

Phase structure of finite density QCD with a histogram method

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for

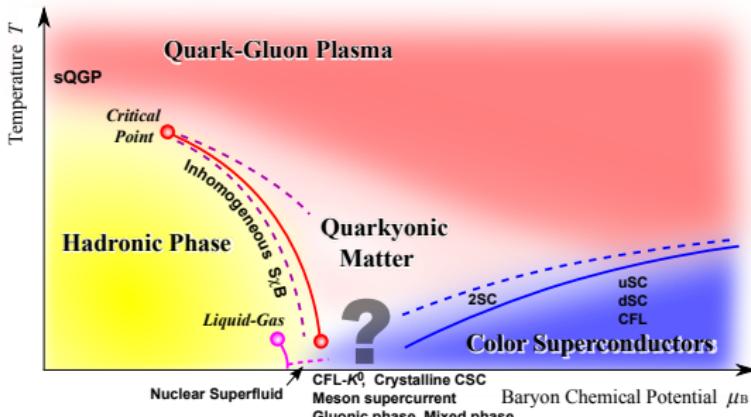
WHOT-QCD collaboration:

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— Finite density lattice QCD —



- Various approaches

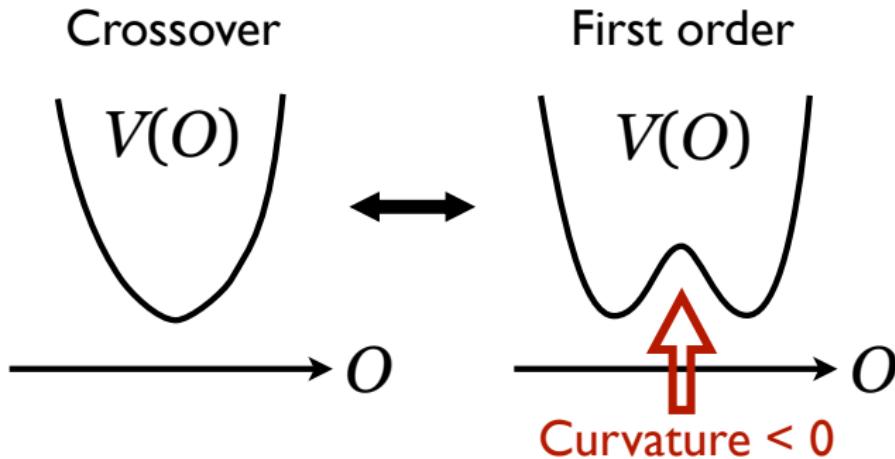
- ✓ Reweighting method
- ✓ Canonical approach
- ✓ Histogram method
- ✓ Taylor expansion
- ✓ Imaginary chemical potential
- ✓ Langevin approach

We explore the phase diagram by
the histogram method
+ phase quenched simulations
+ cumulant expansion for the complex
phase of the quark determinant

— Histogram method and phase transition —

- ✓ Explore the phase diagram by the histogram method

- Label configurations by the value of an operator \hat{O}
- Calculate the histogram $w(O)$ of O and the effective potential $V = -\ln w$
- Change the parameter ($\beta, \kappa, \mu, \dots$) and see if the curvature changes



— Histogram method —

✓ Label gauge configurations by $P = -\frac{S_g}{6\beta N_{\text{site}}}$ and $F(\mu) = N_f \ln \left| \frac{\det M(\mu)}{\det M(0)} \right|$

$$\begin{aligned}\frac{\mathcal{Z}(\beta, \mu)}{\mathcal{Z}(\beta, 0)} &= \frac{1}{\mathcal{Z}(\beta, 0)} \int \mathcal{D}U e^{i\theta(\mu)} |\det M(\mu)|^{N_f} e^{6\beta N_{\text{site}} \hat{P}} \\ &= \int dP dF \textcolor{green}{w(P, F; \beta, \mu)} \left\langle e^{i\theta(\mu)} \right\rangle (P, F; \beta, \mu)\end{aligned}$$

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- ✓ Probability distribution function

$$w(P', F'; \beta, \mu) = \frac{1}{\mathcal{Z}(\beta, 0)} \int \mathcal{D}U \delta(\hat{P} - P') \delta(\hat{F} - F') \underbrace{|\det M(\mu)|^{N_f} e^{6\beta N_{\text{site}} \hat{P}}}_{\text{phase quenched measure}}$$

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- ✓ Complex phase of the quark determinant

$$\langle e^{i\theta(\mu)} \rangle (P', F'; \beta, \mu) = \frac{\langle \langle e^{i\theta(\mu)} \delta(\hat{P} - P') \delta(\hat{F} - F') \rangle \rangle_{(\beta, \mu)}}{\langle \langle \delta(\hat{P} - P') \delta(\hat{F} - F') \rangle \rangle_{(\beta, \mu)}}$$

$\langle \langle \cdot \cdot \cdot \rangle \rangle_{(\beta, \mu)}$: the expectation value with the **phase quenched measure**)

— Sign problem —

- ✓ Poor signal to noise ratio if $e^{i\theta}$ changes sign frequently

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- ✓ Odd terms vanish if the system is invariant under the time reversal
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- ✓ $\mu \rightarrow -\mu$ corresponds to time reversal
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- ✓ The phase factor is **real** and **positive**
- ✓ **No sign problem if the cumulant expansion converges**
- ✓ Need the definition of θ (NOT a Taylor expansion), giving nearly a Gaussian distribution (ideal case)

— Convergence property of the cumulant expansion —

- ✓ We calculate θ as follows:

$$F(\mu) = N_f \ln \left| \frac{\det M(\mu)}{\det M(0)} \right| = N_f \int_0^{\mu/T} \Re \left[\frac{\partial (\ln \det M(\bar{\mu}))}{\partial (\bar{\mu}/T)} \right] d\left(\frac{\bar{\mu}}{T}\right),$$

$$C(\mu, \mu_0) = N_f \ln \left| \frac{\det M(\mu)}{\det M(\mu_0)} \right| = N_f \int_{\mu_0/T}^{\mu/T} \Re \left[\frac{\partial (\ln \det M(\bar{\mu}))}{\partial (\bar{\mu}/T)} \right] d\left(\frac{\bar{\mu}}{T}\right),$$

$$\theta(\mu) = N_f \Im \left[\ln \det M(\mu) \right] = N_f \int_0^{\mu/T} \Im \left[\frac{\partial (\ln \det M(\bar{\mu}))}{\partial (\bar{\mu}/T)} \right] d\left(\frac{\bar{\mu}}{T}\right)$$

- ✓ The complex phase, (the absolute value of) the quark determinant, and the reweighting factor can be obtained as continuous functions of μ .

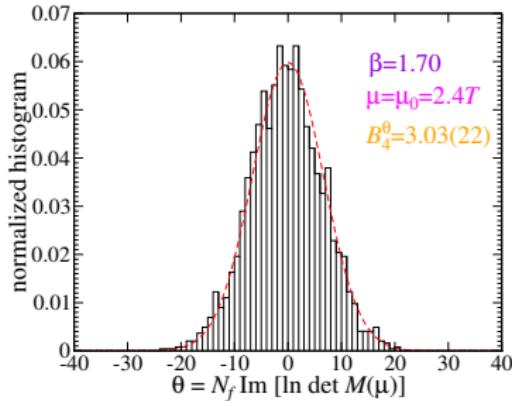
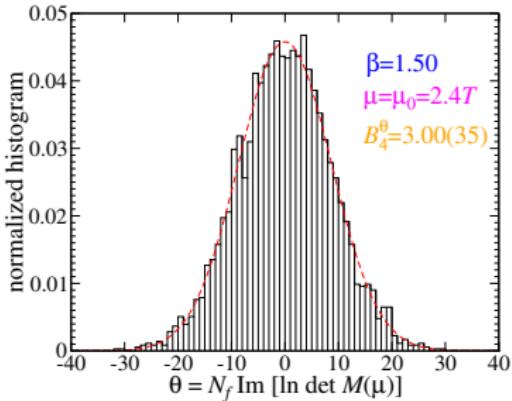
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— Lattice setup —

- ✓ Clover-improved Wilson quark action
- ✓ RG-improved Iwasaki action
- ✓ On $8^3 \times 4$ lattice, $N_f = 2$, $\kappa = 0.141139$ ($m_{\text{PS}}/m_V = 0.8$ for $\beta = 1.80$)
- ✓ $\beta = 1.30, 1.35, 1.40, \dots, 1.90, 1.95, 2.00$
- ✓ $\mu_0/T = 2.4, 2.8, 3.2, 3.6$
- ✓ Measurement every 10 trajectories
- ✓ Random noise method with 50 noises to calculate the derivatives of $\ln \det M$
- ✓ Statistics is 2900

— Curvature of the effective potential —

- ✓ Ratio of the partition function

$$\frac{\mathcal{Z}(\beta, \mu)}{\mathcal{Z}(\beta, 0)} = \int dP dF \textcolor{red}{w}(P, F; \beta, \mu_0) \left\langle e^{i\theta(\mu)} \right\rangle (P, F; \beta, \mu) = \int dP dF e^{-V(P, F; \beta, \mu)}$$

- ✓ Effective potential

$$V(P, F; \beta, \mu) = -\ln w(P, F; \beta, \mu_0) + \frac{1}{2} \left\langle \theta^2 \right\rangle_c (P, F; \beta, \mu, \mu_0)$$

- ✓ Curvature of the effective potential

$$\frac{\partial^2 V}{\partial P^2} = \frac{\partial^2 (-\ln w)}{\partial P^2} + \frac{1}{2} \frac{\partial^2 \left\langle \theta^2 \right\rangle_c}{\partial P^2}, \quad \frac{\partial^2 V}{\partial F^2} = \frac{\partial^2 (-\ln w)}{\partial F^2} + \frac{1}{2} \frac{\partial^2 \left\langle \theta^2 \right\rangle_c}{\partial F^2}$$

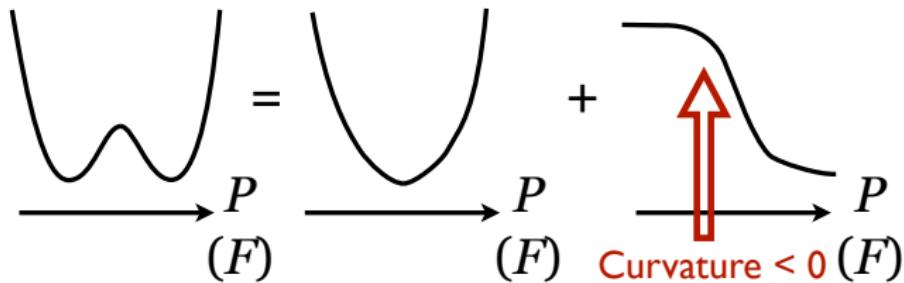
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- ✓ Curvature of the effective potential

$$\frac{\partial^2 V}{\partial P^2} = \frac{\partial^2 (-\ln w)}{\partial P^2} + \frac{1}{2} \frac{\partial^2 \langle \theta^2 \rangle_c}{\partial P^2}, \quad \frac{\partial^2 V}{\partial F^2} = \frac{\partial^2 (-\ln w)}{\partial F^2} + \frac{1}{2} \frac{\partial^2 \langle \theta^2 \rangle_c}{\partial F^2}$$

— Curvature of (- log of) the histogram —

- ✓ Assume a Gaussian distribution for P and F around $\langle P \rangle_{(\beta, \mu)}$ and $\langle F \rangle_{(\beta, \mu)}$,

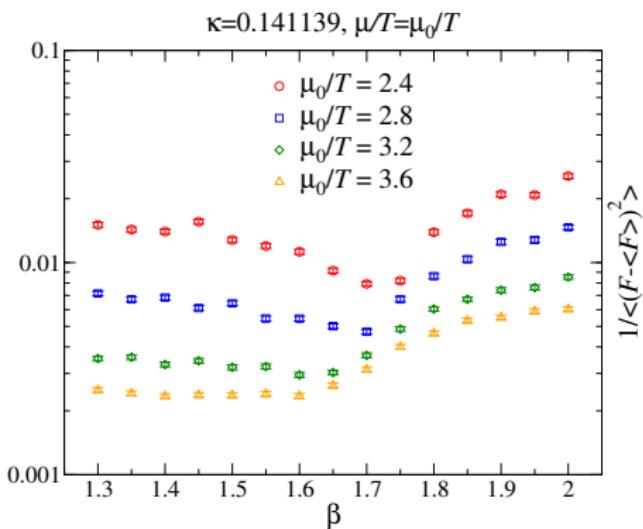
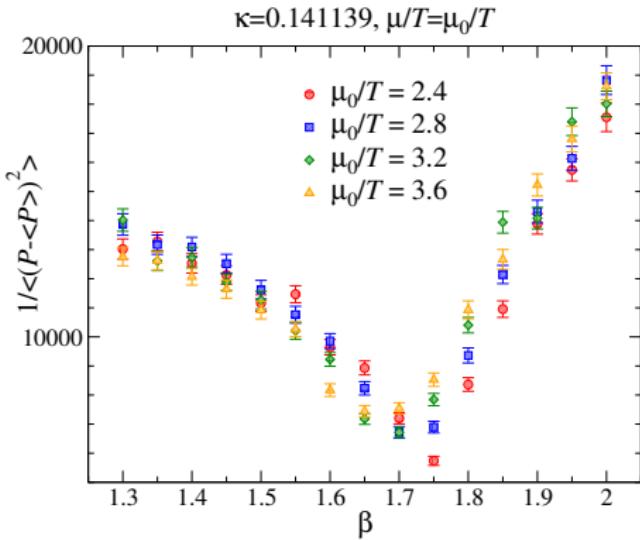
$$w(P, F) = \frac{1}{\sqrt{2\pi}\chi_P} \frac{1}{\sqrt{2\pi}\chi_F} \exp \left[-\frac{1}{2\chi_P}(P - \langle P \rangle)^2 - \frac{1}{2\chi_F}(F - \langle F \rangle)^2 \right]$$

$$\chi_P \equiv \langle (P - \langle P \rangle)^2 \rangle, \quad \chi_F \equiv \langle (F - \langle F \rangle)^2 \rangle$$

- ✓ Curvature of (the log of) the histogram

$$\frac{\partial^2(-\ln w)}{\partial P^2} = \frac{1}{\chi_P} = \frac{1}{\langle (P - \langle P \rangle)^2 \rangle}, \quad \frac{\partial^2(-\ln w)}{\partial F^2} = \frac{1}{\chi_F} = \frac{1}{\langle (F - \langle F \rangle)^2 \rangle}$$

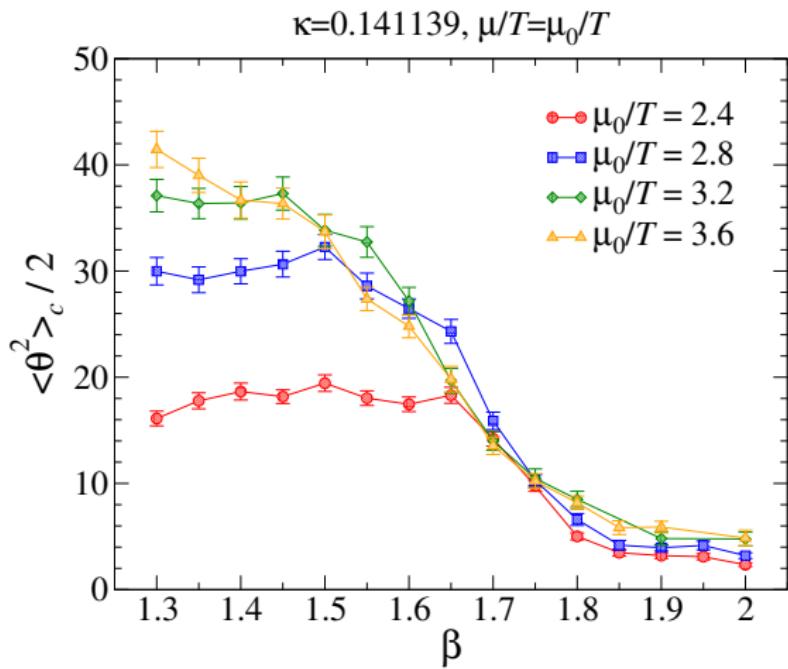
— Curvature of the distribution function —



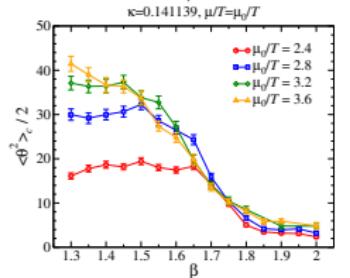
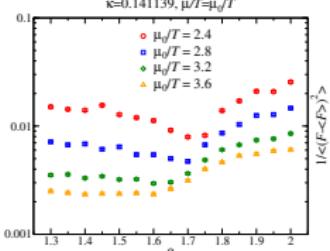
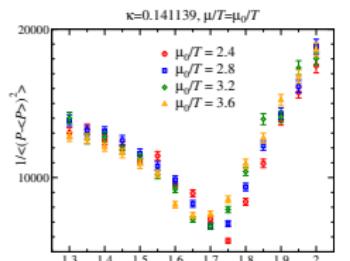
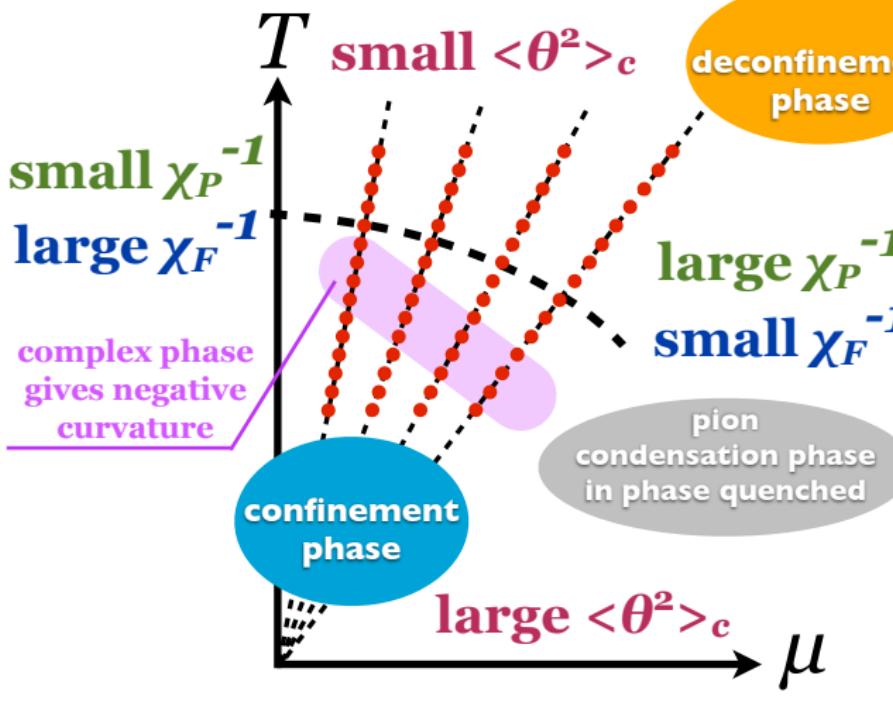
$$\frac{\partial^2(-\ln w)}{\partial P^2} = \frac{1}{\chi_P} = \frac{1}{\langle(P - \langle P \rangle)^2\rangle}$$

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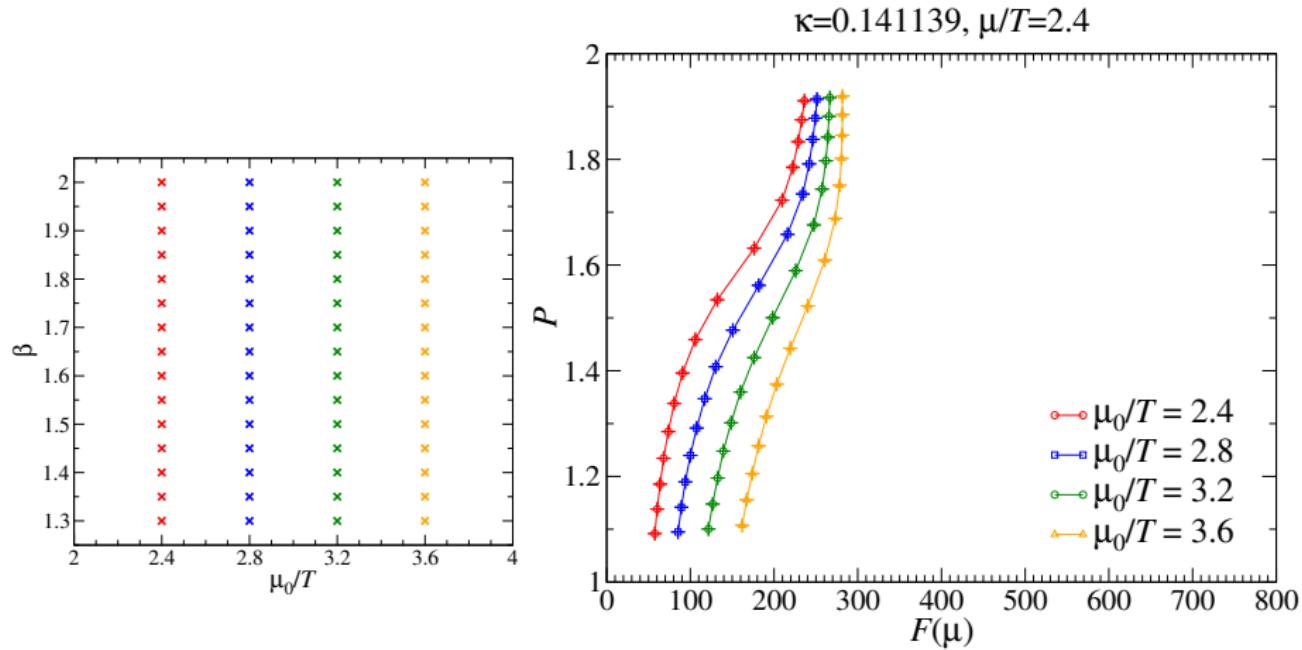
— Second order cumulant of the complex phase —



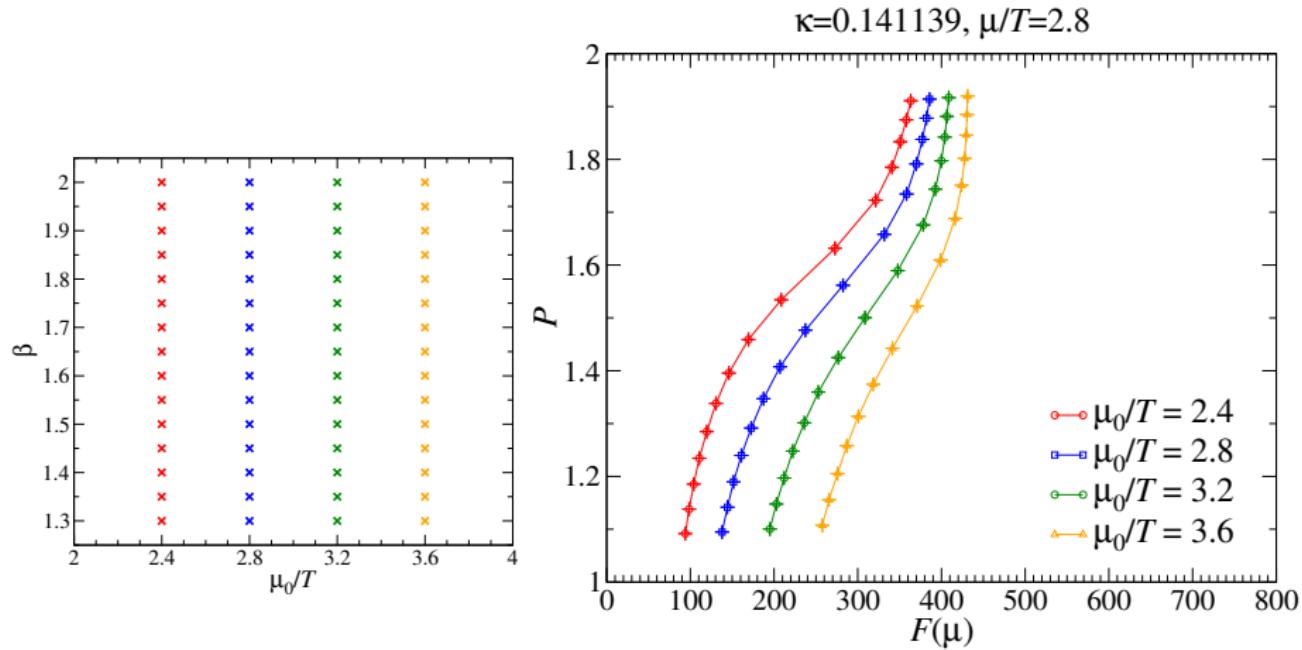
— χ_P , χ_F , and $\langle \theta^2 \rangle_c$ in T - μ plane —



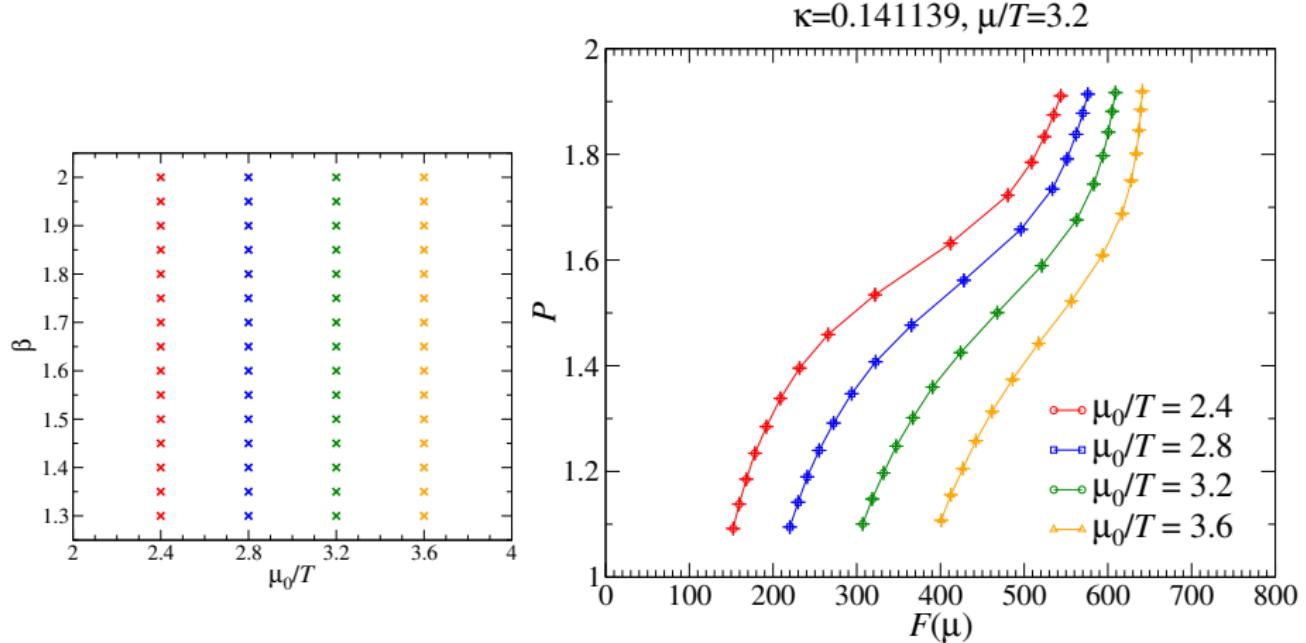
— Average of P and $F(\mu)$ —



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— Reweighting technique —

- ✓ Ratio of the partition function

$$\frac{\mathcal{Z}(\beta, \mu)}{\mathcal{Z}(\beta, 0)} = \int dPdF w(P, F; \beta, \mu) \left\langle e^{i\theta(\mu)} \right\rangle (P, F; \beta, \mu)$$

⇒ We need $w(P, F)$ and $\langle e^{i\theta} \rangle (P, F)$ in a wide range of P and F

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⇒ We need $w(P, F)$ and $\langle e^{i\theta} \rangle (P, F)$ in a wide range of P and F

- ✓ Overlap problem → simulations at various simulation points

$$w(F; \beta, \mu) = R(P, F; \beta, \mu, \mu_0) w(P, F; \beta, \mu_0)$$
$$R(F; \beta, \mu, \mu_0) = \frac{\left\langle \left\langle \left| \frac{\det M(\mu)}{\det M(\mu_0)} \right|^{N_f} \delta(\hat{P} - P') \delta(\hat{F} - F') \right\rangle \right\rangle_{(\beta, \mu_0)}}{\left\langle \left\langle \delta(\hat{P} - P') \delta(\hat{F} - F') \right\rangle \right\rangle_{(\beta, \mu_0)}}$$
$$\langle e^{i\theta(\mu)} \rangle (P', F'; \mu, \mu_0) = \frac{\left\langle \left\langle e^{i\theta(\mu)} \left| \frac{\det M(\mu)}{\det M(\mu_0)} \right|^{N_f} \delta(\hat{P} - P') \delta(\hat{F} - F') \right\rangle \right\rangle_{(\beta, \mu_0)}}{\left\langle \left\langle \left| \frac{\det M(\mu)}{\det M(\mu_0)} \right|^{N_f} \delta(\hat{P} - P') \delta(\hat{F} - F') \right\rangle \right\rangle_{(\beta, \mu_0)}}$$

— Reweighting technique —

- ✓ Complex phase of the quark determinant

$$\left\langle e^{i\theta(\mu)} \right\rangle (P', F'; \mu, \mu_0) = \frac{\left\langle \left\langle e^{i\theta(\mu)} \left| \frac{\det M(\mu)}{\det M(\mu_0)} \right|^{N_f} \delta(\hat{P} - P') \delta(\hat{F} - F') \right\rangle \right\rangle_{(\beta, \mu_0)}}{\left\langle \left\langle \left| \frac{\det M(\mu)}{\det M(\mu_0)} \right|^{N_f} \delta(\hat{P} - P') \delta(\hat{F} - F') \right\rangle \right\rangle_{(\beta, \mu_0)}}$$

- ✓ $F = N_f \ln \left| \frac{\det M(\mu)}{\det M(0)} \right|^{N_f}$ and $C = N_f \ln \left| \frac{\det M(\mu)}{\det M(\mu_0)} \right|^{N_f}$ are strongly correlated, and assume $C \propto F$

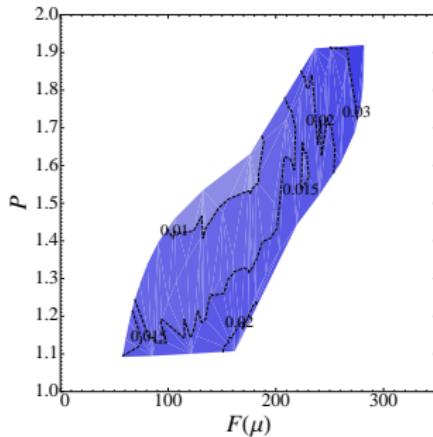
$$\left\langle e^{i\theta(\mu)} \right\rangle (P', F'; \mu, \mu_0) \approx \frac{\left\langle \left\langle e^{i\theta(\mu)} \delta(\hat{P} - P') \delta(\hat{F} - F') \right\rangle \right\rangle_{(\beta, \mu_0)}}{\left\langle \left\langle \delta(\hat{P} - P') \delta(F - F') \right\rangle \right\rangle_{(\beta, \mu_0)}}$$

— Curvature of (- log of) the distribution function —

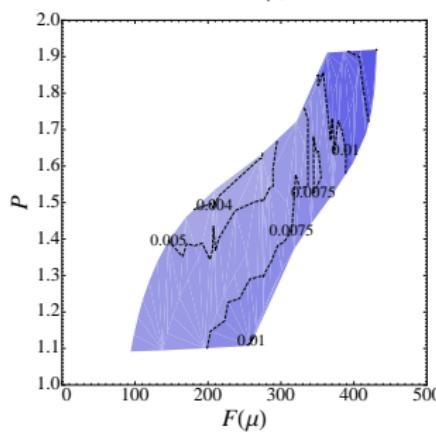
- ✓ Curvature of the distribution function in F direction

$$\frac{\partial^2(-\ln w)}{\partial F^2}$$

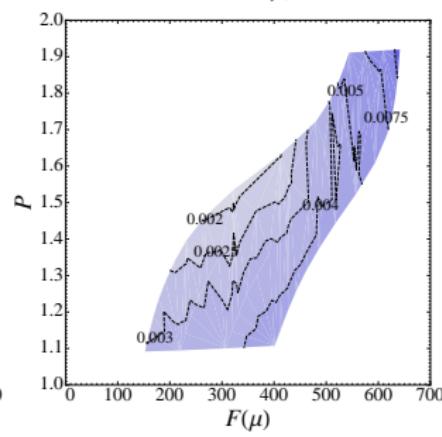
$\kappa=0.144139, \mu/T=2.4$



$\kappa=0.144139, \mu/T=2.8$



$\kappa=0.144139, \mu/T=3.2$



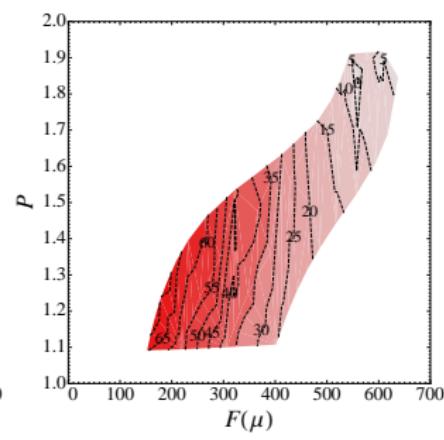
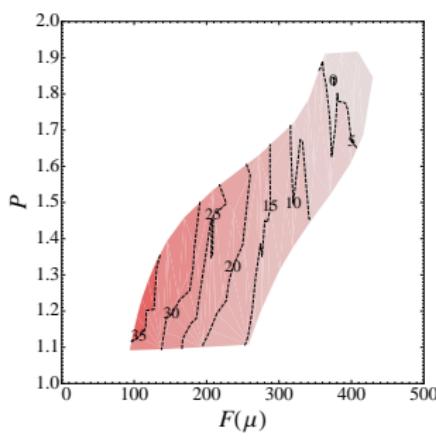
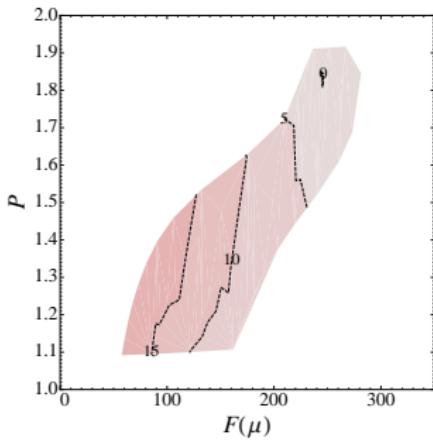
- ✓ Curvature of the distribution function in F direction decreases with increasing μ/T .

— Second order cumulant $\langle \theta^2 \rangle_c$ —

$$\frac{\langle \theta^2 \rangle_c}{2} \text{ at } \kappa=0.144139, \mu/T=2.4$$

$$\frac{\langle \theta^2 \rangle_c}{2} \text{ at } \kappa=0.144139, \mu/T=2.8$$

$$\frac{\langle \theta^2 \rangle_c}{2} \text{ at } \kappa=0.144139, \mu/T=3.2$$

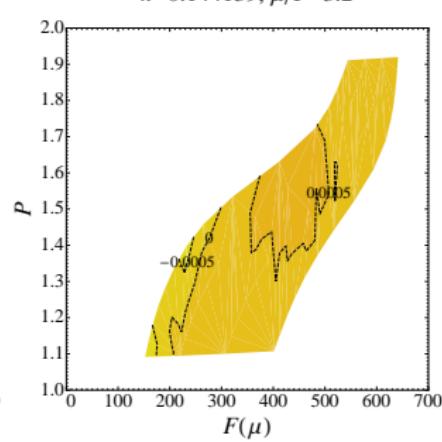
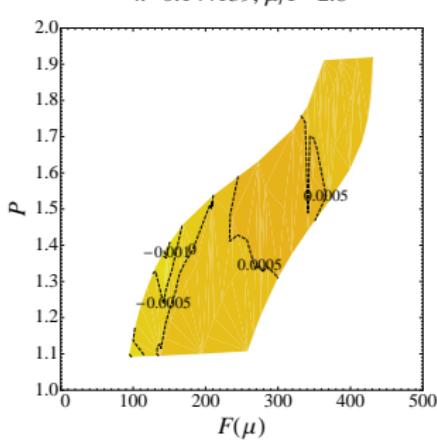
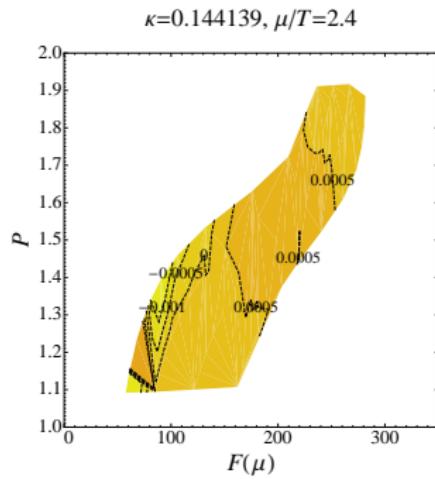


- ✓ The second order cumulant of the complex phase takes large values in small P and F region while small in large P and F region.

— Curvature of the second order cumulant $\langle \theta^2 \rangle_c$ —

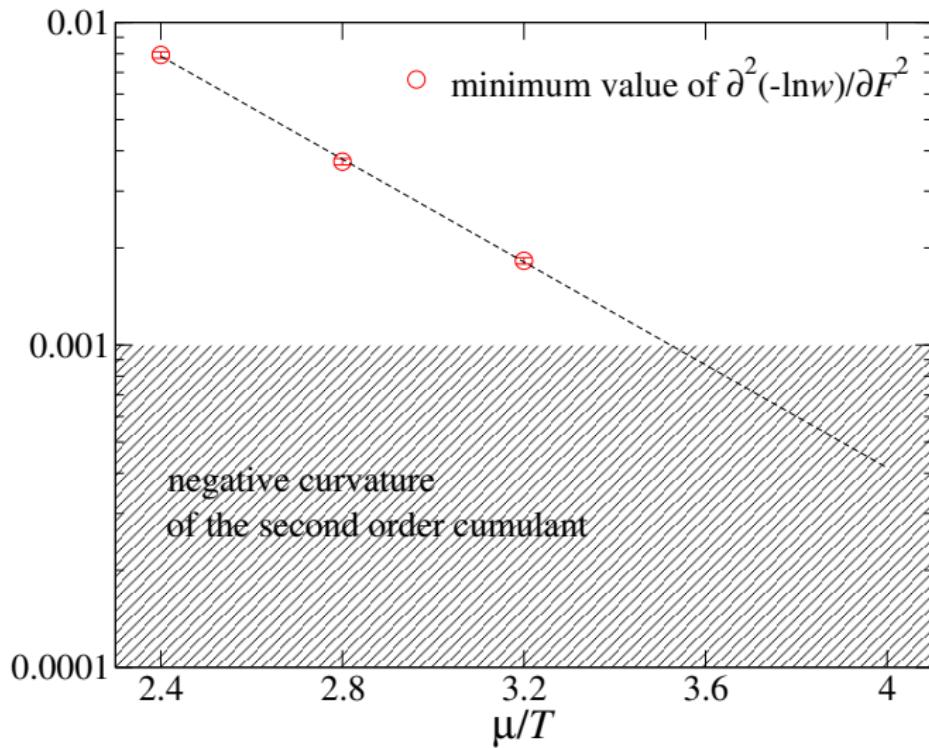
- ✓ Curvature of the second order cumulant $\langle \theta^2 \rangle_c$

$$\frac{1}{2} \frac{\partial^2 \langle \theta^2 \rangle_c}{\partial F^2}$$



- ✓ The second order cumulant has a negative curvature in small P and F region.

— Positive curvature from $(-\ln w)$ and
negative curvature from $\langle \theta^2 \rangle_c$ —



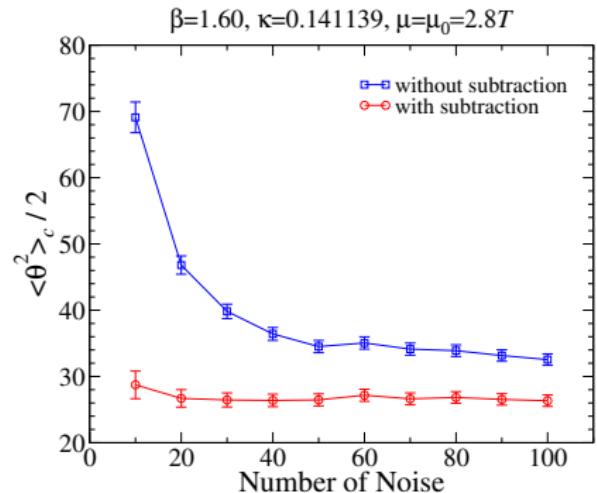
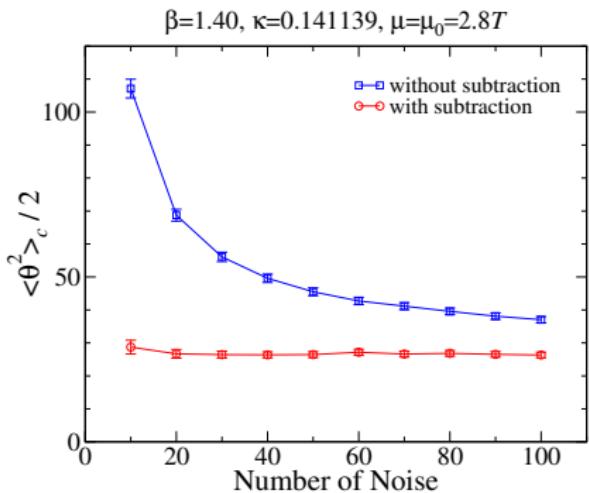
Summary and outlook

- ✓ We investigated the phase structure of finite density QCD by the histogram method
- ✓ We adopted the cumulant expansion for the complex phase
- ✓ The complex phase is calculated by the μ -integration of $\partial \ln \det M(\mu) / \partial \mu$
- ✓ The curvature of the distribution function in F direction decreases with increasing μ/T
- ✓ The curvature of the complex phase in P - F plane has a negative region
- ✓ The first order phase transition at $\mu/T \approx 4.0$?

— Noise error in $(\text{Tr } \theta)^2$ —

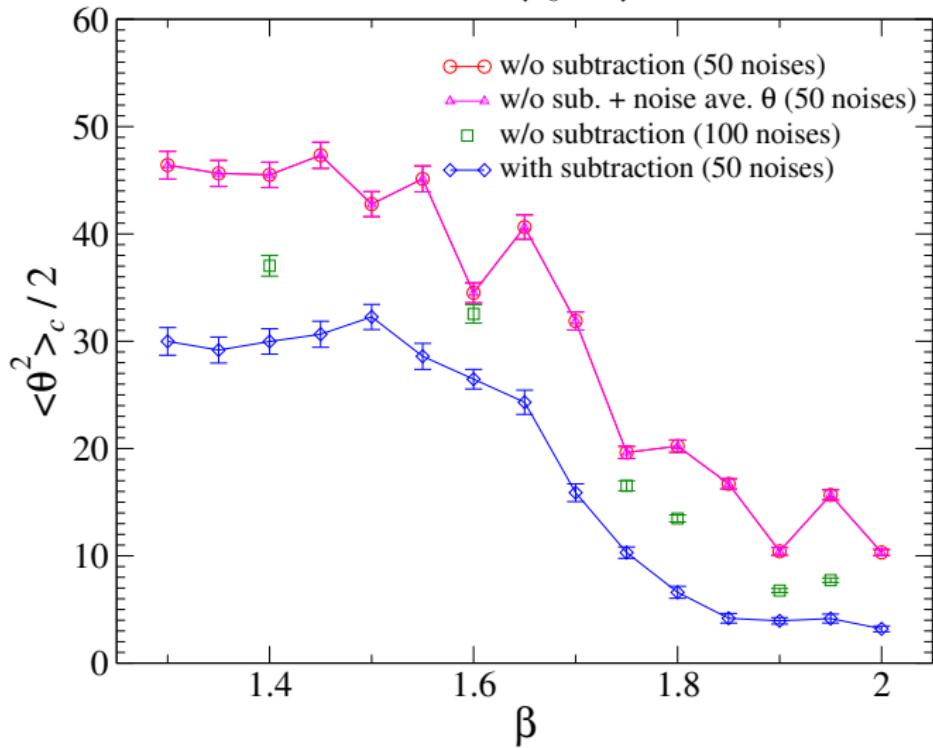
$$(\text{Tr } \theta)^2 = \lim_{N_{\text{noise}} \rightarrow \infty} \left[\left(\frac{1}{N_{\text{noise}}} \sum_{i=1}^{N_{\text{noise}}} \eta_i^\dagger \theta \eta_i \right)^2 - \epsilon^2(\theta) \right]$$

$$\epsilon^2(\theta) = \frac{1}{N_{\text{noise}} - 1} \left[\frac{1}{N_{\text{noise}}} \sum_{i=1}^{N_{\text{noise}}} (\eta_i^\dagger \theta \eta_i)^2 - \left(\frac{1}{N_{\text{noise}}} \sum_{i=1}^{N_{\text{noise}}} \eta_i^\dagger \theta \eta_i \right)^2 \right]$$



— Noise error in $(\text{Tr } \theta)^2$ —

$\kappa=0.141139$, $\mu_0/T=\mu/T=2.8$

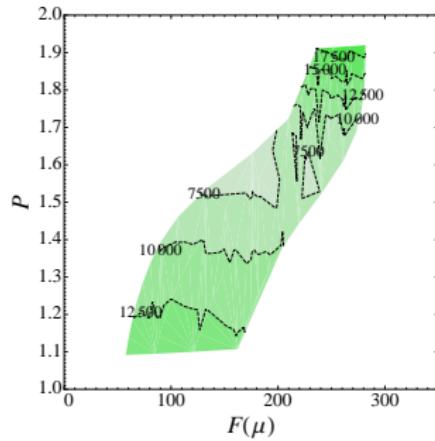


— Curvature of (- log of) the distribution function —

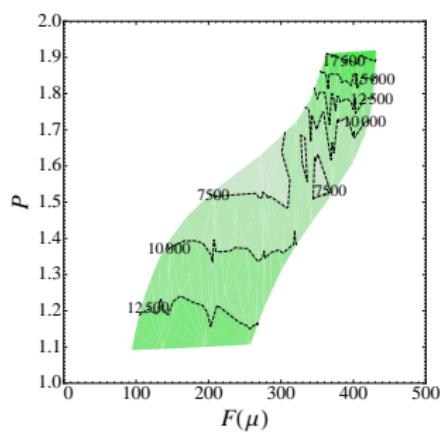
- ✓ Curvature of the distribution function in P direction

$$\frac{\partial^2(-\ln w)}{\partial P^2}$$

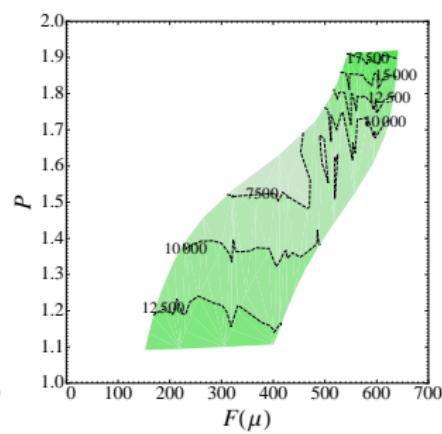
$\kappa=0.144139, \mu/T=2.4$



$\kappa=0.144139, \mu/T=2.8$



$\kappa=0.144139, \mu/T=3.2$

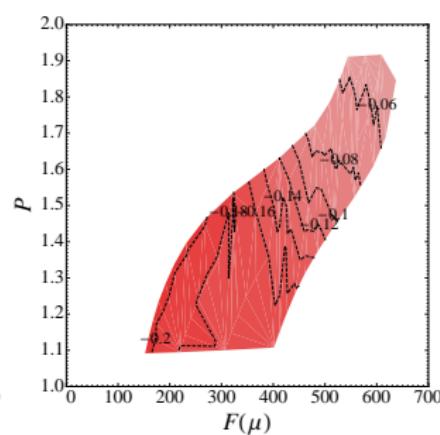
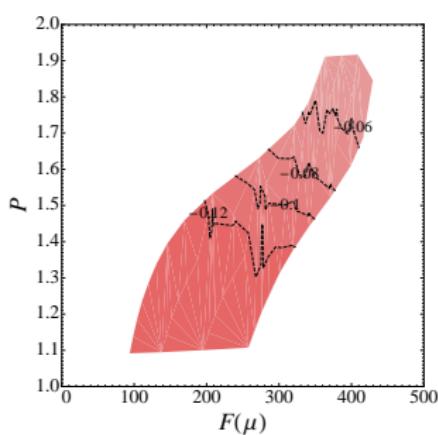
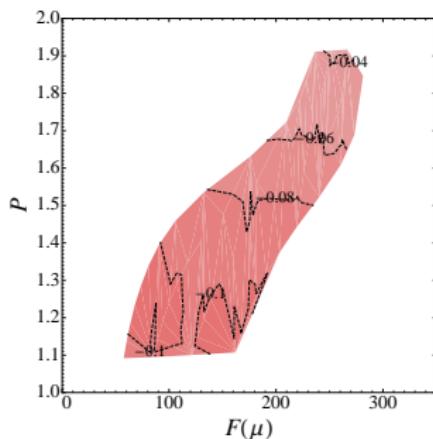


— Slope of the second order cumulant $\langle \theta^2 \rangle_c$ —

$$\frac{1}{2} \frac{\partial \langle \theta^2 \rangle_c}{\partial F} \text{ at } \kappa=0.144139, \mu/T=2.4$$

$$\frac{1}{2} \frac{\partial \langle \theta^2 \rangle_c}{\partial F} \text{ at } \kappa=0.144139, \mu/T=2.8$$

$$\frac{1}{2} \frac{\partial \langle \theta^2 \rangle_c}{\partial F} \text{ at } \kappa=0.144139, \mu/T=3.2$$

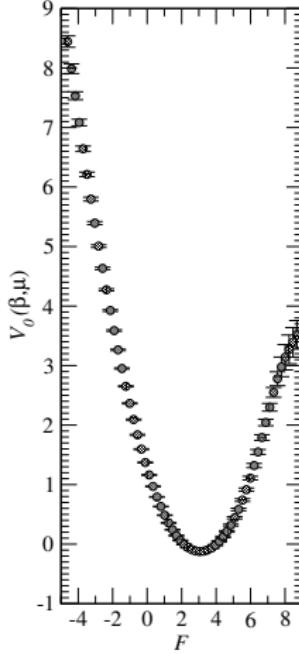
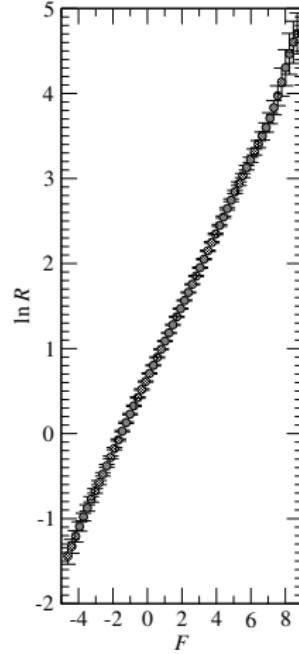
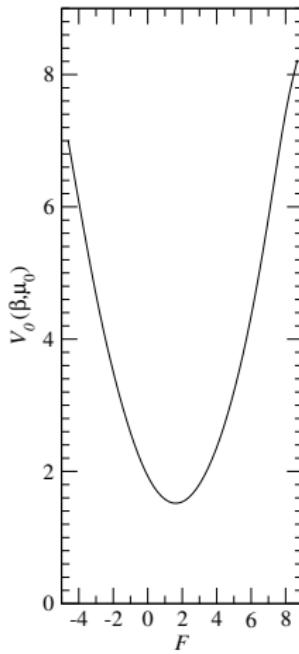


✓ Linear fitting

$$\langle \theta^2 \rangle_c (P, F) \Big|_{(P_0, F_0)} = c_1 + c_2(P - P_0) + c_3(F - F_0)$$

— Distribution function and the reweighting term —

$\beta=1.50, \mu_0/T=0.4, \mu/T=0.8$



— Susceptibility of the plaquette and the Polyakov loop in phased quenched case —

