



# Correlation Matrix Methods for the $\pi$ and $\rho$ Form Factors in Full QCD

#### Benjamin Owen

Waseem Kamleh, Derek Leinweber, M. Selim Mahbub, Ben Menadue & Peter Moran

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#### Outline

- Introduction
  - Correlation Matrix Analysis
  - Extracting Form Factors
- 2 Calculation Details
- Results
  - Meson mass spectrum
  - Meson Form Factors
  - Charge Radii

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• where  $Z_i^{\alpha}$ ,  $Z_j^{\alpha\,\dagger}$  are the couplings of sink operator  $(\chi_i)$  and source operator  $(\chi_j^{\dagger})$  for the state  $\alpha$ 

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use our basis of operators to construct these new operators

$$\phi_{\alpha}^{\dagger}(\vec{p}) = \sum_{i=1}^{N} u_i^{\alpha}(\vec{p}) \chi_i^{\dagger}(\vec{p})$$
$$\phi_{\alpha}(\vec{p}) = \sum_{i=1}^{N} v_i^{\alpha}(\vec{p}) \chi_i(\vec{p})$$

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#### CM Eigenvalue Equation

$$[\mathcal{G}^{-1}(t,\vec{p})\mathcal{G}(t+dt,\vec{p})]_{ij}u_i^{\alpha}(\vec{p}) = \lambda^{\alpha}u_i^{\alpha}(\vec{p})$$
(1a)

$$v_i^{\alpha}(\vec{p})[\mathcal{G}(t+dt,\vec{p})\mathcal{G}^{-1}(t,\vec{p})]_{ij} = \lambda^{\alpha}v_i^{\alpha}(\vec{p})$$
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• Using  $v_i^{\alpha}(\vec{p})$ ,  $u_i^{\alpha}(\vec{p})$  we are able to project out the correlation function for the state  $|\dot{M}_{\alpha}, p\rangle$ 

$$\mathcal{G}_{\alpha}(t,\vec{p}) = v_i^{\alpha}(\vec{p})\mathcal{G}_{ij}(t,\vec{p})u_j^{\alpha}(\vec{p})$$

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- We shall use smearing of the fermion source and sink as a method of increasing our operator basis
- Our basis will comprise standard local operators with some number of sweeps of smearing

 $\{\chi_i\} \to \{\chi_i^n\}, n \text{ number of sweeps of smearing }$ 

 In this work we use of gauge-invariant Gaussian smearing at both source and sink

<sup>&</sup>lt;sup>1</sup>M. S. Mahbub et al., Phys. Rev. D. **80**, 054507 (2009)

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- For further details, refer to paper below

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- Three point function of cross correlators

$$\begin{split} &\mathcal{G}^{\mu}_{ij}(t_{2},t_{1};\vec{p}',\vec{p}) = \sum_{\vec{x}_{1},\vec{x}_{2}} e^{-i\vec{p}'\cdot\vec{x}_{2}} e^{i\vec{q}\cdot\vec{x}_{1}} \langle \Omega | \chi_{i}(x_{2}) J^{\mu}(x_{1}) \chi_{j}^{\dagger}(0) | \Omega \rangle \\ &= \sum_{\alpha=0}^{N-1} \frac{e^{-E_{M_{\alpha}}(\vec{p}')\,t_{2}} e^{-E_{M_{\alpha}}(\vec{q})\,t_{1}}}{2\sqrt{E_{M_{\alpha}}(\vec{p})E_{M_{\alpha}}(\vec{p}')}} Z_{i}^{\alpha}(\vec{p}') Z_{j}^{\alpha\,\dagger}(\vec{p}) \langle M_{\alpha}, p' | J^{\mu}(0) | M_{\alpha}, p \rangle \end{split}$$

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• Crucial thing to note above is that we have source and sink coupling parameters with different momenta

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 In this work we use the following ratio of three- and two-point correlation functions

$$R^{\mu}_{\alpha}(\vec{p}, \vec{p}') = \sqrt{\frac{\mathcal{G}^{\mu}_{\alpha}(t_2, t_1; \vec{p}', \vec{p})\mathcal{G}^{\mu}_{\alpha}(t_2, t_1; \vec{p}, \vec{p}')}{\mathcal{G}_{\alpha}(t_2, \vec{p}')\mathcal{G}_{\alpha}(t_2, \vec{p})}}$$

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- These are  $32^3 \times 64$  lattices with  $\beta=1.9$ , corresponding to a physical lattice spacing of  $0.0907(13)~{\rm fm}$

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- $\bullet$  There are five light quark-masses resulting in pion masses that range from 622~MeV through to 156~MeV
- For our form factor analysis, we will only consider the heaviest quark mass

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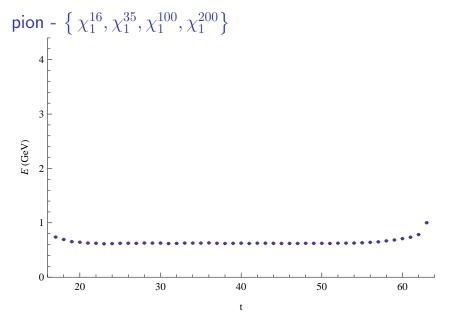
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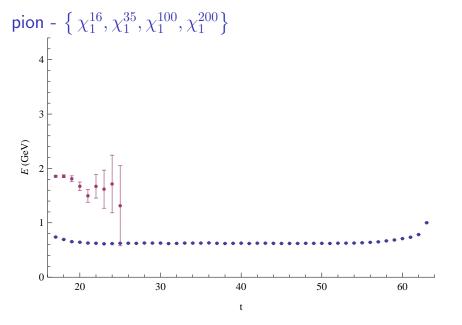
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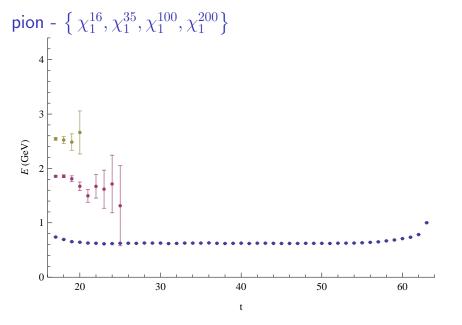
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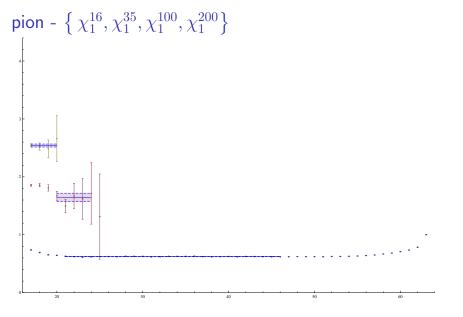
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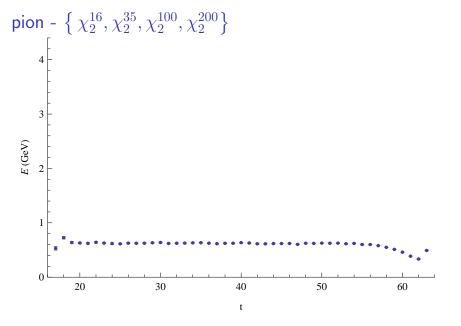
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- ullet Were using  $4 \times 4$  Correlation matrix

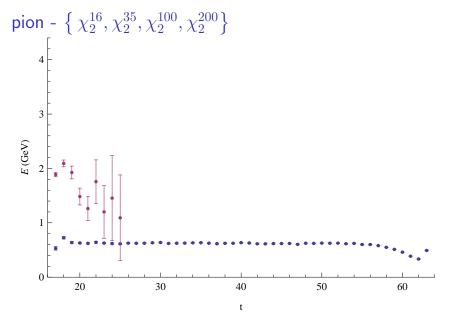


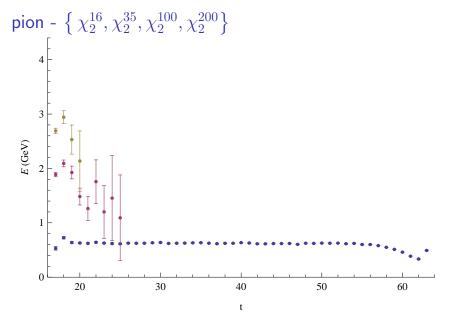


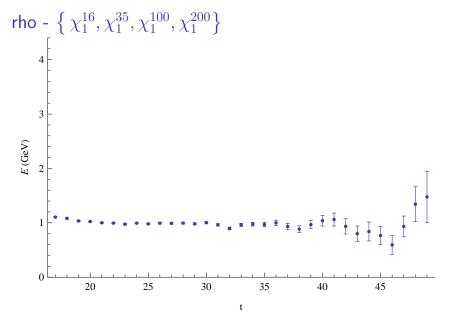


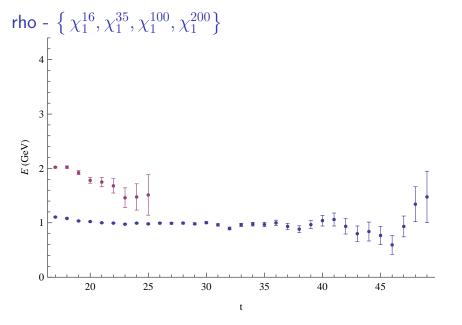


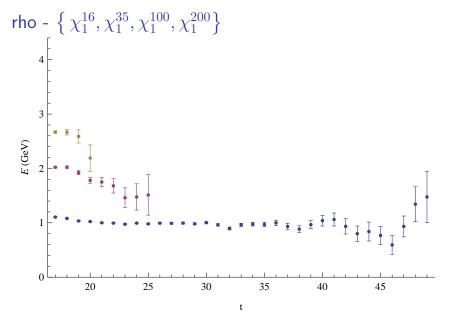


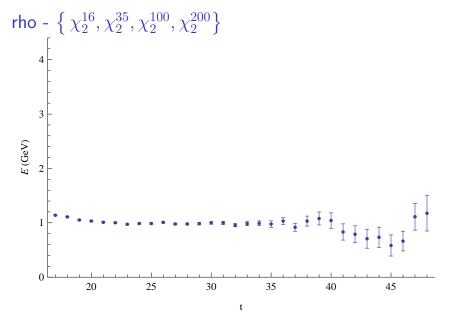


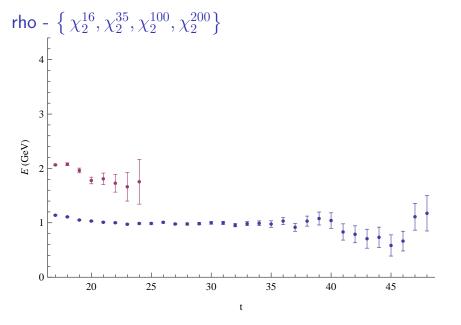


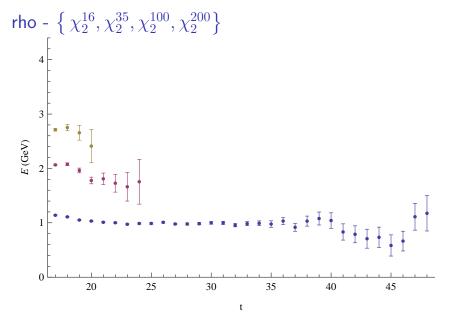


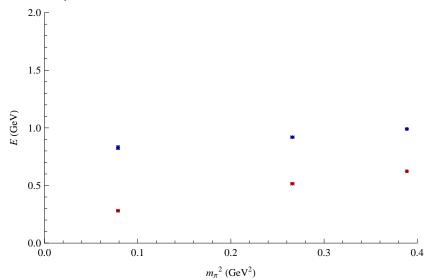


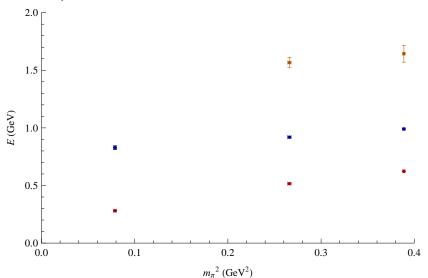


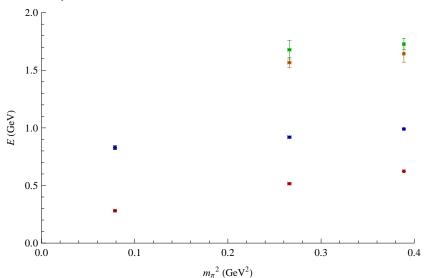


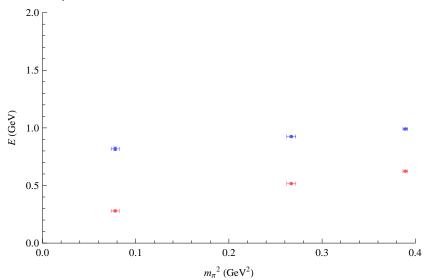


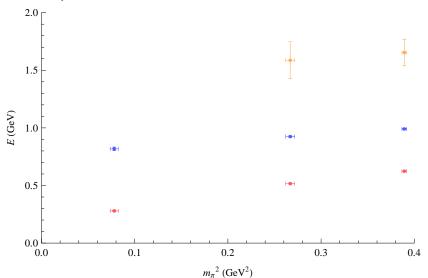


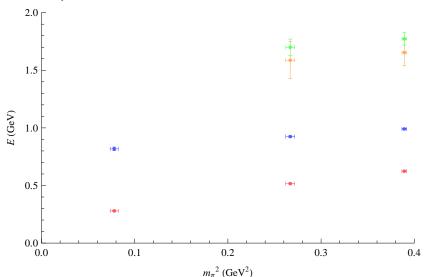


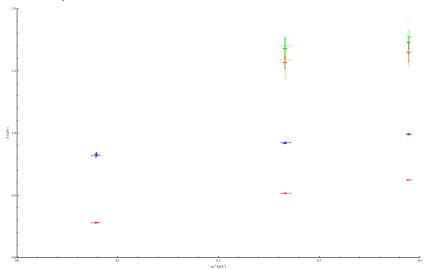












## Form Factor Analysis

 The method and notation in this work closely follows that of a previous study<sup>1</sup> by the CSSM

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#### Vertex Functions

Pion:

$$\langle \pi(\vec{p}') | J^{\mu} | \pi(\vec{p}) \rangle = \frac{1}{\sqrt{E_{\pi}(\vec{p}')E_{\pi}(\vec{p})}} [p'^{\mu} + p^{\mu}] F_{\pi}(Q^2)$$

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Rho:

$$\langle \rho(\vec{p}', s')|J^{\mu}|\rho(\vec{p}, s)\rangle = \frac{1}{\sqrt{E_{\rho}(\vec{p}')E_{\rho}(\vec{p}')}} \epsilon_{\sigma}^{\prime *}(p', s')\Gamma^{\sigma\mu\tau}(p, p')\epsilon_{\tau}(p, s)$$

where

$$\Gamma^{\sigma\mu\tau}(p, p') = -\left\{G_1(Q^2)g^{\sigma\tau}[p'^{\mu} + p^{\mu}] + G_2(Q^2)[g^{\mu\tau}q^{\sigma} - g^{\mu\sigma}q^{\tau}] - G_3(Q^2)q^{\sigma}q^{\tau}\frac{[p'^{\mu} + p^{\mu}]}{2M_{\rho}^2}\right\}$$

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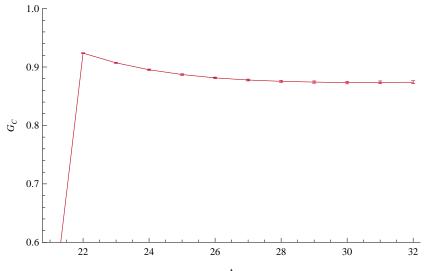
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$$G_C(Q^2) = \frac{2}{3} \frac{\sqrt{E_\rho m_\rho}}{E_\rho + m_\rho} (R_1^{\ 0}_1 + R_2^{\ 0}_2 + R_3^{\ 0}_3)$$

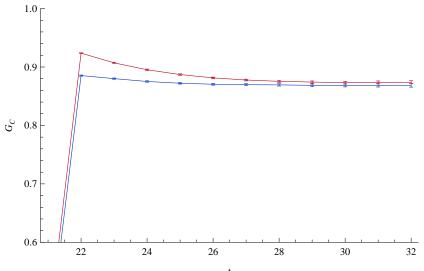
$$G_M(Q^2) = \frac{\sqrt{E_\rho m_\rho}}{p_x} (R_1^{\ 3}_3 + R_3^{\ 3}_1)$$

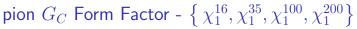
$$G_Q(Q^2) = m_\rho \frac{\sqrt{E_\rho m_\rho}}{p_x^2} (2R_1^{\ 0}_1 - R_2^{\ 0}_2 - R_3^{\ 0}_3)$$

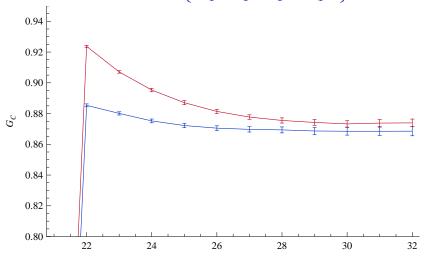




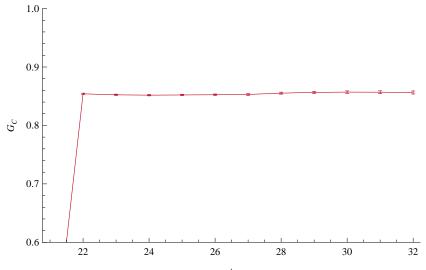




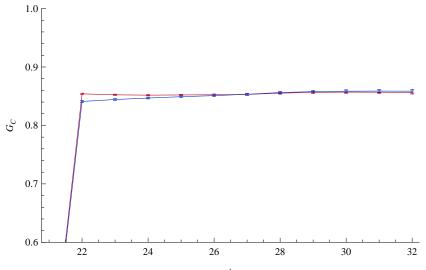




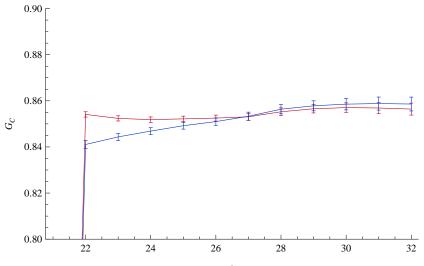


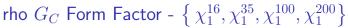


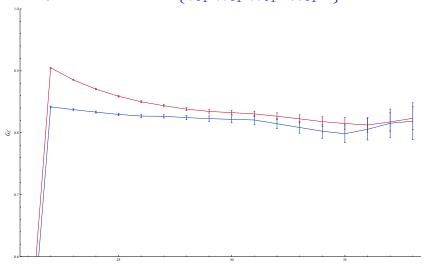


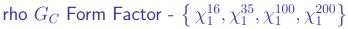


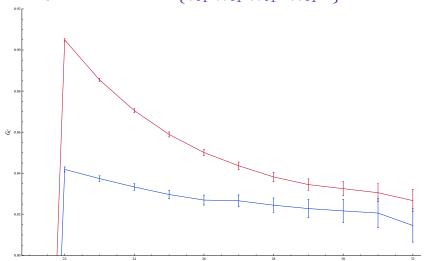


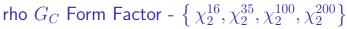


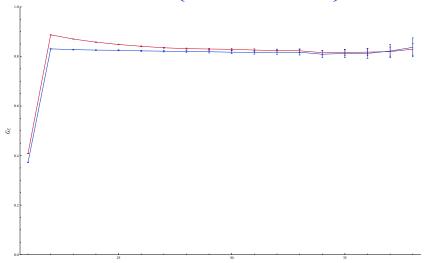


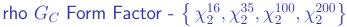


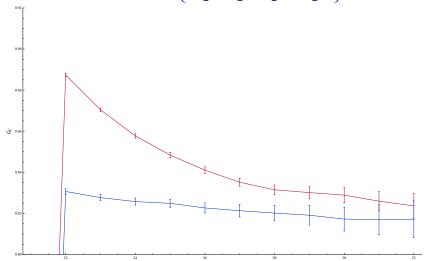


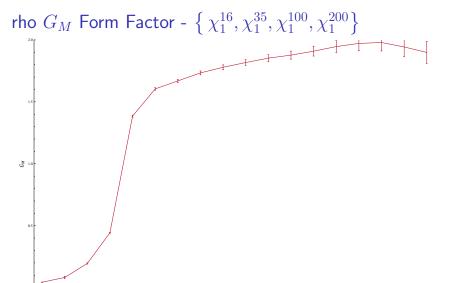


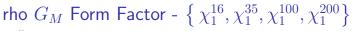


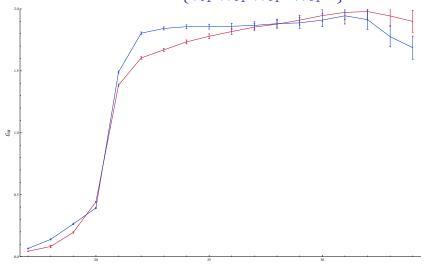


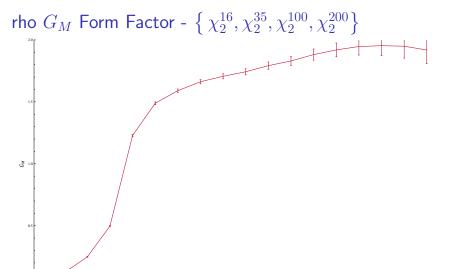


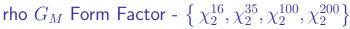


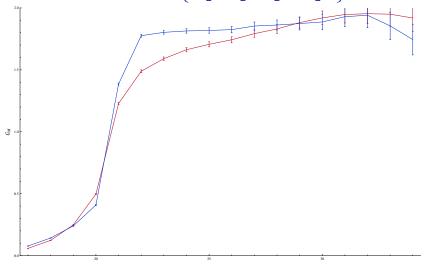




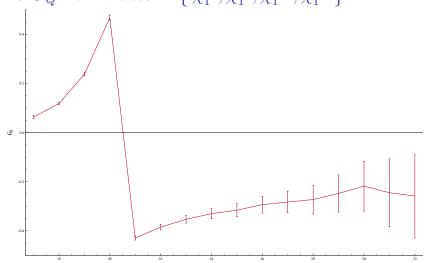


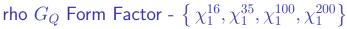


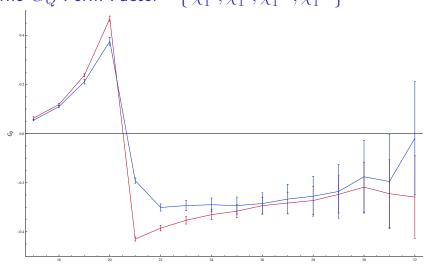


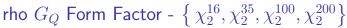


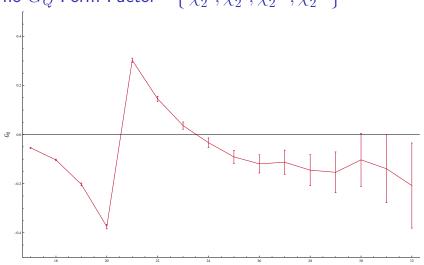


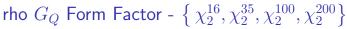


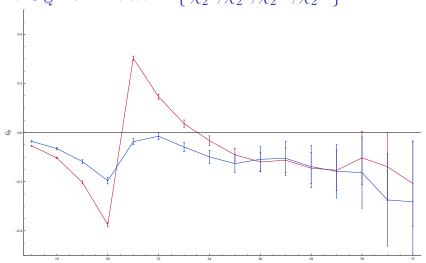












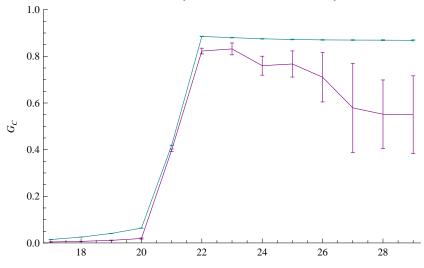
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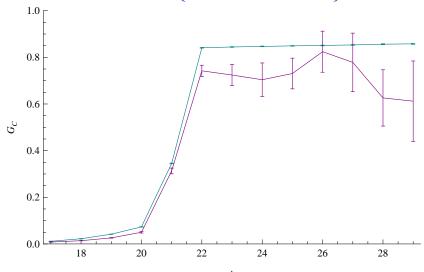
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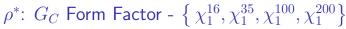
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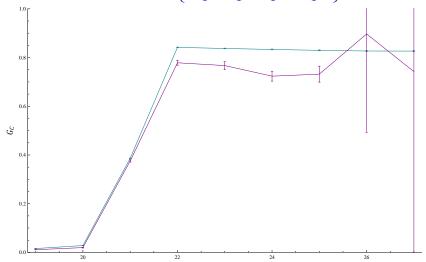
 $\pi^*\colon\thinspace G_C$  Form Factor -  $\left\{\,\chi_1^{16},\chi_1^{35},\chi_1^{100},\chi_1^{200}\right\}$ 

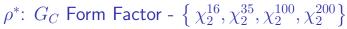


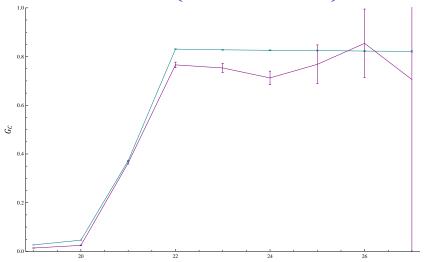
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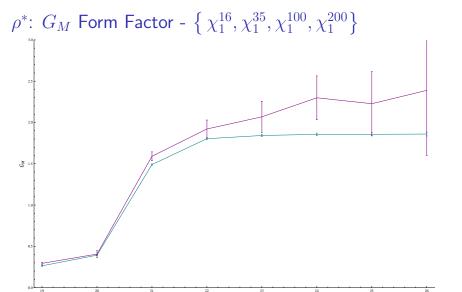


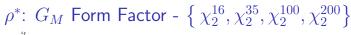


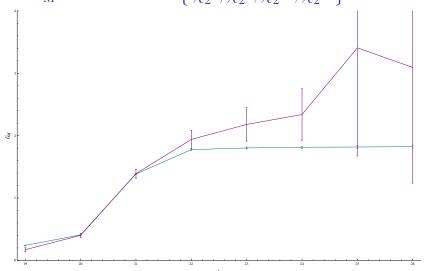












# Charge Radii

Calculate the charge square radius

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#### Ground States

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- $\bullet$  We note that half spatial dimension of our box is  $1.45\,\mathrm{fm}$

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- $\bullet$  Extracted  $\left\langle r^{2}\right\rangle$  for these excited states allowing for comparison with the ground states