

Chiral Polarization Scale at Finite Temperature

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Lattice 2012 - Cairns



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Outline

- Motivation
- Local chirality
- Dynamical polarization
- Chiral polarization scale
- Numerical results
- Conclusions

Local chirality

- Low-lying eigenmodes of the Dirac operator play an important role in chiral symmetry breaking

$$iD\psi = \lambda\psi \quad \psi_R = \frac{1}{2}(1 + \gamma_5)\psi \quad \psi_L = \frac{1}{2}(1 - \gamma_5)\psi$$

- The free-field modes are unpolarized

$$\langle\psi|\gamma_5|\psi\rangle = 0 \quad \Rightarrow \quad \psi(x)^\dagger\gamma_5\psi(x) = \psi_R^\dagger\psi_R - \psi_L^\dagger\psi_L = 0$$

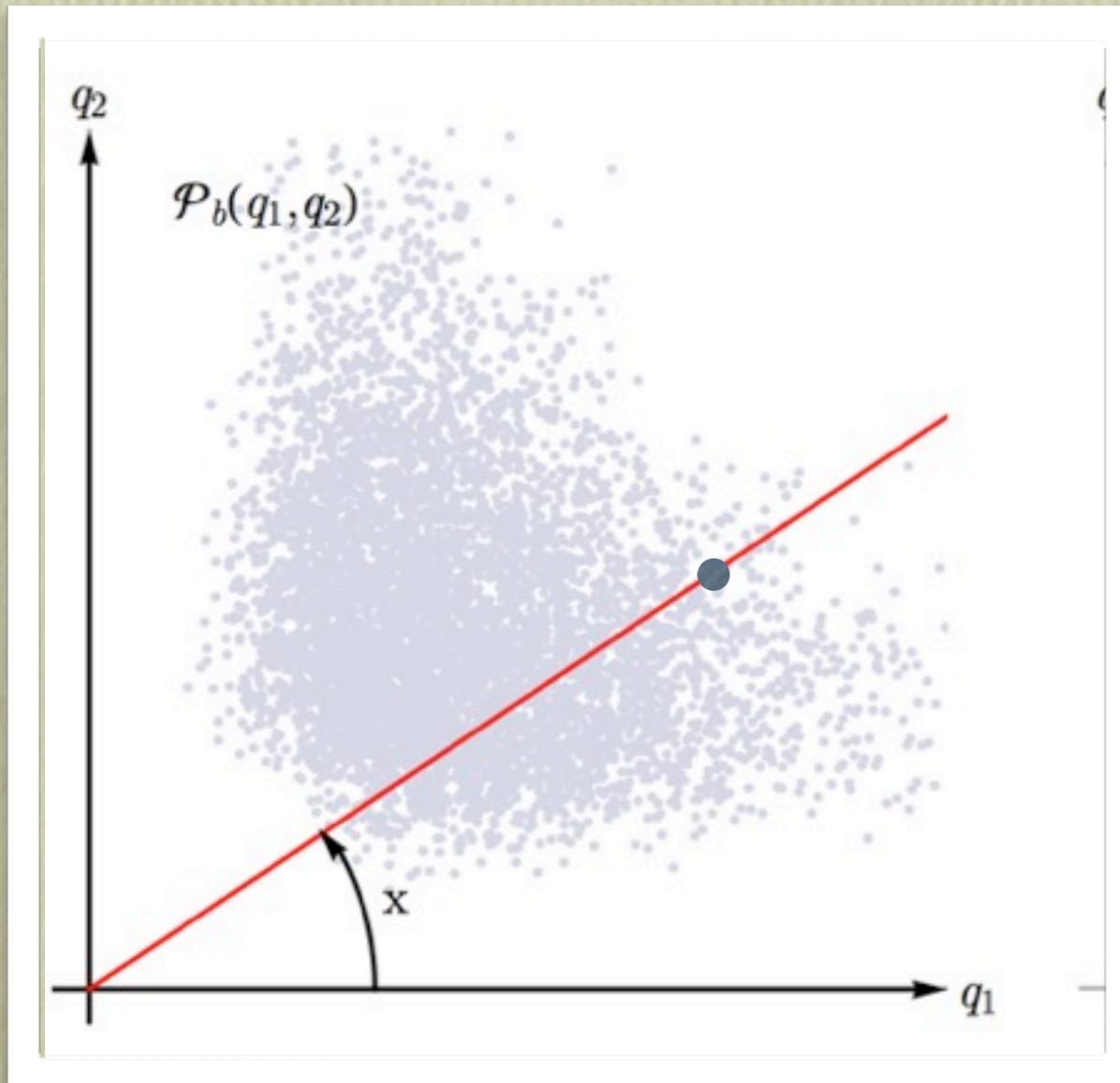
- Eigenmodes for self-dual fields are expected to be polarized

$$\left[-D^2 + \frac{1}{2}\sigma_{\mu\nu}F_{\mu\nu}^S\right]\psi_L = \lambda^2\psi_L \quad F^S = \frac{1}{2}(F + \tilde{F})$$

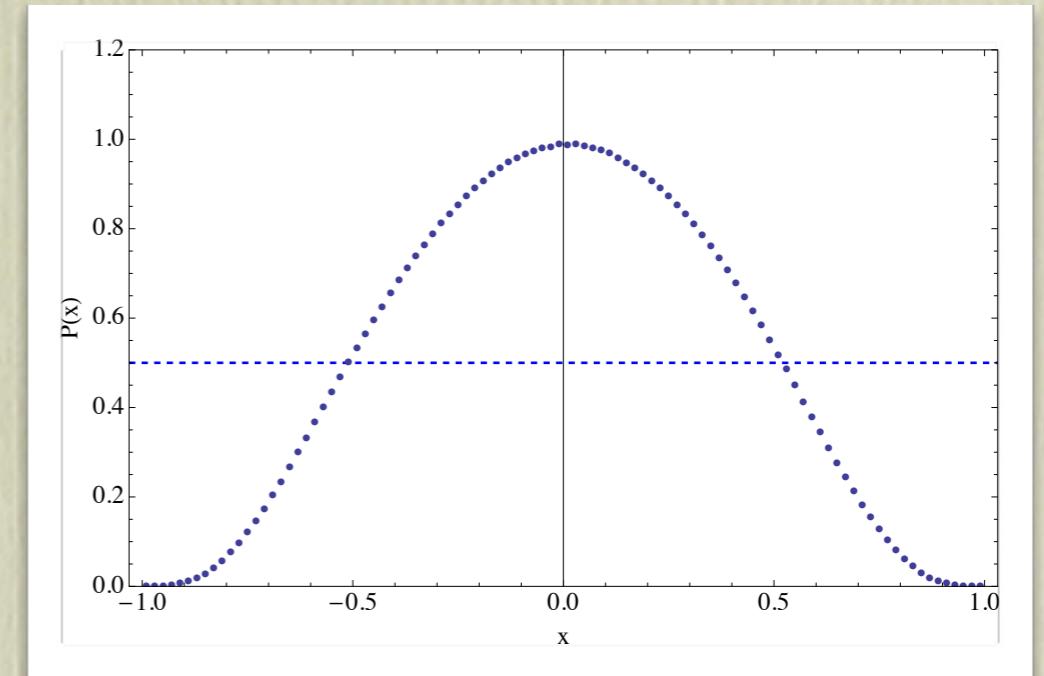
$$\left[-D^2 + \frac{1}{2}\sigma_{\mu\nu}F_{\mu\nu}^A\right]\psi_R = \lambda^2\psi_R \quad F^A = \frac{1}{2}(F - \tilde{F})$$

Dynamical polarization

$$Q_1 = \|\psi_R\| \quad Q_2 = \|\psi_L\|$$

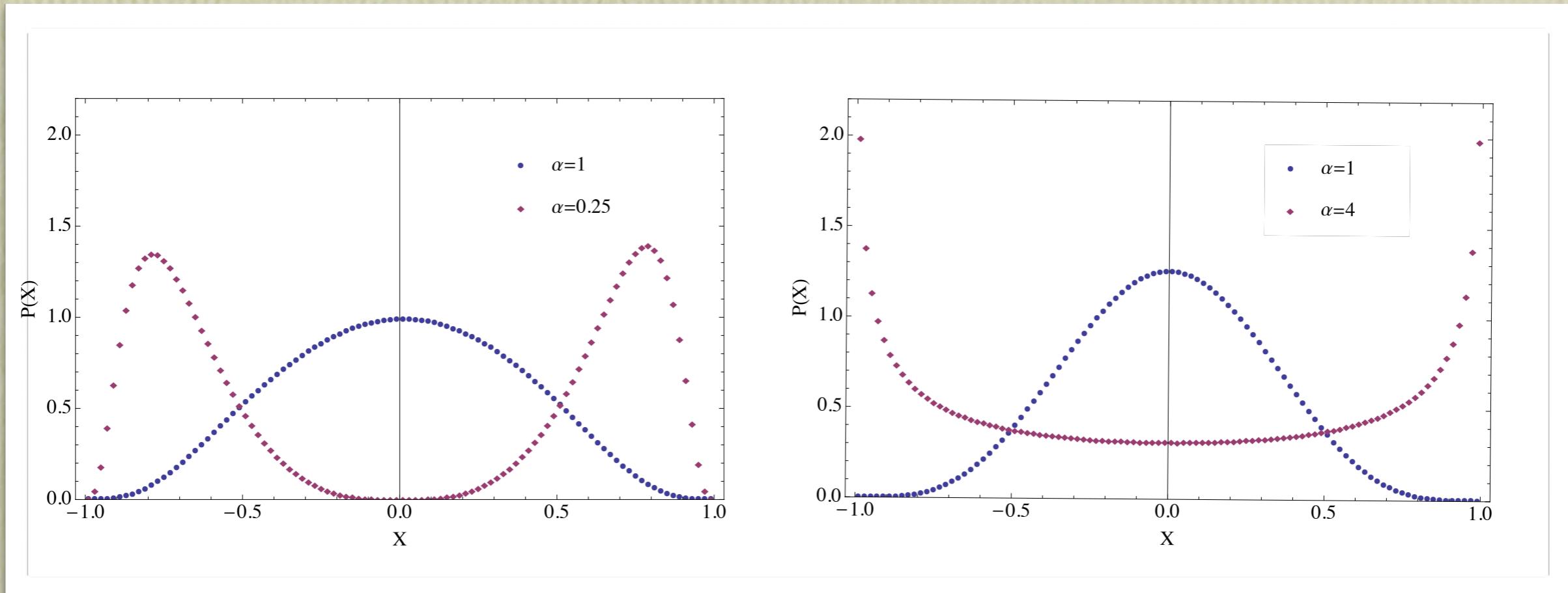


$$x = \frac{4}{\pi} \tan^{-1} \frac{\|Q_2\|}{\|Q_1\|} - 1$$



I. Horvath, N. Isgur, J. McCune, and H. B. Thacker, *Phys. Rev. D* **65** (2002) 014502

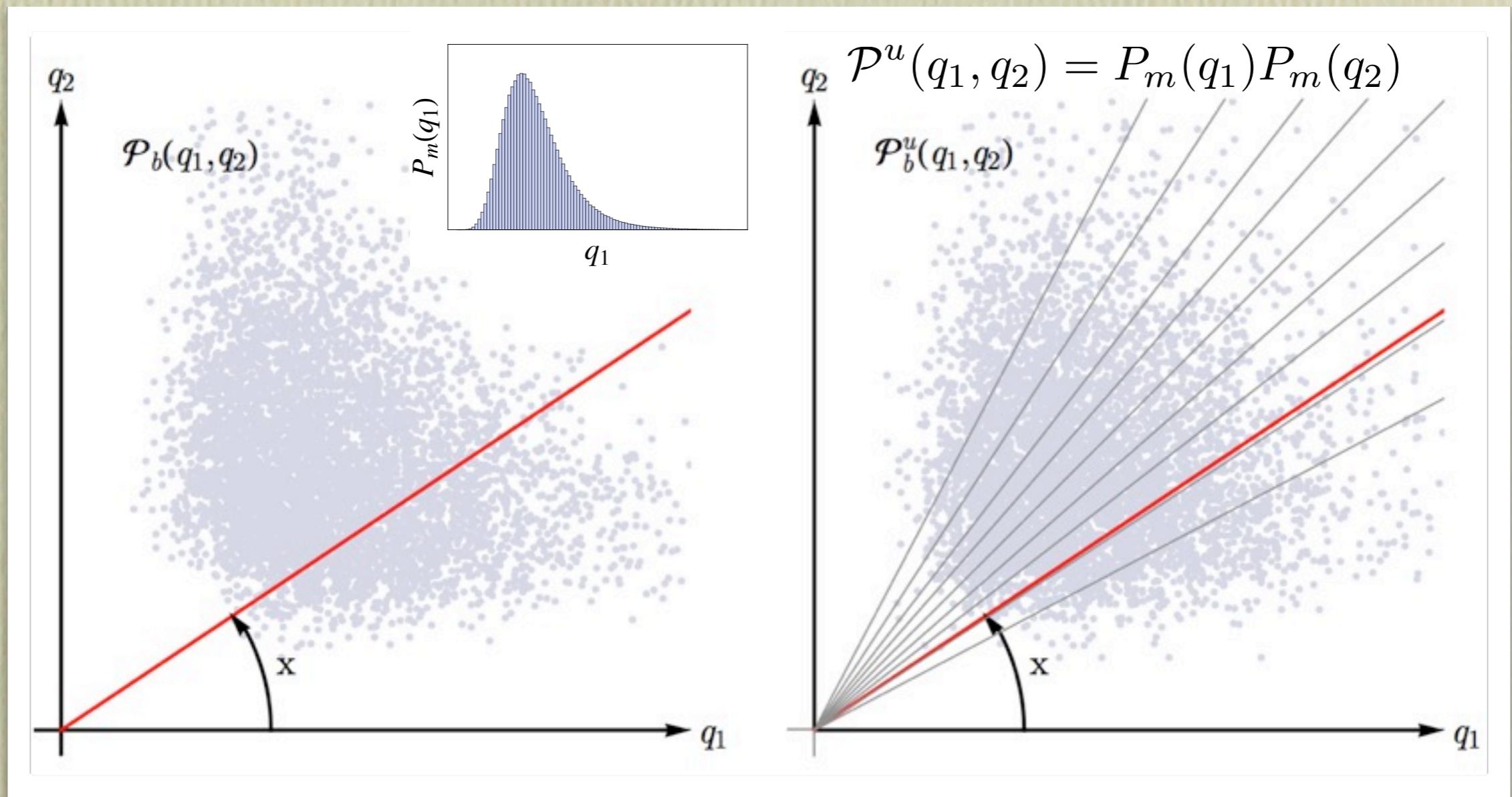
Dynamical polarization



$$X = \operatorname{sgn}(x)|x|^\alpha$$

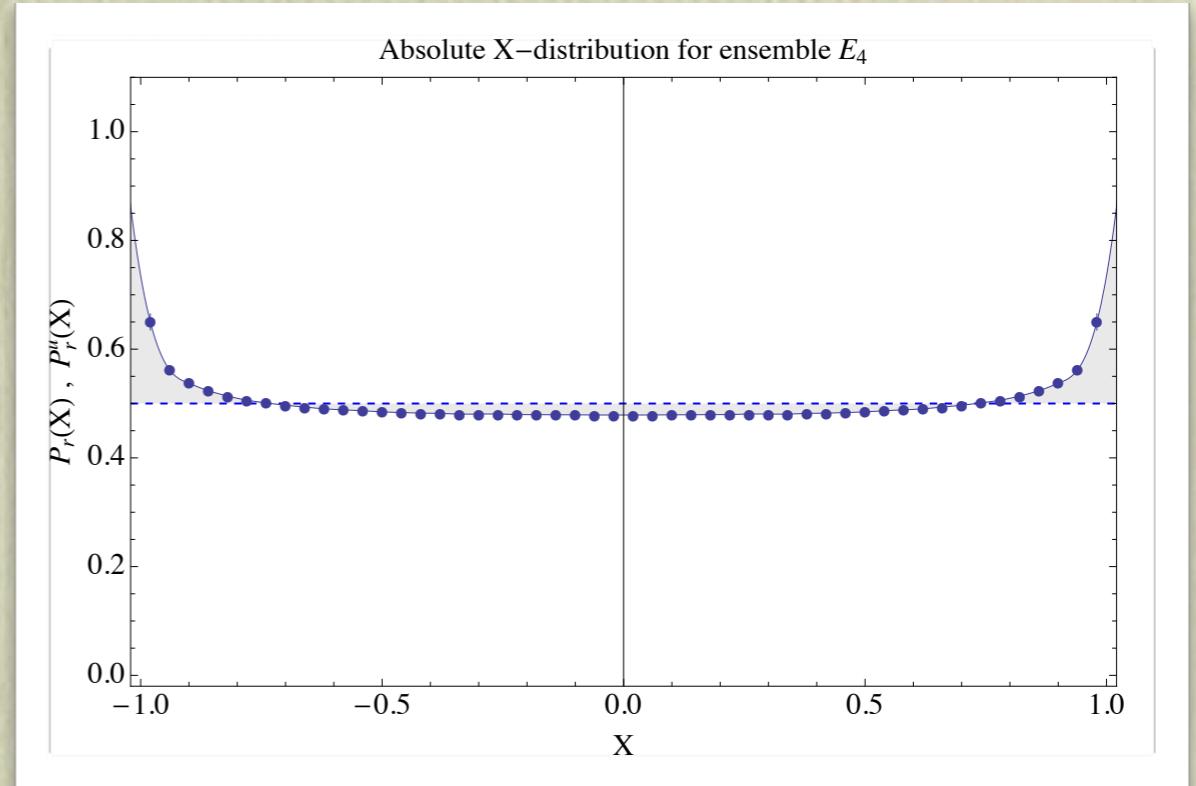
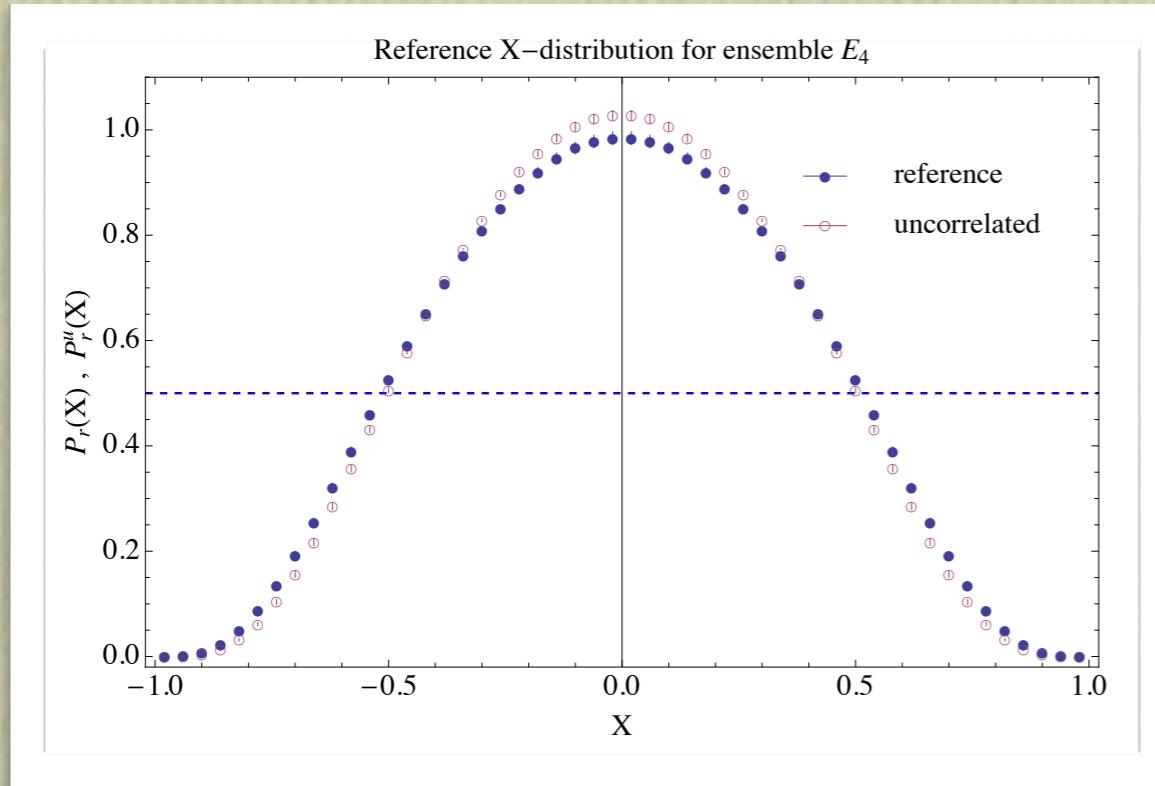
$$X = \frac{\left(\frac{\|Q_2\|}{\|Q_1\|}\right)^\alpha - 1}{\left(\frac{\|Q_2\|}{\|Q_1\|}\right)^\alpha + 1}$$

Dynamical polarization

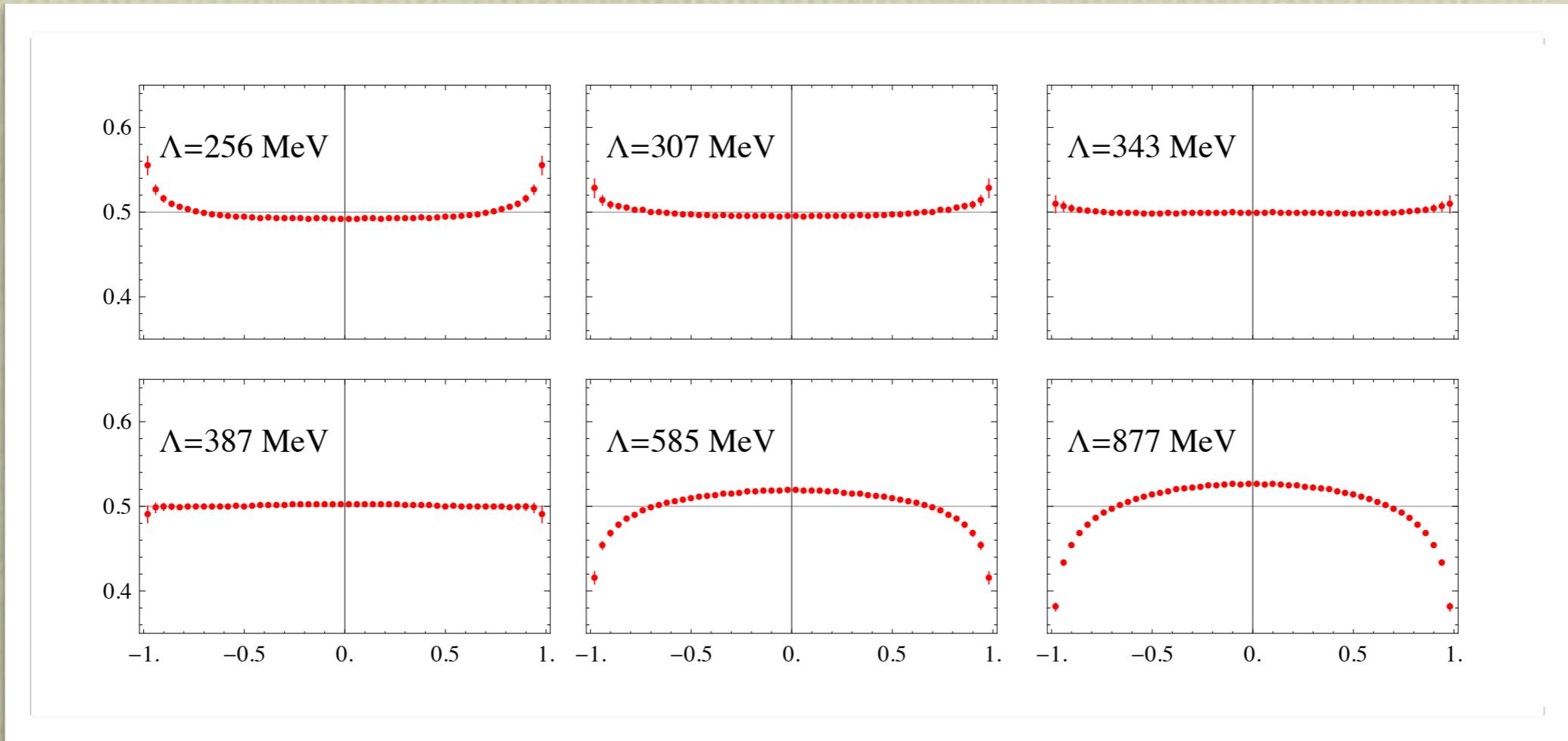


AA, T. Draper, I. Horvath, and T. Streuer, *Annals of Physics* **326** (2011) 1941

Dynamical polarization

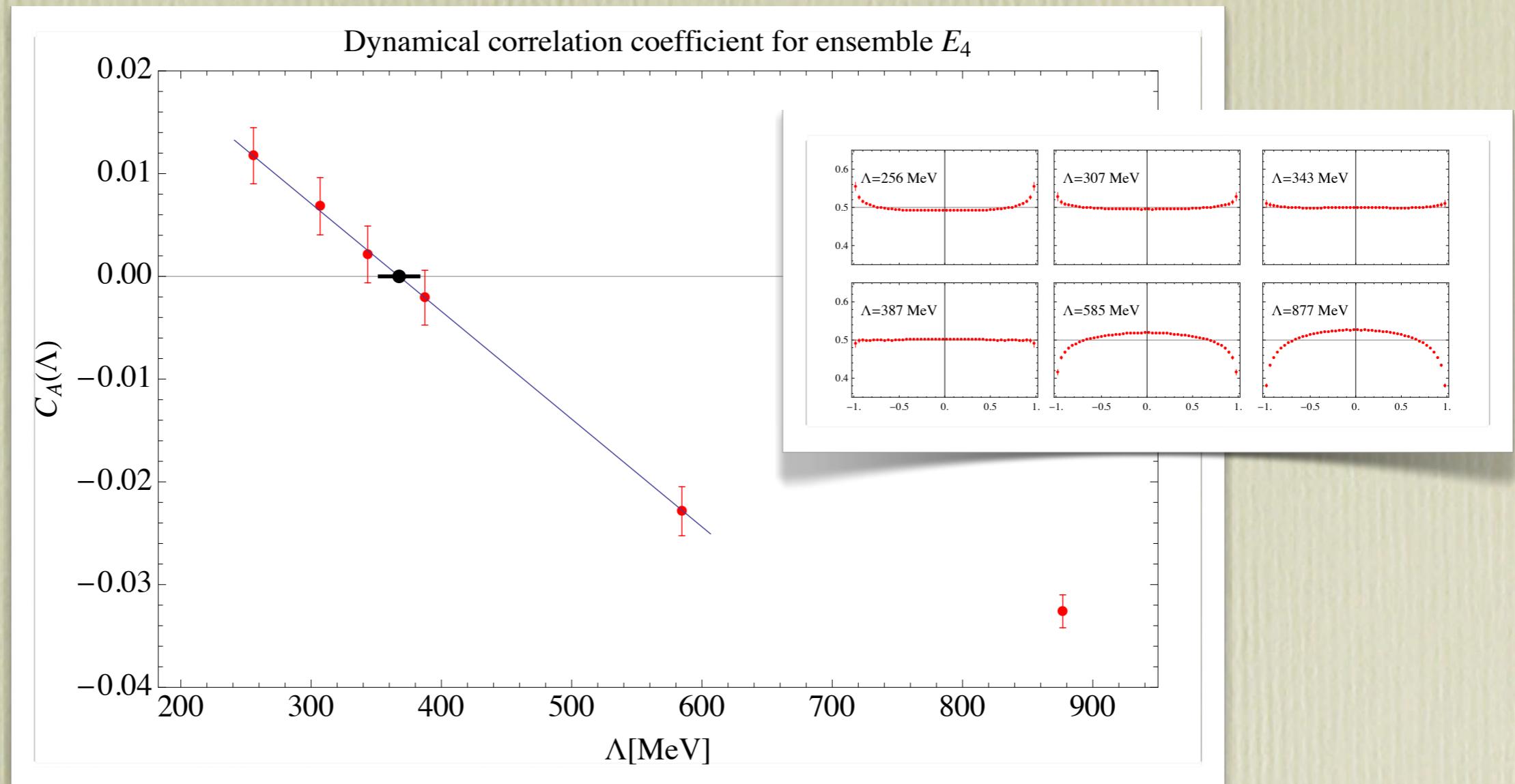


Chiral polarization scale



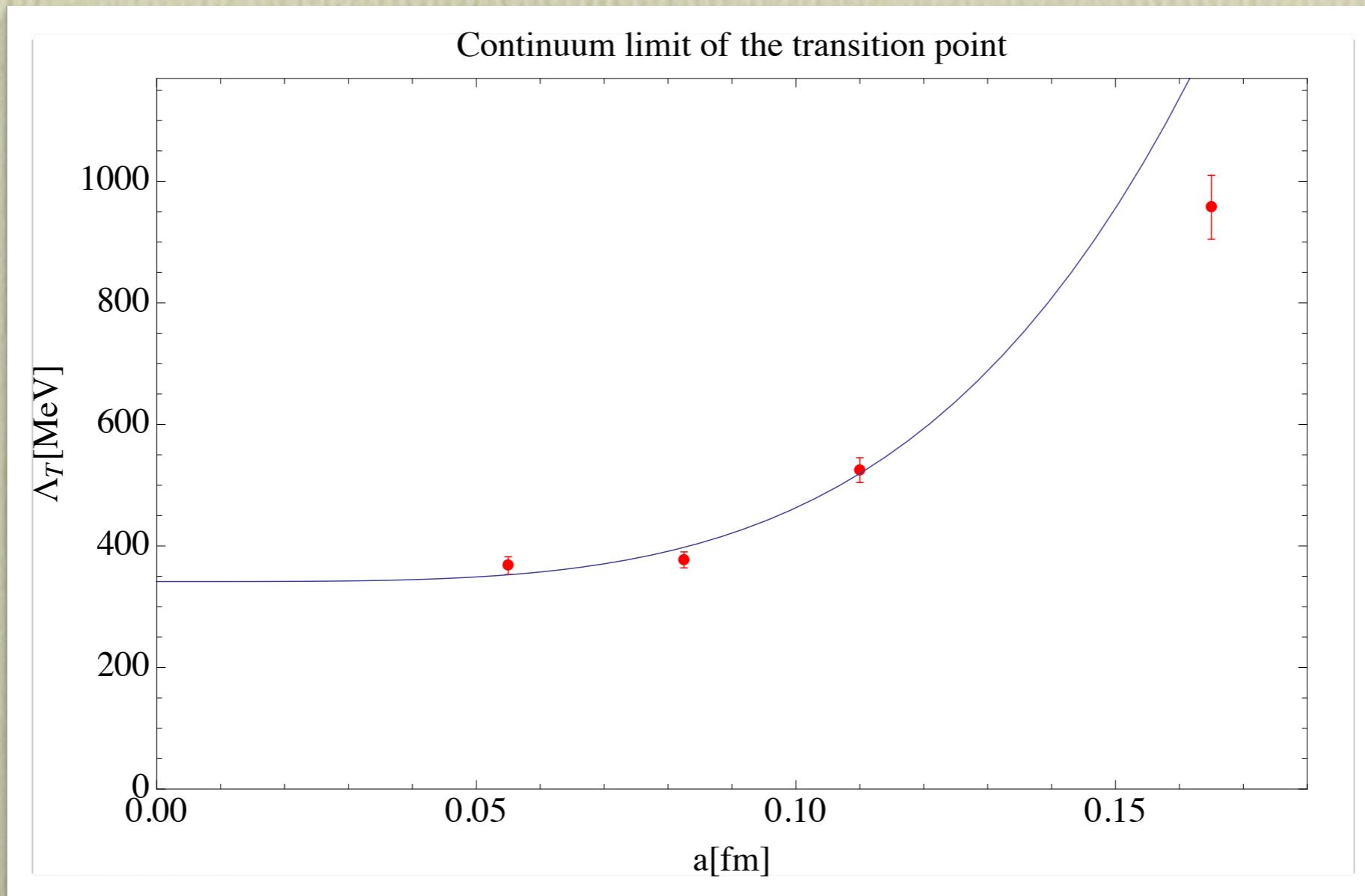
$$D_{\text{ov}}\psi = \lambda\psi \quad Q_1 = \frac{1}{2}(1 + \gamma_5)\psi \quad Q_2 = \frac{1}{2}(1 - \gamma_5)\psi$$

Chiral polarization scale



$$\Gamma \equiv \int_{-1}^1 dX |X| P_A(X) \quad C_A = 2\Gamma - 1$$

Chiral polarization scale



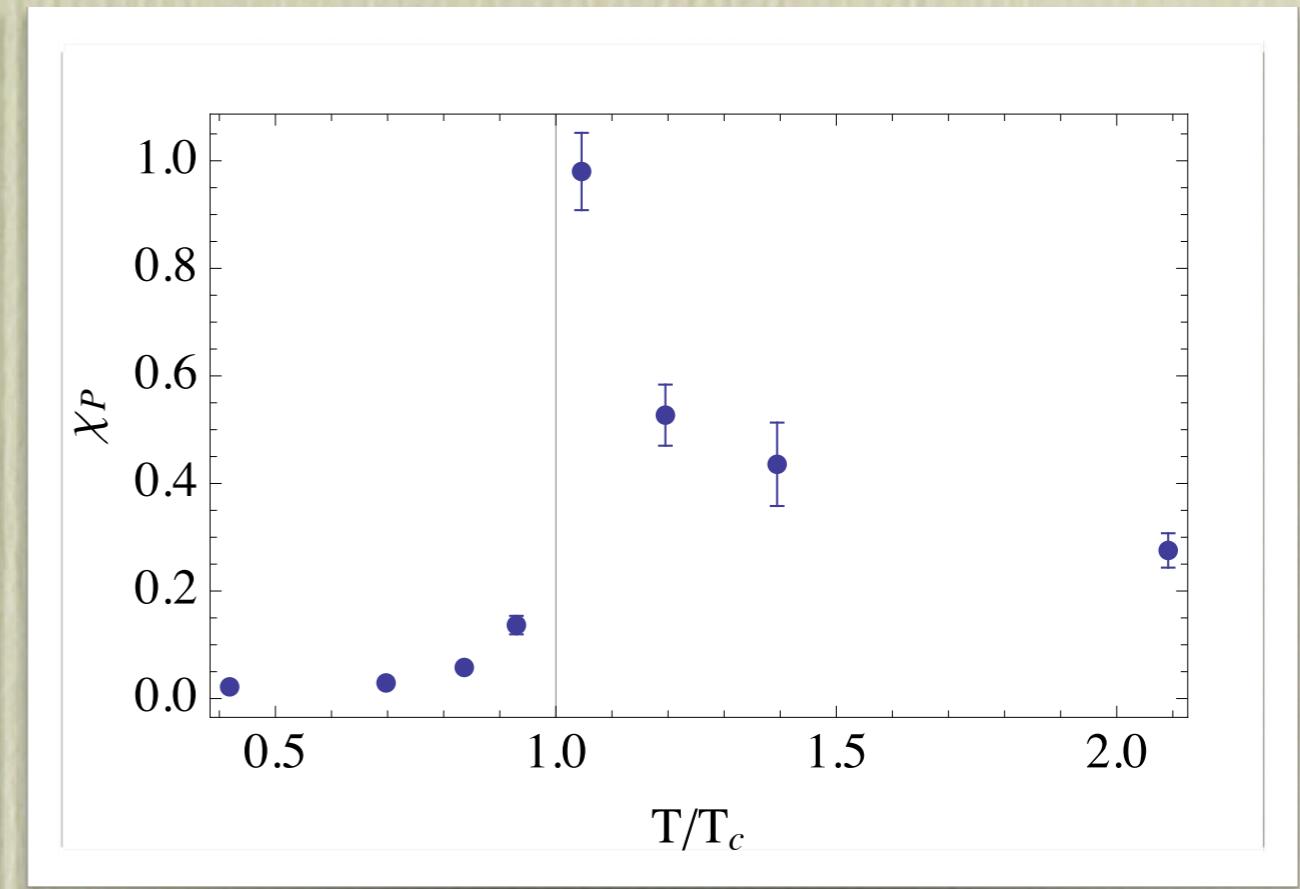
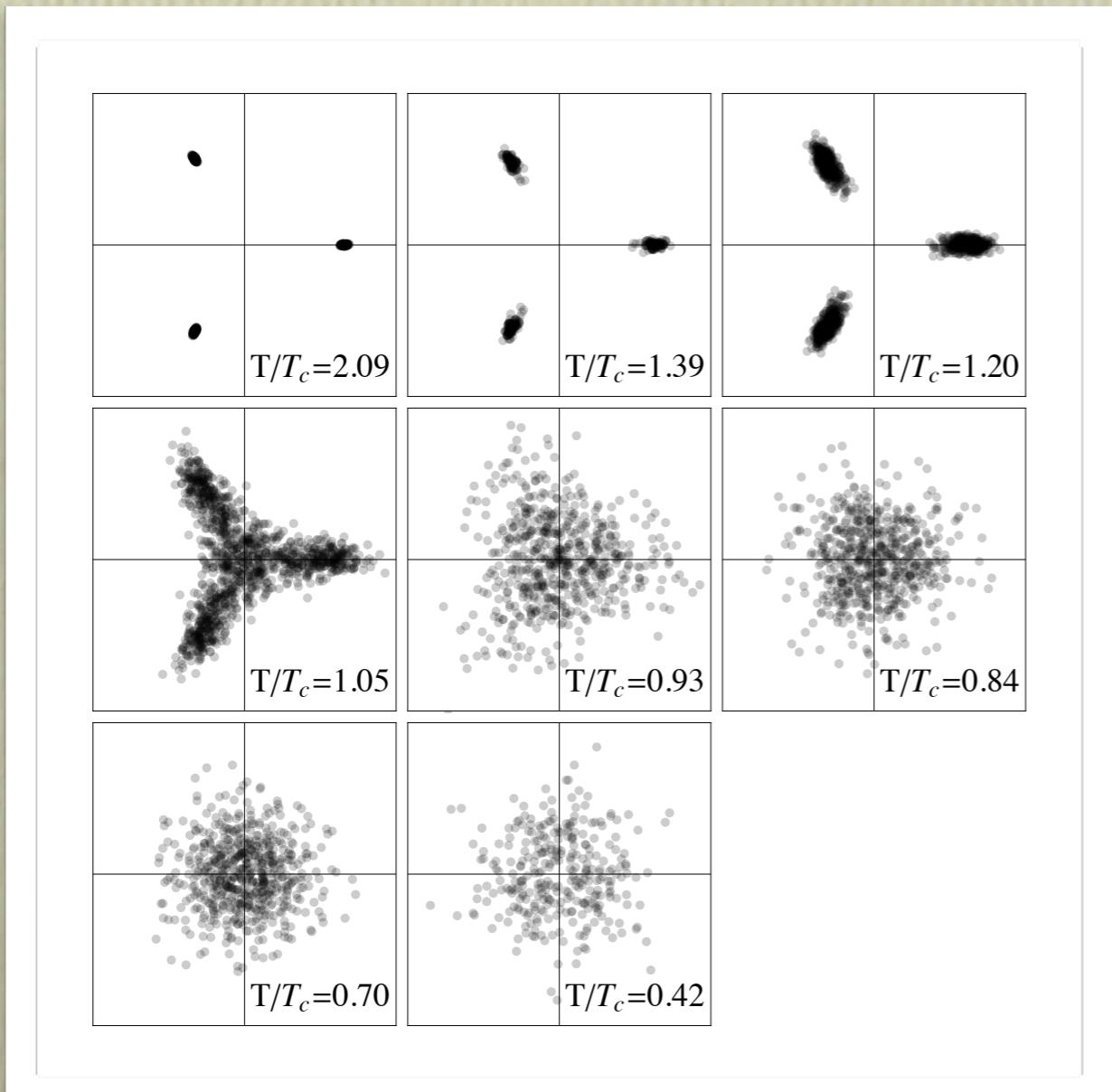
Numerical results

Size	Nconf	T/T _c	a [fm]
4x20 ³	100	2.09	
6x20 ³	100	1.39	
7x20 ³	400	1.20	
8x20 ³	400	1.05	
9x20 ³	200	0.93	
10x20 ³	200	0.84	
12x20 ³	200	0.70	
20x20 ³	100	0.42	0.085

Pure gauge configurations using Wilson action

ALPHA Collaboration, *Nucl. Phys.* **B535** (1998) 389–402.

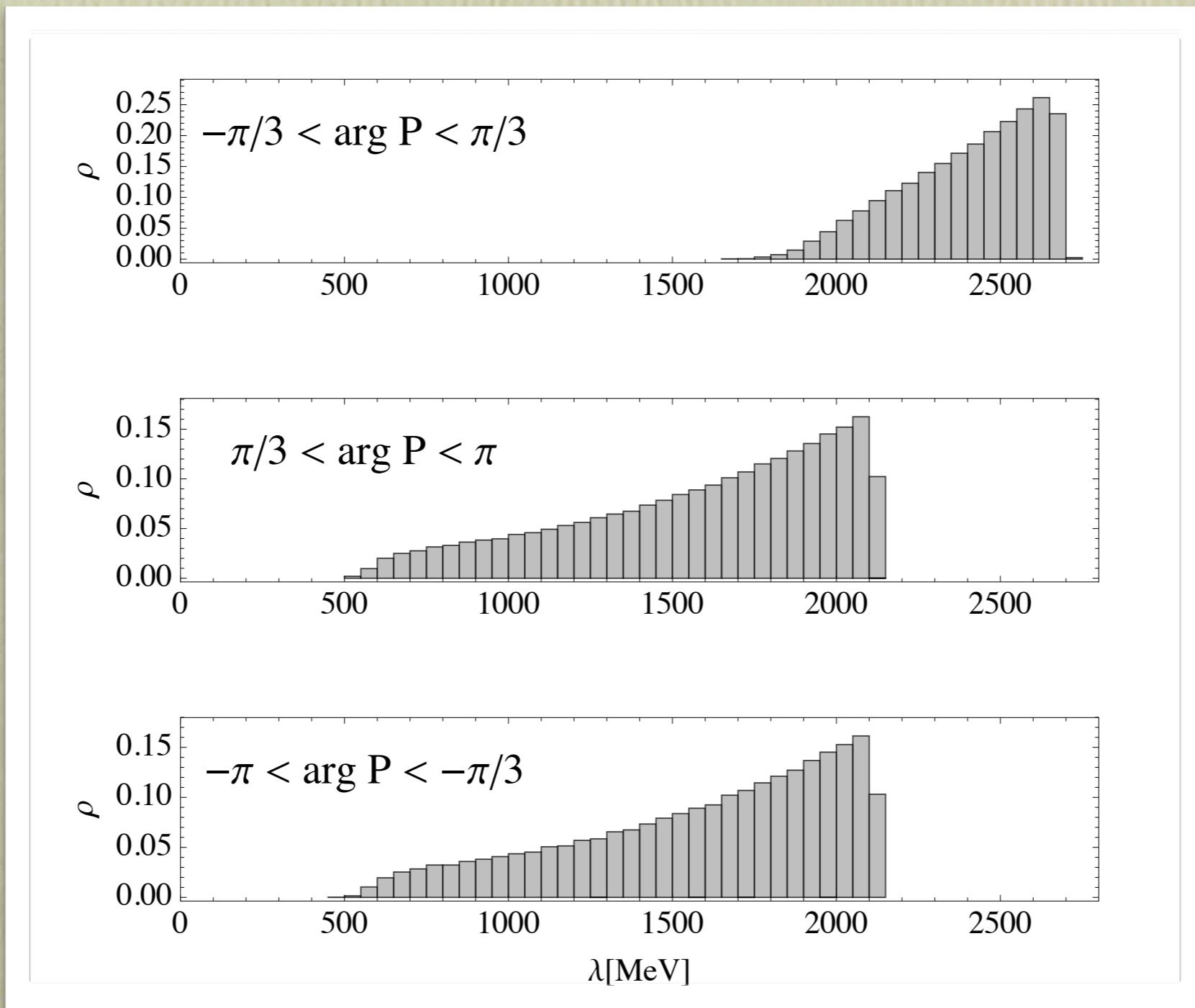
Polyakov loop



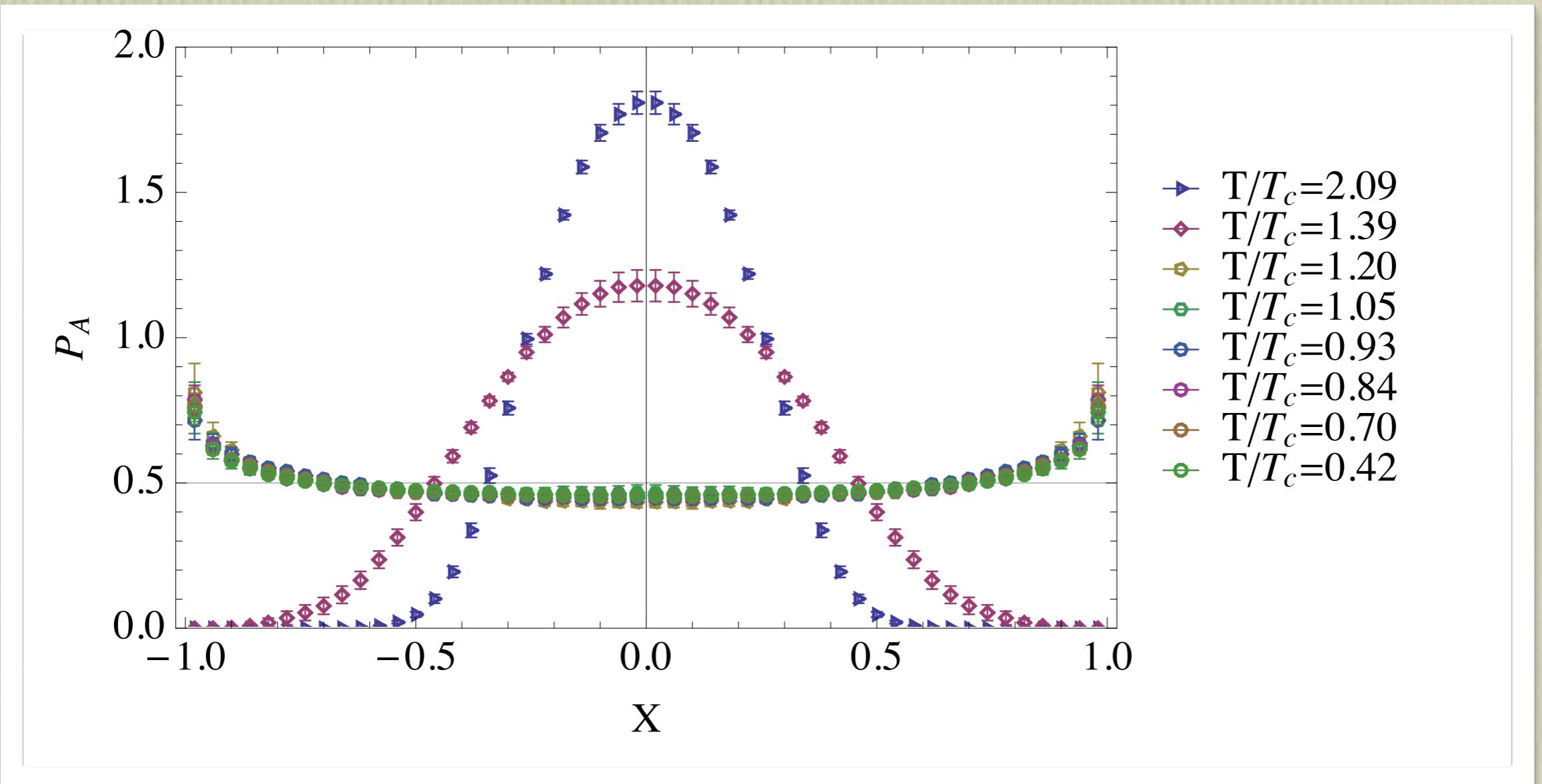
$$P = \frac{1}{3V_s} \sum_{\vec{x}} \text{Tr} \left(\prod_{t=0}^{N_t-1} U_4(\vec{x}, t) \right)$$

$$\chi_P = V_s \left(\langle |P|^2 \rangle - \langle |P| \rangle^2 \right)$$

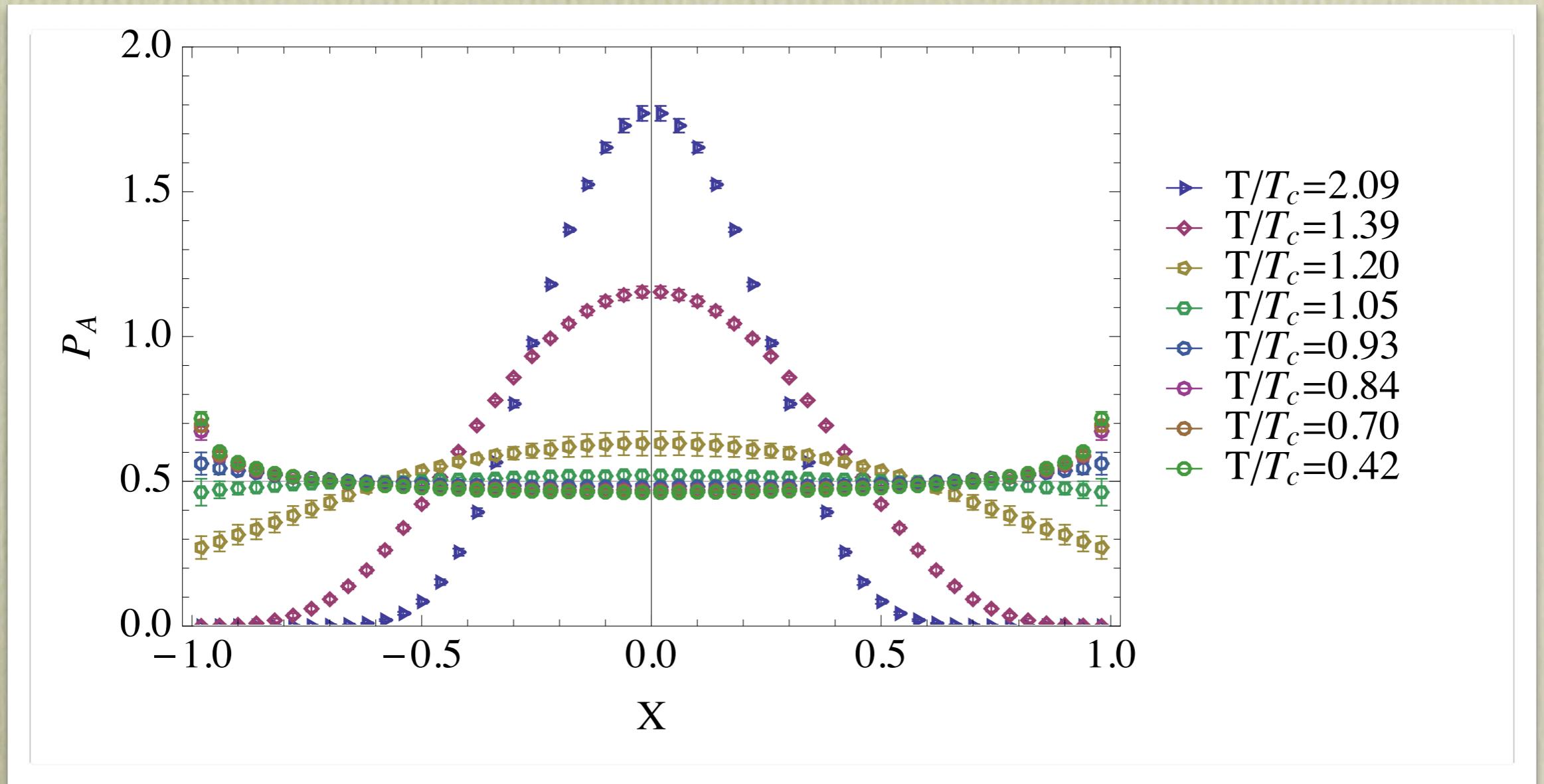
Real sector



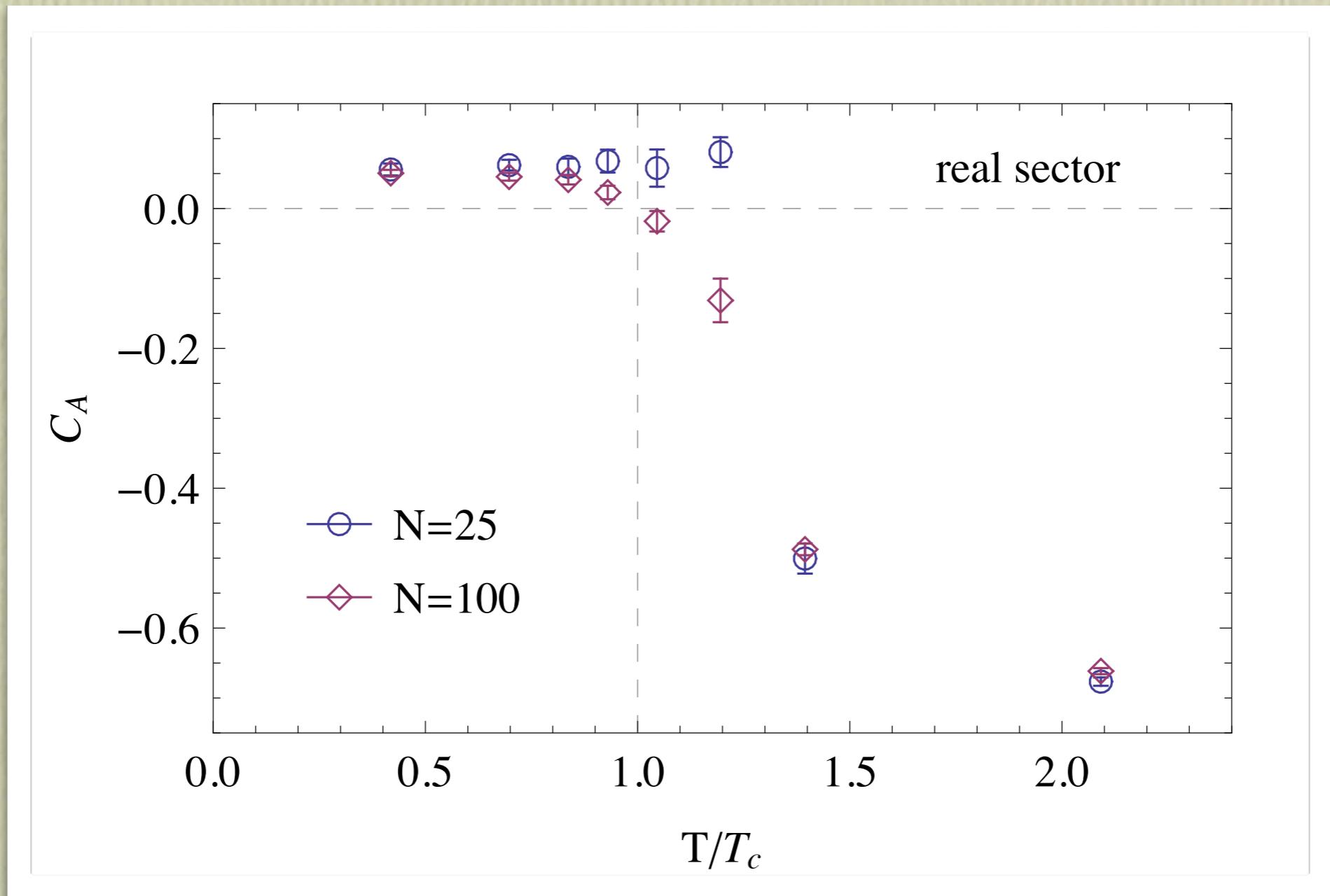
Low modes chiral polarization



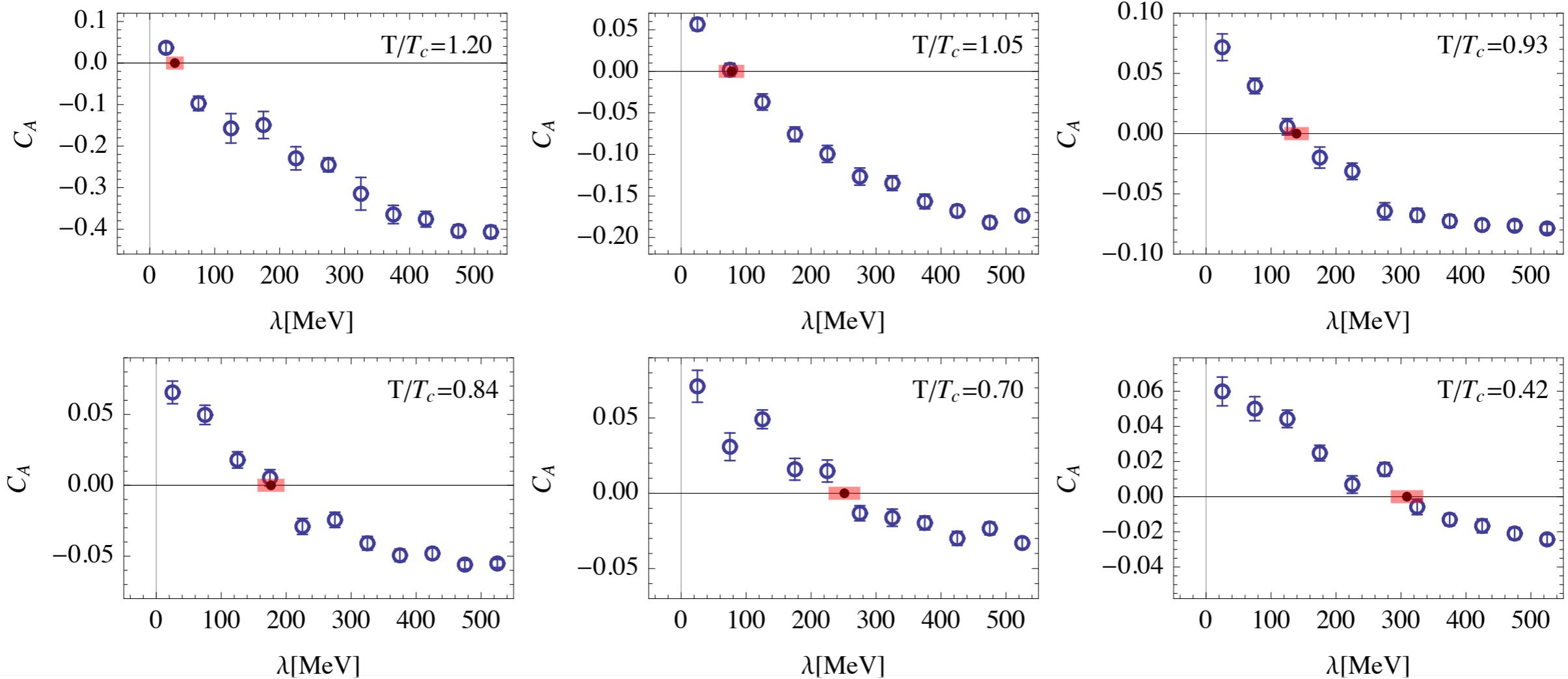
Low modes chiral polarization



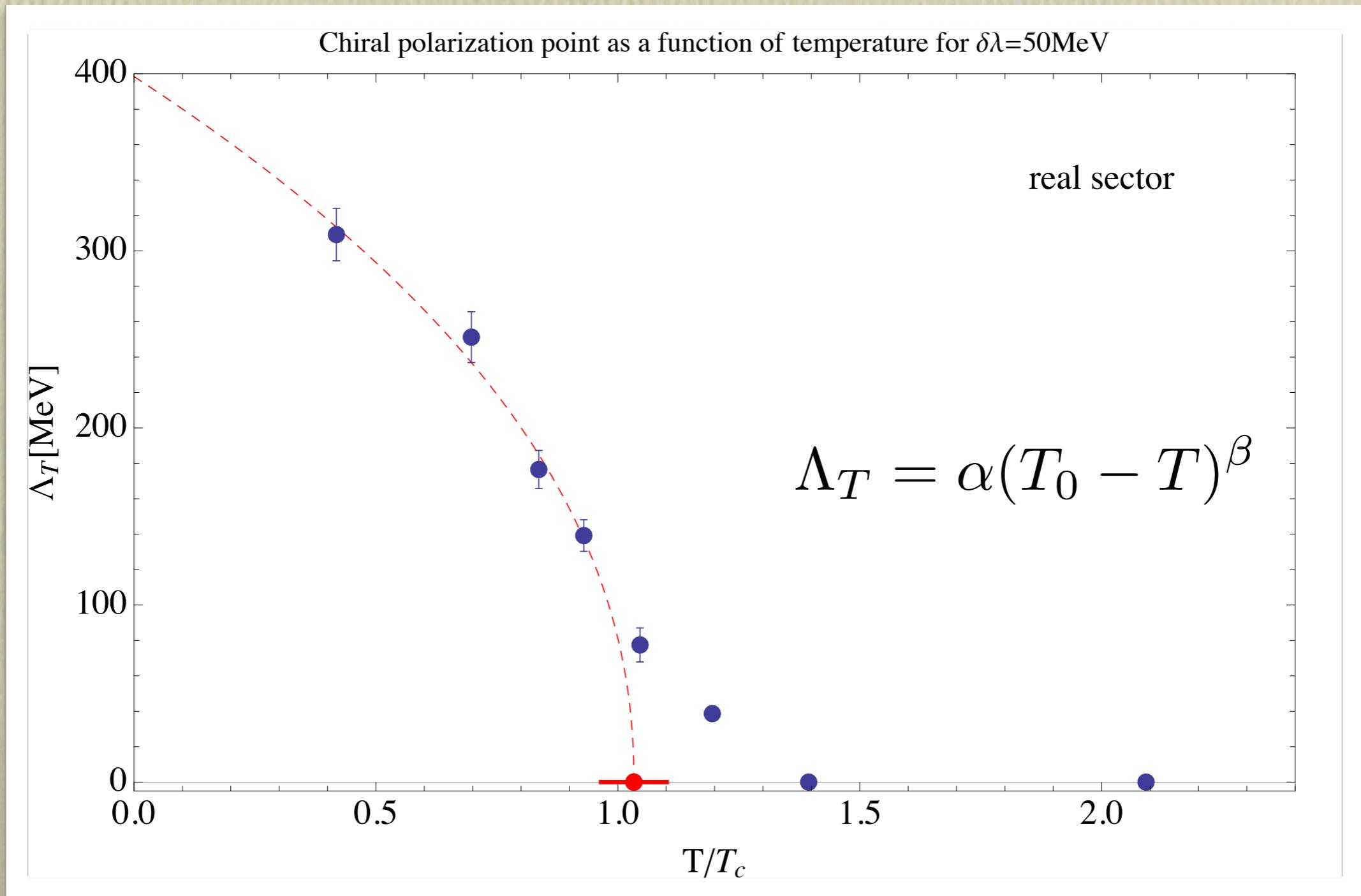
Low modes chiral polarization



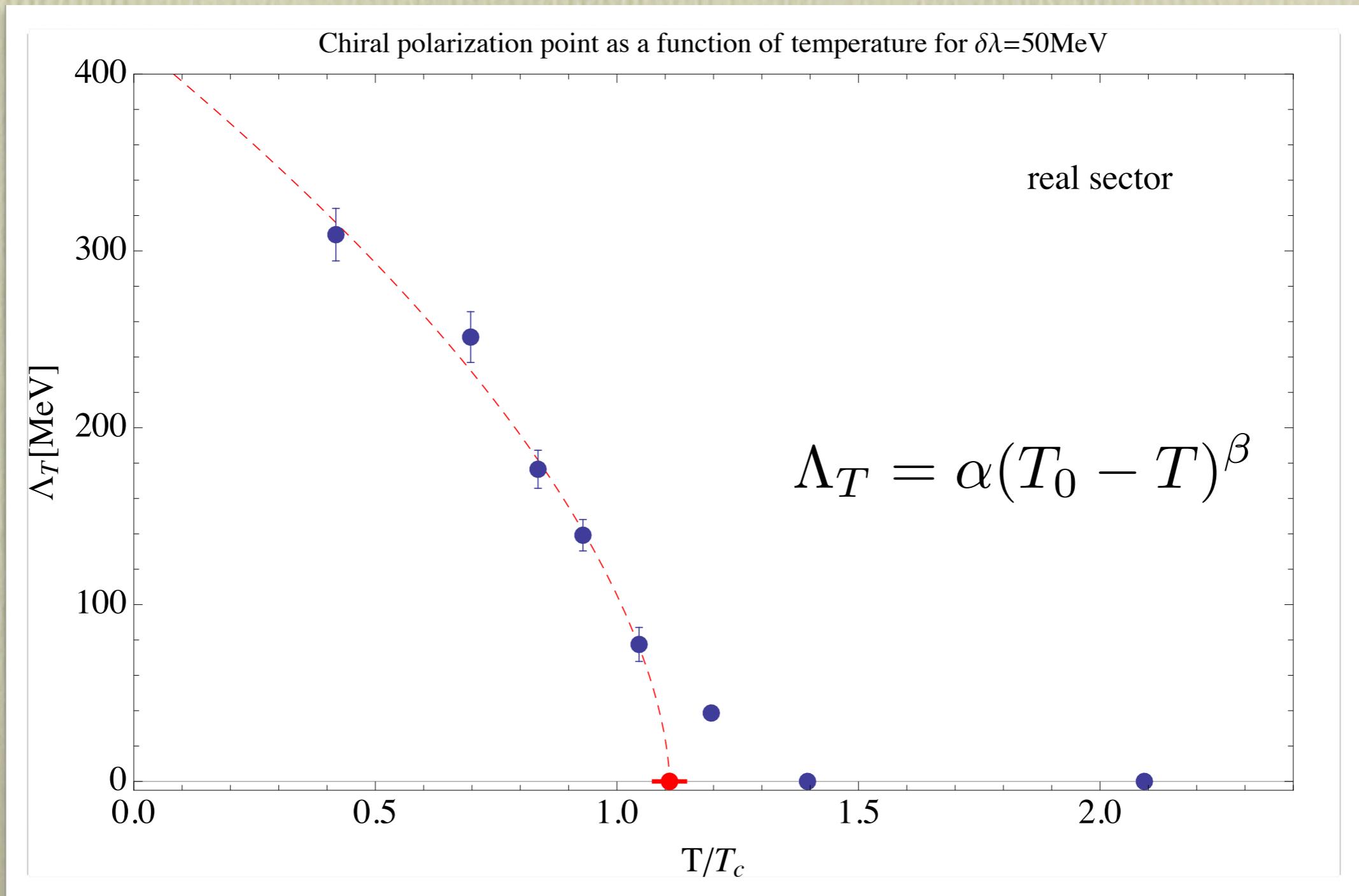
Eigenvalue dependence



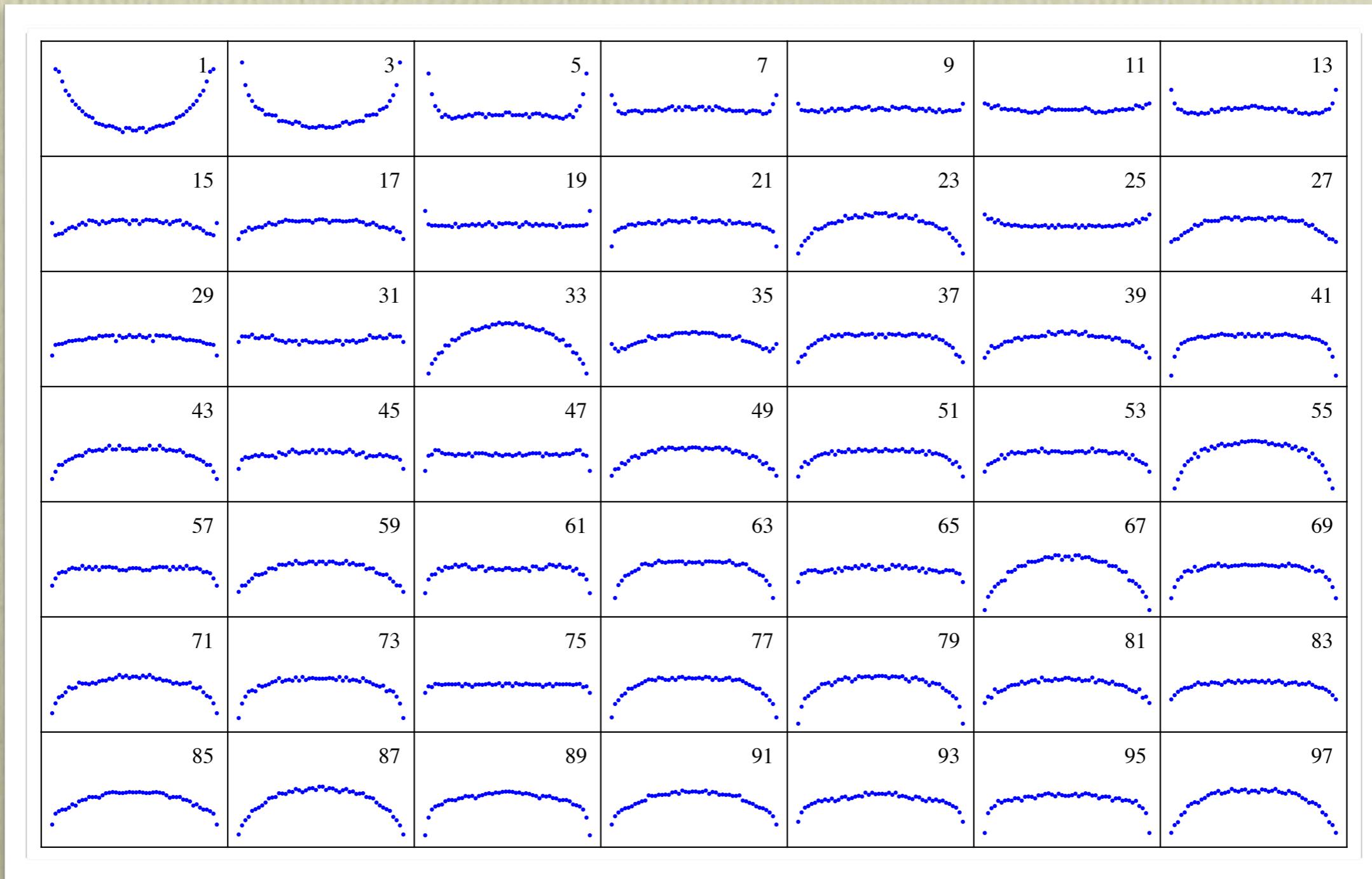
Chiral polarization scale



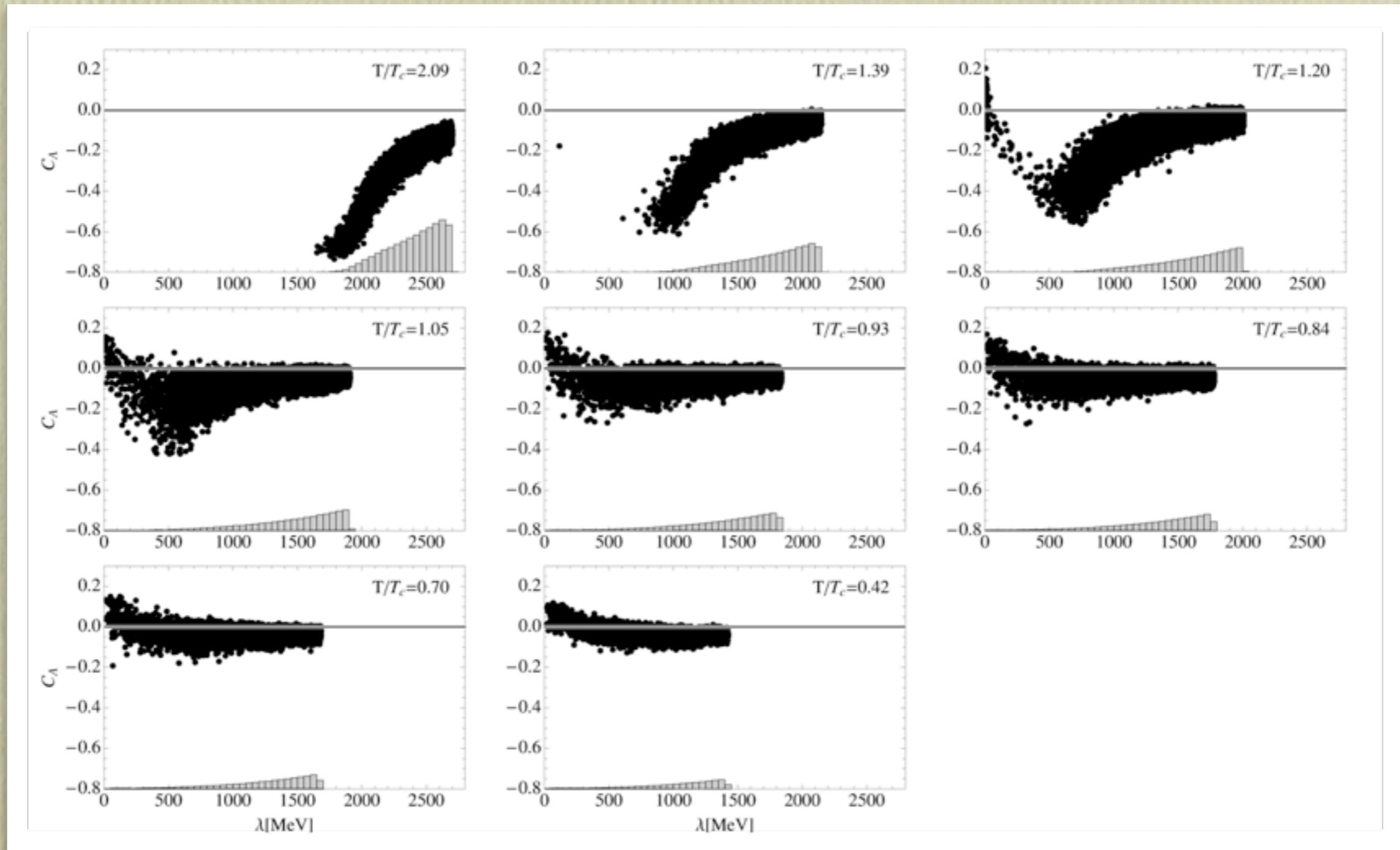
Chiral polarization scale



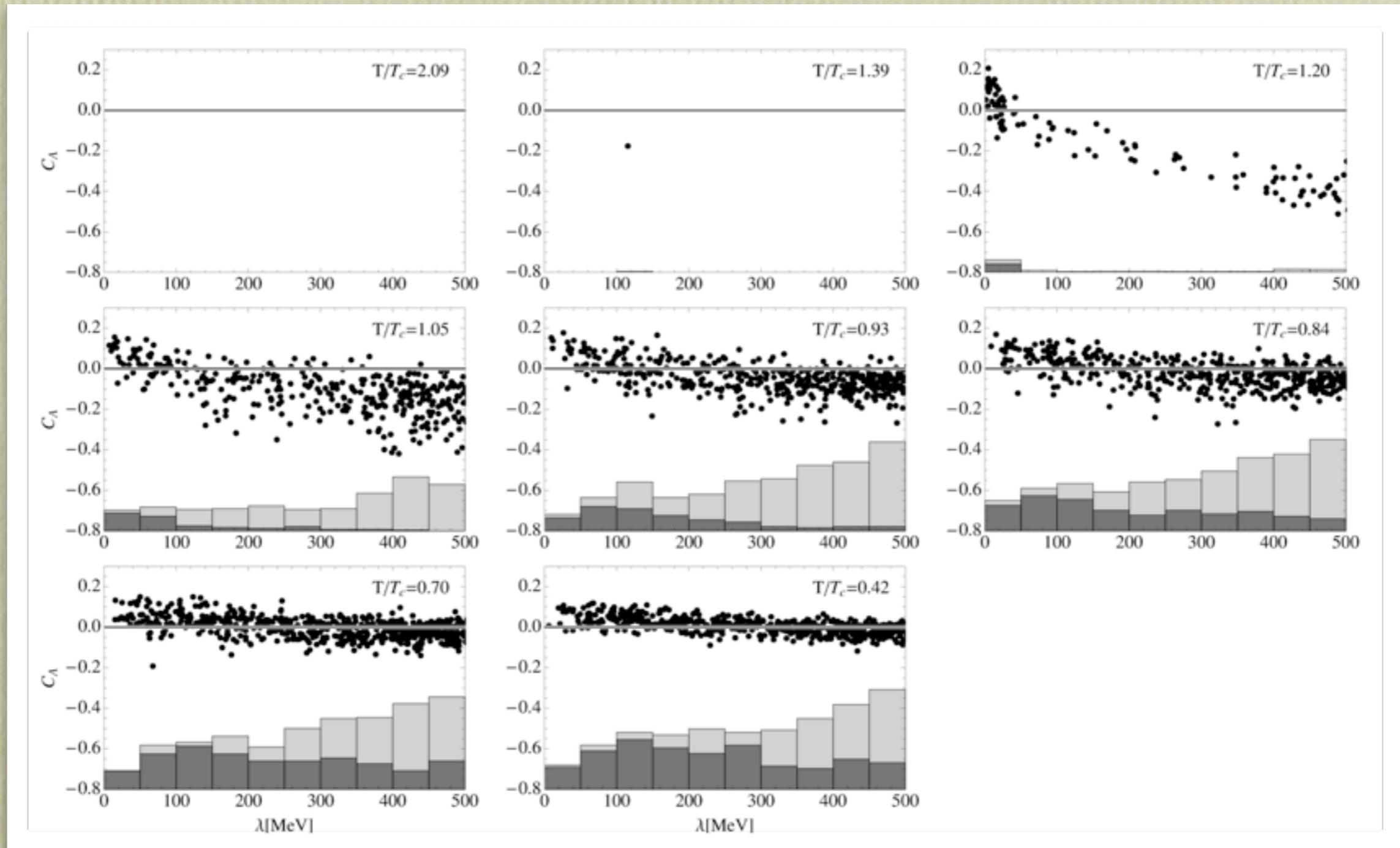
Individual mode polarization



Correlation coefficient



Correlation coefficient



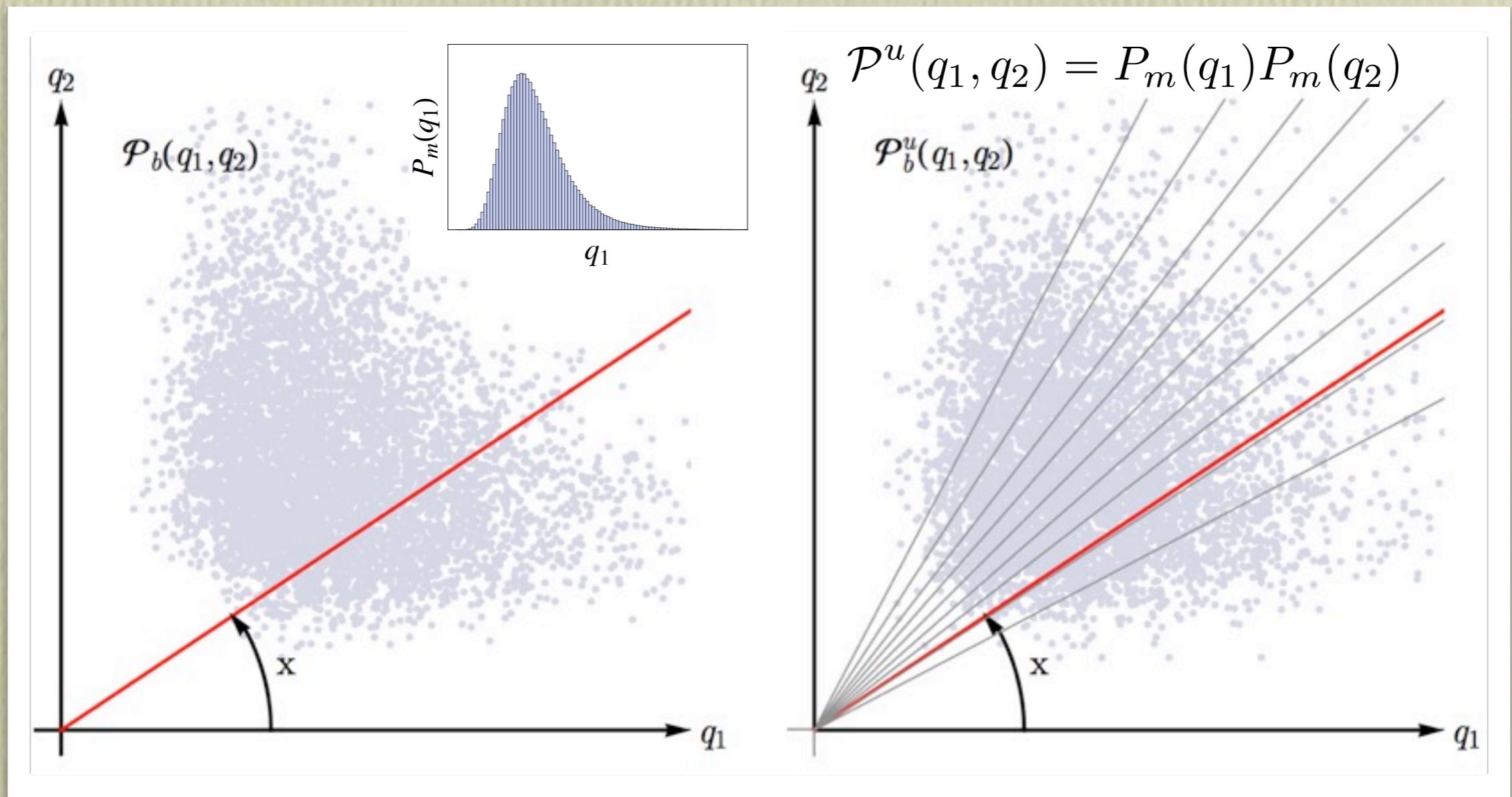
R. G. Edwards, U. M. Heller, J. E. Kiskis, and R. Narayanan, *Phys.Rev. D61* (2000) 074504.

Conclusions and outlook

- We computed the polarization scale for $SU(3)$ pure gauge configurations as a function of temperature
- We find that the polarization scale decreases as we approach the phase transition and vanishes above the phase transition
- The near zero modes, connected to chiral symmetry breaking, are mostly (possibly all) polarized
- Extrapolations from below the phase transition are consistent with a picture where the polarization scale vanishes at T_c
- However, the presence of a small density of near-zero modes at temperatures as high as $1.20 T_c$ modifies this picture
- These are polarized modes and it is not yet clear whether they are a quenching artifact, a finite volume effect or really dynamical

Backup slides

Dynamical polarization



Dynamical polarization

