# Two-Photon Decay of Neutral Pion in Lattice QCD

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Lattice 2012, Cairns, Australia

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### Motivation

- PrimEx@JLab:  $\Gamma_{\pi^0\gamma\gamma} = 7.82(22)$  eV [PrimEx, PRL106, 2011]
- Precision:  $2.8\% \rightarrow 1.4\%$  (projected goal)
- Benchmark test of axial U(1) anomaly in QCD

• 
$$\Gamma_{\pi^0\gamma\gamma} = (\pi/4) \alpha_e^2 m_\pi^3 \mathcal{F}_{\pi^0\gamma\gamma}^2(m_\pi^2, 0, 0)$$
  
•  $\mathcal{F}_{\pi^0\gamma\gamma}(m_\pi^2, p_1^2, p_2^2)$ :  $\pi^0 \to \gamma^* \gamma^*$  transition form factor

•  $m_{\pi}^2 = p_{1,2}^2 = 0$ , ABJ annoaly valid to all orders [Adler, Bardeen, 1969]  $\mathcal{F}_{\pi^0\gamma\gamma}(0,0,0) = \mathcal{F}_{\pi^0\gamma\gamma}^{ABJ} = \frac{1}{4\pi^2 F_{\pi}}$ 

•  $m_\pi^2 
eq 0$ ,  $p_{1,2}^2 
eq 0$ , corrections need to be calculated nonperturbatively

 overlap fermion: exact chiral symmetry ⇒ clean test of chiral anomaly Xu Feng (KEK)
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#### Lattice setup

- $\pi^0 \to \gamma \gamma$  is nontrivial
  - $\gamma$  is not an asympototic state of QCD
  - conventional method to extract the eigenstate fails
  - $1^{--}$  interpolating operator yields vector meson rather than  $\gamma$
- new method is needed
- all-to-all propagators is useful  $\Rightarrow \langle J_{\mu}(t_1) J_{\nu}(t_2) \pi^0(t_{\pi}) \rangle$ 
  - calculate correlator at any time slice of  $t_1$ ,  $t_2$ ,  $t_\pi$
  - disconnected diagram

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### Ensemble information

- four  $m_{u,d}$ :  $m_{\pi} = 540 \rightarrow 290 \text{ MeV} \Rightarrow$  chiral extrapolation
- $m_s$  fixed to be close to its physical value  $\Rightarrow$  dynamical *s*-quark effects
- L/a = 16 and  $24 \Rightarrow$  finite-size effects
- Q = 0 and  $1 \Rightarrow$  fixing-topology effects
- $a = 0.11 \text{ fm} \Rightarrow$  study possible lattice artifacts

### Theoretical setup

• starting point: S-matrix

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# $\langle \gamma(\mathbf{p}_1, \lambda_1) \gamma(\mathbf{p}_2, \lambda_2) | \pi^0(\mathbf{q}) \rangle$

• transition form factor is defined by matrix element

 $\int d^4x e^{ip_1x} \langle \Omega | T\{J_{\mu}(x)J_{\nu}(0)\} | \pi^0(q) \rangle = \epsilon_{\mu\nu\alpha\beta} p_1^{\alpha} p_2^{\beta} F_{\pi^0\gamma\gamma}(m_{\pi}^2, p_1^2, p_2^2)$ 

 ϵ<sub>μναβ</sub> p<sub>1</sub><sup>α</sup> p<sub>2</sub><sup>β</sup>: induced by the negative parity of the π<sup>0</sup>

 analytical continuation from Minkowski to Euclidean space-time
 [Ji, Jung, 2001; Dudek, Edwards, 2006]

$$\int dt e^{\omega t} \int d^3ec x e^{-iec p_1\cdotec x} \langle \Omega | \, \mathcal{T}\{J_\mu(x)J_
u(0)\} | \pi^0(q) 
angle$$

• Euclidean space-time, 3-point correlation function

 $\langle \Omega | J_\mu(ec{p}_1,t_1) J_
u(ec{p}_2,t_2) \pi^0(-ec{q},t_\pi) | \Omega 
angle$ 

# Amplitude $A_{\pi}(\tau)$

• extract  $|\pi^0(q)\rangle$ : set  $\tau = t_1 - t_2$  and  $t = \min\{t_1, t_2\}$ 



# Time dependence of $A_{\pi}(\tau)$

• VMD model:  $\mathcal{F}_{\pi^0\gamma\gamma}^{\text{VMD}}(m_{\pi}^2, p_1^2, p_2^2) = c_V G_V(p_1^2) G_V(p_2^2)$ •  $G_V(p^2) = M_V^2/(M_V^2 - p^2)$  is the (lightest) vector meson propagator



#### Fit ansatz

- lowest vector meson effects should be accounted for first
- corrected by including excited-state effects
- includes contributions from excited states as a polynomial of  $p_{1,2}^2$

$$\begin{aligned} \mathcal{F}_{\pi^{0}\gamma\gamma}(m_{\pi}^{2},p_{1}^{2},p_{2}^{2}) &= c_{V}G_{V}(p_{1}^{2})G_{V}(p_{2}^{2}) \\ &+ \sum_{m}c_{m}\left((p_{2}^{2})^{m}G_{V}(p_{1}^{2}) + (p_{1}^{2})^{m}G_{V}(p_{2}^{2})\right) \\ &+ \sum_{m,n}c_{m,n}(p_{1}^{2})^{m}(p_{2}^{2})^{n} \end{aligned}$$

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# Momentum domain



- $A_{\pi}(\tau) \Rightarrow \langle \Omega | T\{J_{\mu}(\vec{p}_1,t)J_{\nu}(\vec{p}_2,0)\} | \pi^0(q) \rangle$
- perform integration:  $\int dt \ e^{\omega t} \langle \Omega | T\{J_{\mu}(t)J_{\nu}(0)\} | \pi^{0} \rangle$   $\Rightarrow \mathcal{F}_{\pi^{0}\gamma\gamma}(m_{\pi}^{2}, p_{1}^{2}, p_{2}^{2})$ 
  - $\blacktriangleright \ \omega$  is input by hand
  - form factor relies on  $p_1^2 \& p_2^2$

• tunning 
$$\omega$$
, fixing  $\vec{p}_{1,2}$ :  
 $p_1^2 = \omega^2 - \vec{p}_1^2$   
 $p_2^2 = (E_{\pi} - \omega)^2 - \vec{p}_2^2$ 

 diff. spatial momentum setup, diff. contour of (p<sub>1</sub><sup>2</sup>, p<sub>2</sub><sup>2</sup>)

# Form factor

• along the contour of  $(p_1^2, p_2^2)$  (left),  $\mathcal{F}_{\pi^0\gamma\gamma}(m_{\pi}^2, p_1^2, p_2^2)$  (right) is determined



### Combined fit



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# On-shell photon limit



•  $F(m_{\pi}^2,0,0)\equiv \mathcal{F}_{\pi^0\gamma\gamma}(m_{\pi}^2,0,0)/\mathcal{F}_{\pi^0\gamma\gamma}^{\mathrm{ABJ}}$ 

- data with  $m_{\pi}L \ge 4$ : consistent with ABJ and PrimEx
- L = 16: smallest two quark mass,  $m_{\pi}L < 4$ , big FS effects
- FS effects checked at topological sector Q = 0 and 1

#### Finite-size effects

• expand the correlator into three hadronic matrix elements:

 $\langle J_{\mu}J_{\nu}\pi^{0}
angle
ightarrow \langle \Omega|J_{\mu}|V
angle\langle V|J_{
u}|\pi^{0}
angle\langle \pi^{0}|\pi^{0}|\Omega
angle$ 



• FS effects in  $g_{\nu}$ ,  $g_{V\pi\gamma}$ ,  $F_{\pi}$  accumulate and add up to a large effect

### Finite-size corrections



• FS corrections  $R_{\mathcal{O}} \equiv \mathcal{O}(\infty)/\mathcal{O}(L)$ , assume  $R_{F(m_{\pi}^2,0,0)} = R_{g_{\rho}}R_{g_{\rho\pi\gamma}}R_{F_{\pi}}$ 

# Disconnected-diagram effects



- all-to-all propagator: control error of disc. contribution
- although not significant, conn+disc systematically shift down
- precision level (3% for form factor): disc. diagram should be included

### Lattice artifacts

discrete data v.s. continuum case?



 $\bullet$  disc. effects in VMD model: less than 5  $\times$  10  $^{-4},$  neglegiable

#### Numerical results

• after examining possible systematic effects

$$egin{array}{rll} F(0,0,0)&=&1.009(22)(29)\ F(m_{\pi,\mathrm{phy}}^2,0,0)&=&1.005(20)(30)\ \Gamma_{\pi^0\gamma\gamma}&=&7.83(31)(49)~\mathrm{eV} \end{array}$$

• ABJ anomaly and PrimEx measurement

$$egin{array}{rcl} F(0,0,0)&=&1\ F(m_{\pi,\mathrm{phy}}^2,0,0)&=&1.004(14)\ \Gamma_{\pi^0\gamma\gamma}&=&7.82(22)~\mathrm{eV} \end{array}$$

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# Conclusions

- $\pi^0 \to \gamma \gamma$  calculation is done successfully
  - by analytic continuation
  - using all-to-all propagators
- ABJ anomaly confirmed in the chiral limit
  - with the overlap fermion, it is satisfied at finite lattice spacing
- correction to ABJ can also be calculated precisely
- remaining major effect would be isospin breaking
- open a possibility to calculate amplitudes with no QCD asymptotic state
  - $\blacktriangleright \pi^+ \to I^+ \nu_I \gamma$
  - ▶ g 2, light-by-light

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### Motivated by muon $g_{\mu} - 2$

- Exp. determination of  $g_{\mu}$  2 to 0.54 ppm [E821@BNL, PRD73, 2006]
- S.M. prediction of  $g_{\mu}$  2 to 0.51 ppm [Jegerlehner, EPJC71, 2011]
- Discrepancy:  $3.3\sigma \Rightarrow$  New Physics ??
- HLbL is predicted to be dominant error in the next round



- Difficult: HLbL involves  $\langle J_{\mu}J_{\nu}J_{\rho}J_{\sigma}\rangle$
- Better to start with  $\pi^0(\eta, \eta') \rightarrow \gamma^* \gamma^*$

# $\pi^{\rm 0}$ contribution

Contribution	BPP	HKS	KN	MV	BP	PdRV	N/JN
$\pi^0,\eta,\eta^\prime$	85±13	$82.7{\pm}6.4$	$83{\pm}12$	114±10	-	$114{\pm}13$	99±16
$\pi, K$ loops	$-19{\pm}13$	$-4.5{\pm}8.1$	_	_	-	$-19{\pm}19$	$-19{\pm}13$
$\pi, K$ loops + other subleading in $N_c$	-	-	_	$0{\pm}10$	-	_	_
axial vectors	$2.5{\pm}1.0$	$1.7{\pm}1.7$	_	$22\pm5$	-	$15{\pm}10$	$22\pm5$
scalars	$-6.8{\pm}2.0$	_	_	_	-	$-7\pm7$	$-7\pm2$
quark loops	$21{\pm}3$	$9.7{\pm}11.1$	-	-	-	2.3	$21{\pm}3$
total	83±32	$89.6 \pm 15.4$	80±40	$136\pm25$	110±40	$105 \pm 26$	$116\pm39$

• summary table [Jegerlehner, Nyffeler, Phys.Rept.477:1-110,2009]

- $\pi^0(\eta, \eta') \rightarrow \gamma^* \gamma^*$  are consistent to total contributions
- Among three PS mesons,  $\pi^0$  takes about  $\tilde{7}0\%$
- ▶ calulation on the  $\pi^0 \to \gamma^* \gamma^*$  can be duplicated to the  $\eta$ ,  $\eta'$  sector

# Non-perturbative nature

• Invariant mass spectrum for two-photon



- Three spikes presents three bound states:  $\pi^0$ ,  $\eta$ ,  $\eta'$
- Bound states  $\rightarrow$  confinement  $\rightarrow$  LQCD

### Rho mass



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