The string tension at large N from smeared Wilson loops

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Large N gauge theory

- Large N gauge theories are very interesting: Simpler, close to phenom., connected to string picture, etc.
- Interesting simplifying properties: planar PT, factorization, stable resonances, etc.
- LGT provides first-principles approach to gauge theories.
- Combining both is numerically challenging ($\times N^2$ degrees of freedom).
- An early (80's) look at the lattice loop eqs. at large N led to an interesting observation

Schwinger-Dyson equations for Wilson loops adopt a simple form on the lattice

Expectation values of Wilson loops become volume independent in the large N limit (Assuming Z_N^4 symmetry is unbroken) \Rightarrow Eguchi-Kawai model: LGT in one point

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If NOT it would be possible to modify lattice model.

- Quenched EK model (Bhanot, Heller, Neuberger)
- Twisted EK model (A.G-A, Okawa)
- Double trace deformations (Shifman-Unsal-Yaffe)
- Cut-off scale adjoint fermions (Kovtun-Unsal-Yaffe)

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Alternatively, it could happen that reduction only applies beyond a certain length scale $l > l_c$ (Narayanan-Neuberger). Reduction could still work for other large N theories

TEK Model

Use twisted boundary conditions (magnetic flux):

$$\mathcal{S}=-b \mathcal{N}\sum_{\mu,
u} z^{(\mathcal{N})}_{\mu
u} ext{Tr}(U_{\mu}U_{
u}U^{\dagger}_{\mu}U^{\dagger}_{
u})$$

With appropriate choice of $z_{\mu\nu}^{(N)}$ symmetry subgroup unbroken at weak coupling; $N = L^2$ $z_{\mu\nu}^{(N)} = \exp\{i2\pi k\epsilon_{\mu\nu}/L\}$ $Z_N^4 \longrightarrow Z_N^2 = Z_L^4 \Leftrightarrow \text{Reduction } (N^2 = L^4)$

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- Teper and Vairinhos found breaking for k=1 N > 100 Explanation: Other minima dominate (First order transition)
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- Teper and Vairinhos found breaking for k=1 N > 100 Explanation: Other minima dominate (First order transition)
- Azeyanagi et al showed \Rightarrow No continuum reduction
- AGA-Okawa 2010: No breaking if $k/L \longrightarrow \text{const} \neq 0$.

conclusions

Checking our prescription

- Complicated semiclassical analysis. Ongoing 2+1 dim. analysis a first step (M. Garcia Perez talk)
- Check of symmetry breaking up to $N = 37^2$
- Check the physics (This study arXiV:1206.0049)
 Extract String tension from SU(N) LGT ⇒ Take continuum
 limit ⇒ Extrapolate to large N ⇒ Compare with TEK

$$\begin{array}{ll} N=3 & b \in [0.3278, 0.3611] \mbox{ (7 values)} & 260 \mbox{ confs. } 32^4 \\ N=5 & b \in [0.3513, 0.3772] \mbox{ (5 values)} & 260 \mbox{ confs. } 32^4 \\ N=6 & b \in [0.3541, 0.3792] \mbox{ (5 values)} & 260 \mbox{ confs. } 32^4 \\ N=8 & b \in [0.3569, 0.3815] \mbox{ (5 values)} & 260 \mbox{ confs. } 32^4 \\ N=4 & b \in [0.3464, 0.3725] \mbox{ (5 values)} & 260 \mbox{ confs. } 32^4 \\ N=29^2 & b \in [0.36, 0.385] \mbox{ (5+1 values)} & 5400 \mbox{ confs TEK} \\ \end{array}$$

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Observables

- Must use the same method for LGT and TEK (single loops, symmetric lattice, ...)
- Use rectangular Wilson loops W(R, T) for $R \approx T$
- To eliminate constant + perimeter dependence: Creutz ratios

$$\chi(T,R) = -\log \frac{W(T+0.5, R+0.5)W(T-0.5, R-0.5)}{W(T+0.5, R-0.5)W(T-0.5, R+0.5)}$$

 One must use 4d APE Smearing for noise reduction We eliminate the effect of smearing on χ(T, R) by extrapolating back.

Results

- For string tension we use square loops $R = T \in [3.5, 8.5]$
- At each N and b the 6 values are well-fitted to a 3-param formula

$$\chi(R,R) = \kappa + \frac{2\gamma}{R^2} + \frac{\eta}{R^4}$$

Good Fits Click

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Continuum extrapolation

- Data display scaling with 1% accuracy for top 5 b values and intermediate size loops
- To fix the scale we use 2 perturbative prescriptions with improved coupling definitions (Edwards et al and Allton et al.)
- Non-perturbative prescription (similar to Sommer scale)

$$-RT\chi(R,T)|_{R=T=r_0a(b)}=1.65$$

Large N limit

• Using Allton et al a(b) we made a linear fit in $1/N^2$ Click

 $\Lambda_{\overline{\mathrm{MS}}}/\sqrt{\sigma} = 0.525(2)$ for TEK 0.523(5)

- \bullet Agrees with N dependence by Allton (slope $\approx 0.3)$
- Value is off by 4-5% (compatible within systematic errors) (See Lohmayer talk)
- Most of the N-dependence is in $r_0 \Lambda_{\overline{\mathrm{MS}}}$ Click

Approach to large sizes

The parameter γ defining approach to large R is very interesting: Our value $\gamma=0.272(5)$

• Slope is correlated with s.t. value

$$\delta\sigma\approx-\frac{\delta\gamma}{20a^2}$$

- These parameters are connected with the string picture of the chromoelectric flux-tube (Luscher-Symanzik-Weisz)
- This can be generalized to rectangular loops. We have

$$\chi(R,T) = \kappa + \phi(z)(\frac{1}{R^2} + \frac{1}{T^2}) + \dots$$

with z = R/T (aspect ratio).

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with z = R/T (aspect ratio).

Our values do not coincide with Nambu-Goto open string parameters.

Aspect ratio dependence similar to I.o. perturbation theory **PLOT**

▶ FORMULAS

Conclusions

- TEK model works up to 1-2% precision in the physical range of params.
- We obtained a very precise determination of the large N string tension
- N-dependence is small and lies mostly in the scale
- String tension value is lower by 4-5% than Allton et al. (Systematics should be analysed further)
- Interesting results were obtained for the parameters describing approach to large sizes. Our results differ from the predictions of the Nambu-Goto string.

A more detailed account of our results including volume-dependence and additional data (N = 4, 10 and TEK at $N = 37^2$), is in progress.



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$$\phi_{NG} = -\frac{1}{1+1/z^2} \frac{\partial}{\partial z} \left(z \frac{\partial \log(\eta(iz))}{\partial z} \right)$$
$$\phi_{PT} = \frac{(N^2 - 1)g^2}{4\pi^2 N} \frac{1 + z \operatorname{atan}(z) + \operatorname{atan}(1/z)/z}{z + 1/z}$$



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