  with only known vs with fitted HO coefficients



– dominant overall scale-setting error, residual HO  $D = 0$  issues hence difficult to improve significantly

- The  $\tau$  decay analysis case:

Significant improvement requires better understanding of  $D = 0$  truncation uncertainty and residual duality violation (if any)

– Theory error currently dominant ( $\sim 2.5$  times expt'l)

–  $D = 0$  truncation largest theory error source (for  $|FOPT - CIPT| \oplus O(a^5)$  estimate  $\sim 0.010$  of  $0.012$  total)  $\Rightarrow$  important bottleneck for future improvements (though reducible by suitable weight choice)

- Recent explorations *a la* Beneke-Jamin, Caprini-Fischer (taking into account divergent nature of  $D = 0$  series) promising for eventual reduction of  $D = 0$  truncation uncertainty
- Constraining duality violation, and related issues:
  - \* Further constraints on Regge-inspired Cata, Goltermann, Peris DV model now demonstrated [KM+CGP], with simultaneous  $\langle aG^2 \rangle$  determination
  - \*  $\langle aG^2 \rangle$  determination relevant because
    - $\langle aG^2 \rangle$  determination proved NOT feasible using  $w(y)$  with integrated DV sufficiently suppressed to neglect

- Second largest theory error source in MY08 from input  $\langle aG^2 \rangle$
- $\langle aG^2 \rangle$  renormalon ambiguity  $\Rightarrow$  should be determined simultaneously with corresponding truncated  $D = 0$  series
- \* Preliminary results [KM+CGP] show additional constraints push maximum DV impact on  $\alpha_s$  determination to low end of former CGP range ( $\sim 0.0004$  on  $\alpha_s(M_Z)$ )
- Non-trivial error reduction thus appears feasible, though not yet explicitly demonstrated

*HAPPY BIRTHDAY TONY!*

(Canadian English version)

*HAPPY BIRTHDAY TOYNY!*

(SA English version, to the Canadian ear)

**GREAT TO SEE YOU BACK HERE AT THE CSSM!**