

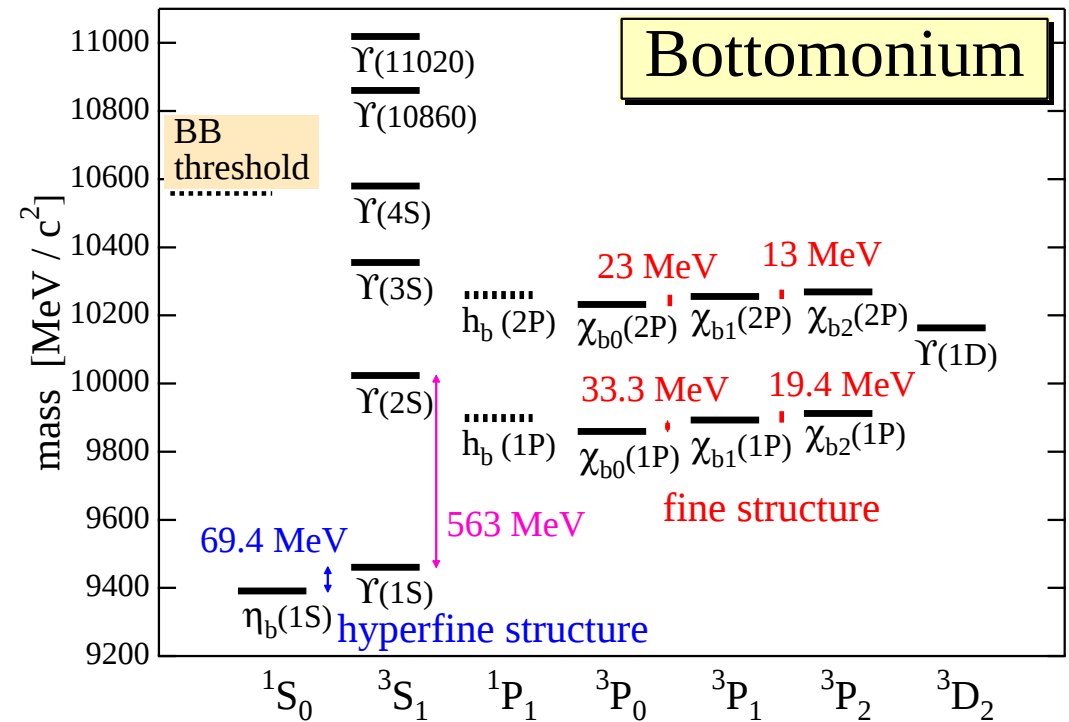
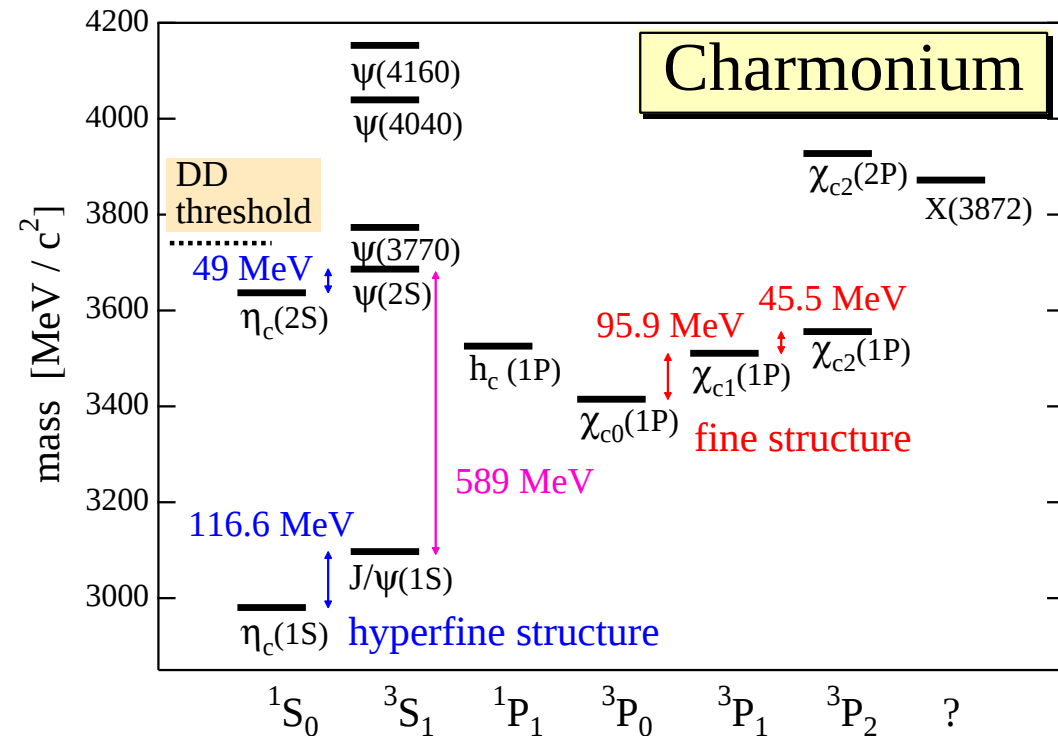
Heavy quarkonium spectroscopy in pNRQCD with lattice QCD input

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— Lattice 2012 —

Quarkonium

► Bound state of heavy quark and anti-quark



- Systematic patterns, HFS and FS (remnant of degeneracy)
 - ⇒ The challenge in **QCD**
 - ⇒ A promising approach: **pNRQCD with lattice QCD input**

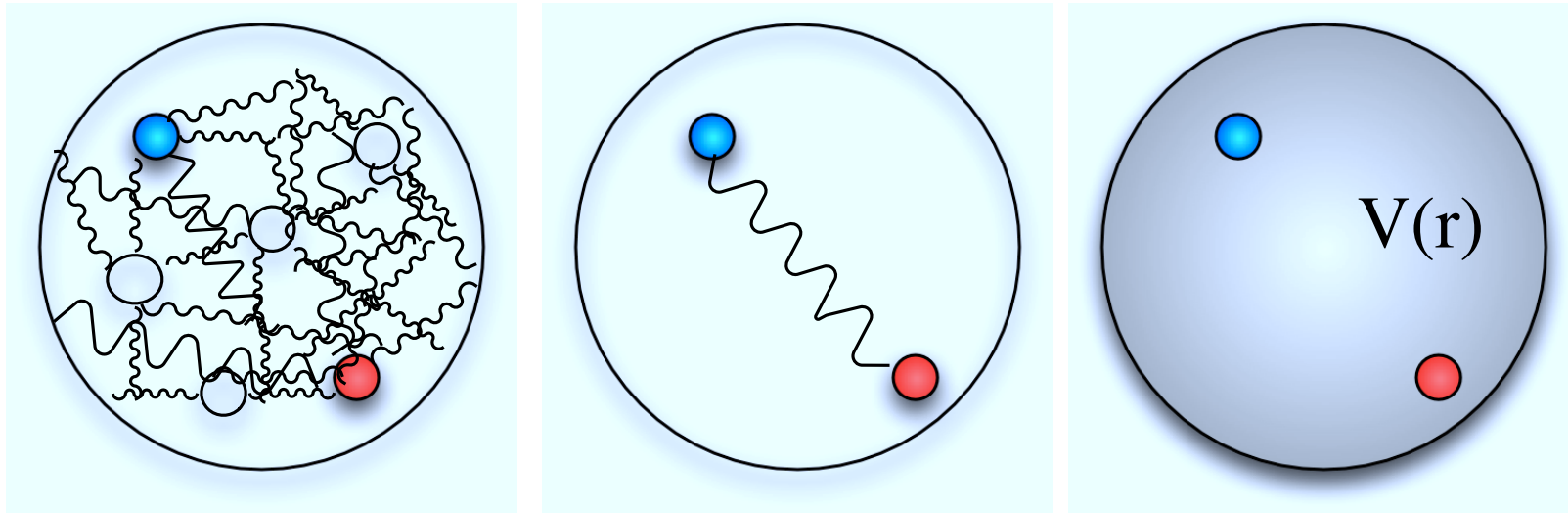
Contents

- ▶ pNRQCD ?
- ▶ Lattice QCD input ?
- ▶ Quarkonium spectroscopy ?

pNRQCD ?

- ▶ **Potential NRQCD** is proposed to make nonrelativistic nature of the quarkonium manifest by paying attention to the hierarchy of energy scales $m \gg mv \gg mv^2$ [Brambilla, Pineda, Soto&Vairo('99ff)]
- ▶ integrate out the scale above m (NRQCD) and mv (**pNRQCD**)

QCD $\Rightarrow$ NRQCD $\Rightarrow$ pNRQCD



- High energy contribution is inherited in the effective interaction in the resulting effective field theories

pNRQCD ?

► NRQCD [Caswell&Lepage('86), Manohar('97), Pineda&Soto('98), ...]

$$\begin{aligned}\mathcal{L} = & \psi^\dagger \left(\mathbf{c}_F g \frac{\boldsymbol{\sigma} \cdot \mathbf{B}}{2m} + \mathbf{c}_D g \frac{D \cdot E - E \cdot D}{8m^2} + i \mathbf{c}_S g \frac{\boldsymbol{\sigma} \cdot (\mathbf{D} \times \mathbf{E} - \mathbf{E} \times \mathbf{D})}{8m^2} + \dots \right) \psi \\ & + (\psi \rightarrow \chi) \\ & + \frac{d_{ss}}{m^2} \psi^\dagger \psi \chi^\dagger \chi + \frac{d_{sv}}{m^2} \psi^\dagger \boldsymbol{\sigma} \psi \cdot \chi^\dagger \boldsymbol{\sigma} \chi \\ & + \frac{d_{vs}}{m^2} \psi^\dagger T^a \psi \cdot \chi^\dagger T^a \chi + \frac{d_{vv}}{m^2} \psi^\dagger T^a \boldsymbol{\sigma} \psi \cdot \chi^\dagger T^a \boldsymbol{\sigma} \chi + \dots\end{aligned}$$

- Strength of interaction is controlled by effective couplings called the **matching coefficients**, which are determined to reproduce the same scattering amplitude both in QCD and NRQCD
⇒ **no need to use bound state information**

pNRQCD ?

► Matching coefficients in NRQCD (leading order)

- Bilinear term c_F, c_D, c_S [Manohar ('97)]

$$c_F = 1 + \frac{\alpha_s}{2\pi} (C_F + C_A) + O(\alpha_s^2)$$

$$c_D = 1 + \frac{\alpha_s}{2\pi} C_A + O(\alpha_s^2)$$

$$c_S = 2c_F - 1 \quad (C_F \text{ and } C_A \text{ are Casimir factor})$$

- Contact term $d_{ss}, d_{sv}, d_{vs}, d_{vv}$ [Pineda & Soto ('98)]

$$d_{ss} = \frac{2}{3}\pi\alpha_s + O(\alpha_s^2), \quad d_{sv} = -\frac{2}{9}\pi\alpha_s + O(\alpha_s^2)$$

$$d_{vs} = -\frac{1}{2}\pi\alpha_s + O(\alpha_s^2), \quad d_{vv} = \frac{1}{6}\pi\alpha_s + O(\alpha_s^2)$$

pNRQCD ?

- Since mv can be $O(\Lambda_{\text{QCD}})$, the matching between NRQCD and pNRQCD must be performed **nonperturbatively**
 $\Rightarrow$ **quantum mechanical matching** [Brambilla et al. ('99ff)]

1. write NRQCD hamiltonian

$$\mathcal{L}_{\text{NRQCD}} \rightarrow H = H^{(0)} + \frac{1}{m}H^{(1)} + \frac{1}{m^2}H^{(2)} + \dots$$

2. compute expectation values of $H^{(i \geq 1)}$ with the eigenstate of $H^{(0)}$ for $q-\bar{q}$ system and match them to **pNRQCD matching coefficients**

$$V_{\text{pNRQCD}}(r) = V^{(0)}(r) + \frac{1}{m}V^{(1)}(r) + \frac{1}{m^2}V^{(2)}(r) + \dots$$

$\Rightarrow$ **local in time but nonlocal in space (potential)**

$\Rightarrow$ **typical scale of binding energy mv^2**

Lattice QCD input ?

► pNRQCD potential

$$V(r) = V^{(0)}(r) + \frac{1}{m} V^{(1)}(r) + \frac{1}{m^2} V^{(2)}(r) + \dots$$

- $V^{(0)}$ is the static potential
- $V^{(1)}$ is typical in QCD which originates from 3-gluon vertex
- $V^{(2)}$ contains spin-dependent and spin-independent corrections

► Corrections are represented using the FSCs on the q - $\bar{q}$ source:

$$V^{(1)} = - \int dt \, t \langle\langle E(0) \cdot E(0) \rangle\rangle$$

$$V_{ls}^{(2)} = \frac{\mathbf{c}_S}{2r} \frac{dV^{(0)}}{dr} + \frac{\mathbf{c}_F}{r} \left(\underbrace{\epsilon_{ijk} \hat{r}_i \int dt \, t \langle\langle B^j(0) E^k(0) \rangle\rangle}_{V'_1} + \underbrace{\epsilon_{ijk} \hat{r}_i \int dt \, t \langle\langle B^j(0) E^k(r) \rangle\rangle}_{V'_2} \right)$$

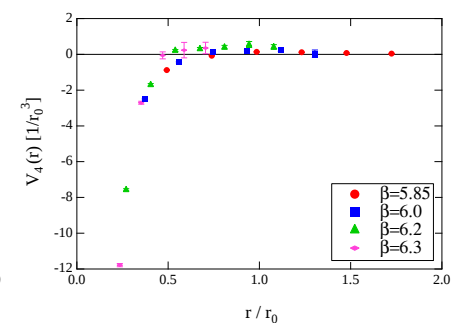
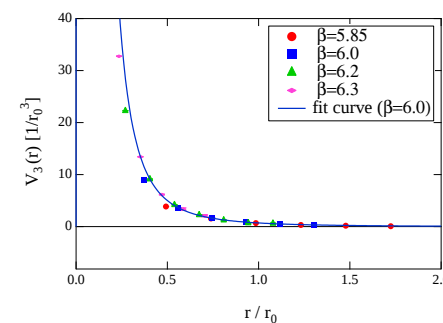
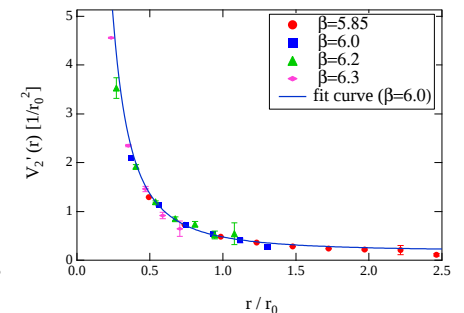
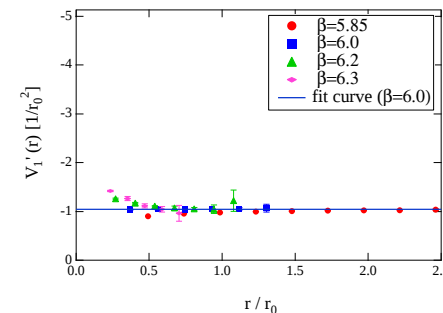
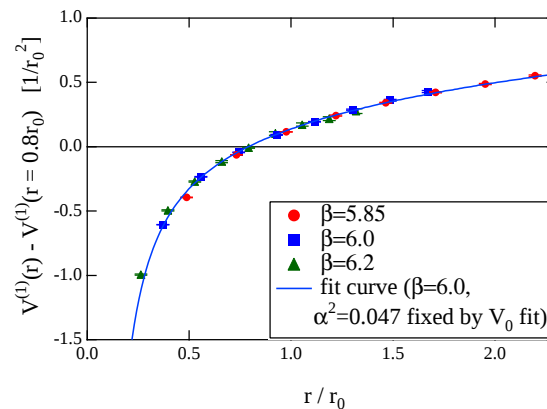
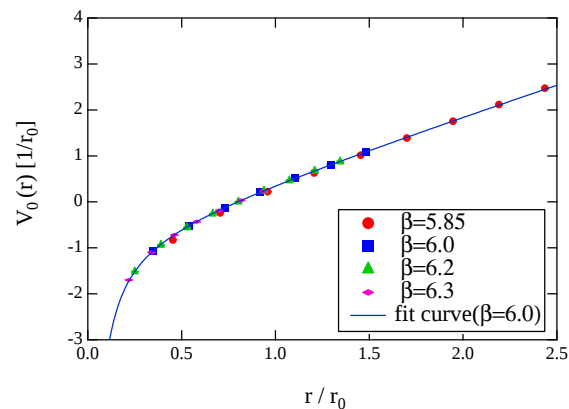
⇒ compute by **lattice QCD simulations**

Lattice QCD input ?

- Lattice QCD potentials up to $O(1/m^2)$ [Koma, Koma & Wittig('06ff)]

$$V(r) = V^{(0)} + \frac{1}{m}V^{(1)} + \frac{1}{m^2}V^{(2)}$$

⇒ compute the FSCs on the PLCF with the multilevel algorithm



Quarkonium spectroscopy ?

- ▶ Solve the Schrödinger equation

$$H^{(0)}\psi = \left(\frac{p^2}{m} + V^{(0)} \right) \psi = E^{(0)}\psi$$

- ▶ Add the corrections in the first order perturbation theory

$$E = E^{(0)} + \delta E = E^{(0)} + \langle \psi | \delta V | \psi \rangle$$

- ▶ Investigate the effects of $V^{(1)}$ and $V_{ls}^{(2)}$ (others are to come...)

$$\delta E^{(1)} = \frac{1}{m} \langle \psi | V^{(1)} | \psi \rangle$$

$$\delta E_{ls}^{(2)} = \frac{1}{m^2} \langle \psi | \left(\frac{\mathbf{c}_s}{2r} \frac{dV^{(0)}}{dr} + \frac{\mathbf{c}_F}{r} (V_1' + V_2') \right) l \cdot s | \psi \rangle$$

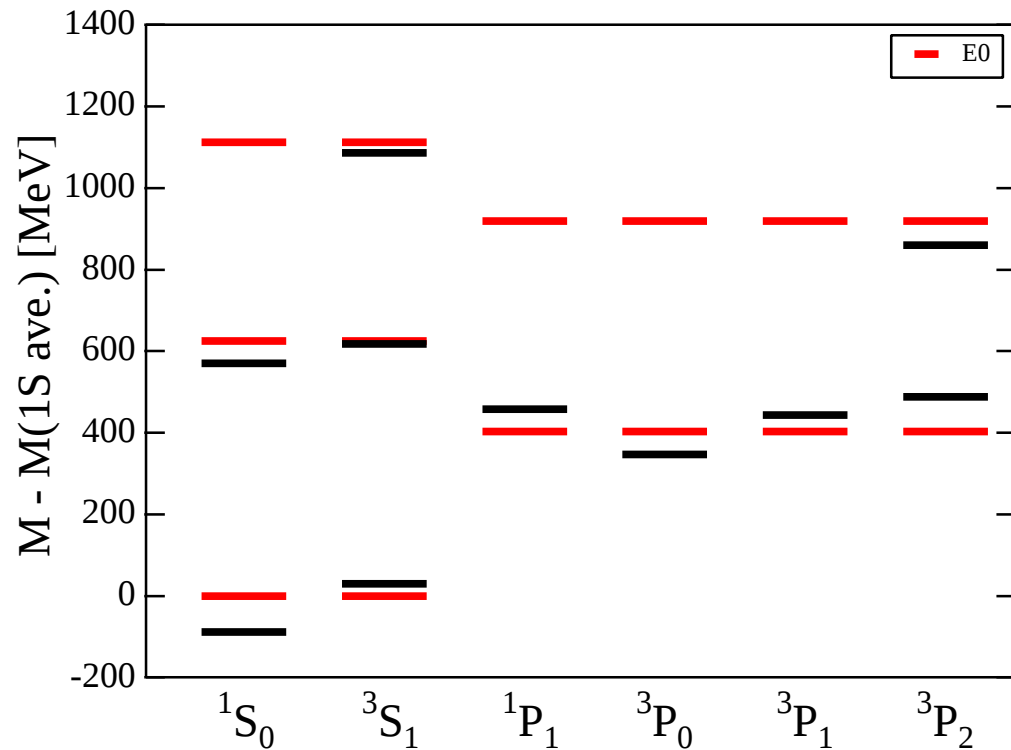
Quarkonium spectroscopy ?

- ▶ Parameters are fixed by lattice QCD except for the quark masses (and constant shift)
 - $r_0 = 0.50$ fm to fix lattice spacing
 - $V^{(0)} = -\frac{\alpha}{r} + \sigma r$ ($\alpha = 0.297$, $\sigma = 1.06$ GeV/fm)
 - $\delta E^{(1)} = \frac{1}{m} \langle \psi | -\frac{9\alpha^2}{8r^2} + \sigma^{(1)} \ln r | \psi \rangle$ ($\alpha = 0.297$, $\sigma^{(1)} = 0.142$ GeV²)
 - $\delta E_{ls}^{(2)} = \frac{1}{m^2} \langle \psi | \left(\frac{3\alpha}{2r^3} - \frac{(1-4\epsilon)\sigma}{2r} \right) l \cdot s | \psi \rangle$ ($\mathbf{c_F} = \mathbf{c_S} = 1$, $\epsilon = 0.2$)
 $\Rightarrow \epsilon$ characterizes long-range contribution in V_2'
- $m_c = 1.27$ GeV (charm), $m_b = 4.20$ GeV (bottom)

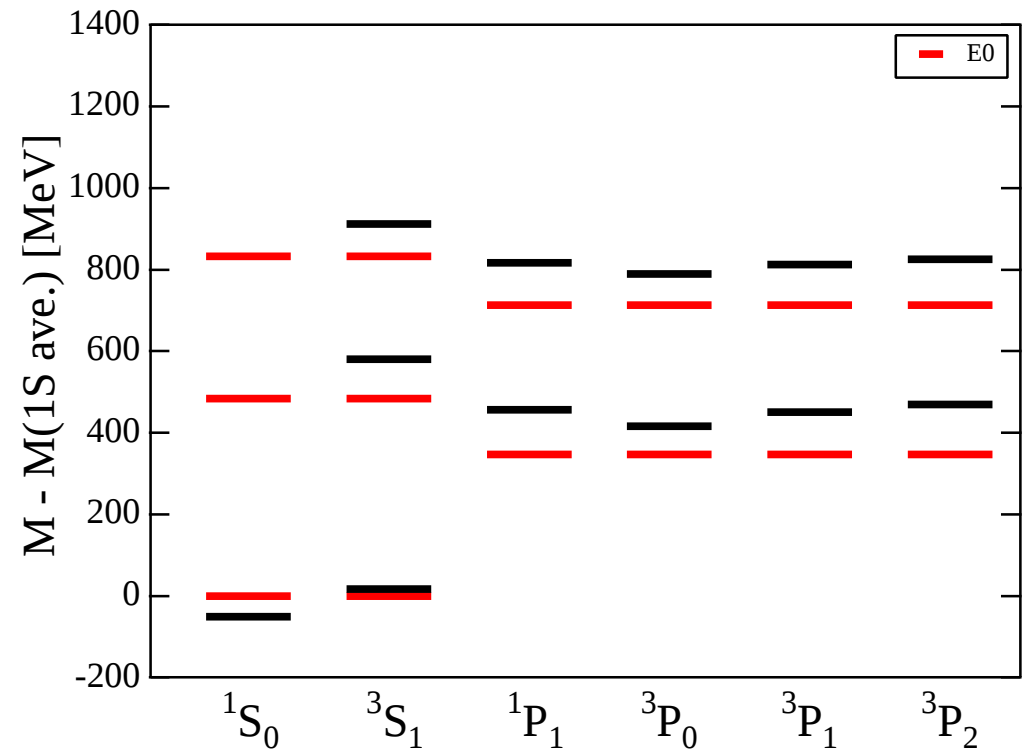
Quarkonium spectroscopy ?

► $E = E^{(0)}$

charmonium



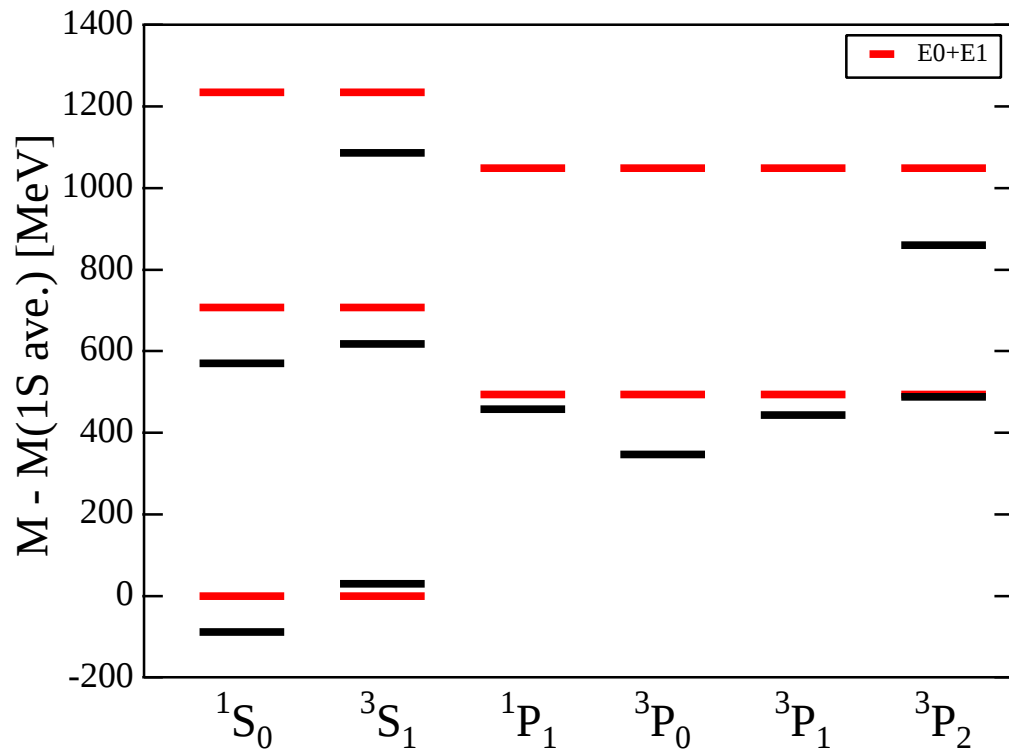
bottomonium



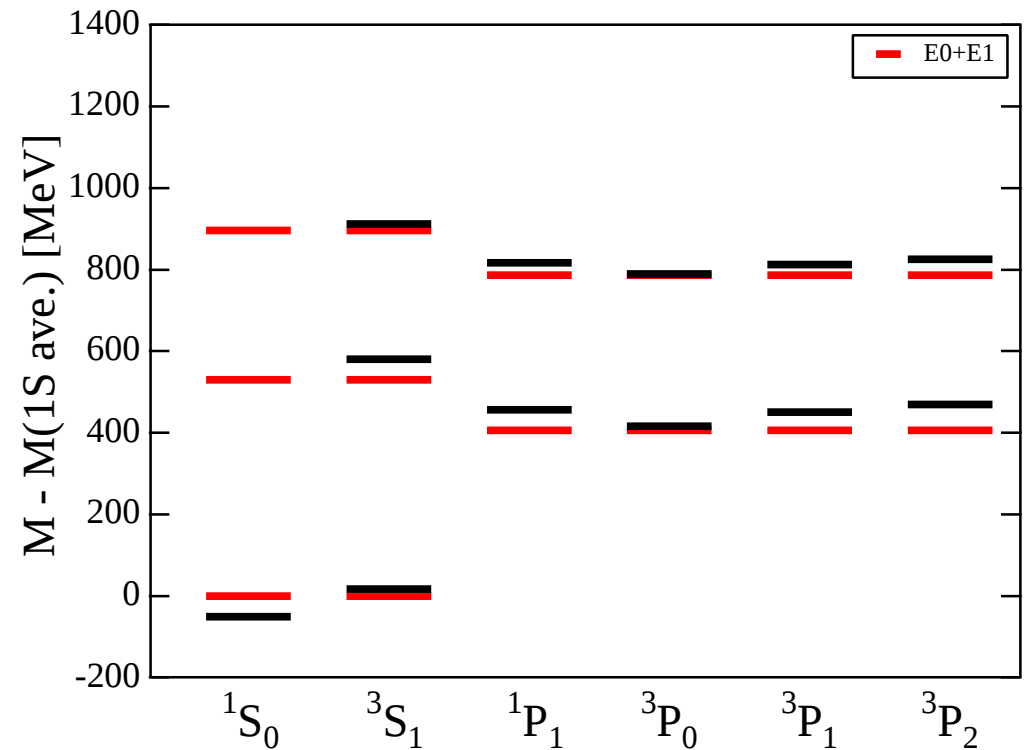
Quarkonium spectroscopy ?

► $E = E^{(0)} + \delta E^{(1)}$

charmonium



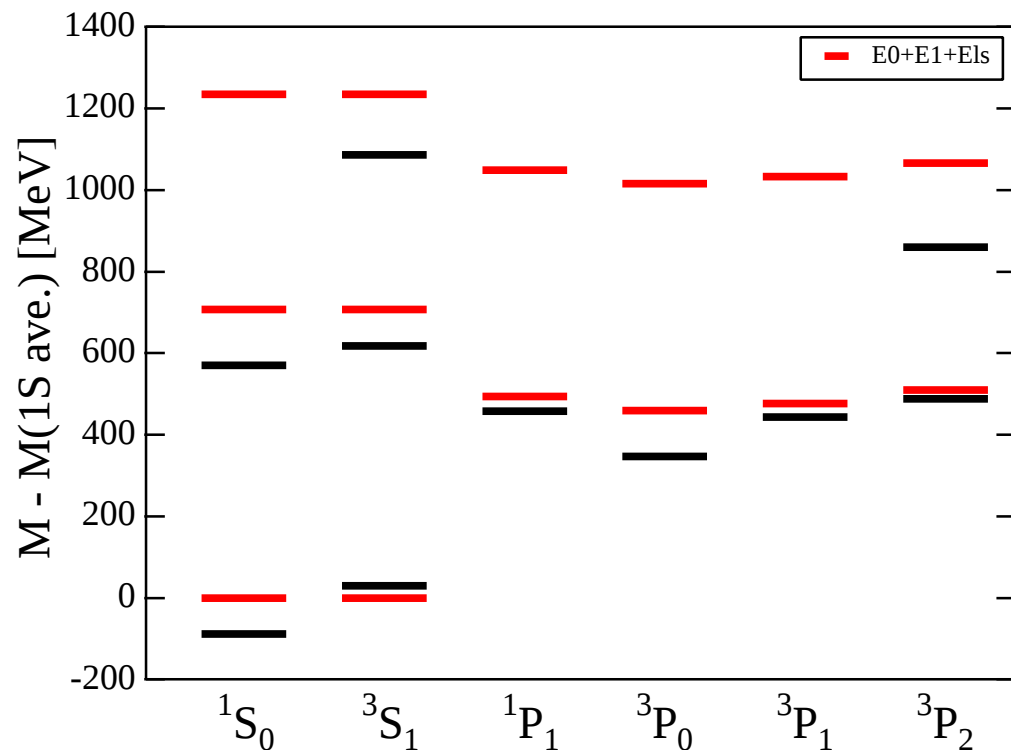
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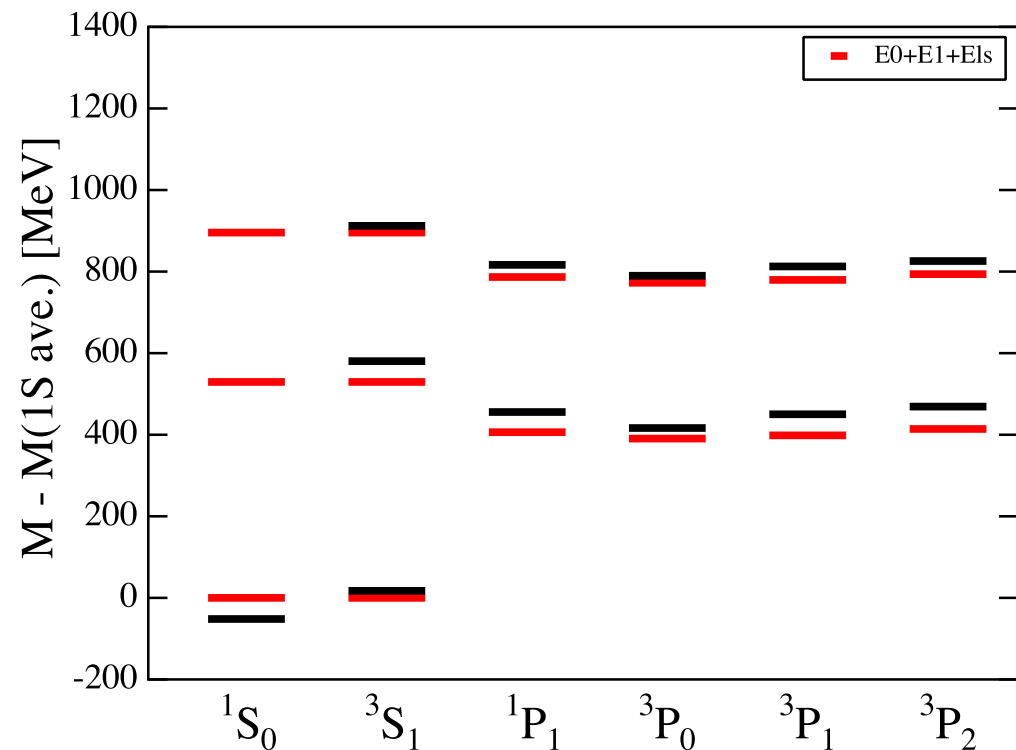
Quarkonium spectroscopy ?

► $E = E^{(0)} + \delta E^{(1)} + \delta E_{ls}^{(2)}$

charmonium



bottomonium



Summary and outlook

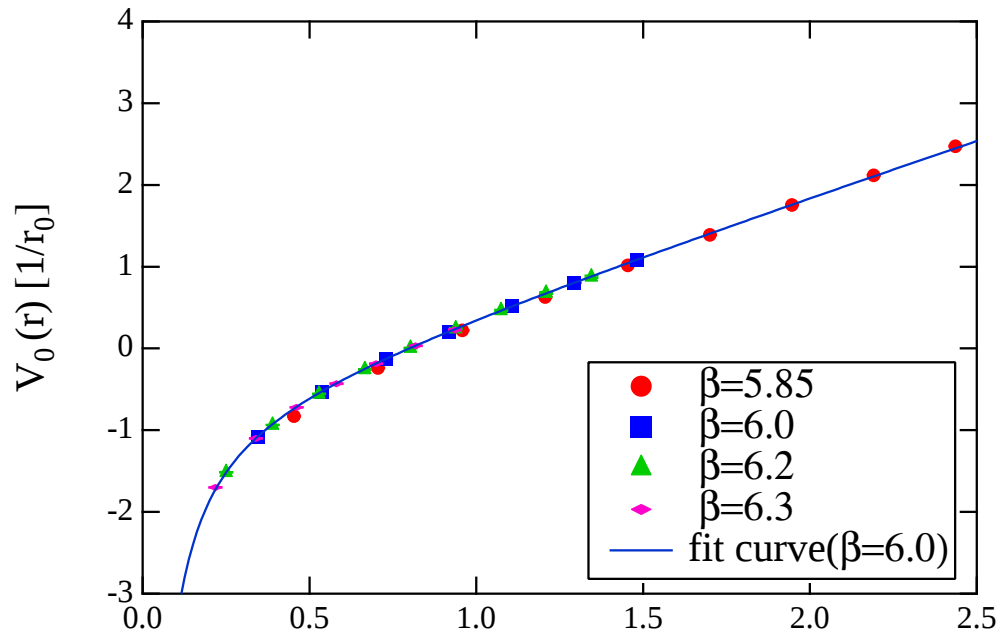
- ▶ We have studied heavy quarkonium spectra systematically based on pNRQCD with lattice QCD input
 - pNRQCD is derived by paying attention to the hierarchy of energy scales in the heavy quarkonium system $m \gg mv \gg mv^2$
 - Nonperturbative input in pNRQCD (potentials) can be computed by using lattice QCD
[cf. LECs in χ PT]
- ▶ Compute the spectra with all $O(1/m^2)$ corrections
 - Short distance corrections to the lattice potentials
 - NLO (and further) matching coefficients in NRQCD

Appendix

- ▶ $V^{(0)}$
- ▶ $V^{(1)}$
- ▶ $V_{ls}^{(2)}$
- ▶ **Effect of $V^{(1)}$**
- ▶ **Effect of $V_{ls}^{(2)}$**
- ▶ **Effect of $V_{ls}^{(2)}$ with tensor term**

$$V^{(0)}$$

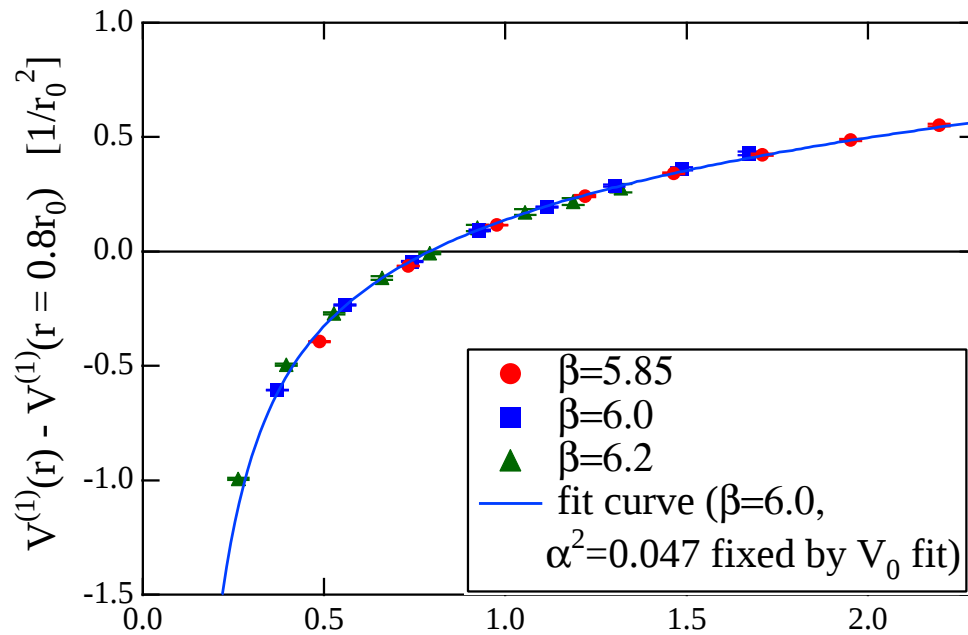
► $V(r) = V^{(0)}(r) + \delta V(r)$



Pert.: $V^{(0)} = -\frac{\alpha}{r} \quad (\alpha \equiv C_F \alpha_s) \quad \Rightarrow \quad \text{NP: } V^{(0)} = -\frac{\alpha}{r} + \sigma r + c^{(0)}$

$$V^{(1)}$$

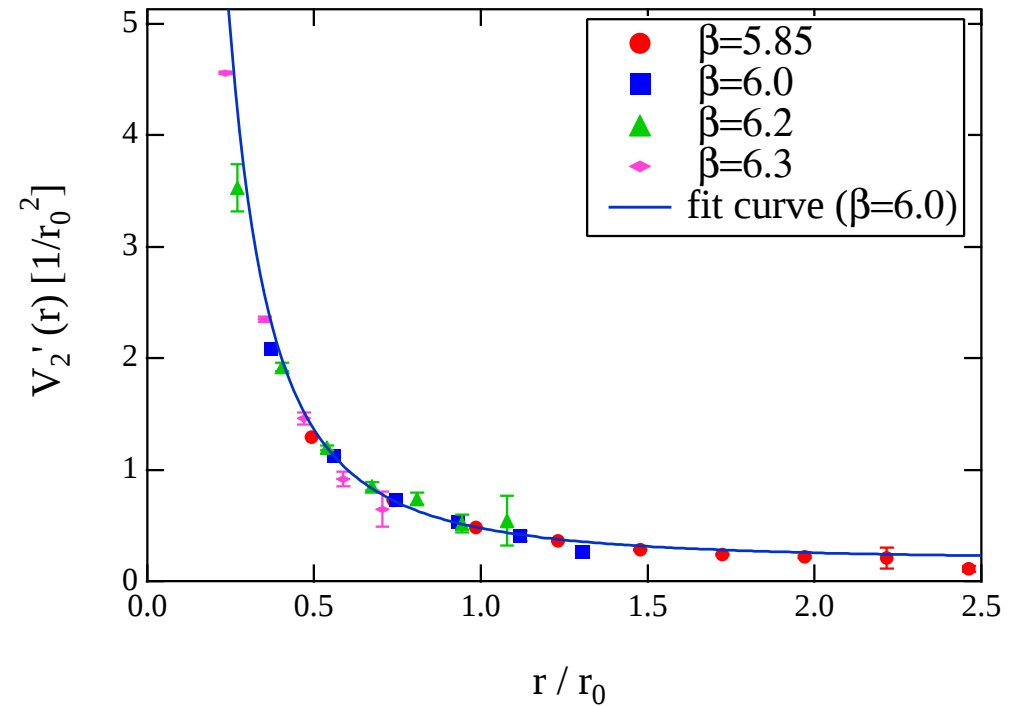
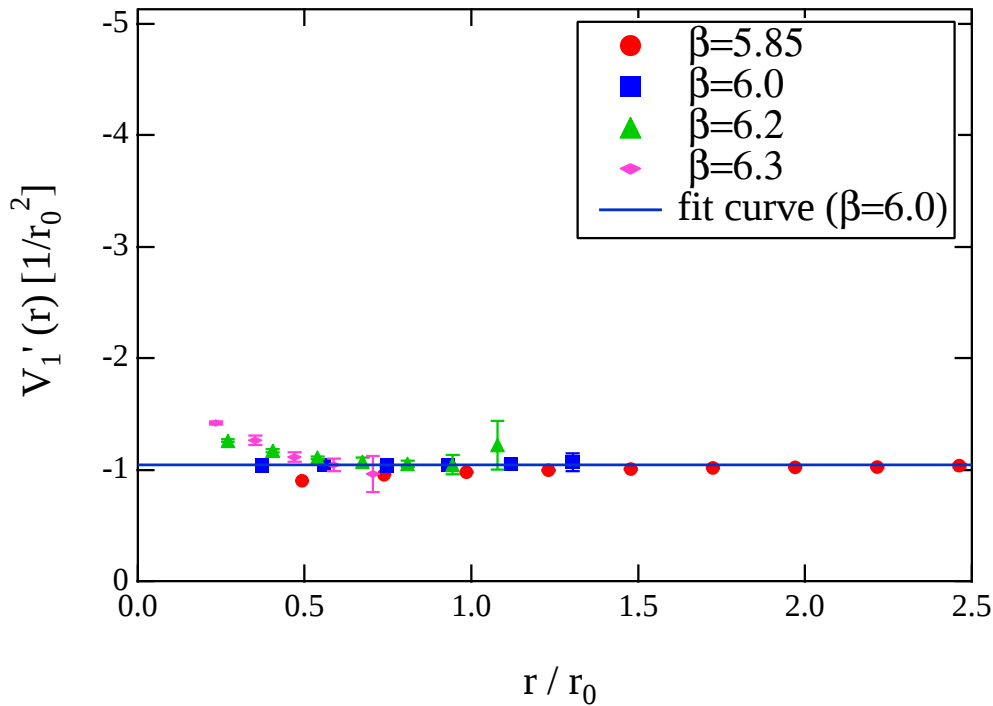
► $\delta V(r) = \frac{1}{m} V^{(1)}(r)$ (*y-axis of fig is $V^{(1)}/2$ in this definition*)



$$\text{Pert.: } V^{(1)} = -\frac{9\alpha^2}{8r^2} \quad \Rightarrow \quad \text{NP: } V^{(1)} = -\frac{9\alpha^2}{8r^2} + \sigma^{(1)} \ln r + c^{(1)}$$

$V_{ls}^{(2)}$

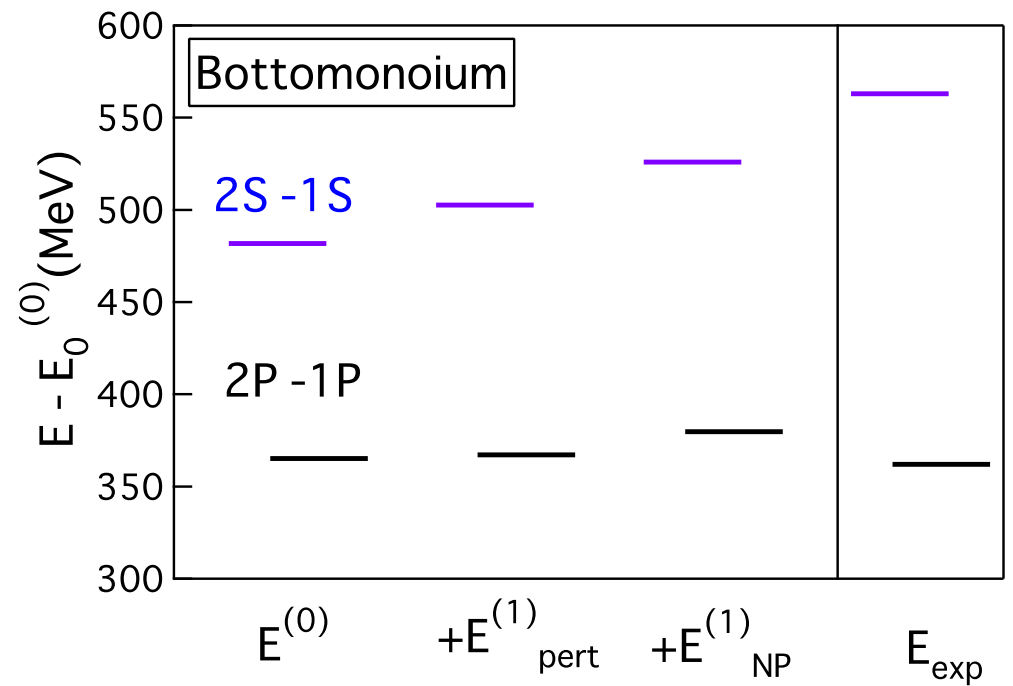
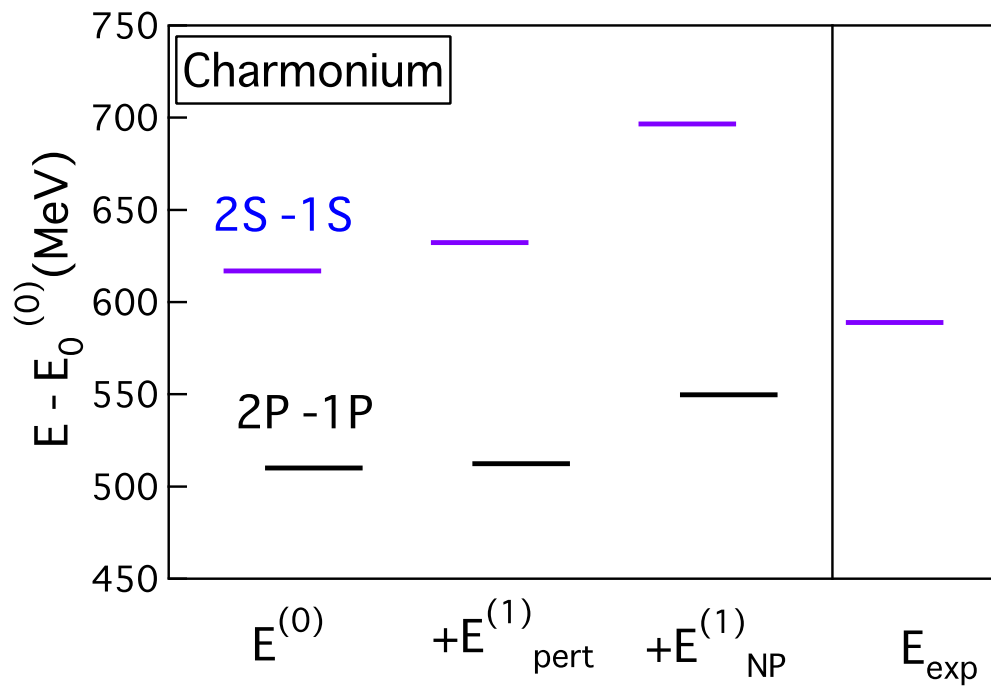
$$\blacktriangleright \delta V(r) = \frac{1}{m^2} \left(\frac{\mathbf{c}_s}{2r} \frac{dV^{(0)}(r)}{dr} + \frac{\mathbf{c}_F}{r} (V_1'(r) + V_2'(r)) \right) l \cdot s$$



Pert.: $V_1' = 0, V_2' = \alpha/r^2 \Rightarrow$ **NP:** $V_1' = -(1-\epsilon)\sigma, V_2' = \alpha/r^2 + \epsilon\sigma$

Effect of $V^{(1)}$

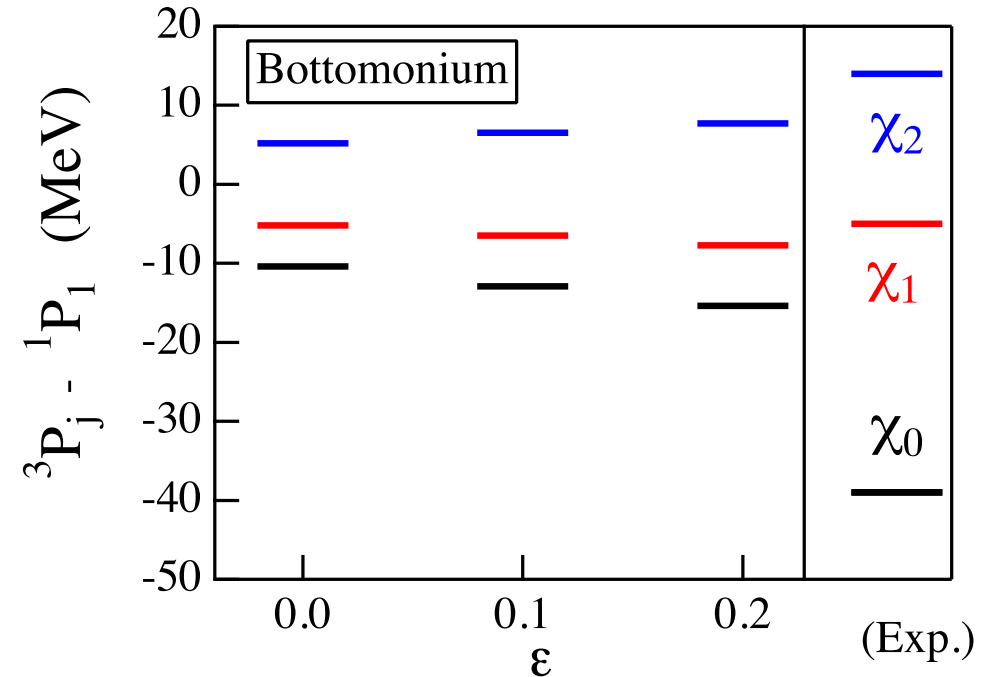
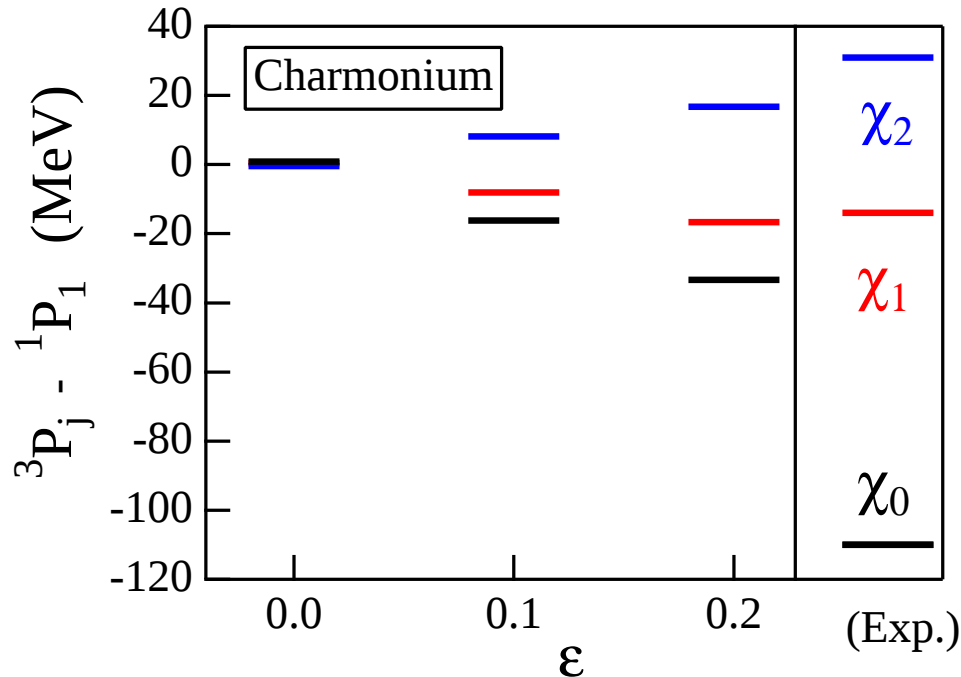
$$\blacktriangleright \delta E^{(1)} = \frac{1}{m} \langle \psi | -\frac{9\alpha^2}{8r^2} + \sigma^{(1)} \ln r | \psi \rangle = \delta E_{\text{pert}}^{(1)} + \delta E_{\text{NP}}^{(1)}$$



Effect of $V_{ls}^{(2)}$

► $\delta E_{ls}^{(2)} = \frac{1}{m^2} \langle \psi | \left(\frac{3\alpha}{2r^3} - \frac{(1-4\epsilon)\sigma}{2r} \right) l \cdot s | \psi \rangle, \quad \langle {}^3P_J | l \cdot s | {}^3P_{J=0,1,2} \rangle = -2, -1, 1$

► ϵ dependence of the level

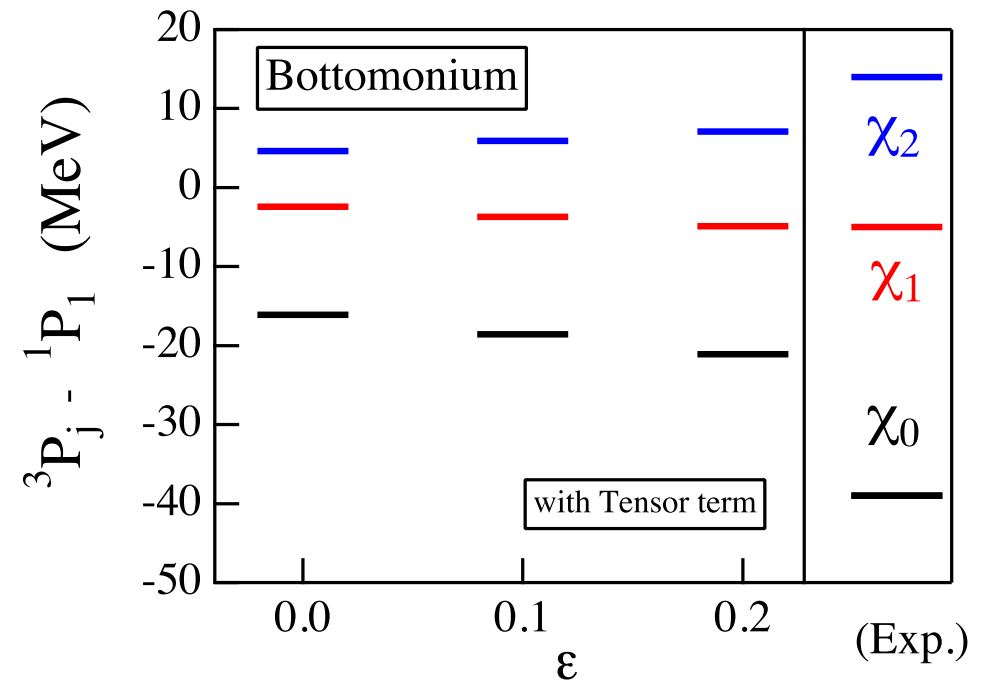
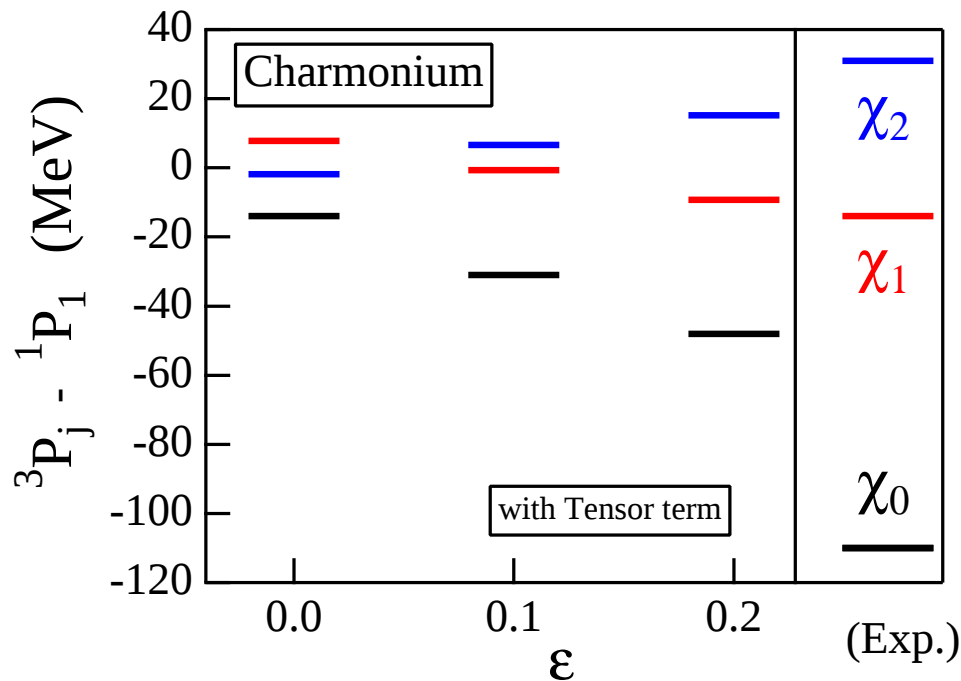


• from lattice QCD $\Rightarrow \epsilon \sim 0.2$

Effect of $V_{ls}^{(2)}$

► $\delta E_{ls}^{(2)} = \frac{1}{m^2} \langle \psi | \left(\frac{3\alpha}{2r^3} - \frac{(1-4\epsilon)\sigma}{2r} \right) l \cdot s | \psi \rangle, \quad \langle {}^3P_J | l \cdot s | {}^3P_{J=0,1,2} \rangle = -2, -1, 1$

► ϵ dependence of the level (include tensor term)



• $V_3 = 3\alpha/r^3, \quad \langle {}^3P_J | s_{12} | {}^3P_{J=0,1,2} \rangle = -1/3, 1/6, -1/30$