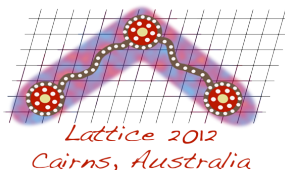


Lattice Determination of the Hadronic Vacuum Polarisation for $(g - 2)_\mu$

Benjamin Jäger



Institute for Nuclear Physics and Helmholtz Institute Mainz,
University of Mainz



In collaboration with A. Jüttner, M. Della Morte and H. Wittig

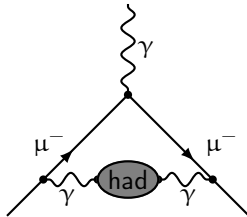
based on arXiv:1112.2894

1 Introduction

- Lattice Calculations

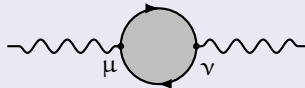
2 Results

- Vacuum Polarization $\Pi(q^2)$
- Hadronic Contribution to a_μ



In the continuum

- Vacuum polarization tensor defined as current-current correlator



$$\Pi_{\mu\nu}(q) = i \int d^4x e^{iqx} \langle J_\mu(x) J_\nu(0) \rangle = (q_\mu q_\nu - q^2 g_{\mu\nu}) \Pi(q^2)$$

- Currently only connected diagram considered
 - Two flavour χ PT $\Rightarrow$ Disconnected diagram $\approx -10\%$ [Jüttner, Della Morte, 2010]
 - Lattice study $\Rightarrow$ Disconnected diagram compatible with 0 (large error bars) [ETMC, 2011]

On the lattice

- $\Pi_{\mu\nu}$ can be expressed in terms of gauge links $U_\mu(n)$ and propagators D_{lat}^{-1}

$$\Pi_{\mu\nu}(q) = a^4 \sum_{n \in \Lambda} e^{iq(n + a\hat{\mu}/2)} \left\langle J_\mu^c(n) J_\nu^l(0) \right\rangle = (q_\mu q_\nu - q^2 \delta_{\mu\nu}) \Pi(q^2)$$

- Use local current J^l at the source and conserved point-split current J^c at sink:
→ **Only 1 inversion needed**, but $\Pi_{\mu\nu}$ needs to be renormalized. [Boyle, et al, 2011]
- Twisted boundary conditions applied to valence quarks

$$\psi(x + L) = \exp\left(i \frac{\Theta_i}{L} x_i\right) \psi(x)$$

⇒ Momentum becomes tunable by Θ_i : $q_i = \frac{2\pi n_i}{L} - \frac{\Theta_i}{L}$ [Sachrajda, Villadoro, 2005]

- Determine α_μ^{had} by convolution integral: $4\alpha^2 \int_0^\infty F\left(\frac{q^2}{m_\mu^2}\right) (\Pi(0) - \Pi(q^2)) dq^2$

Simulation details

- $\mathcal{O}(a)$ improved Wilson fermions (Wilson clover)
- $N_f = 2$ and $N_f = 2 +$ quenched strange
- CLS ensembles:

β	a [fm]	lattice	L [fm]	m_π [MeV]	$m_\pi L$	Labels
5.20	0.079	64×32^3	2.5	473, 363, 312	6.0, 4.7, 4.0	A3, A4, A5
5.30	0.063	64×32^3	2.0	606, 451	6.2, 4.7	E4, E5
5.30	0.063	96×48^3	3.0	324, 277	5.0, 4.2	F6, F7
5.30	0.063	128×64^3	4.0	195	4.0	G8
5.50	0.050	96×48^3	2.4	536, 430, 340	6.5, 5.2, 4.1	N4, N5, N6
5.50	0.050	128×64^3	3.2	270	4.4	O7

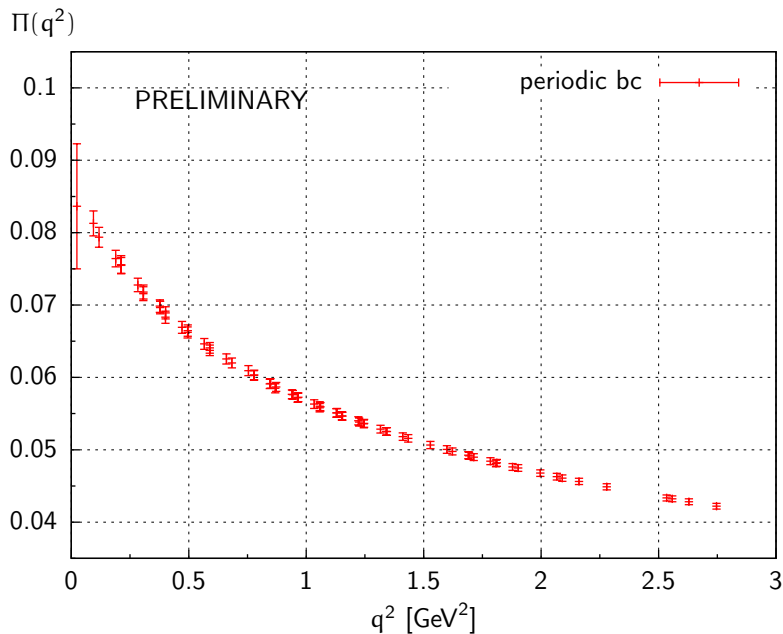
1 Introduction

- Lattice Calculations

2 Results

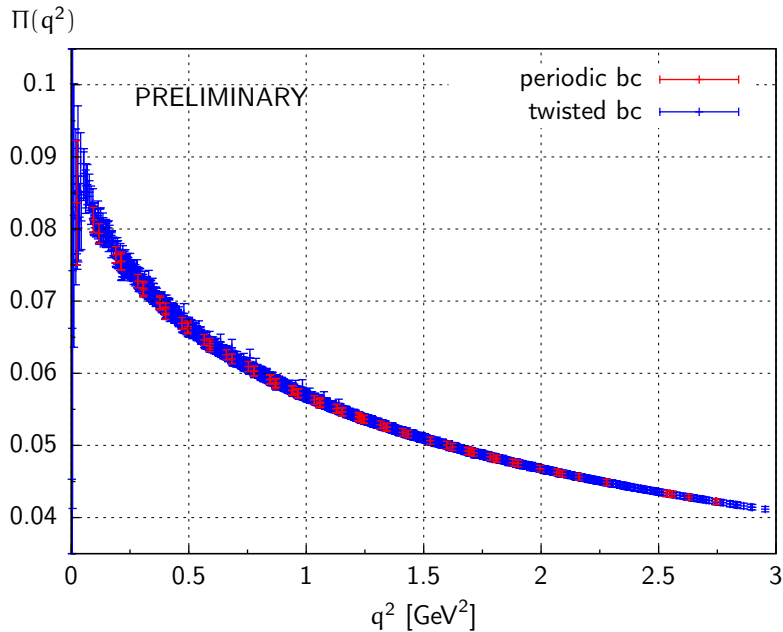
- Vacuum Polarization $\Pi(q^2)$
- Hadronic Contribution to α_μ

Vacuum Polarization $\Pi(q^2)$



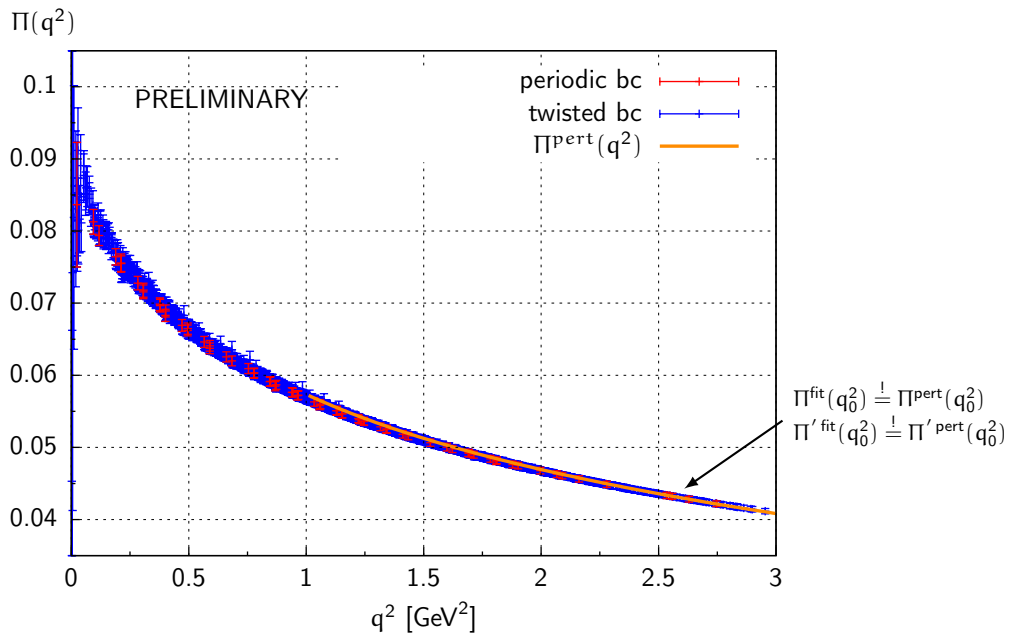
- G8 ensemble: $\beta = 5.3$, $m_{\pi} = 195$ MeV, $L = 4.0$ fm

Vacuum Polarization $\Pi(q^2)$



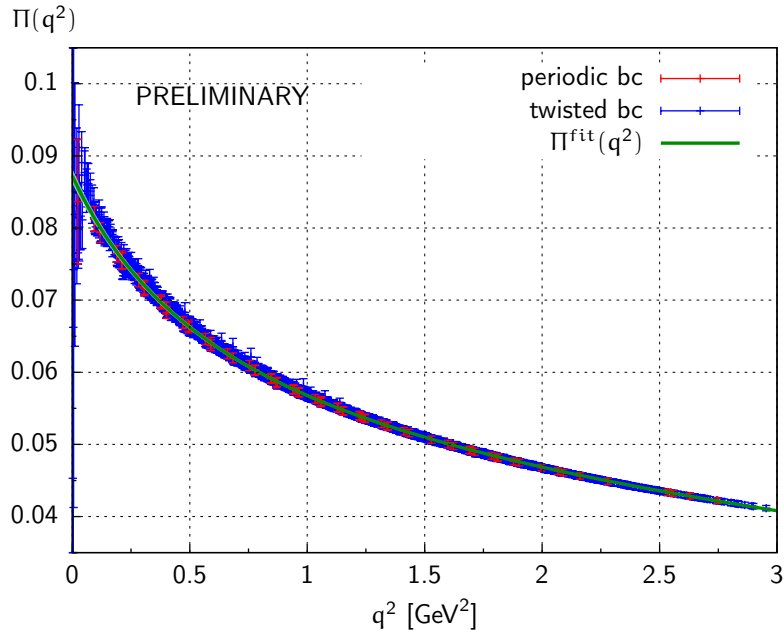
● G8 ensemble: $\beta = 5.3$, $m_{\pi} = 195$ MeV, $L = 4.0$ fm

Vacuum Polarization $\Pi(q^2)$



- 2-loop perturbation theory matched to lattice data at $q_0^2 \approx 2.6$ GeV²

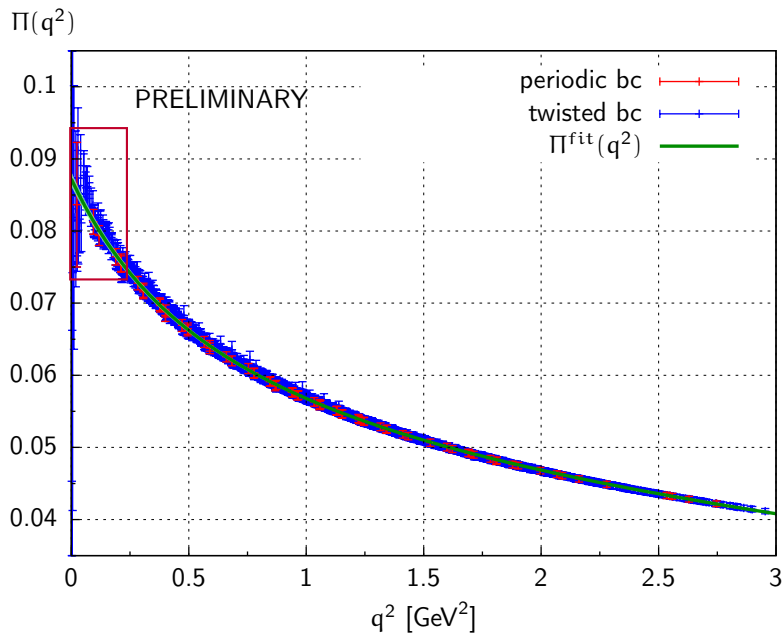
Vacuum Polarization $\Pi(q^2)$



- Fit $\Pi(q^2)$ to well-behaved functions (e.g. Padé)

[Blum, et al, 2012]

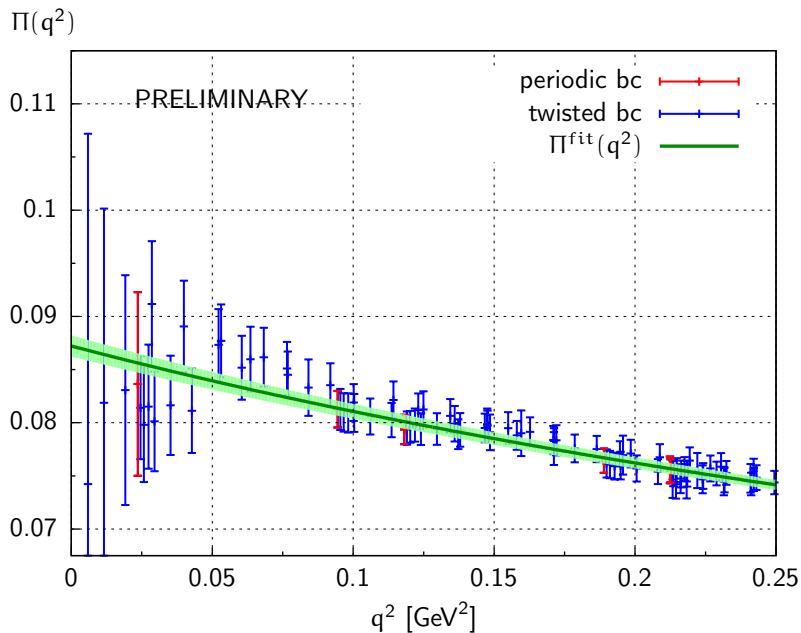
Vacuum Polarization $\Pi(q^2)$



- Fit $\Pi(q^2)$ to well-behaved functions (e.g. Padé)

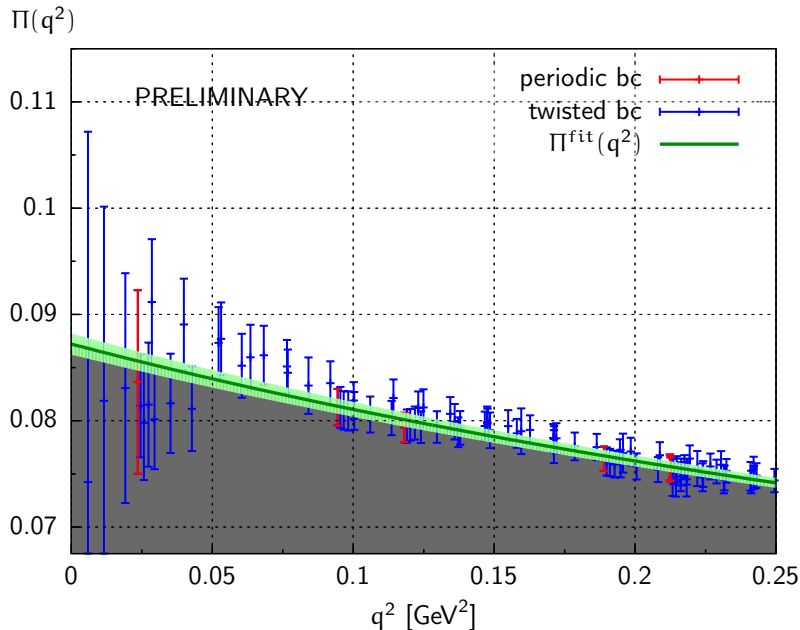
[Blum, et al, 2012]

Vacuum Polarization $\Pi(q^2)$



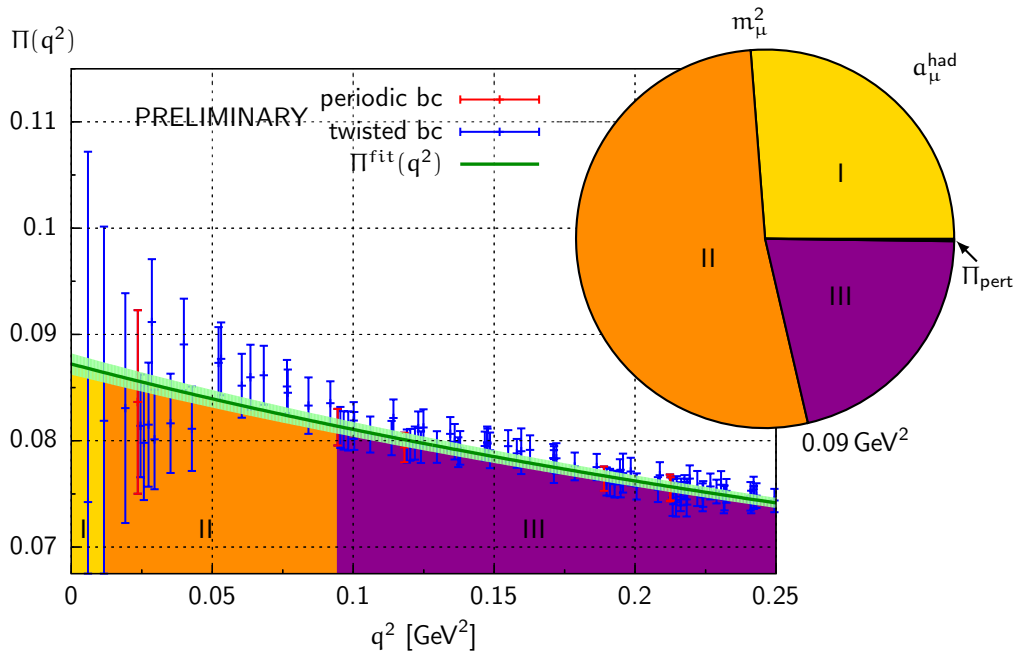
- Data from twisted boundary conditions improve stability of the fit

Vacuum Polarization $\Pi(q^2)$



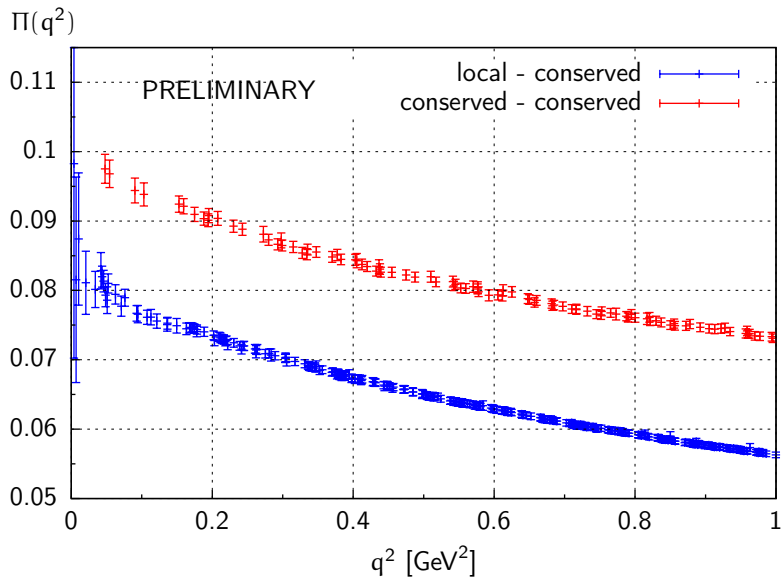
- Determine a_μ^{had} by convolution integral: $a_\mu^{\text{had}} = 4\alpha^2 \int_0^\infty F\left(\frac{q^2}{m_\mu^2}\right) (\Pi(0) - \Pi(q^2)) dq^2$

Vacuum Polarization $\Pi(q^2)$



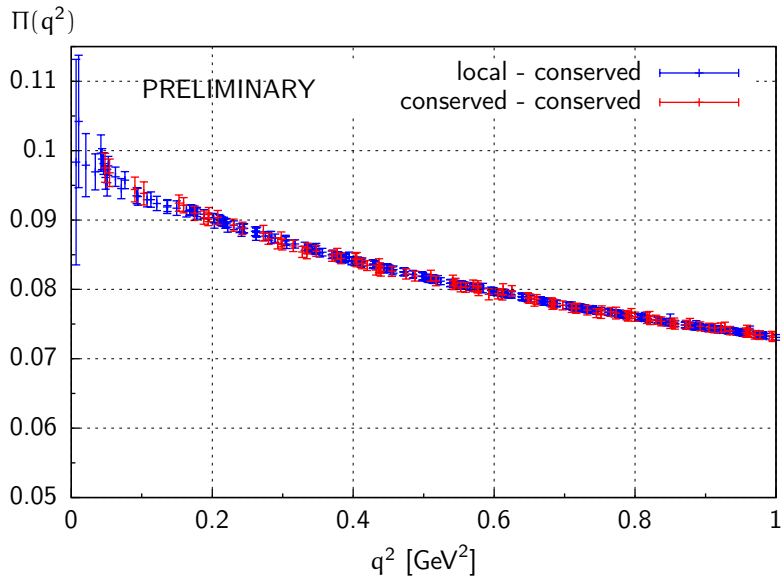
- Twisted boundary conditions improve the crucial low momentum behaviour

Vacuum Polarization $\Pi(q^2)$



● F6 ensemble: $\beta = 5.3$, $m_\pi = 324$ MeV, $L = 3.0$ fm

Vacuum Polarization $\Pi(q^2)$



- Subtracted vacuum polarisation $\hat{\Pi}(q^2)$ is unchanged $\rightarrow \alpha_\mu^{\text{had}}$ remains unchanged

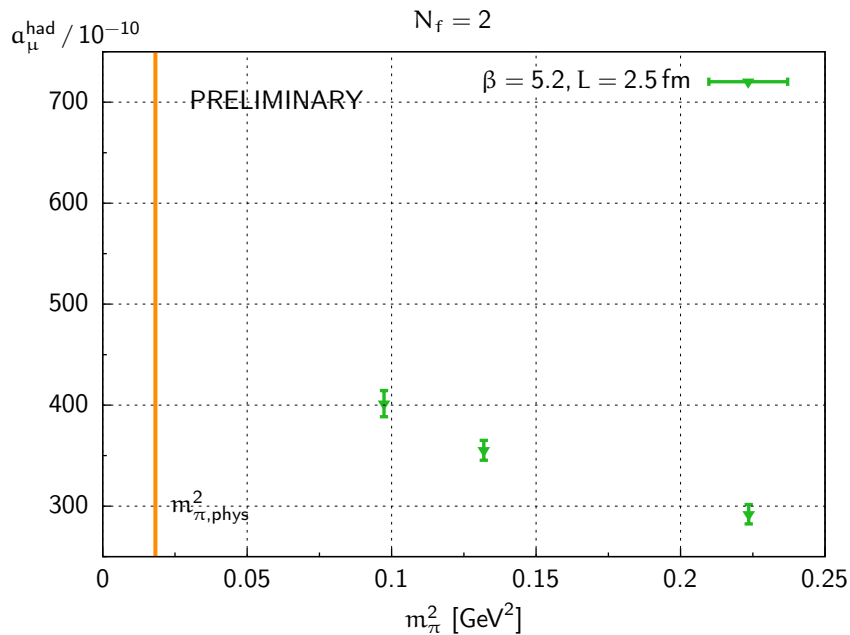
1 Introduction

- Lattice Calculations

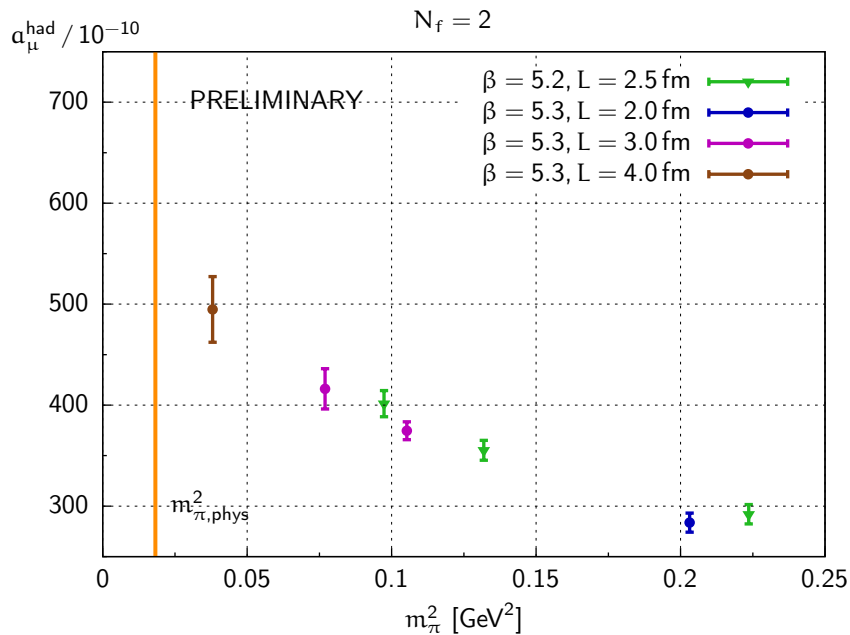
2 Results

- Vacuum Polarization $\Pi(q^2)$
- Hadronic Contribution to a_μ

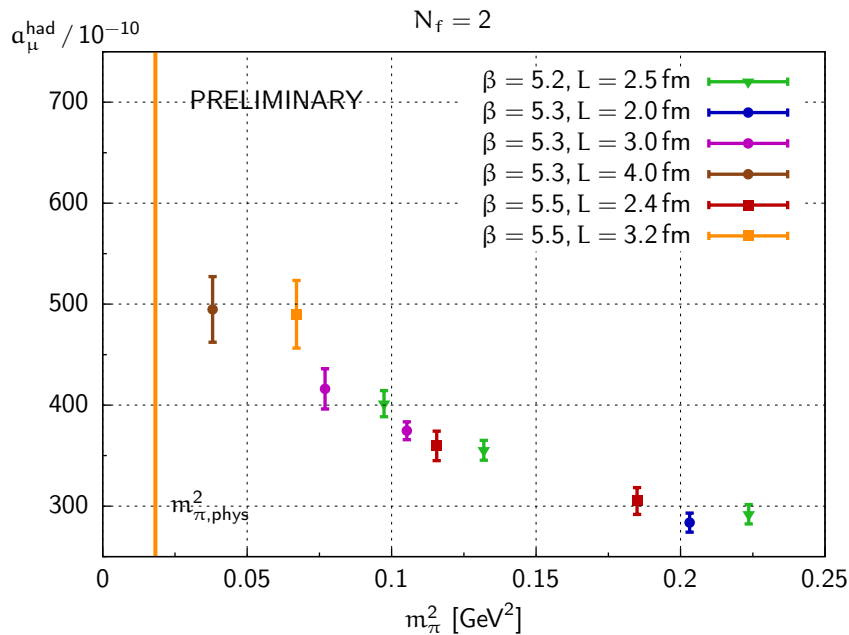
Hadronic Contribution to a_μ for $N_f = 2$



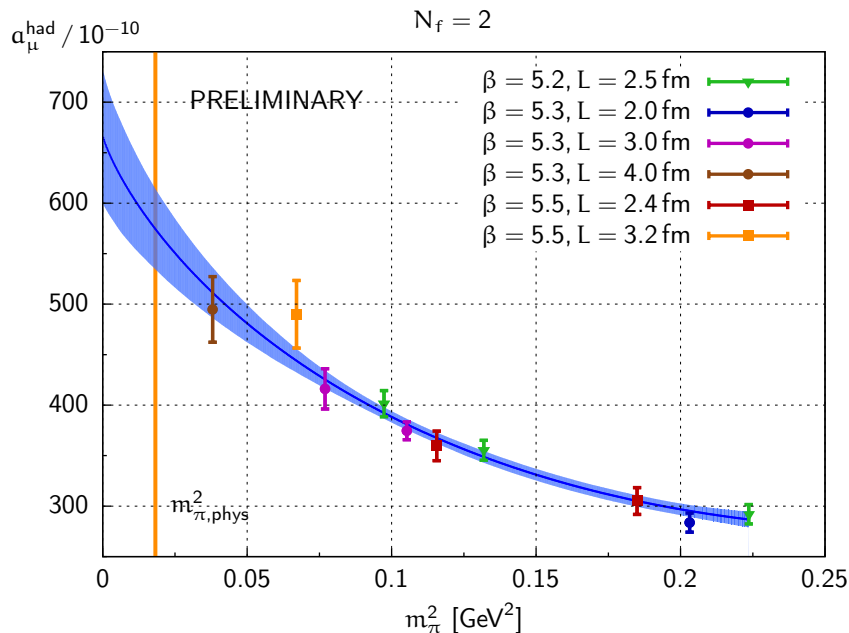
Hadronic Contribution to a_μ for $N_f = 2$



Hadronic Contribution to a_μ for $N_f = 2$

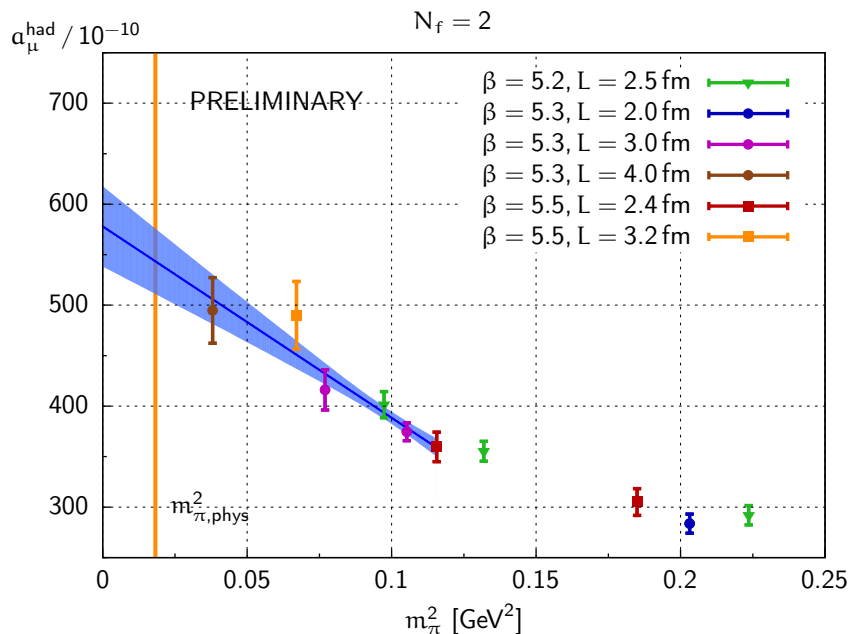


Hadronic Contribution to a_μ for $N_f = 2$



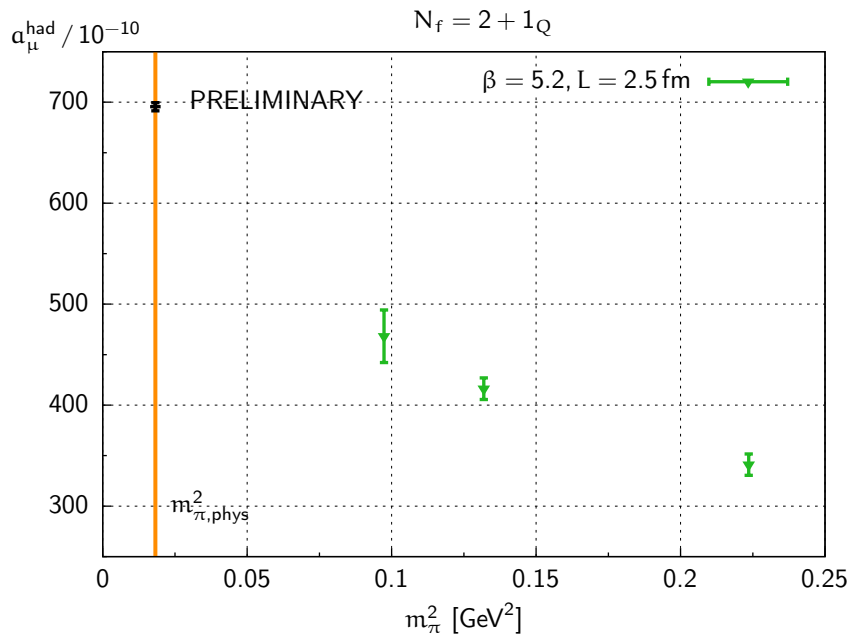
- Chiral behavior unknown: χ PT inspired fit : $A + Bm_\pi^2 + Cm_\pi^2 \ln(m_\pi^2)$

Hadronic Contribution to α_μ for $N_f = 2$

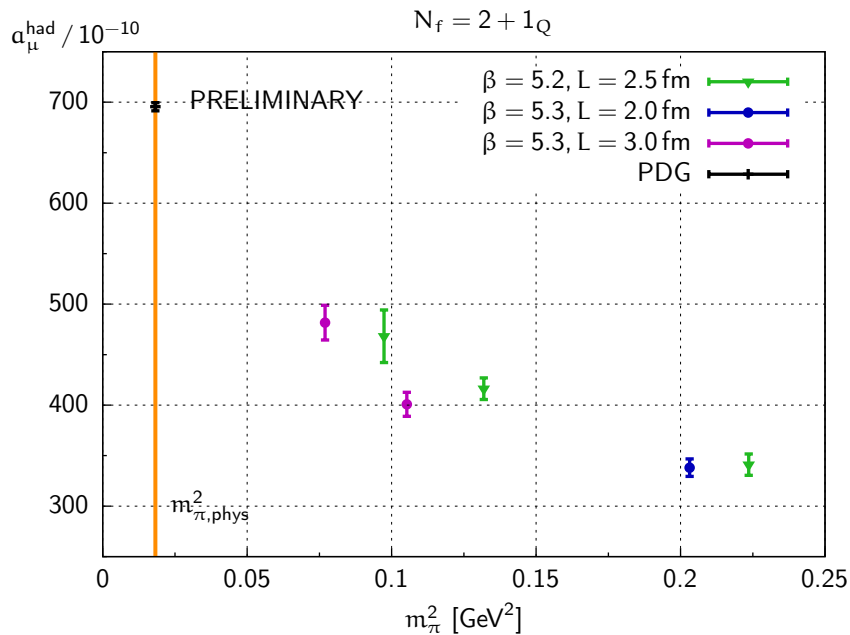


- Chiral behaviour unknown: Try linear extrapolation on most chiral points

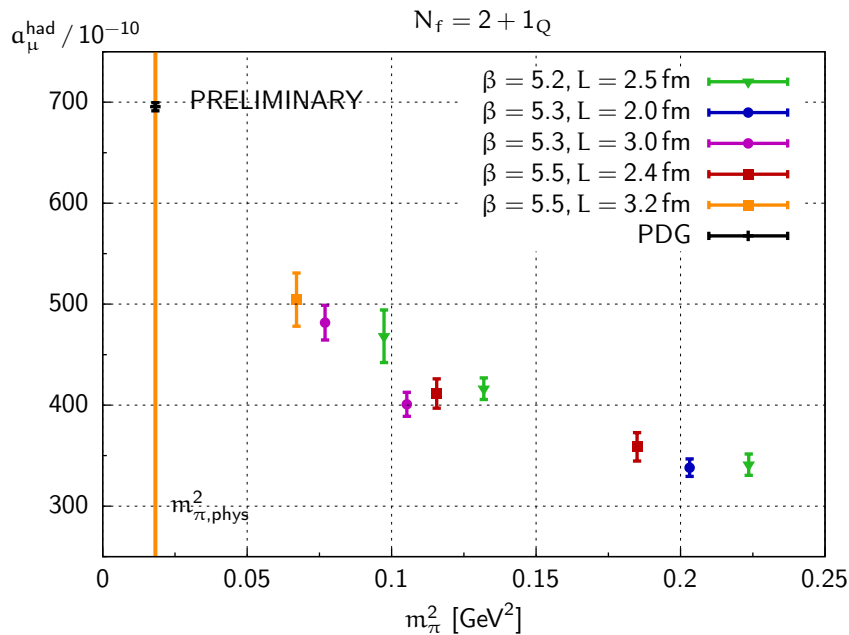
Hadronic Contribution to a_μ for $N_f = 2 + 1_Q$



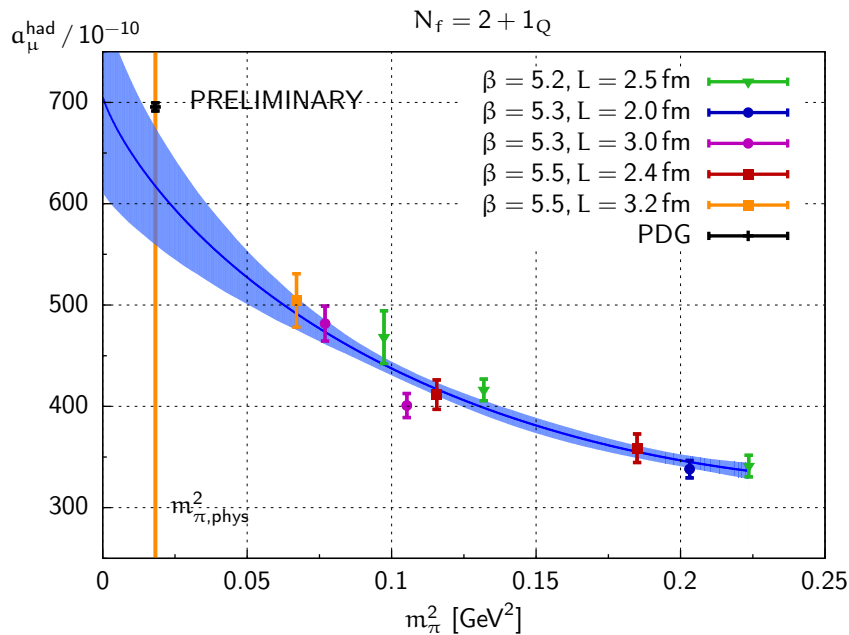
Hadronic Contribution to a_μ for $N_f = 2 + 1_Q$



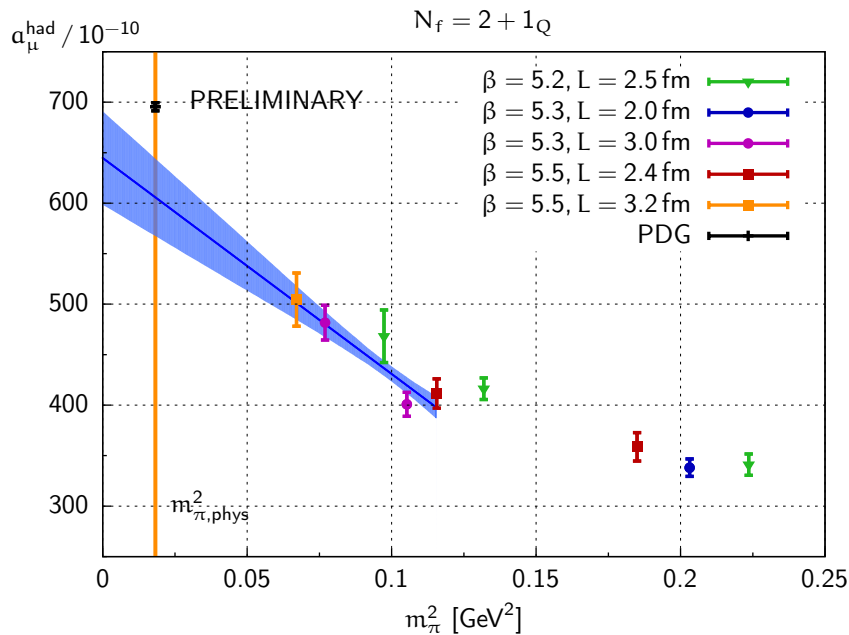
Hadronic Contribution to a_μ for $N_f = 2 + 1_Q$



Hadronic Contribution to a_μ for $N_f = 2 + 1_Q$



Hadronic Contribution to a_μ for $N_f = 2 + 1_Q$



Conclusion

- Lattice QCD can calculate α_μ^{had} from first principles
- Twisted boundary conditions improve momentum dependence of $\Pi(q^2)$ and help to control the systematic uncertainties of α_μ
- New currents (local and conserved) reduce numerical cost by factor 5
- Chiral extrapolation improved by additional ensembles ($m_\pi^2 < 200 \text{ MeV}$)

Outlook

- Further improvements necessary to compete with phenomenological approach
 - Improve statistics (e.g. by multiple sources)
 - Study finite size and volume effects
 - Dynamical strange quark (and charm quark)
 - Disconnected diagrams (e.g. by hopping parameter expansion) → talk by V. Gülpers
 - Simulations at the physical pion mass
 - Isospin breaking

Thank you for your attention!