

Lattice gauge theory and physics beyond the standard model

Joel Giedt
Rensselaer Polytechnic Institute

Outline

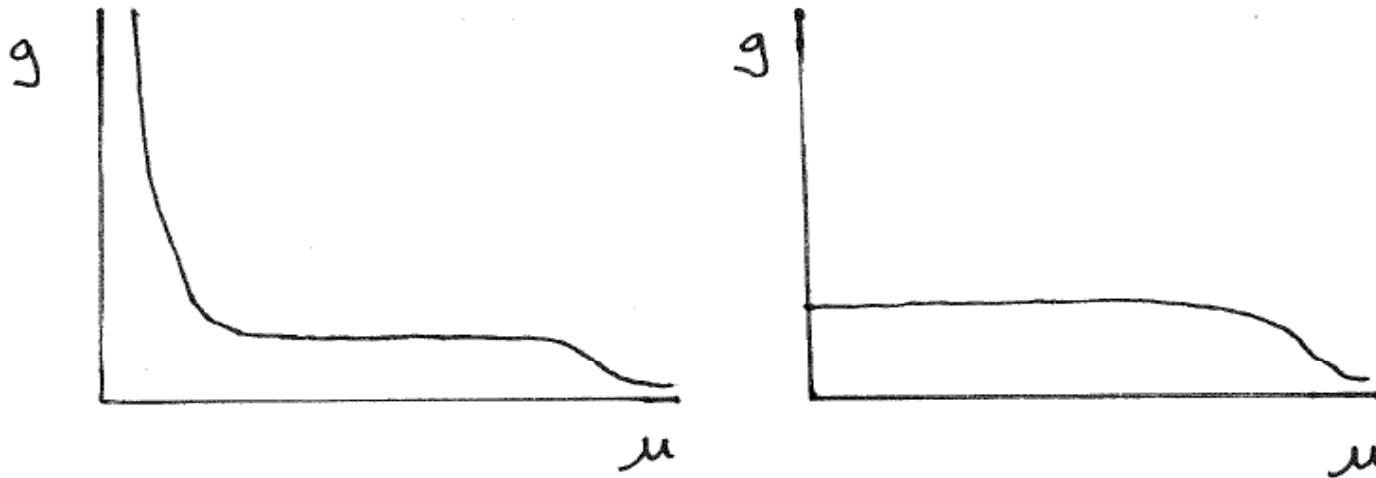
- Dynamical electroweak symmetry breaking/Technicolor
 - Conformal/confining
 - Measurements
- Not included (w/ apologies):
 - Supersymmetry
 - Condensed matter
 - Matrix models
 - Quantum gravity
 - Extra dimensions
 - Weak interaction
 - Higgs physics
 - Standard model tests
 - Generic quantum field theory

Technicolor

- Standard model Higgs suffers from hierarchy problem and triviality problem
- Replace Higgs vev with condensate of new fermions under new color group (technicolor)
- To avoid FCNCs, need to push up scale associated with exchange particle (ETC) that generates four-fermion terms from which quark/lepton masses are derived.
- Need large condensate at ETC scale to get viable quark masses: walking with $\gamma=1$

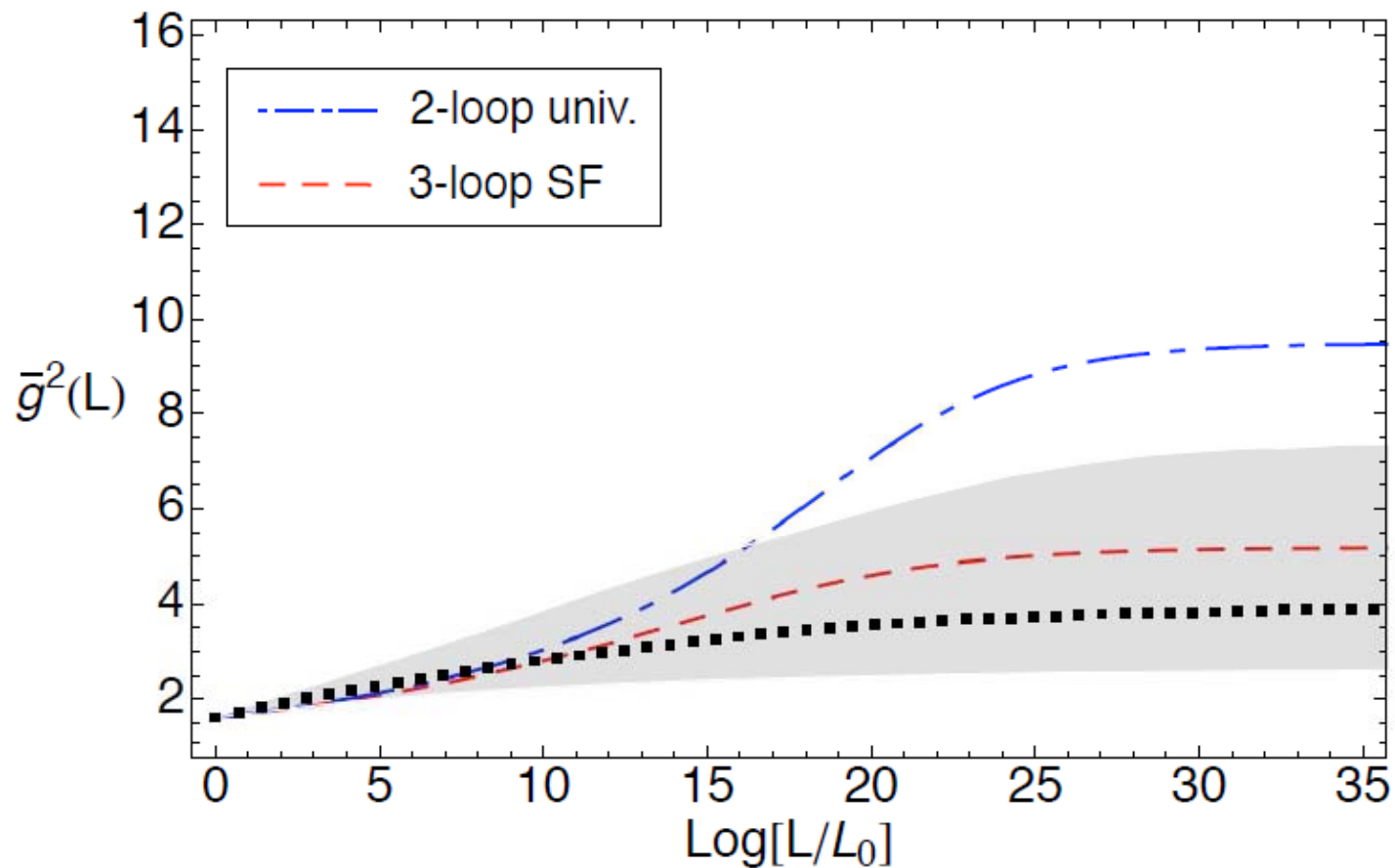
$$\langle \bar{\psi}\psi \rangle_{ETC} = \langle \bar{\psi}\psi \rangle_{TC} \left(\frac{\Lambda_{ETC}}{\Lambda_{TC}} \right)^\gamma$$

- Much effort has gone into attempting to distinguish between the two pictures below.
- It's not easy given that on a single lattice we only see a small range of μ and it's necessary to keep bare \mathbf{g} small to have small \mathbf{a} .
- In addition, we often have a nonzero mass, which drives the theory away from the fixed point at long distances (RHS \rightarrow LHS).

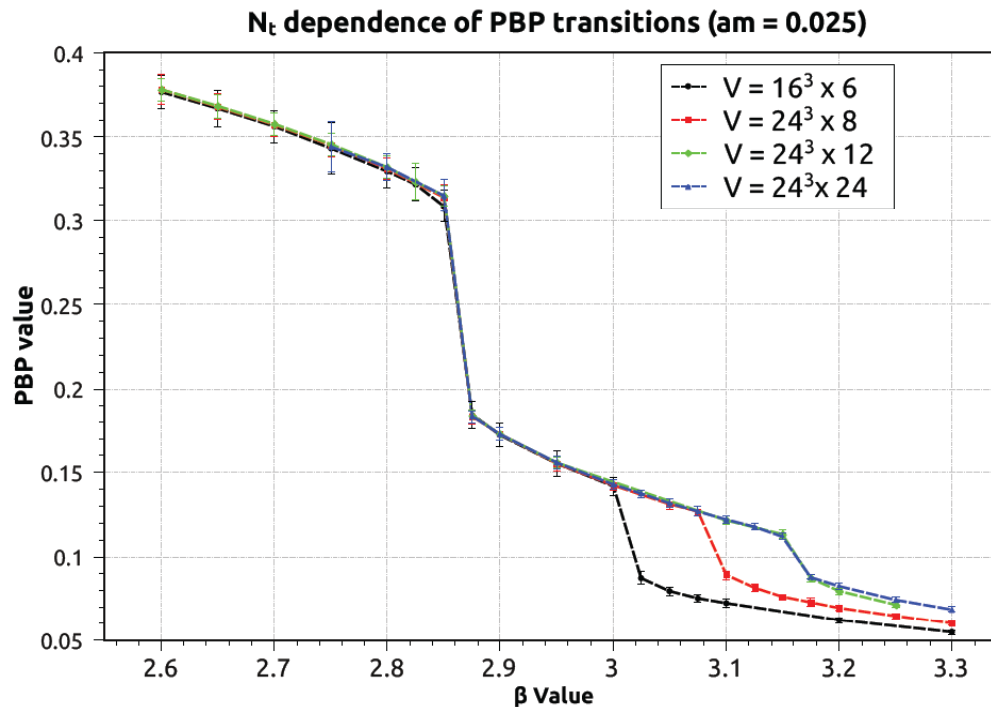


12 flavors of fundamentals in SU(3)

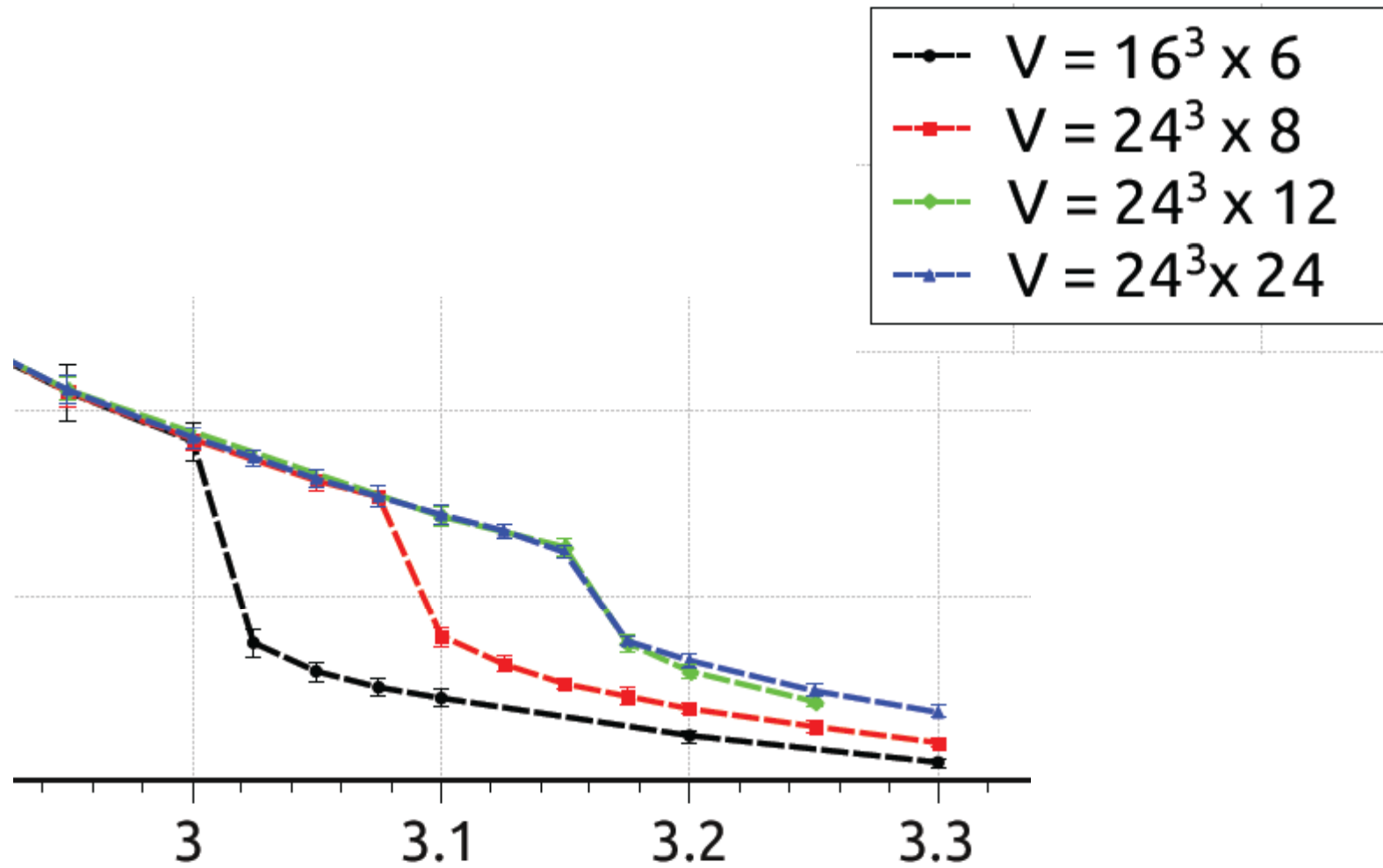
- Of course you have the Appelquist, Fleming, Neil result...



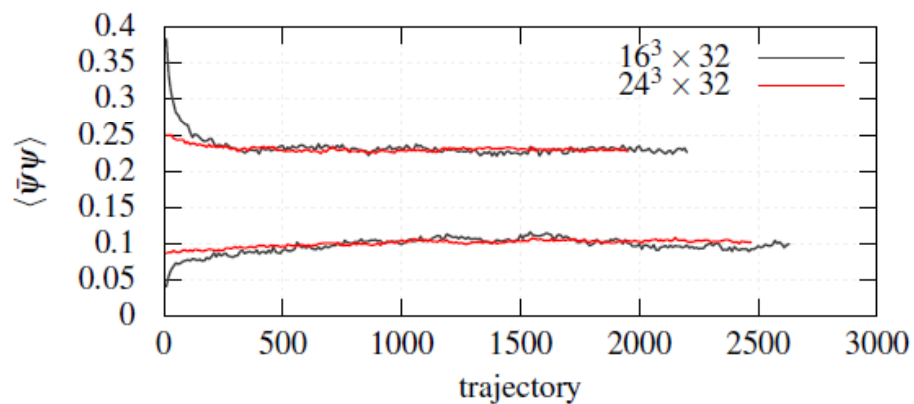
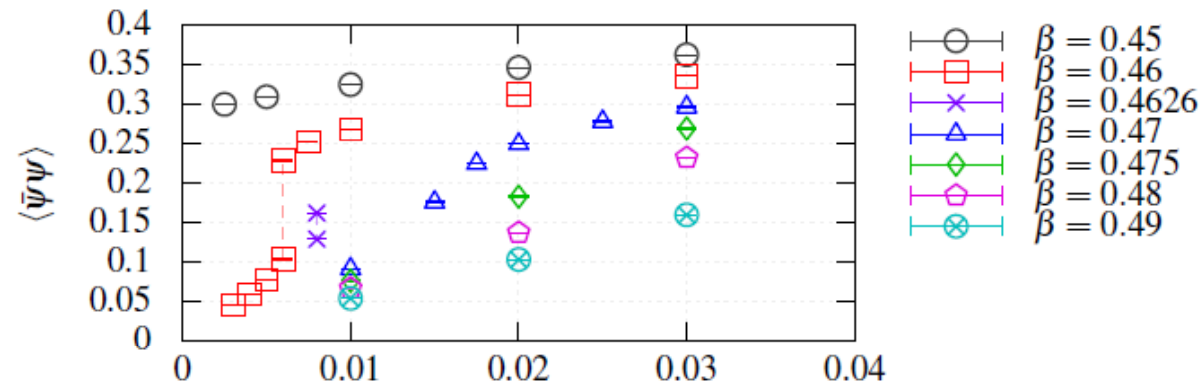
- More recently, Deuzeman, Lombardo, Nunez da Silva & Pallante observe two transitions, one of which is clearly bulk and the other is argued to be bulk:



Lattice 2011 proceedings,
1111.2590

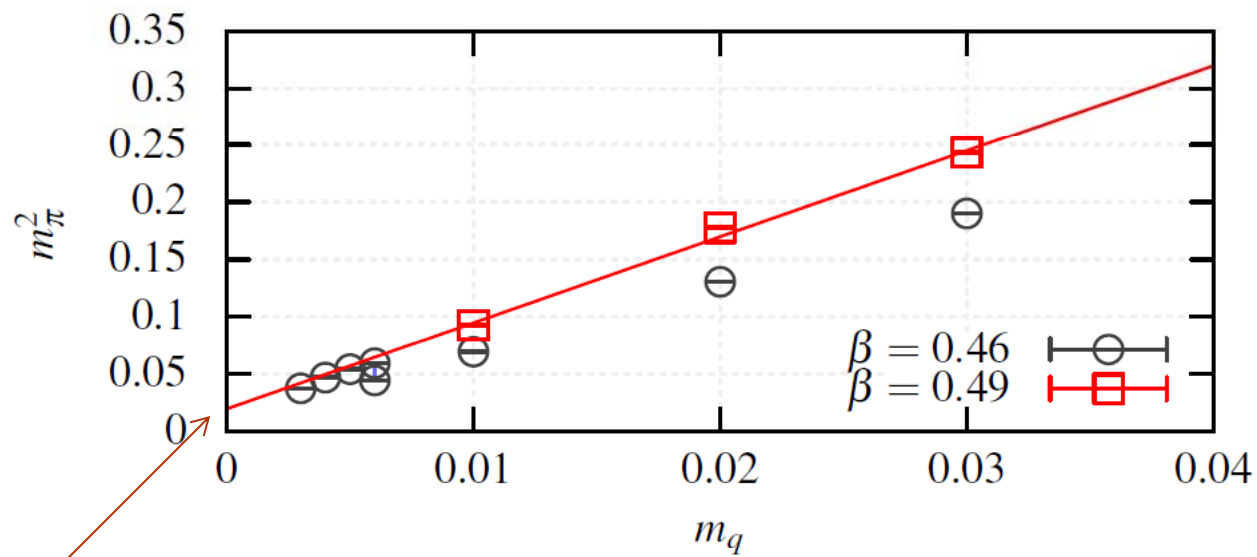


- Jin & Mawhinney observe a bulk transition with two different actions



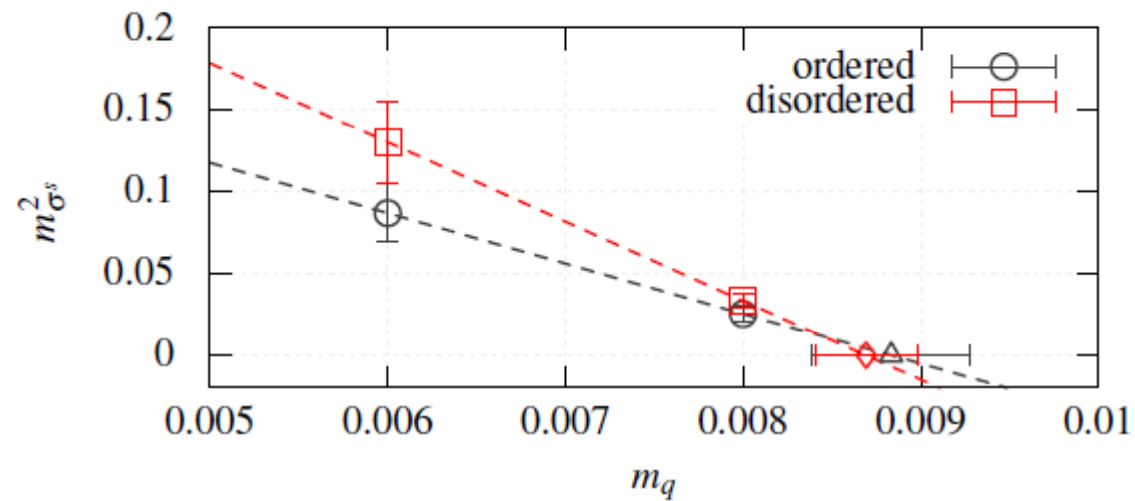
Lattice 2011 proc.,
1203.5855

- On the weak side they see behavior consistent (?) with χ SB



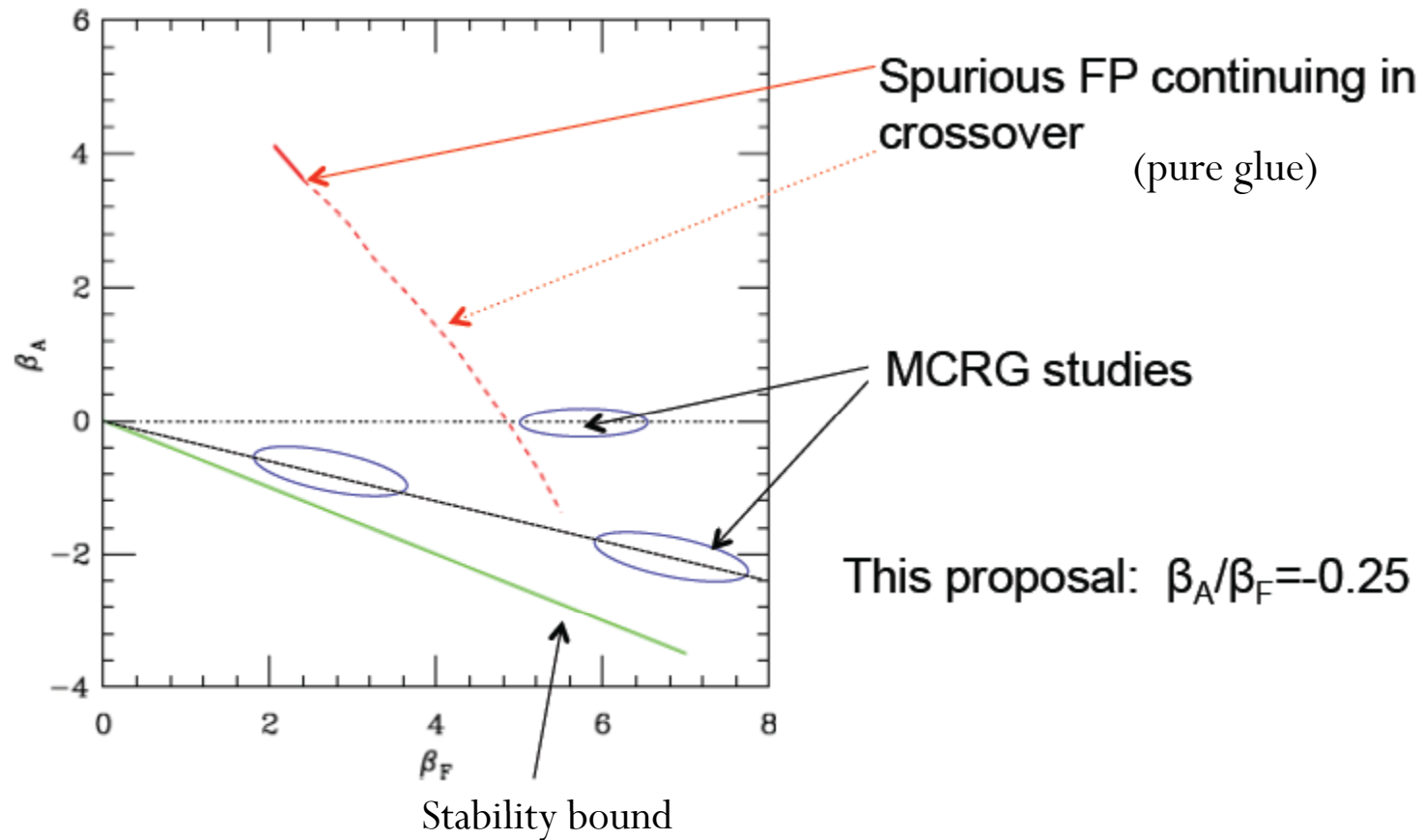
Nonzero intercept?

- They also find a massless scalar at the end of the first order line and interpret it as a UVFP



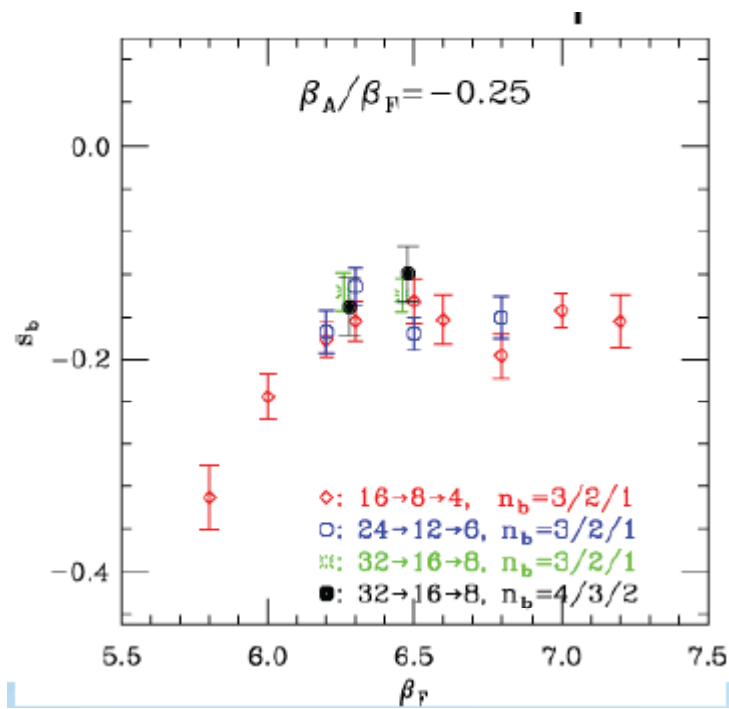
The negative adjoint explorations

A. Hasenfratz

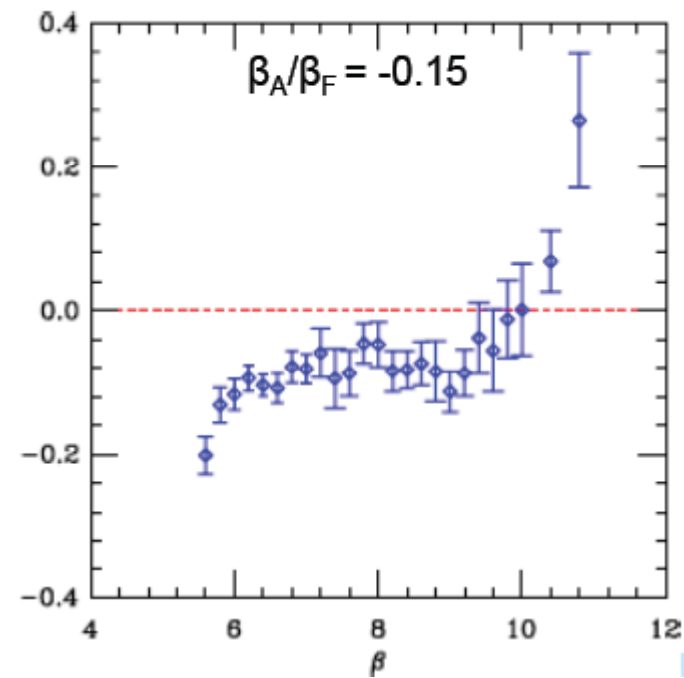


12 flavor results

MCRG (See talk Thurs. by Petropoulos)



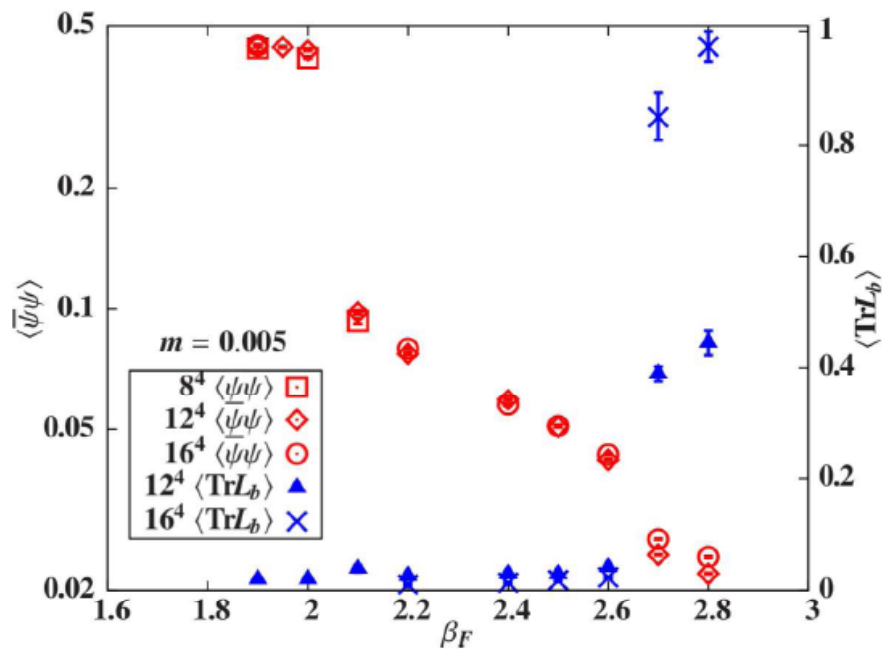
Backward flow



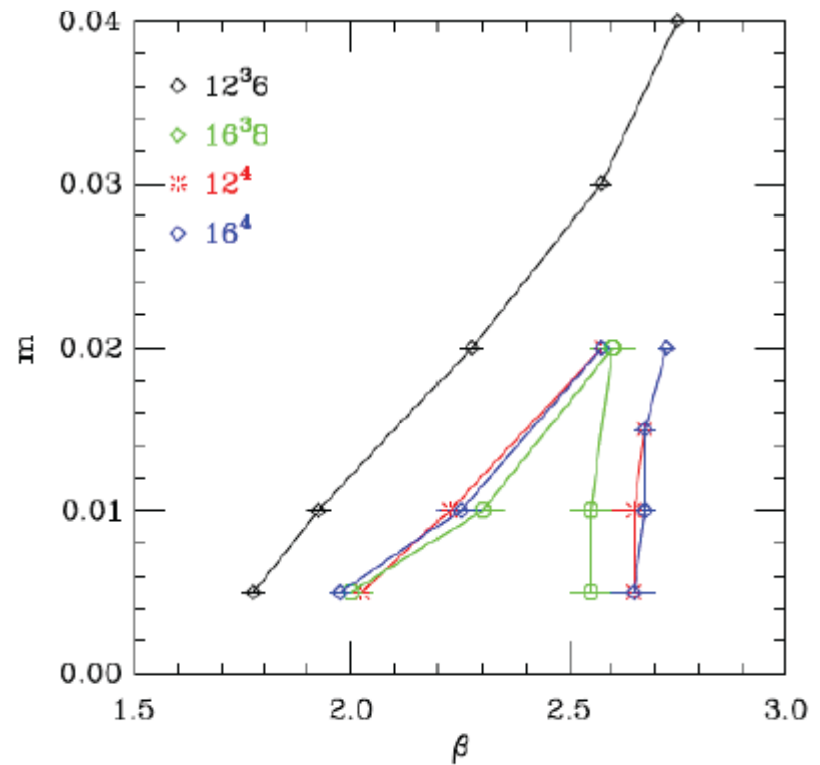
IRFP

Cheng, Hasenfratz & Schaich

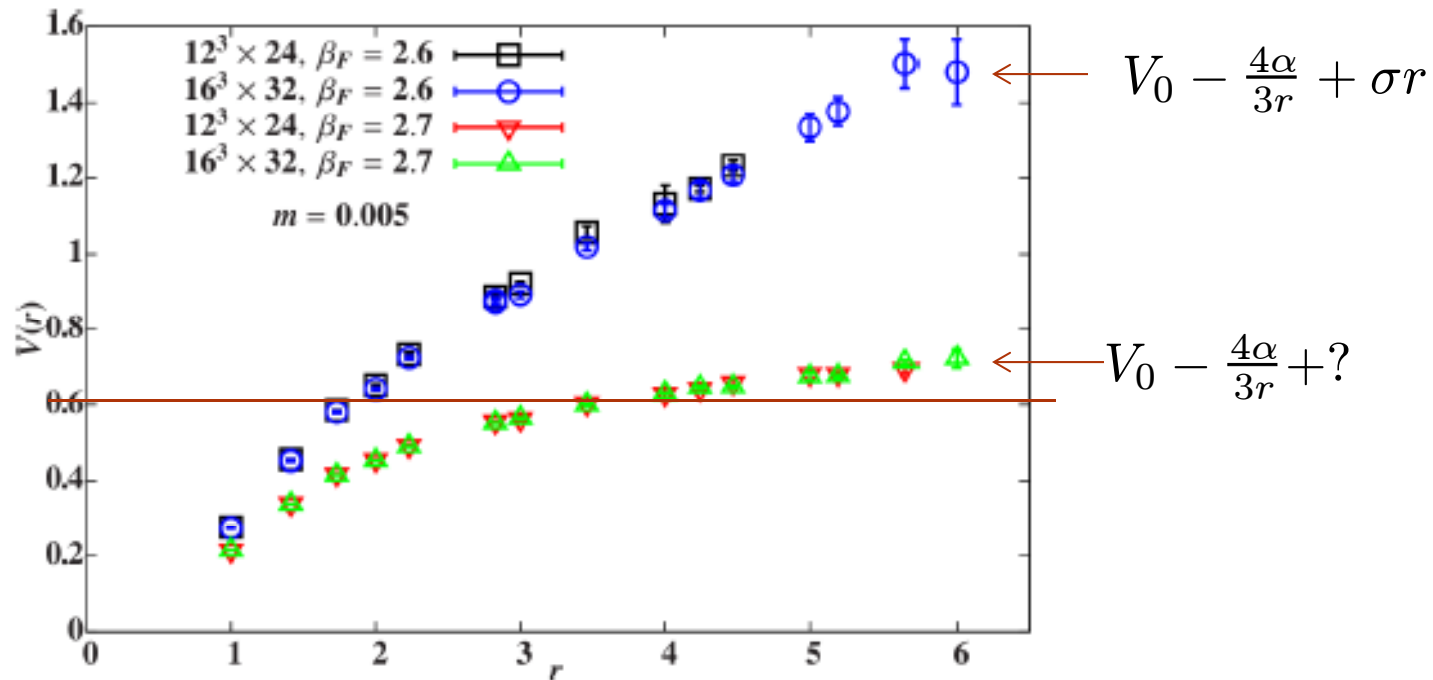
See talk by Schaich from Mon.



$ma = 0.005$

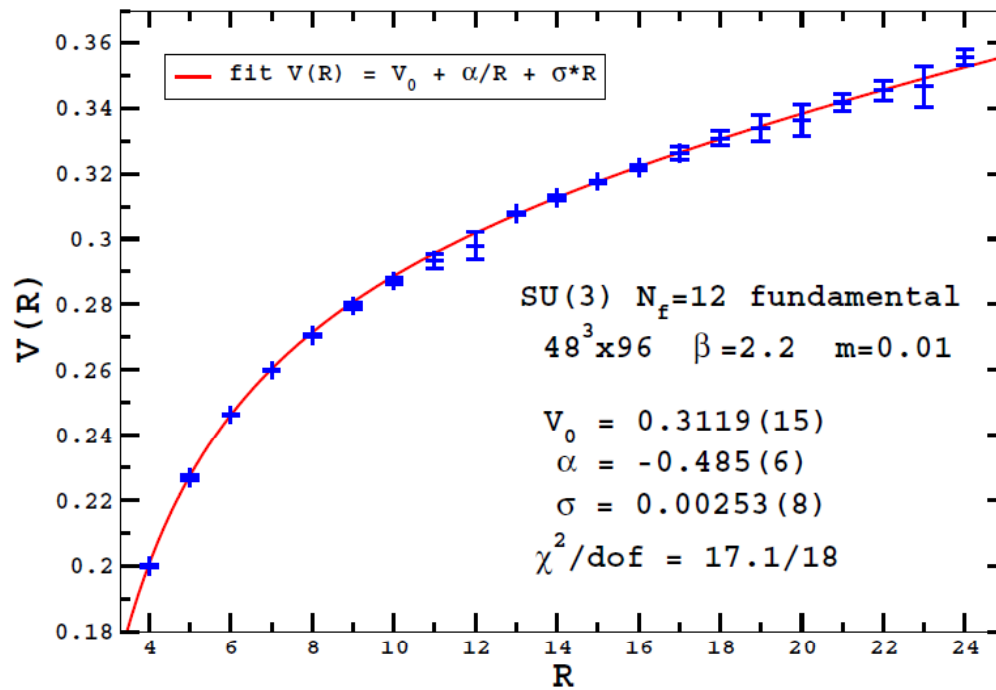


- One thing that differs significantly between the two phases is the static potential...

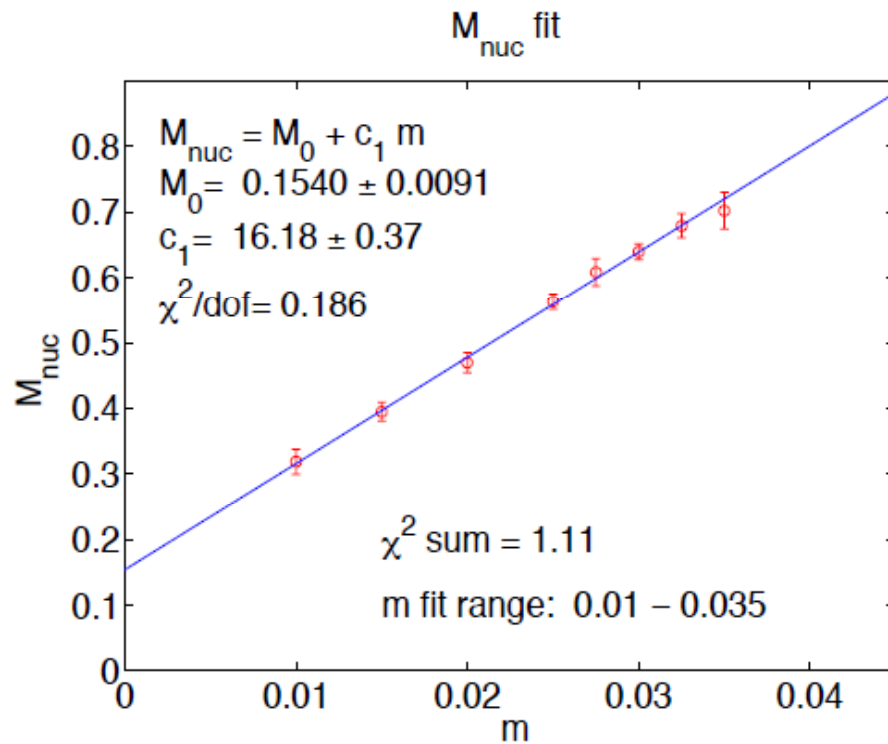


- This is to be compared with the results of LHC (Fodor, Holland, Kuti, Nogradi & Schroeder 2010) :

See talk by Holland from Mon.

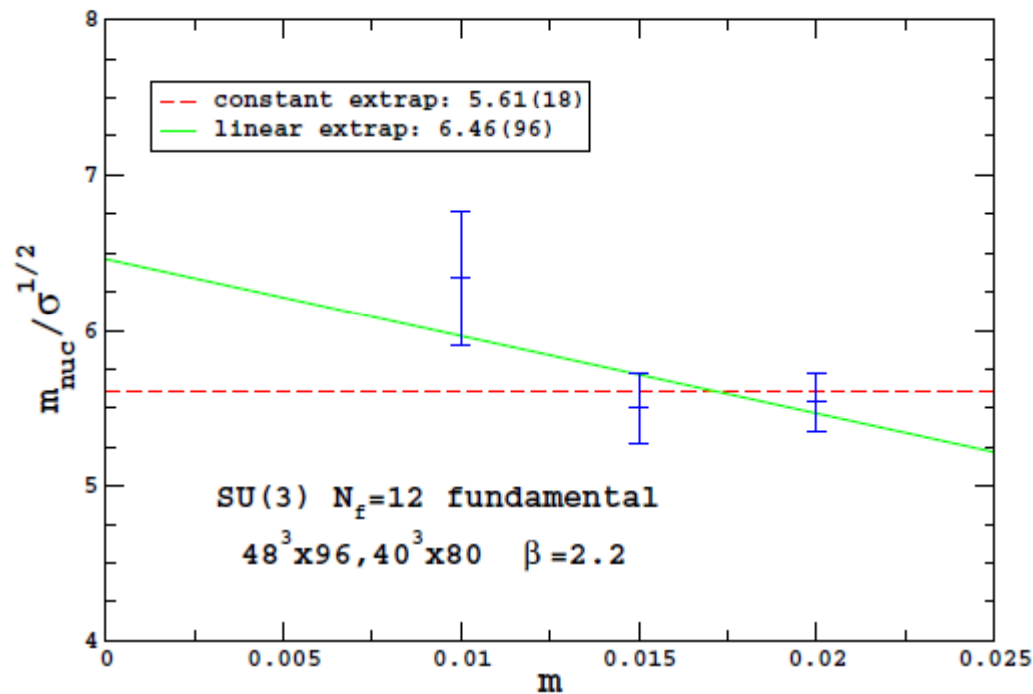


- LHC find a nonzero nucleon gap in the chiral limit:



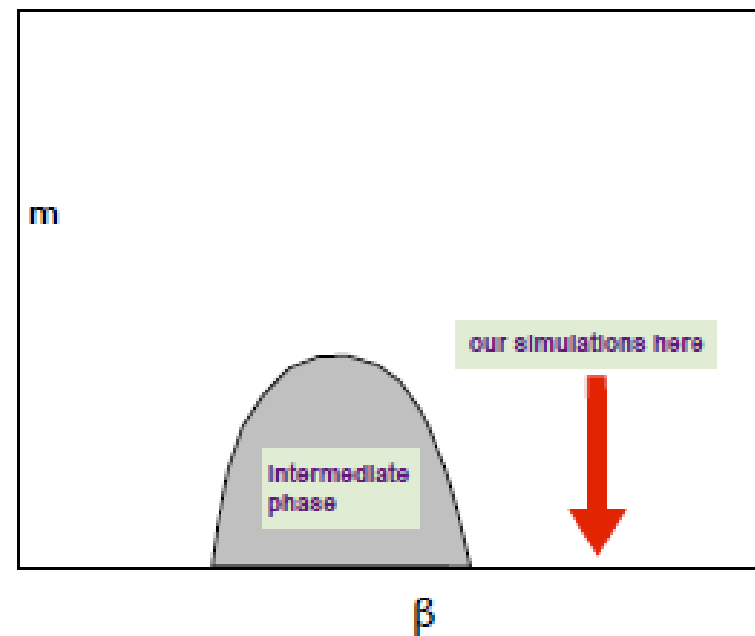
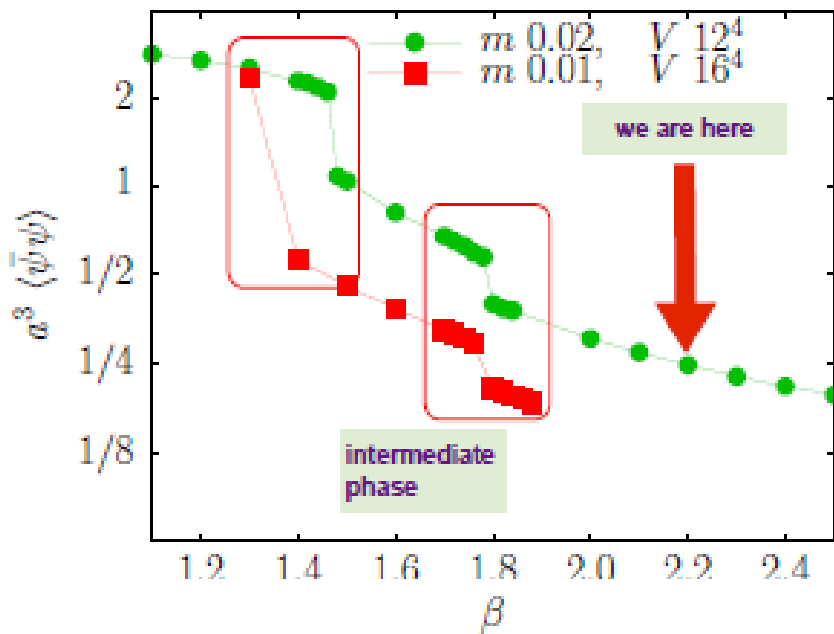
$$\frac{M_0}{\sigma_{M_0}} = 17$$

- When they measure the string tension in these units, extrapolating to the chiral limit, they have a nonzero result:

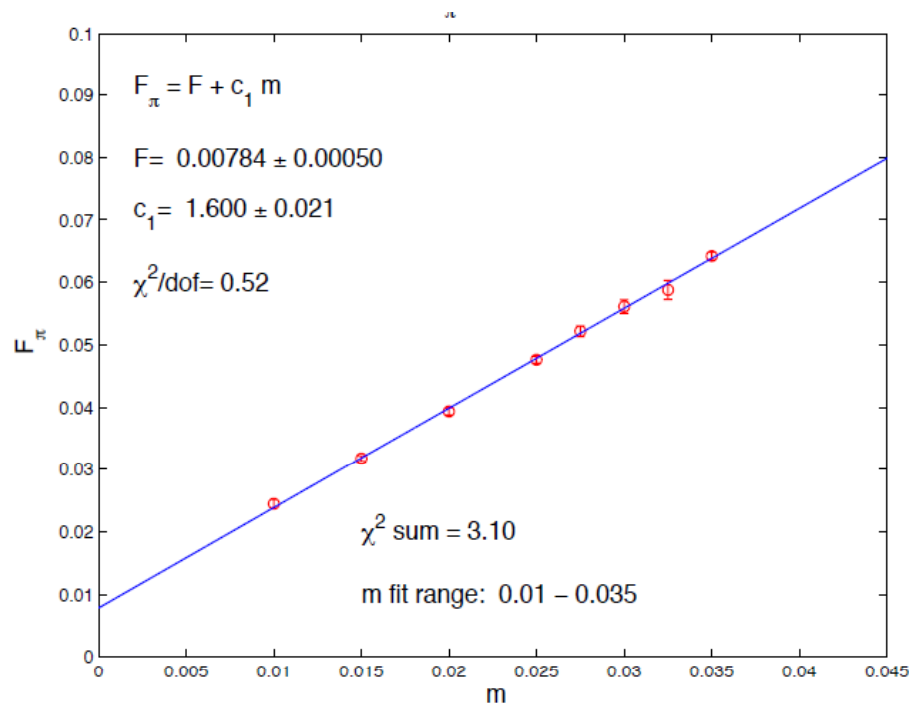


This ought to address the issue that the massive theory is always confining at long enough distance scale.

- LHC has identified the two phase transitions relative to their principal simulation point:

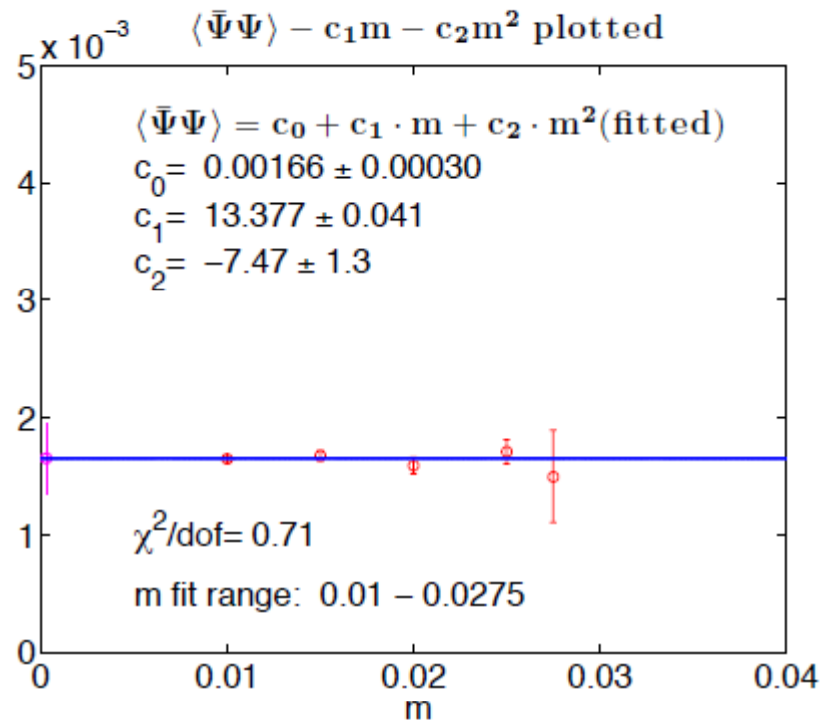


- F extrapolates to a nonzero value, indicating spontaneous chiral symmetry breaking.



L=32,40,48

- The chiral condensate comes out nonzero



- Similar results are found for the nucleon, f_0 , rho and the a_1 .
- It looks like the 12-flavor theory has a hadronic scale given by $Ma=0.1$ for this lattice spacing.
- The existence of a dynamical scale is forbidden by the IRCFT hypothesis.

Conformal fits

- The Lattice 2011 proceedings generalizes the function and includes finite volume (FSS):

$$ML = f(x), \quad x = m^{1/y_m} L$$

$$f(x) = c_1 x + c_{\text{exp}} (c_1 x)^{-1/2} \exp(-c_1 x), \quad x > x_{\text{cut}}$$

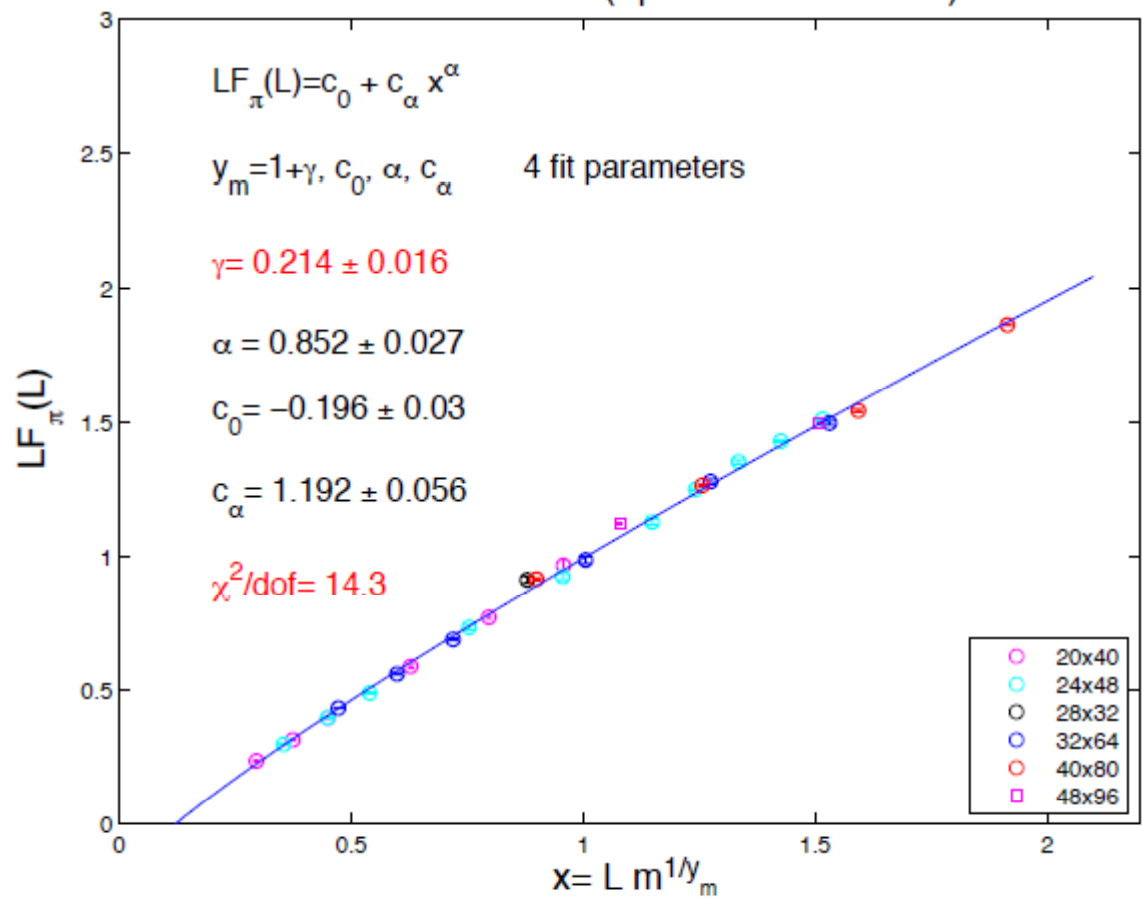
$$f(x) = c_0 + c_\alpha x^\alpha, \quad x < x_{\text{cut}}$$

with qualifications.

- See talk by Wong from Mon.

channel	γ	χ^2/dof
pion	0.393(8)	2.83
F	0.214(16)	14.3
rho	0.300(17)	1.51
nucleon	0.288(27)	1.45

conformal FSS (Fpi - PCAC channel)



- The γ values don't agree within the errors derived from the fit.
- But how large is the systematic error associated with choice of fitting function?
- For example, could I bring the 0.39 value from the pion down to the 0.30 value from the rho by choosing a different function?

- DeGrand approached this problem by extracting γ from the FHKNS results using an approach that doesn't assume a specific form for the finite size scaling function $f(x)$.

$$P(y_m) = \frac{1}{N_{over}} \sum_p \sum_{j \neq p} \sum_{i, over} \left(\frac{\xi_L(m_{i,j})}{L_j} - f_p(L_j^{y_m} m_{i,j}) \right)^2$$

TABLE I: Exponent y_m from various hadronic channels. Errors are from a single-elimination jackknife.

channel	y_m
pseudoscalar	1.35(23)
nucleon	1.43(26)
f_π	1.23(31)
vector	1.33(22)
axial vector	1.32(12)

- The LHC collaboration has responded with spline based general B-form fits

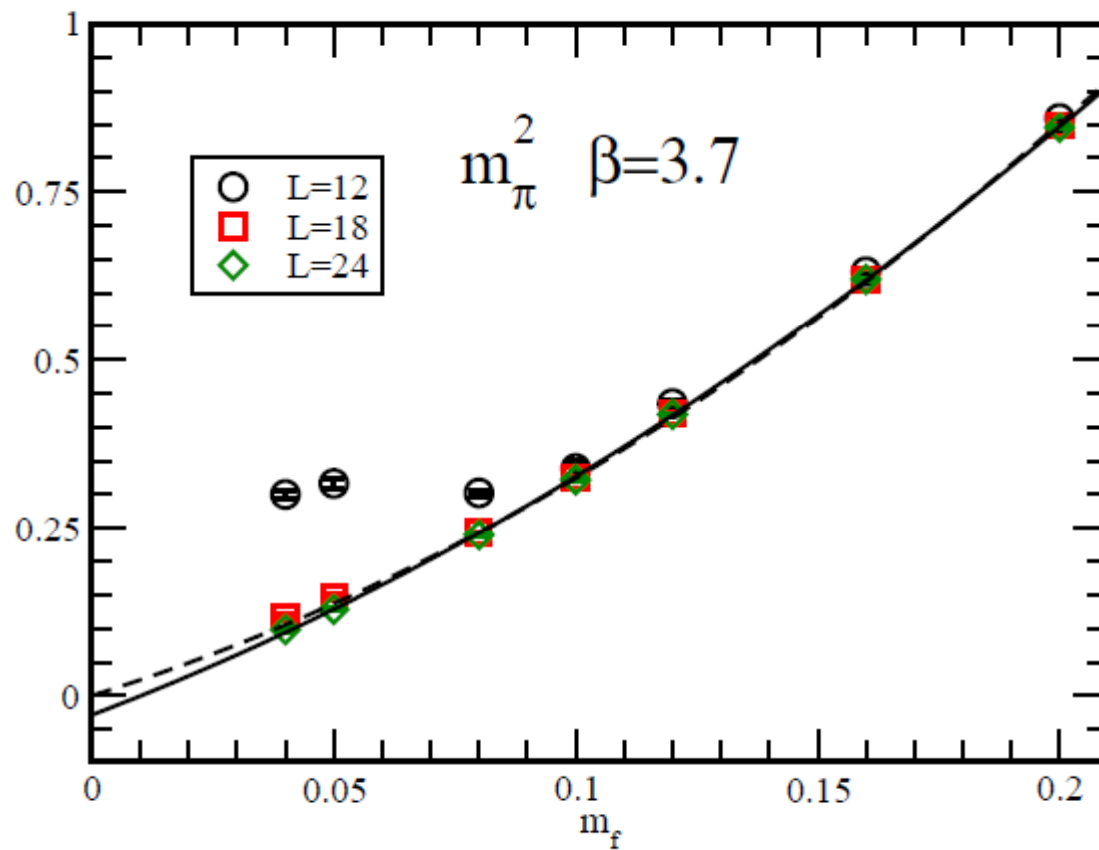
channel	γ	χ^2/dof
pion	0.405(21)	1.47
rho	0.315(75)	1.02
F	0.23(2)	8.05

- Cheng, Hasenfratz and Schaich have been looking at the mode number of the Dirac operator (following A. Patella) to determine the anomalous mass dimension.

$$\nu(\lambda) - \nu(\lambda_0) = cV(\lambda^{4/y_m} - \lambda_0^{4/y_m})$$

- They have added many more eigenvalues (x10) and now perform fits on separate volumes (Talk: Hasenfratz, Tues. afternoon).
- They find that it is necessary to go to fairly large β (well past the bulk phase transition) to see y_m volume independence.

- The LatKMI collaboration (Aoki et al.) have studied spectra.



$$m_\pi^2 = c_0 + c_1 m_f + c_2 m_f^2$$

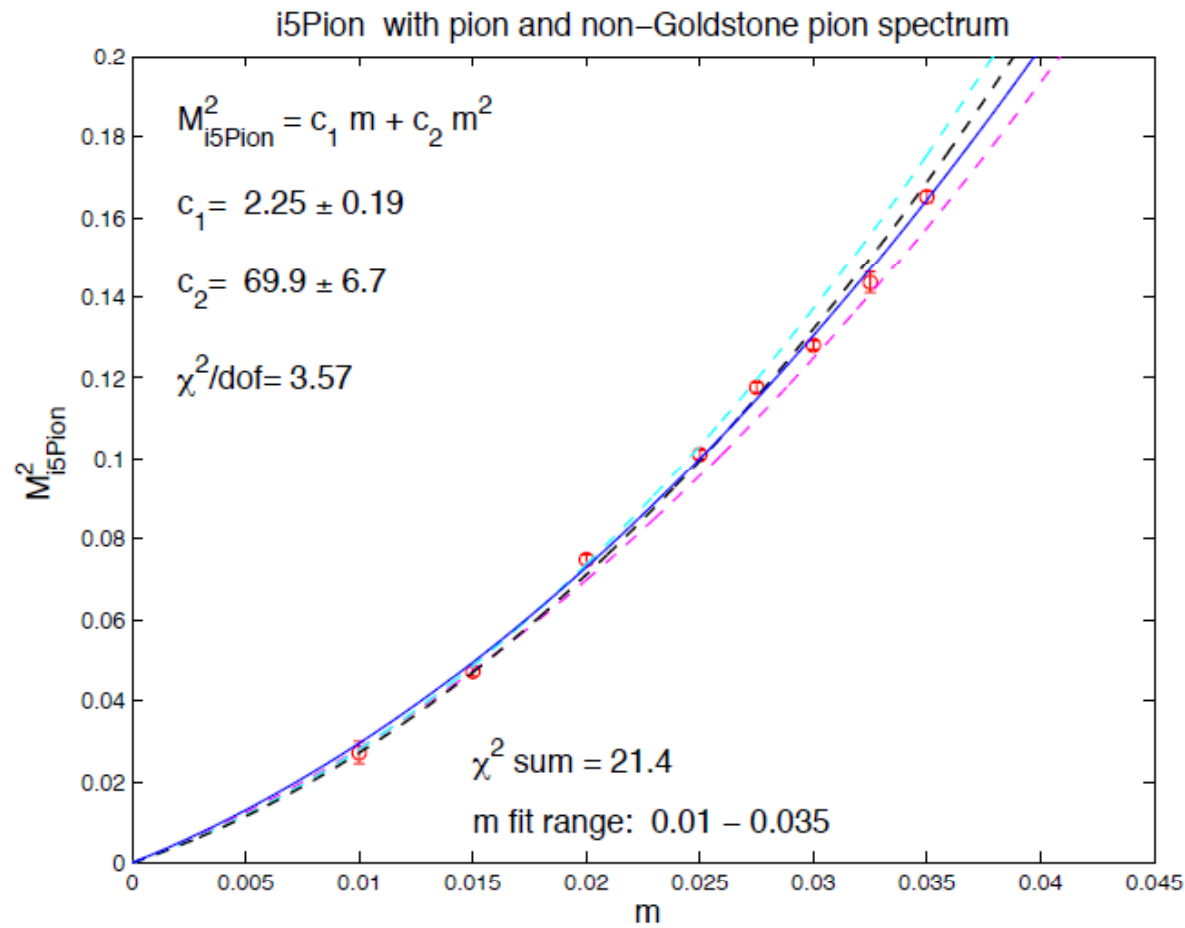
$$\chi^2/\text{dof} = 4.1$$

$$c_0 = 0$$

$$\chi^2/\text{dof} = 28$$

Impossible to reconcile
with LHC (Fodor et al.)

LHC Pions



- They also do a FSS fit with $f(x) = c_0 + c_1 x$
- The constant term is a guess.

$$\gamma = 0.44 \pm ?, \quad \chi^2/\text{dof} = 4$$

- All results “preliminary” (Lattice 2011 proc.)
- Follow-up work and further results were presented Mon. (Ohki)

- Appelquist, Fleming, Lin, Neil & Schaich [1106.2148] have considered mass corrected hyperscaling:

$$M_X = C_X m^{1/(1+\gamma)} + D_X m$$

- It is sensible that corrections should be analytic in m .
- Similar eq. for F .
- Condensate: $\langle \bar{\psi}\psi \rangle = A_C m + B_C m^{(3-\gamma)/(1+\gamma)}$
- SD inspired generalization:

$$\langle \bar{\psi}\psi \rangle = A_C m + B_C m^{(3-\gamma)/(1+\gamma)} + C_C m^{3/(1+\gamma)} + D_C m^3$$

- Without the D-terms,

$$\chi^2/\text{dof} = 133/53$$

- With the D-terms and finite volume corrections,

$$\chi^2/\text{dof} = 42/44$$

Sextet

- In the Lattice 2010 proceedings, LHC (FHKNS) report that a collective fit to the pion, F and the chiral condensate produces:
 - $\chi^2 / \text{dof} = 1.24$ for χ_{SB}
 - $\chi^2 / \text{dof} = 6.96$ for IRCFT

- Lattice 2011 proc. LHC results for chiral fit:

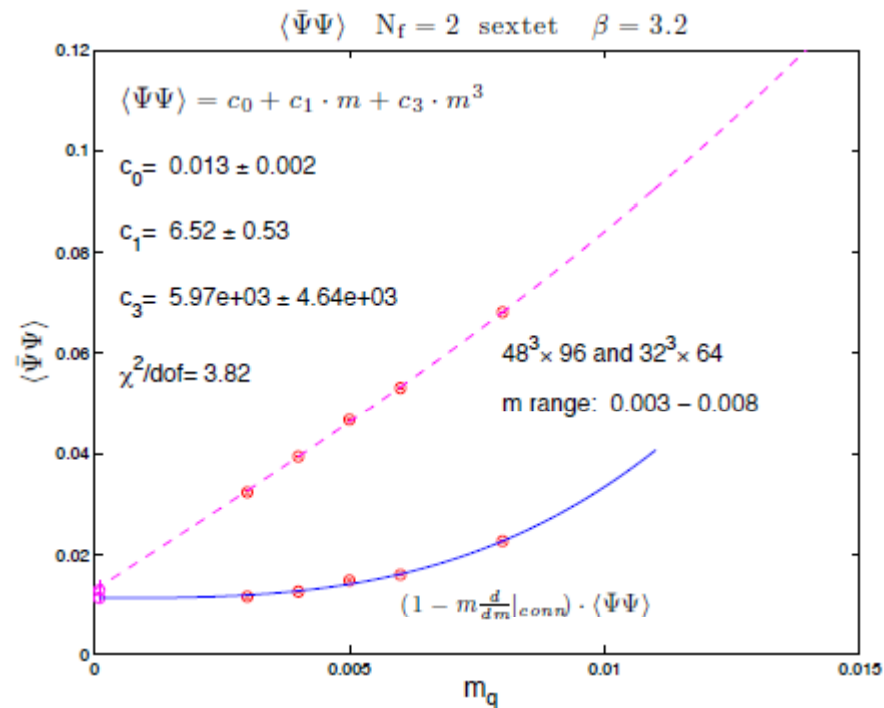
channel	χ^2/dof
pion	1.6
F	0.87
rho	0.56
f0 (connected)	0.48

- Lattice 2011 proc. LHC results for conformal fit

channel	γ	χ^2/dof
pion	1.091(34)	2.0
F	2.13(18)	2.0

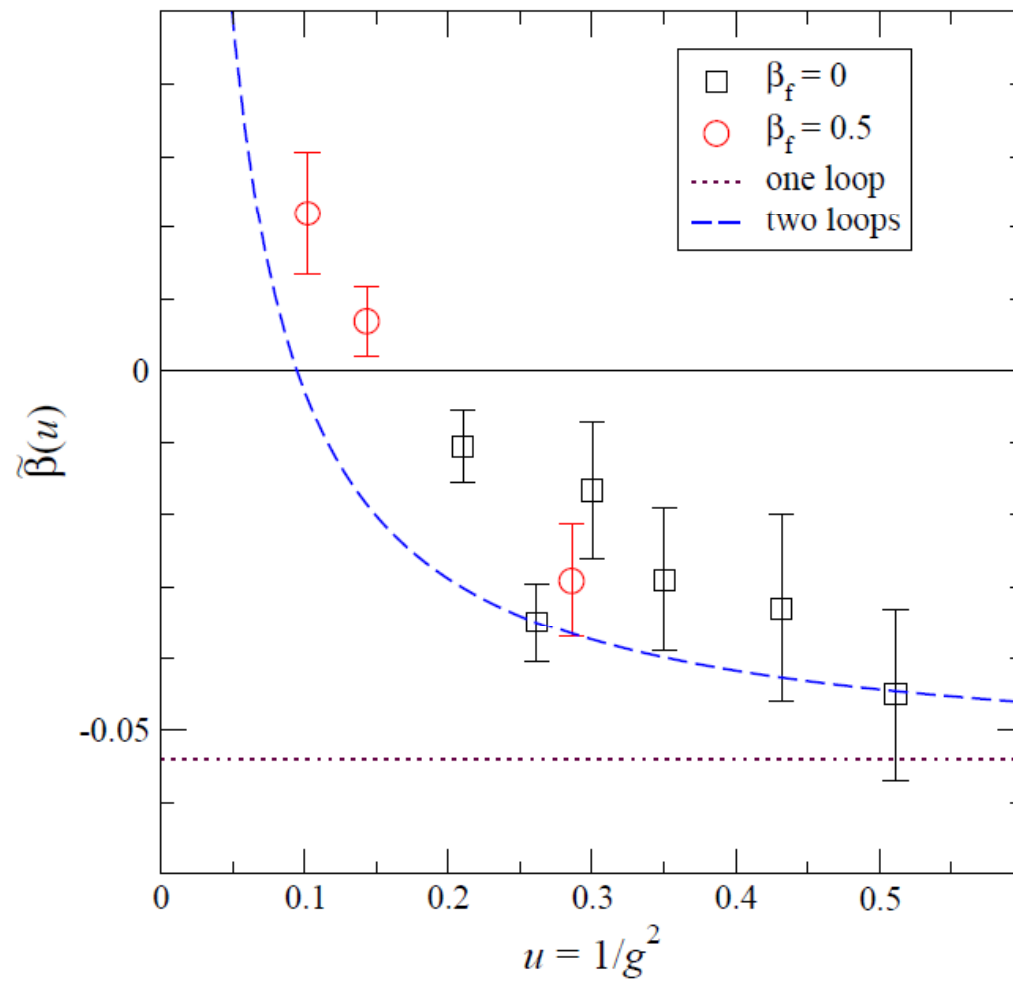
$$M, F \sim m^{1/(1+\gamma)}$$

- And the chiral condensate:

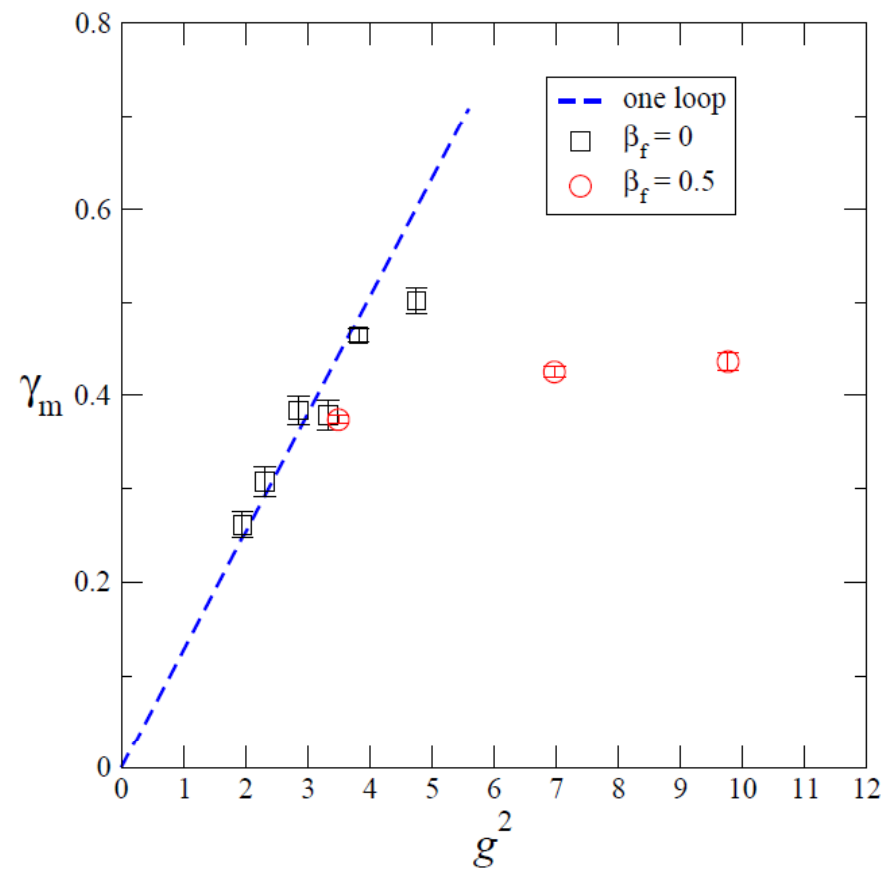


For most recent results, see Kuti talk from Mon.

- DeGrand, Shamir and Svetitsky have added a fat link sextet repr. gauge action term which allows them to push into stronger coupling w/o hitting the bulk transition.
- They compute a discrete beta function from differences between two lattices.

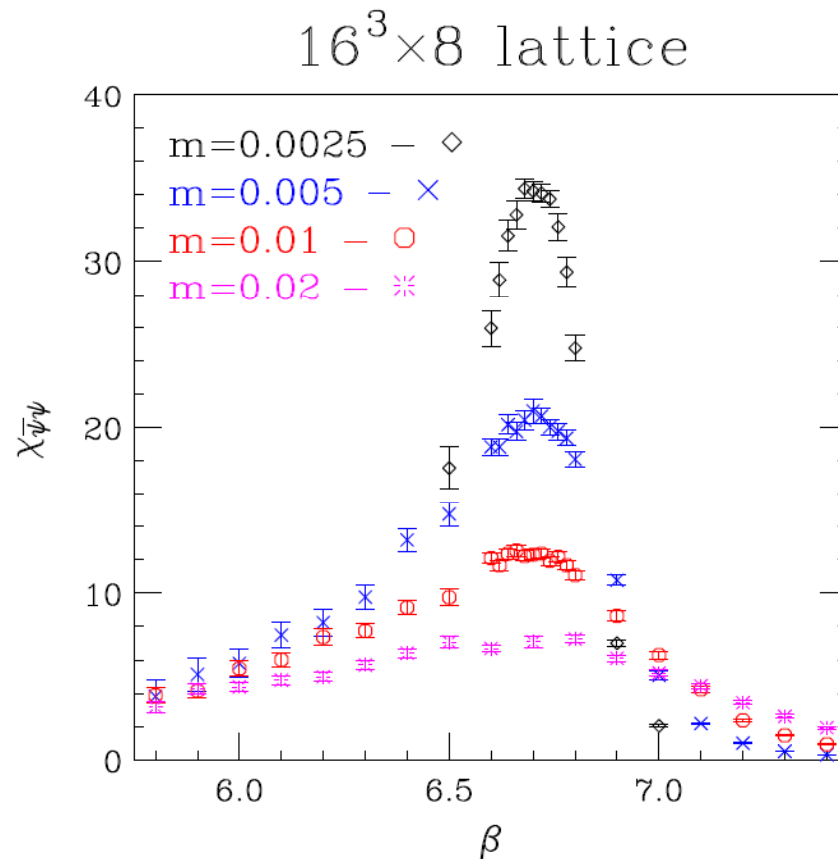


- They also measure γ



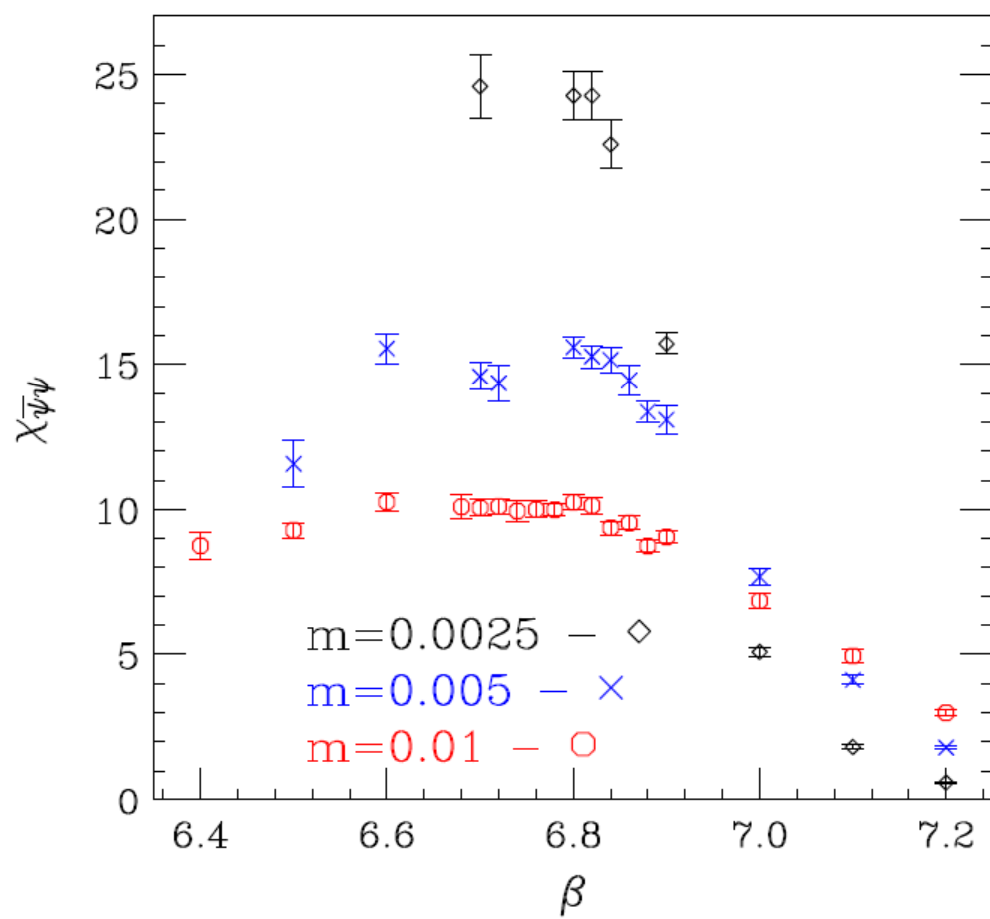
- Similar results by same group for $SU(4)$.

- Kogut and Sinclair have been studying the finite temperature transition in this theory.



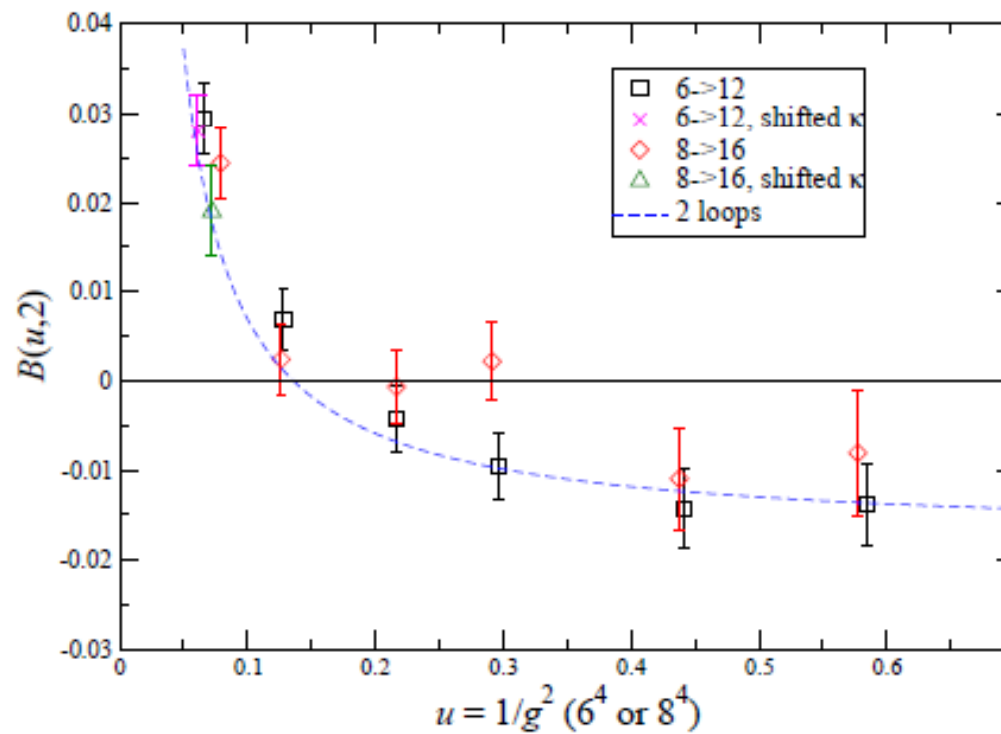
- The scaling of the peak with mass is consistent with a second order finite temperature phase transition and inconsistent with a first order bulk transition.
- Results are too preliminary (smaller mass may be needed) for an accurate estimate of the peak location for $N_t = 12$ (however, see Sinclair's talk from Mon).
- Ultimately $N_t = 18$ will also be needed.

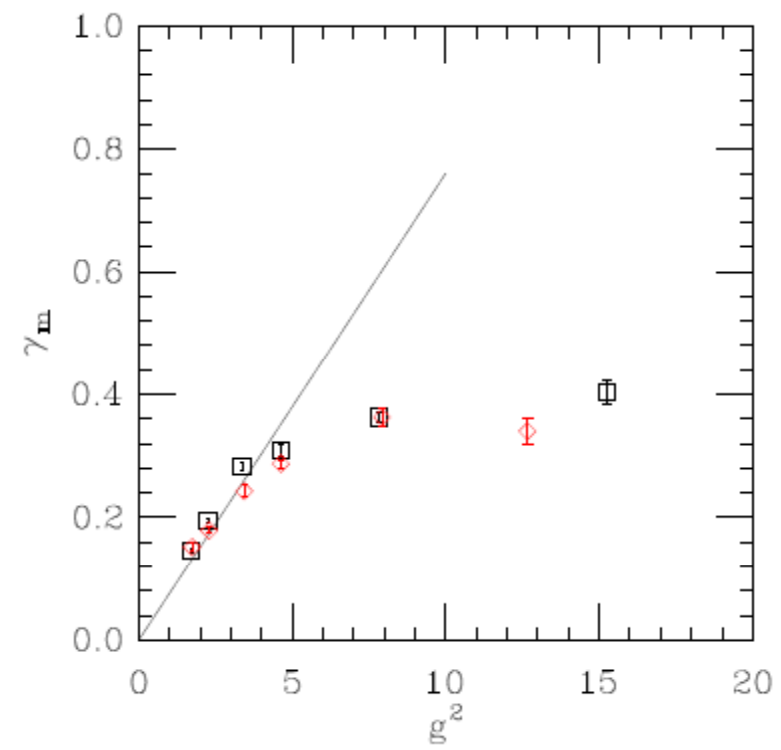
$24^3 \times 12$ lattice



SU(2) $N_f=2$ Adjoint

- DeGrand, Shamir, Svetitsky results for beta function and γ





- JG & Weinberg estimate from FSS
- Using approach advocated by DeGrand

Channels and fits

Observable	Quadratic	Log Quad	PWL	Combined
m_π	1.67(93)	1.26(54)	1.51(33)	1.46(27)
m_ρ	1.67(88)	1.37(39)	1.56(31)	1.50(23)
m_{a_1}	1.40(52)	1.42(27)	1.41(22)	1.41(16)
f_π	1.65(22)	1.49(54)	1.60(29)	1.62(17)

$$\gamma = 0.51 \pm 0.16$$

- Patella introduced Dirac mode number approach

$$\nu(\lambda) - \nu(\lambda_0) = cV(\lambda^{4/y_m} - \lambda_0^{4/y_m})$$

- He obtains

$$\gamma = 0.371(20)$$

- (Orthogonal: large-N reduction, Keegan talk Wed., along the lines of Hietanen & Narayanan, Koren, Okawa talks Wed.)

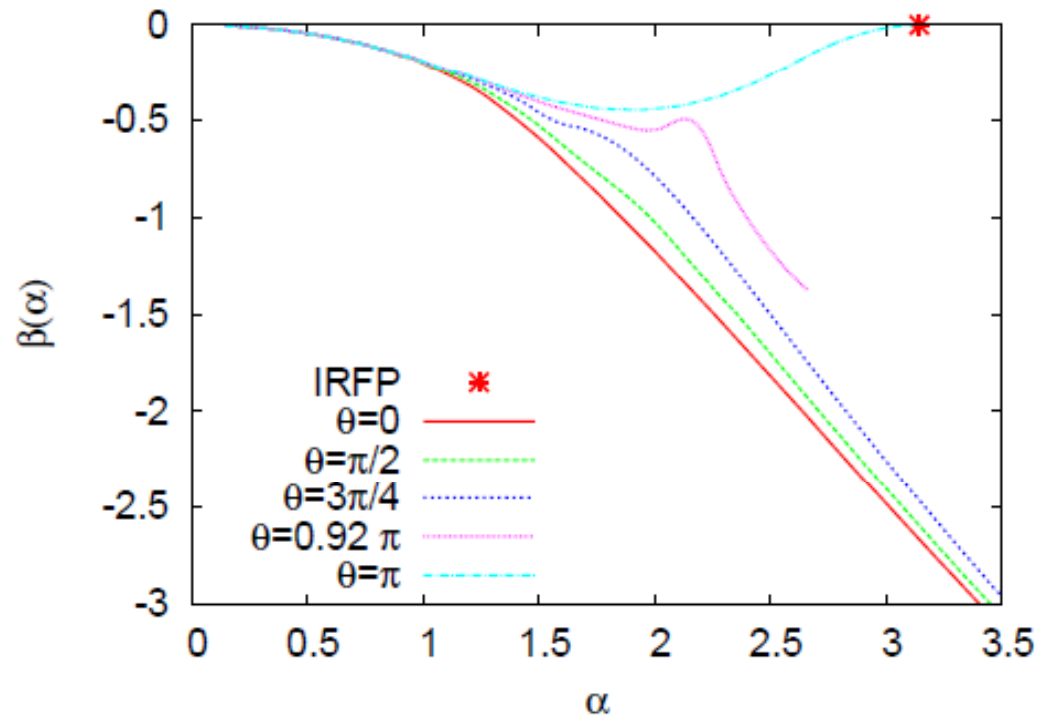
Comparison

Method	γ
SF [Bursa et al. 2009]	$0.05 < \gamma < 0.56$
SF [Degrand et al. 2011]	0.31 ± 0.06
Perturbative 4-loop [Pica & Sannino 2010]	0.500
Schwinger-Dyson [Ryttov & Shrock 2010]	0.653
All-orders hypothesis [Pica & Sannino 2010]	0.46
MCRG [Catterall et al. 2011]	$ \gamma < 0.6$
FSS [Del Debbio et al. 2010]	$0.05 < \gamma < 0.20$
FSS [Del Debbio et al. 2010]	0.22 ± 0.06
FSS [JG & Weinberg 2012]	0.51 ± 0.16
Mode number [Patella 2012]	0.371 ± 0.020

- de Forcrand, Pepe and Wiese [1204.4913] have looked at the 2d $O(3)$ spin model with vacuum angle $\theta \approx \pi$
- This theory was suggested as a proxy by Nogradi [1202.4616]
- This allows them to go arbitrarily close to a CFT, and so have walking behavior (Pepe talk Wed).
- This is an asymptotically free theory, so it is an IRCFT when $\theta = \pi$

$$\alpha(L) = \alpha_* - \frac{1}{C \ln(L/L_0)}$$

- They are able to compute the beta function with high precision:



SU(2) with fundamental flavors

- Finish group [Karavirta (talk Mon.), Rantaharju (talk Tues.), Rummukainen, Tuominen] have recently studied $N_f = 4, 6$ and 10 .
- Schrödinger functional with clover fermions.
- $N_f = 4$ appears confining and similar to QCD.
- $N_f = 10$ appears conformal, with Banks-Zaks FP.
- $N_f = 6$ is inconclusive probably because it is right near the bottom of the conformal window.
- Previous work (2010) by Bursa, Del Debbio, Keegan, Pica, Pickup found zero of beta function for constant extrapolation and inconclusive for linear extrapolation.
- Voronov, Hayakawa Wed.?

SU(3) with $N_f=6$ fundamentals

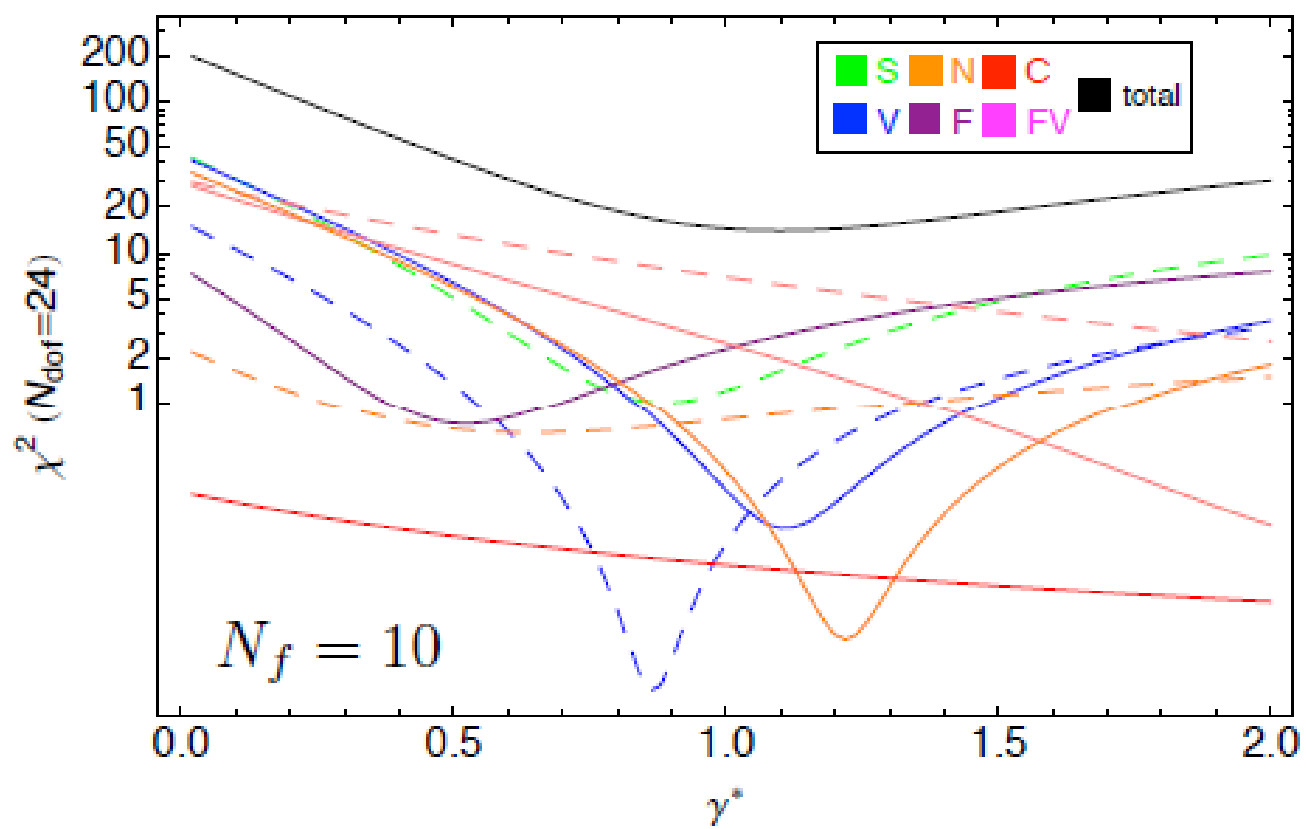
- Recent work of LSD collaboration calculating scattering length.
- Decreased compared to $N_f=2$.

SU(3) with $N_f=10$ fundamentals

- Recent work by LSD collaboration using DWF (talk by Fleming Tues afternoon).
- Based on the behavior of m_π/f_π with m_q they don't expect to be able to fit to χ PT.
- Hard to reconcile with good chiral fits for $N_f=12$.
- Hyperscaling + fits similar to 12 flavor fits give

$$\chi^2/\text{dof} \sim 1, \quad \gamma \approx 1$$

- More data on more volumes needs to be added: is there an F problem with FSS?



Four fermion couplings

- It is of interest to see whether these can push a theory out of the conformal window.
- The four fermion coupling could provide a tunable way to be arbitrarily close to that window: as much walking as desired.
- Catterall (Mon talk) has looked at gauged NJL model on lattice.
- No evidence of second order critical line.
- Further work, on models inside the conformal window, needs to be performed.

Conclusions

- Two theories have been studied in depth:
 - SU(3) 12-flavor fundamental
 - SU(2) two-flavor adjoint
- In the first case the controversy has not been resolved and there are claims that cannot be reconciled: more large lattice studies needed, more efforts to understand discrepancies.
- In the second case there is consensus but large lattice spectral studies like LHC should be performed to see how different it is from 12 flavors.
- One theory (10 flavors) may have a γ that is large enough for WTC

- Hyperscaling $M \sim m^{1/y_M}$ is based on crude approximations such as $\gamma = \gamma_*$
- It would be good to have predictions for corrections to these formulas and to be able to fit these corrections from lattice data: different γ 's for different channels? (Kurachi talk Thurs.?)
- Similarly, in FSS we should attempt fits to the scaling violations.
- Schrödinger functional studies need to be repeated on larger lattices to reduce lattice artifacts.
- Finite temperature transition needs to be pushed harder.