

Two-Color Schrödinger Functional with Six Flavors of Stout Smeared Wilson Fermions

Gennady Voronov

Yale University

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Lattice Strong Dynamics (LSD) Collaboration



Heechang Na
James Osborn



Michael Buchoff
Chris Schroeder
Pavlos Vranas



Richard Bower
Michael Cheng
Claudio Rebbi
Oliver Witzel



Mike Clark
Ron Babich



David Schaich



Joe Kiskis



Ethan Neil



Saul Cohen



Sergey Syritsyn



Tom Appelquist
George Fleming
Meifeng Lin
Gennady Voronov

Motivation

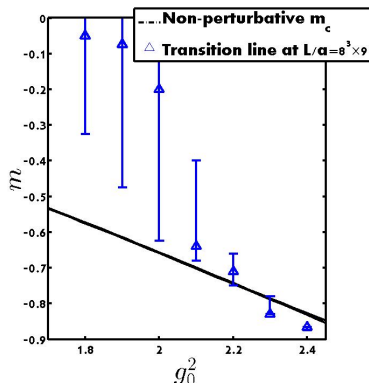
- ▶ A technicolor model based on a walking gauge theory may explain electroweak symmetry breaking while avoiding experimental constraints for SM fermion masses and flavor-changing neutral currents. Such a theory is expected to reside just below the conformal window.
- ▶ $SU(2)$ gauge theories are special in that there is an enhanced $SU(2N_f)$ global symmetry.
- ▶ Even color theories might be particularly interesting for dark matter model building.
- ▶ Evidence for IRFP at 8^1 and 10 flavors and χ SB at 4 flavors². Several inconclusive 6 flavors calculations.

¹Itou et al.

²Karavirta et al

Smearing Motivations

- ▶ Previously were unable to reach sufficiently large SF \bar{g}^2 with six flavors of Wilson fermions due to presence of bulk phase transition along m_c line.
- ▶ Try to get to stronger renormalized coupling by smearing the gauge field in the fermion part of the action.



Stout Smearing³

- ▶ Smearing technique which is analytic and therefore can be implemented in tandem with HMC algorithm.

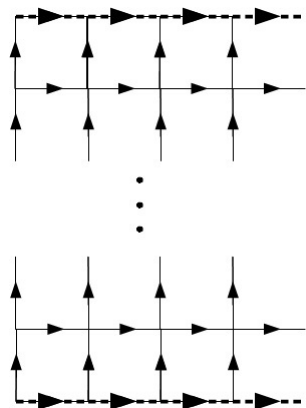
- ▶
$$\Omega_\mu(x) = \left[\sum_{\nu \neq \mu} \rho_{\mu\nu} (U_\nu(x) U_\mu(x + \hat{\nu}) U_\nu^\dagger(x + \hat{\mu}) + U_\nu^\dagger(x - \hat{\nu}) U_\mu(x - \hat{\nu}) U_\nu(x - \hat{\nu} + \hat{\mu})) \right] U_\mu^\dagger(x)$$

$$= \sum_{\nu \neq \mu} \rho_{\mu\nu} \left(\begin{array}{c} \uparrow \rightarrow \uparrow \\ \downarrow \leftarrow \downarrow \end{array} + \begin{array}{c} \downarrow \rightarrow \downarrow \\ \uparrow \leftarrow \uparrow \end{array} \right) \leftarrow \begin{array}{c} \nu \\ \uparrow \\ \mu \end{array}$$

- ▶
$$Q_\mu(x) = \frac{i}{2} \left(\Omega_\mu^\dagger(x) - \Omega_\mu(x) \right) - \frac{i}{2N} \text{Tr} \left(\Omega_\mu^\dagger(x) - \Omega_\mu(x) \right)$$
- ▶
$$U_\mu^{(n+1)}(x) = \exp \left(i Q_\mu^{(n)}(x) \right) U_\mu^{(n)}(x)$$

Stout Smearing with Dirichlet BC

- Care must be taken to implement smearing in such a way that the boundary gauge field is not smeared by the bulk gauge field.



Smearing of:

- adjacent link

$$\uparrow = \left(\begin{array}{c} \text{---} \rightarrow \\ \downarrow \\ \leftarrow \end{array} + \begin{array}{c} \text{---} \leftarrow \\ \downarrow \\ \rightarrow \end{array} \right) \downarrow$$

- link near boundary

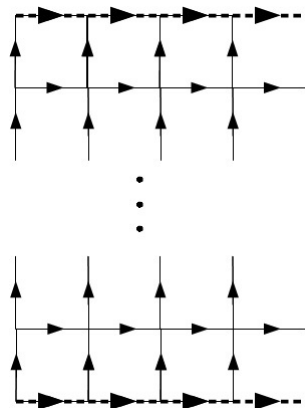
$$\rightarrow = \left(\begin{array}{c} \uparrow \\ \text{---} \rightarrow \\ \downarrow \end{array} + \begin{array}{c} \downarrow \\ \leftarrow \\ \uparrow \end{array} \right) \leftarrow$$

- boundary link

$$\text{---} \rightarrow \text{---} = 0$$

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Smearing of:

- adjacent link

$$\uparrow = \left(\begin{array}{c} \text{---} \rightarrow \\ \downarrow \\ \leftarrow \text{---} \end{array} \right) \downarrow$$

- link near boundary

$$\rightarrow = \left(\begin{array}{c} \text{---} \rightarrow \\ \uparrow \\ \downarrow \end{array} \right) \leftarrow$$

- boundary link

$$\text{---} \rightarrow \text{---} = 0$$

Critical Mass I

- ▶ Wilson fermion mass is additively renormalized, no chiral symmetry at $m_0 = 0$.
- ▶ Need to determine the critical mass $m_c(g_0^2)$, i.e. the input bare mass that result in zero renormalized quark mass.
- ▶ Define⁴ a renormalized quark mass using the lattice PCAC relation:

$$\left\langle \frac{1}{2} (\partial_\mu^* + \partial_\mu) (A_R)_\mu^a(x) \mathcal{O} \right\rangle = 2m_R \langle (P_R)^a(x) \mathcal{O} \rangle + O(a^2).$$

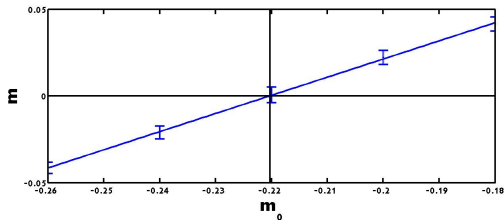
- ▶ Calculate m_R using Schrödinger Functional. For details see ⁵

⁴M. Lüscher, S. Sint, R. Sommer, and P. Weisz

⁵M. Lüscher, S. Sint, R. Sommer, P. Weisz, and U. Wolff (1996)

Critical Mass II

- ▶ At fixed g_0^2 and L/a calculate m at a variety of m_0 , around $m(m_0) = 0$, interpolate with a linear polynomial to determine



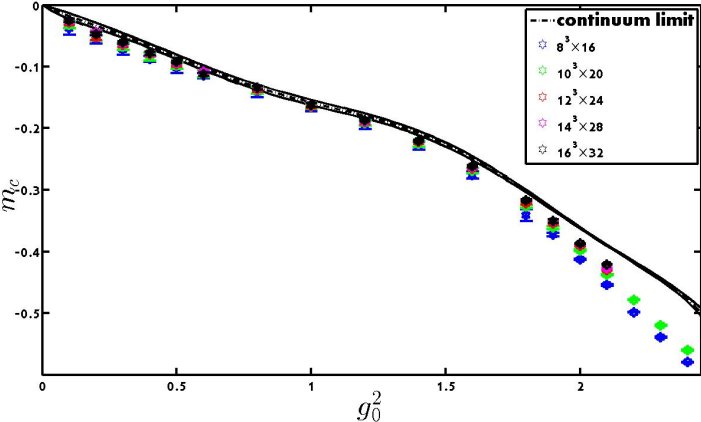
m_c :

- ▶ Do this for variety of g_0^2 and $L/a = 8, 10, 12, 14,$ and 16 . Fit all $m_c(g_0^2, a/L)$ to polynomial

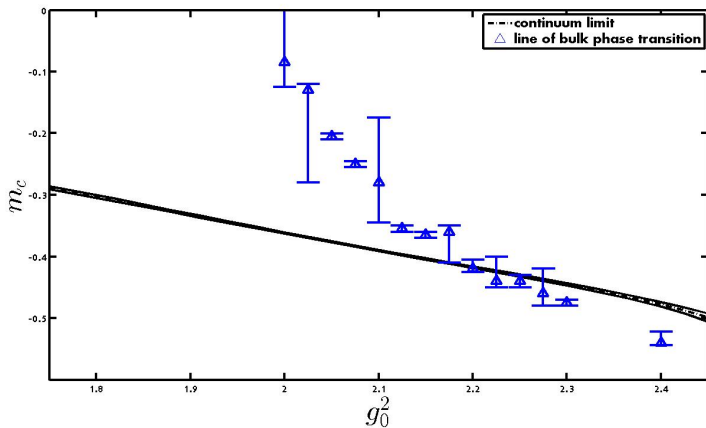
$$m_c^{\text{fit}} \left(g_0^2, \frac{a}{L} \right) = \sum_{i=1}^n \sum_{j=0}^1 a_{ij} g_0^{2i} \left(\frac{a}{L} \right)^j.$$

- ▶ Finally use $m_c^{\text{fit}}(g_0^2, 0)$, as the critical mass to use in current and future $SF \bar{g}^2$ calculations.

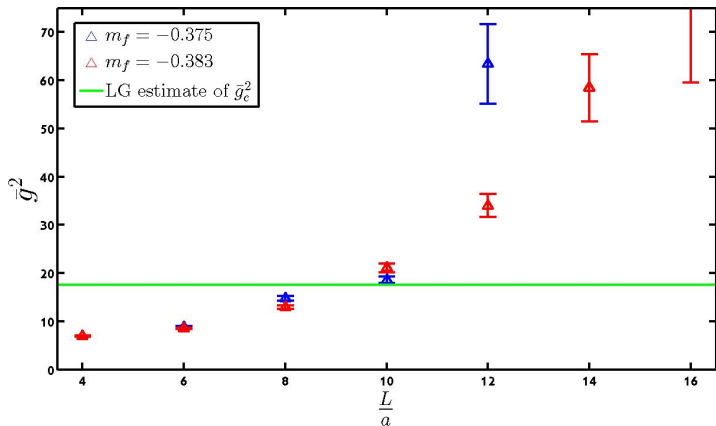
Critical Mass III



Critical Mass and Bulk Phase Transition



SF Renormalized Coupling at fixed $g_0^2 = 2.1$

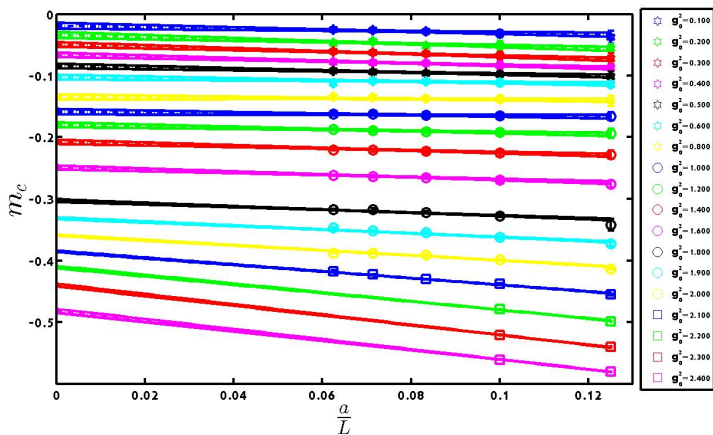


Conclusions

- ▶ Preliminary indication that $N_f = 6$ is chirally broken in the IR.
- ▶ This is not definitive. Much more running at other values of g_0^2 and a full step-scaling analysis with a continuum extrapolation is required and in progress.
- ▶ Can certainly reach a \bar{g}^2 sufficiently strong (according to ladder gap equations⁶) to break chiral symmetry on moderate lattice volumes without running past a lattice bulk phase transition.

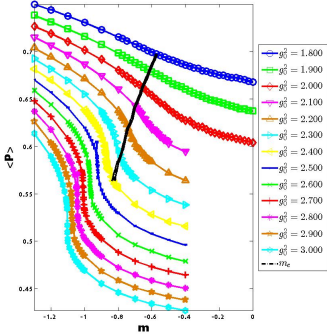
⁶K. Higashijima, A. Cohen, and H. Georgi

Backup Slide 1



Backup Slide II

▶ Not Stout Smeared



▶ Stout Smeared

